# **Results from the adaptive optics coronagraph at the William Herschel Telescope**

S. J. Thompson, <sup>1</sup><sup>\*</sup> A. P. Doel,<sup>1</sup> R. G. Bingham,<sup>1</sup> A. Charalambous,<sup>1</sup> R. M. Myers,<sup>2</sup> N. Bissonauth, $2$  P. Clark<sup>2</sup> and G. Talbot<sup>3</sup>

<sup>1</sup>*Department of Physics and Astronomy, University College London, Gower Street, London WC1E 6BT* <sup>2</sup>*Centre for Advanced Instrumentation, Durham University, South Road, Durham DH1 3LE* <sup>3</sup>*Isaac Newton Group of Telescopes, Apartado de Correos 321, E-38700, Santa Cruz de la Palma, Tenerife, Spain*

Accepted 2005 September 14. Received 2005 September 1; in original form 2005 June 17

## **ABSTRACT**

Described here is the design and commissioning of a coronagraph facility for the 4.2-m William Herschel Telescope (WHT) and its Nasmyth Adaptive Optics for Multi-purpose Instrumentation (NAOMI). The use of the NAOMI system gives an improved image resolution of∼0.15 arcsec at a wavelength of  $2.2 \mu m$ . This enables the Optimised Stellar Coronagraph for Adaptive optics (OSCA) to suppress stellar light using smaller occulting masks and thus allows regions closer to bright astronomical objects to be imaged. OSCA provides a selection of 10 different occulting masks with sizes of 0.25–2.0 arcsec in diameter, including two with full grey-scale Gaussian profiles. There is also a choice of different sized and shaped Lyot stops (pupil plane masks). Computer simulations of the different coronagraphic options with the NAOMI segmented mirror have relevance for the next generation of highly segmented extremely large telescopes.

**Key words:** instrumentation: adaptive optics – instrumentation: miscellaneous.

#### **1 INTRODUCTION**

Compared to a standard coronagraphic imaging system, one with adaptive optics (AO) gives much higher spatial resolution and a high dynamic range allowing the environments of bright objects to be studied closer in than ever before (Malbet 1996). Additional instrumentation used in combination with a coronagraph allows other new areas of research to be pursued. There are few coronagraphs that have this facility, for example, Coronagraphic Imager with Adaptive Optics (CIAO) on Subaru, which has a choice of linear polarimeters (Murakawa et al. 2003) and Optimised Stellar Coronagraph for Adaptive optics (OSCA) on the William Herschel Telescope (WHT), which has a spectroscopic capability [with the Optically Adaptive System for Imaging Spectroscopy (OASIS) integral-field unit].

The AO system at the WHT, Nasmyth Adaptive Optics for Multipurpose Instrumentation (NAOMI), consists of a single Shack– Hartmann wavefront sensor normally using  $8 \times 8$  sub-apertures and a 76-element segmented mirror. This is a reasonably high-order AO system and can offer partial correction for wavelengths down to 700 nm. NAOMI has been at the telescope since 2000 and during early 2003 was moved to a new environment-controlled laboratory at the opposite Nasmyth platform. This provides a more dust-free and thermally stable environment and so should improve the performance of the system (Myers et al. 2003).

OSCA is a fully deployable instrument which when in use leaves the focus of the NAOMI beam unchanged. This enables OSCA to be used in conjunction with a number of instruments that have already been commissioned at the WHT (Fig. 1). The main imaging camera used with OSCA is the Isaac Newton Group Red Imaging Device (INGRID); a  $1024 \times 1024$ -element HgCdTe-cooled near-infrared (near-IR) detector at the NAOMI focus. The pixel scale when used with NAOMI is ~0.04 arcsec pixel<sup>-1</sup>, hence Nyquist sampling is only obtained down to  $H$  band (1.6  $\mu$ m). Prior to the detector but within the camera cryo-chamber are a set of three wheels containing broad-band filters (*Z*–*K*), narrow-band filters and a choice of pupil stops, respectively. OSCA also has the option of being used in conjunction with an integral-field spectrograph (OASIS) for imaging and spectral analysis at visible wavelengths.

