

# Interlayer tunnelling in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+d}$ single crystals

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We present measurements of the intrinsic quasi-particle conductivity along the  $c$ -axis of 2212-BSCCO single-crystal mesa structures in the superconducting and normal states. Direct measurement of the mesa temperature enables corrections to be made for self-heating and permits the acquisition of reliable  $I$ - $V$  characteristics over a wide range of temperatures and voltages. Unlike a conventional superconductor, there is no evidence for any change in the quasiparticle conductivity at  $T_c$ , consistent with precursor pairing of electrons in the normal state. At low temperatures the initial low-voltage linear conductivity exhibits a  $T^2$  dependence, approaching a limiting value at zero temperature.

Keywords:  $c$ -axis transport; Intrinsic Josephson effects; Bi-2212; Out-of-plane conductivity

The  $c$ -axis, interlayer quasi-particle DC conductivity has been measured in the normal and superconducting states using small mesas lithographically patterned on the surface of a number of single crystals of 2212-BSCCO with a range of doping. Here we concentrate on the limiting behaviour of the interlayer conductivity at low temperatures and its temperature dependence on passing through  $T_c$ .

Mesa structures typically  $30 \times 30 \mu\text{m}^2 \times 15\text{-}30 \text{ nm}$  were lithographically patterned [1] on as-grown and oxygen annealed 2212-BSCCO single crystals of various doping grown using a thermal gradient method [2].

Below  $T_c$ , the voltage-dependent conductivity was measured in zero magnetic field by biasing junctions in the resistive phase-slip state. A novel three-contact sample geometry was employed, allowing four-probe measurements to be performed whilst simultaneously monitoring the mesa temperature via the temperature-dependent contact resistance (see inset to Fig. 1). The normal state, low-bias ( $\sim 1 \mu\text{A}$ ) conductivity was measured down to  $\sim 50 \text{ K}$  in a field of 6.6 T, overlapping in temperature with zero field measurements in the superconducting state.

Very similar multi-branched  $I$ - $V$  characteristics were observed to those reported elsewhere [3]. The temperature-corrected  $I$ - $V$  curves are well described by a tunnel current given by

$$I = \alpha(V + \beta V^3) \quad (1)$$

(see dashed line in Fig. 1). The intrinsic linear conductance,  $\alpha$ , is inconsistent with models involving incoherent interlayer tunnelling between layers with a  $d(x^2+y^2)$  order parameter symmetry and isotropic scattering around the 2D Fermi surface, as proposed previously [4].

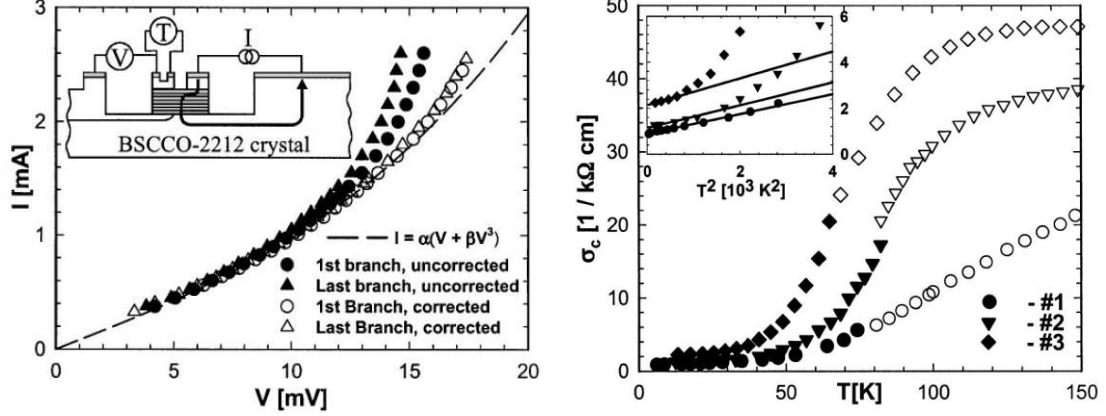


Fig. 1 (left).  $I$ - $V$  characteristics of 1st and 11th (scaled by 1/11) phase-slip branches for sample 2 at 24 K, before and after correction for sample heating. Note that the temperature-corrected data lie on the same line.

Fig. 2 (right). Quasiparticle conductivity plotted for three samples spanning optimal doping ((#1)  $T_c = 75 \pm 3.5$  K, (#2)  $T_c = 87 \pm 0.9$  K, (#3)  $T_c = 86 \pm 0.5$  K). Open symbols: normal state. Solid symbols: from fits to  $I$ - $V$  curves below  $T_c$ . Inset shows low-temperature behaviour.

The derived linear component of the conductivity,  $\sigma_c$ , is plotted as a function of temperature in Fig. 2 in addition to the associated normal state conductivity. At low temperatures, the conductivity is given by

$$\sigma_c(T) = \text{const}(1 + \gamma T^2), \quad (2)$$

consistent with impurity-assisted interlayer hopping (see e.g. Ref. [5]), approaching a doping-dependent limiting value at  $T=0$  which is in good agreement with other measurements [6]. When corrected for heating,  $\sigma_c(T)$  for samples #1 and #2 (underdoped) exhibit a  $T^2$  dependence over a larger temperature range (up to  $\sim 45$  K) than previously reported, whilst sample #3 (overdoped) shows a  $T^2$  variation below  $\sim 20$  K.  $\sigma_c$  is continuous at  $T_c$ , with no evidence for the discontinuity in slope expected from any model involving the onset of pairing at  $T_c$ . Because it is continuous on passing through  $T_c$ , changes in conductivity above  $T_c$  are unlikely to be associated with superconducting fluctuations. For all samples,  $\sigma_c(T)$  decreases monotonically from temperatures well above  $T_c$ , consistent with the onset of a pseudo-gap, possibly due to precursor pairing in the normal state.

Latyshev et al. [6] predict a universal relationship between  $\beta$  (1) and  $\gamma$  (2). Within experimental error, the ratio  $\gamma/\beta$  for all of our samples is consistent with their model, which assumes coherent interlayer transport between low-lying impurity induced states localised in the d-wave gap nodes.

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