An important criterion in creating a high-contrast imaging system is keeping scattered light to a minimum. Compared to other AO coronagraphs, the system at the WHT allows the insertion of an occulting mask in a focal plane before the AO system as well as those in the focal plane within the AO. Ideally this 'pre-AO' mask would be made of a dichroic material that is transparent to the wavelengths used for wavefront sensing and opaque to the science/observation wavelength. The many square segments in the NAOMI adaptive mirror contribute more scattered light than a similar-sized continuous face-sheet mirror and the gaps between the segments also have a higher emissivity in the IR, so this additional pre-AO stop for the NAOMI–OSCA system could be extremely useful in cutting down scattered light at the science wavelengths and thus improving

<sup>-</sup>E-mail: sjt@star.ucl.ac.uk (SJT)

sensitivity in the final image. An additional means of reducing the diffraction effects of the segmented mirror subsequent to the application of a focal occulting mask is the use of a Lyot (pupil) mask that is matched to the segmented mirror. Such a mask is shown in Fig. 2 (inset top, left-hand mask) and discussed in Section 3.2.

OSCA was commissioned at the WHT in mid-2002 and became available for general use in 2003. The OSCA  $+$  OASIS option is currently unavailable although work is being planned to rectify the problem.

#### **2 DESIGN OVERVIEW**

OSCA (Thompson et al. 2003) is based on the classic Lyot (Lyot 1939) coronagraph design. The sizes of the focal-plane masks (0.25– 2.00 arcsec) were chosen to take advantage of the improved point spread function (PSF) offered by NAOMI. Ten focal-plane masks are available for selection on the instrument. One is for alignment purposes only and two are of a more novel Gaussian profile compared to the standard hard-edged circular discs.

The standard focal-plane (occulting) masks are made of chromium, deposited on the substrates by a contactphotolithography process. Computer simulations, both here (in Section 3) and by others (Nakajima 1994), have shown that occulting masks of a Gaussian profile give better suppression than the standard solid disc design. However, creating such a mask is not trivial. Canyon Materials Ltd. has a patented glass formula, High Energy Beam Sensitive (HEBS) glass, which behaves in a similar way to photographic film. HEBS glass is sensitive towards electron-beam exposure whereby exposure with a certain electron-beam dosage changes the optical density of the material. In this way the HEBS glass can act as a mask material and by varying the electron-beam exposure different grey-scale (optical density) levels can be written into the glass (Walter et al. 1996). To create the Gaussian-profile masks for OSCA (Fig. 3) 669 different grey-scale levels were written in ring steps of 0.5-µm width. One restriction on their usage is that the glass is only suitable for the wavelength range  $0.4-0.8$  µm. There is a visible camera on OASIS, so these masks were included in OSCA with this use in mind.

The pupil-plane masks are manufactured from 0.25-mm thick, hard stainless steel by a process of photochemical machining. Since OSCA is located at a Nasmyth focus of the telescope, the telescope field de-rotator (which ensures the observation object does not rotate as the telescope tracks on its alt-az mounting) has the effect of rotating the telescope pupil. Since the standard Lyot mask used in OSCA includes vane masking, the Lyot mask must be rotated to maintain the alignment with the telescope pupil image. The standard Lyot mask mounted on a Newport SR50 step motor rotation stage is shown in Fig. 2. Additional pupil stops (non-rotating) were manufactured at the Isaac Newton Group (ING) to go in the IN-GRID pupil filter wheel. These provide cooled masks for OSCA which are particularly important for observing in *K*-band otherwise the thermal background noise is very high.

An important requirement of the optical design was to ensure that the focal position was unaltered with OSCA deployed in the beam path and since space is very limited on the optical table, to fit this design into a very small spatial envelope (see Fig. 1). The optical specifications were that it should be diffraction limited at a wavelength of  $2.2 \mu m$  over the full field of view and have less than 0.1-arcsec distortion over the field for the full wavelength range. Additionally, the system has an exit pupil in the same position as that for NAOMI so providing an unchanged optical interface to subsequent optics which were designed to work with NAOMI. The field



**Figure 1.** OSCA installed at the GRACE at the WHT. The light from NAOMI enters from the left through the dichroic and exits to the INGRID detector shown at the right-hand side of the picture. For scale, the hole spacing on the optical table is 25 mm.



**Figure 2.** Standard OSCA Lyot mask mounted on a step-motor rotation stage. Other available Lyot masks are shown inset.

of view for OSCA is designed to be 20 arcsec, but ghosting within the INGRID camera optics has currently reduced this to 15 arcsec. The system also has a 1:1 magnification to leave the NAOMI fieldscale unchanged. To accommodate the required wavelength range of 0.4–2.4 µm all mirrors in the system were coated with protected silver on Zerodur. Similarly, all transmissive optics were made from an IR-grade fused silica (Spectrosil WF) which gives over 90 per cent transmission across the entire specified wavelength range and can be polished to a very high surface quality.

To prevent ghosting within OSCA the substrates for the focalplane masks have a slight wedge angle. The wedges are also tilted slightly in the design, to ensure the ghost-beam misses the pupil. This is achieved by aligning the whole occulting-mask filter-wheel



**Figure 3.** White light images of the OSCA Gaussian-profiled HEBS occulting masks. Speckles are dust in the CCD camera.



**Figure 4.** OSCA ray-trace. The sequence of optical components is numbered and the original NAOMI beam (without OSCA) is overlaid to show the focus at INGRID (circled) unchanged.

assembly at a slight angle to the beam and checking that the ghostbeam is thrown completely out of the optical path.

The path of light through OSCA is shown in Fig. 4 and is numbered sequentially along the beam direction. When deployed in the NAOMI beam, flat mirror (1) picks up the beam. The occulting masks are located at the focal plane (2). The Lyot stop is located at the pupil plane (4). The mirrors at (3) and (5) are a pair of off-axis paraboloids, and finally the flats at (6) and (7) position the focus to the original NAOMI focal position.

The important mechanical requirements of the OSCA system were for it to be thermally stable from −5◦C to 25◦C (i.e. positional accuracy is held within the specified tolerances and that the components will function correctly), vibrationally stable, for none of the components to infringe upon the beam path within or outside of OSCA, the entire system to be deployable in and out of the beam with a repeatability of  $\pm 50 \,\mu$ m, the centre of the focal-plane masks to be automatically positioned to  $\pm 10 \,\mu m$ , 5 yr component lifetime (standard ING requirement) and for the heat output to be minimized so as not to interfere with INGRID.

The unit can be automatically deployed to interrupt the beam passing through to INGRID and the OASIS pick-off mirror. A vertical deployment was implemented due to the space restrictions. From Fig. 1 it can be seen that OSCA comprises a base plate which is rigidly clamped to the bench and a top plate on which all the optomechanical components are seated. Deployment is achieved by the top plate pivoting about a groove and cone of hardened stainless steel at the first off-axis paraboloid end of OSCA with a pneumatic actuator at the second off-axis paraboloid end to bring the plate up into the beam.

The focal masks are located within a wheel that can be automatically selected from the control room. Each focal mask is re-located to high precision by use of notches around the mounting wheel and a pneumatically deployed arm which locks into the notches (detentarm). This facility along with the capability to deploy OSCA into and out of the beam accurately and automatically allows for a flexible and versatile observing programme over the course of a night.

#### **3 SYSTEM SIMULATIONS**

Computer simulations were written and run for a coronagraphic system with a segmented AO corrector (as in NAOMI). A program that generates Kolmolgorov-type turbulence phase-screens (Lane, Glindemann & Dainty 1992) was used as the input to the simulation. This program does not incorporate the outer scale of atmospheric turbulence, although this effect is negligible over a 4-m aperture.

Unless stated otherwise, the simulations discussed here used a phase-screen (aperture) size of  $256 \times 256$ , padded to four times its size with zeroes for the fast Fourier transforms (FFTs), giving a 1024  $\times$  1024 focal plane. In the final image this results in a higher resolution than that achieved by INGRID. The AO and coronagraph code were developed in MATLAB (Thompson 2004) based on an original C-code by A. P. Doel. The simulations include photon noise in the wavefront sensor, mirror hysteresis and pixel noise in the detector.

### **3.1 Effect of focal-mask properties on performance**

The chromium disc focal-plane masks in OSCA were measured to transmit more light at longer wavelengths. To model the effect this might have on the suppression performance, a set of simulations were performed using a 1-arcsec disc mask function with transmission values of 0, 0.1, 1 and 10 per cent. These transmission levels are equivalent to neutral density (ND) values of  $\infty$ , 3, 2 and 1. A 2000 phase-screen AO simulation with the coronagraph using an 80, 20 (80 per cent primary masking, 120 per cent secondary masking) Lyot stop was run for each of these disc transmissions and the results plotted in Fig. 5.



**Figure 5.** Change in coronagraph suppression performance for disc masks 1 arcsec in diameter with different transmission.



**Figure 6.** Simulated PSF profiles for the different shaped coronagraphic focal masks, the standard (80, 20) Lyot stop was used for each. Characteristic width for each mask is 1.0 arcsec. Phase-screens corrected with a NAOMIlike simulated system – 2000, 256  $\times$  256 phase-screens with a four times FFT padding factor. Peak for no-coronagraph lies at  $7 \times 10^8$ .

The suppression factors (no coronagraph/with coronagraph) for these masks measured at a radius of 1 arcsec (i.e. 0.5 arcsec from the edge of the mask), starting with the 0 per cent transmission are: 4.80, 4.80, 4.65 and 3.83, respectively. Therefore, for small amounts of transmission, 1 per cent and less, there is negligible loss of suppression performance compared to a totally opaque mask. For a 10-per cent transmission mask the difference becomes significant, with the factor loss in suppression measured at 1-arcsec radius being 1.25 compared to the opaque mask.

Simulations were also performed for a wide variety of different shaped focal masks, Fig. 6 shows the simulated suppression profiles of those mask types most relevant for OSCA. As expected, the best suppression close in to the mask (0.5–1.0 arcsec) is achieved by using a Gaussian-shaped focal mask. For distances >2 arcsec the disc-mask with a pre-AO stop gives the best suppression. The segmented adaptive mirror contributes a significant amount to the diffracted/scattered light in the image, so reducing the unwanted light falling on the mirror acts to reduce the background. The effect of smoothing the edges of sharp-edged masks and using Gaussianprofiled masks improves the suppression by reducing additional diffracted light in the pupil plane (see the result of diffraction in the pupil plane by a hard-edged mask in Fig. 9b). This then increases the efficiency of the Lyot mask in removing the light in the remaining wings of the PSF. The light from remainder of the PSF, when viewed at the next pupil plane is then mostly concentrated about the edges of this aperture image (primary, secondary, vanes, segment edges, etc.), that is, the parts of the aperture image that contain the most high spatial frequencies.

#### **3.2 Effect of a segmented mirror on performance**

There are a number of additional factors which need to be taken into consideration when designing a coronagraph with an adaptive segmented mirror. The segmentation pattern will cause additional unwanted diffraction effects which distribute light away from the



Figure 7. (a) Simulated WHT aperture function; (b) as in (a) but with an approximate NAOMI segmented mirror pattern applied.

central maximum, this includes the gap between the segments and any phase effects due to step mismatching between the segments. The pupil in OSCA (no focal or Lyot mask) is illustrated in Fig. 7. The normalized peak intensity ratio of the PSF produced by this aperture compared to the same aperture without segmentation (with no phase mismatching) is 0.97. From this simple calculation it is seen that a segmented AO system is of lower contrast by design compared to a continuous face-sheet mirror.

However, the loss in performance when using a coronagraph with a segmented mirror is much more than 3 per cent. The reason for this can be seen by examining the distribution of light in the pupil plane after the application of the focal stop, an example is given in Fig. 9(a). For the case of a Gaussian focal-plane mask applied to the PSF obtained by using the aperture in Fig. 7(b), it is found that 45 per cent of the total light in the pupil is distributed about the segment edges. So if a normal Lyot mask is used here (primary, secondary and vane masking) a significant proportion of the light from that remaining of the masked star will stay in the final image, thus reducing the suppression performance.

The use of a more complex Lyot mask which masks the individual mirror segments as well as the telescope primary and secondary gives improved suppression performance compared to a Lyot mask which only masks the telescope primary and secondary. The individual segments in NAOMI are 7.6 mm across with a ∼0.1-mm gap between each one. The ratio of gap-to-segment size is the same order of magnitude to those proposed for future segmented extremely large telescopes (ELTs), that is, 1-m segments with 10-mm gaps. Hence, these results have relevance for high-contrast imaging with ELTs. To model the effect of the gaps between the NAOMI segments more pixels are required across the aperture. The results shown in Fig. 8 were obtained using 1024 pixels across the aperture diameter with a 2-pixel gap between the mirror segments. Due to the size of this array this was a static simulation (i.e. not an AO simulation), a two times padding factor was used in the FFTs. The lines show the mean trends (a convolution filter has been applied to flatten out the high-frequency periodicity) that the segmented aperture produces. In the high-contrast direction ( $45°$  to image axes) the benefit is most evident at distance >5 arcsec from the centre, reducing counts by two orders of magnitude. For the radial averaged lines (which include the bright axial diffraction peaks) the benefit of the grid Lyot mask is notable from 1 arcsec.

For the full AO simulations (with 256 pixels across the aperture, four times padding and no gaps), the effect of phase mismatching between the segments can be seen as an over-intensity about the segments in the pupil plane subsequent to the application of the occulting mask, as shown in Fig. 9(d). Segments with the greatest intensity are those that are 'turned off' and so have the greatest phase step between them and adjacent segments. Taking gaps and



**Figure 8.** Plot showing the benefit of using a grid-type Lyot mask matched to the NAOMI segment pattern compared to a standard Lyot mask (primary and secondary masking only) – the intensity has been scaled to the same throughput for both the masks. Simulations used a 1024-pixel aperture diameter padded to 2048 for the FFTs, eight segments across the aperture with a 2-pixel gap width. A 2-arcsec solid disc was used as the occulting mask.

phase errors between segments into consideration the benefit of a Lyot mask which masks the individual mirror segments becomes apparent.

A Lyot mask matched to the NAOMI mirror segments (80 per cent undersizing of segments) was created with OSCA [shown in Fig. 2 (inset top, left-hand mask)] but is as yet untested on-sky. The mask requires very careful alignment and there has been insufficient commissioning time to trial this new mask. Telescope schedules allowing, trials may be performed towards the end of 2005.

## **4 LABORATORY TESTING THE FOCAL-PLANE MASKS**

A spatially filtered collimated 613-nm laser beam and a series of lenses and masks were arranged in the laboratory to simulate an ideal coronagraphic system. A simple iris was used for the entrance aperture and another one at 80-per cent diameter to act as the Lyot mask. Images were recorded at the final focus for a variety of different occulting spots using a Santa Barbara Instrument Group (SBIG) camera, this consists of a  $375 \times 242$ -pixel, Peltier-cooled CCD. The usual calibrations were taken – dark frames and background images between every image.

These tests were performed after OSCA had been shipped to the WHT. The Gaussian masks were commissioned at a later date



**Figure 10.** Results of laboratory testing of focal masks in an ideal coronagraphic system.

and a method to compare them to the standard masks on OSCA was devised. The lithography template plate was used in place of a 0.5-arcsec disc occulting mask, approximately the same as the full width at half-maximum of the Gaussian masks. The ND value of this mask was also closely matched to the maximum ND level at the peak of the Gaussian mask so offered a good comparison between the two different shapes of the mask. The 1.0-arcsec mask tested here was a spare from OSCA and had an ND level of ∼5.5 (compared to ∼2.5 for the 0.5-arcsec mask and Gaussian) at this wavelength.

Images were taken at the focus for all three different occulting masks, both with and without the Lyot mask. The images were reduced (dark and background subtracted and scaled for exposure differences) and then radial averages were plotted about the PSF peak. Fig. 10 shows the radial averages of the three different masks using the same Lyot mask. The 1.0-arcsec mask performs best of all, entirely due to its much larger size (covers  $4 \times$  the area thus removing much more of the PSF) and greater opacity. Comparing the other two masks which differ mainly in their shape rather than any other factors it can be seen that as the simulations predicted, the Gaussian-shaped mask provides greater suppression closer in to the centre than the disc mask does.

Fig. 11 shows the significant effect adding a Lyot stop has on the suppression with the Gaussian mask, the CCD count at 10 pixels from the centre is  $5 \times 10^3$ , compared to  $1.2 \times 10^4$  for the 0.50-arcsec disc. Without a Lyot stop the value at 10 pixels from the centre for the Gaussian mask is  $8 \times 10^4$  and for the 0.50-arcsec disc is 5.5  $\times$ 104. This shows that for a Gaussian mask at the focus compared to a solid disc of comparable size, an undersized pupil stop further along the beam works to greater advantage in suppressing the final



**Figure 9.** Pupil-plane images after the application of a 2.0-arcsec hard circular stop. (a) Aperture with segment gap representation of 2 pixels, image has been scaled to minimum and maximum levels; (b) same image as (a) but cut at a low intensity level to bring out faint structures; (c) as in (b) but with the grid-type Lyot mask applied; (d) no gaps, but full AO simulation showing the effect of phase errors between the segments, the secondary vane pattern based on the WHT aperture is included. The outer four and inner four segments are turned off in the AO simulation, the increased intensity about these indicates the greater phase mismatching. The simulation matches the mid-point of the segments in piston and so the corners are also at higher intensity due to phase errors.



**Figure 11.** Radially averaged line profiles taken about the peak count, for different coronagraph configurations with the Gaussian focal mask.



**Figure 12.** Plot of 'Lyot factors' for the Gaussian and 0.50-arcsec mask, where Lyot factor denotes ratio between the suppression curves with and without a Lyot stop. The *x*-axis starts just outside the edge of the 0.50-arcsec disc mask.

image. In terms of attenuation factors, the effect of adding the Lyot stop (measured at 10 pixels) with the Gaussian mask compared to no Lyot stop brings about a factor of 16 drop in measured CCD counts, for the 0.50-arcsec disc this factor is only 4.6. At 60 pixels, for the Gaussian the factor is 1.83 and for the 0.50-arcsec disc it is 1.75, approximately the same. The extra effectiveness of the Lyot stop with the Gaussian mask is therefore greatest close in to the mask, with the 'Lyot factor' dropping as the distance from the mask increases until it converges with the 0.50-arcsec mask values. This trend is shown in Fig. 12.

#### **5 COMMISSIONING RESULTS**

OSCA and its electronics were shipped to La Palma at the beginning of 2002 May and it underwent first commissioning soon after. Only two nights on-sky were allocated for the OSCA commissioning so time was very limited. The first night was plagued by extremely bad seeing throughout (5 arcsec recorded at the worst) and the AO system could not be used. Since the largest occulting mask in OSCA is 2.0 arcsec it was also not feasible to do any performance testing.

The last night (24th May) saw very variable seeing over the course of the night and sky-location and the presence of high cirrus cloud also caused problems on occasion. NAOMI was used although cen-



Figure 13. Cross-sections taken in the high-contrast (45<sup>°</sup>) direction from *H*-band images (a) without OSCA (solid line) and (b) with 2.0-arcsec mask (dashed line). The peak at zero is due to the OSCA mask having a small throughput. The jagged drop to zero  $\left( \langle 25 \rangle \right)$  with no coronagraph is due to detector saturation. The INGRID pixel scale is 0.04 arcsec pixel−1.

tring objects on the OSCA occulting masks was difficult when the seeing was bad and OSCA performance was degraded. The average seeing was 1.5 arcsec, and with the AO system an average corrected PSF width of 0.5 arcsec was obtained. Fig. 13 shows the suppression obtained using OSCA in *H*-band using the 2.0-arcsec mask during these conditions. From the graph it can be seen that just outside the edge of the occulting mask the photon count has been reduced by a factor of 3.5. These values have already been adjusted to account for the loss in throughput due to the Lyot mask.

Attempts were made to observe science targets during the course of the night. The objects were chosen based on their need for coronagraphic observations, that is, features that would not otherwise be easily observable and would demonstrate the benefits of using a coronagraph. They also had to have a *V*-band magnitude of less than 12 (i.e. brighter) due to the sensitivity of the NAOMI wavefront sensor and be observable during the night at an altitude greater than  $z = 40^\circ$ ; below this the turbulence is generally higher due to the high air-mass. Additionally, two possible subtraction stars were found for each target. These were selected to be as close to the target in all respects – sky position, colour/spectral type and *V* magnitude. The attempt was to try and select single stars (no known companions) with a 10-arcsec field about them that is free of any other (particularly bright) stars. Finding good subtraction stars is difficult as there are no comprehensive catalogues for this and information can be incomplete or incorrect in the ones that are available, the ones discussed here were chosen with the aid of the SIMBAD data base.

The reduced data for one of the commissioning targets – HD150451AB – revealed a faint detection (see Fig. 14) of the recently identified cool white dwarf companion HD150451C (Carson 2005). At the time a potential brown dwarf companion was suspected. Since no field rotations were performed to confirm this was not an AO artefact and the signal-to-noise ratio for the dwarf was extremely low, no specific conclusions could be made as of 2002 May. The data collected over 2 years by Carson (2005) have confirmed this to be a companion to HD150451AB, and although initial data suggested the companion to be a methane brown dwarf, recent spectroscopic measurements favour a cool white dwarf classification.

Sky frames were taken for all data (every 10 min in *K* and every 15 min in *J*) to allow more accurate background subtractions. Two different PSF subtraction stars were observed for each target, although with much longer intervals than is ideal due to time



**Figure 14.** Detection of HD150451C (white dwarf) using NAOMI+OSCA. The central star has had a scaled calibration star PSF subtracted from it to increase contrast.



**Figure 15.** A series of same object subtractions over time in *J* band to illustrate the changing PSF due to NAOMI. (a) After 1 min, (b) 5 min, (c) 15 min and (d) 40 min. Atmospheric seeing conditions were average (0.7 arcsec in *V*) and each of the images used in the subtraction had an integration time of 30 s. All images are scaled to same level.

constraints. It became apparent that the PSF stability was poor so that subtractions across time-scales greater than 5 min were contaminated with many AO residuals. Fig. 15 shows this PSF change over time, taken from *J*-band images of HD141569.

Due to the limited time, nothing of interest was uncovered in data taken on MWC297. However, the PSF star chosen for MWC297 revealed a group of four stars in close proximity (see Fig. 16a), with the closest being 1.6 arcsec from the central PSF star. No off-mask images were taken for this object since it was originally only to be used as a PSF subtraction star, so only rough estimates of the *K*-magnitude were possible. The occulted profile was compared to that of an unmasked IR standard star and scaled so that the wings closely matched. This scalefactor (corrected for integration time) was then used directly with the IR standard star magnitude and



**Figure 16.** *Ks*-band image of BD-04 4476 (SAO 142339). (a) There are four stars about the central (occulted) star at distances of 1.59, 2.72, 8.26 and 9.67 arcsec. (b) 2MASS image of the same region. The two stars marked with crosses are the outermost stars in the OSCA image.

applied to the usual magnitude–flux relation to obtain a value of 8.9. The *K*-magnitude estimate for the closest star is  $12.2 \pm 0.3$ – a value for the flux contribution from the central star has been subtracted and was estimated by plotting radial averages about the star. The central star and the outermost two stars in this image are listed in the latest Two Micron All Sky Survey (2MASS) catalogue of point sources. This lists BD-04 4476 as having a *K* magnitude of 8.5 and the outer two stars 12.6 and 11.3, respectively. Taking into account the estimates involved with the focal stop and that 2MASS cannot resolve the inner two stars, our measurements appear to be reasonable.

An upgrade to OSCA was carried out in 2003 April. The Gaussian occulting masks were installed along with a razor-edged antiscatter mask in front of the focal-plane substrates (to reduce scattering from the edges of substrates on which the occulting spots are deposited). The whole of OSCA was also moved and installed at GRound based Adaptive optics Controlled Environment (GRACE) (a new, environment-controlled laboratory on one of the Nasmyth platforms on the WHT). This should alleviate previous problems of dust contamination and temperature problems for NAOMI.

### **6 CONCLUSIONS**

OSCA is a high-precision stellar coronagraph, produced on a low budget, over a short time-scale and meeting all the design specifications. However, due to the limited time assigned for the commissioning and the overlap with NAOMI engineering schedules, a complete and thorough testing of OSCA has not been possible and as a result there have been no performance tests done with OSCA in optimum seeing conditions *and* optimum alignment (both of OSCA and other instrumentation).

The mechanics and electronics of OSCA have operated consistently to date and succeeded in maintaining the required positioning and alignment, including that for the Lyot stop rotation. The deployment mechanism for OSCA has proven to be very successful; it allows OSCA to be raised in to (and out of) the beam very quickly and with consistently accurate positioning, allowing for a flexible and varied observing programme over the course of a night.

Simulations of the NAOMI and OSCA systems also have relevance for the next generation of ELTs where high-contrast imaging is essential in the search for extrasolar planets. When designing a coronagraphic system for a highly segmented aperture suitable masking in the Lyot (pupil) plane must be devised to counteract the strong diffraction pattern that will result otherwise – reducing suppression performance and confusing the data.

A number of interesting astronomy targets have been observed during the OSCA commissioning runs. The most positive results were the very faint detection of the cool white dwarf in HD150451 and the discovery of a potential companion (∼150 au) in BD-04 4476. More data are required to draw any firm conclusions on the nature of these objects.

Information regarding the current status of OSCA and instructions for observing with the system can be found on the ING website.<sup>1</sup>

#### **ACKNOWLEDGMENTS**

In memory of Richard Bingham, who passed away during the publication of this paper. Richard was a great colleague and a brilliant optical designer and will be sadly missed by us all.

<sup>1</sup> http://www.ing.iac.es/Astronomy/instruments/osca

# 1210 *S. J. Thompson et al.*

We would like to thank the ING at the WHT for their support during the commissioning of OSCA. We also thank the astronomers at University College London and elsewhere that gave us help and suggestions for observing and objects of scientific interest. The funding for OSCA was provided by the UK Particle Physics and Astronomy Research Council (PPARC). SJT acknowledges PPARC for providing the PhD studentship, during which the research for this paper was completed.

This research has made use of the SIMBAD data base, operated at CDS, Strasbourg, France

# **REFERENCES**

Carson J. C., 2005, PhD thesis, Cornell Univ., p. 107

Lane R. G., Glindemann A., Dainty J. C., 1992, Waves in Random Media. IOP Publishing, Bristol, p. 209

Lyot M. B., 1939, MNRAS, 99, 580

- Malbet F., 1996, A&AS, 115, 161
- Murakawa K. et al., 2003, SPIE, 4841, 881
- Myers R. M., Longmore A. J., Benn C. R., Buscher D. F., Clark P., Dipper N. A., Doble N., Doel A. P., 2003, SPIE, 4839, 647
- Nakajima T., 1994, ApJ, 425, 348
- Thompson S. J., 2004, PhD thesis, Univ. College London

Thompson S., Doel A. P., Bingham R. G., Charalambous A., Bissonauth N., Clark P., Myers R. M., Talbot G., 2003, SPIE, 4839, 1085

Walter D., Stein R. D., Long P., Wu C., Lee S. H., 1996, SPIE, 2689, 153

This paper has been typeset from a TEX/LATEX file prepared by the author.