

# InAs/InAlGaAs Quantum Dot Lasers on InP and Si

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(Invited Paper)

**Abstract**—We report the development of InAs/InAlGaAs quantum-dot (QD) lasers grown on both InP and Si substrates. A modified indium flush technique was employed to control dot-height distribution and tailor the emission wavelength by using a strained partial capping layer. Using this approach, 7-stack InAs/InAlGaAs QD lasers on InP substrates exhibit a low threshold current density ( $J_{th}$ ) of 63 A/cm<sup>2</sup> per QD layer and high-temperature operation up to 140 °C. Furthermore, electrically pumped InAs/InAlGaAs QD lasers directly grown on Si are also demonstrated, with a low  $J_{th}$  of 1.35 kA/cm<sup>2</sup> and a maximum operating temperature of 100 °C. This work highlights the effectiveness of the modified indium flush in achieving high-performance InAs/InAlGaAs QD lasers. These results represent a significant step forward in the development of high-performance C-/L-band QD lasers in the InAs/InAlGaAs/InP material system for Si photonics.

**Index Terms**—Quantum dots, semiconductor lasers, indium flush, InAs/InAlGaAs, molecular beam epitaxy.

This paper was submitted at xx. This work was supported in part by the UK Engineering and Physical Sciences Research Council under Grants EP/V029606/1, EP/V029681/1, EP/Z532848/1, EP/X015300/1, EP/T028475/1, EP/S024441/1, EP/X035123/1, EP/W021080/1, and EP/P006973/1; European Union's Horizon 2020 program (101129904). (Corresponding author: Mingchu Tang; Huiyun Liu). Hui Jia and Jae-Seong Park contributed equally to this work.

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## I. INTRODUCTION

High-performance semiconductor lasers operating in the C- and L-band wavelengths are of great importance for modern optical communication systems applied in data center and Si photonics [1-3]. Furthermore, amid the rise of artificial intelligence and big data, light sources compatible with large-scale photonic integration are urgently required to tackle the surging bandwidth and energy consumption [4]. In particular, light sources that can be monolithically integrated onto Si platforms are considered key elements for next-generation electronic and photonic integrated circuits for artificial intelligence and big data [5].

Quantum-dot (QD) lasers have emerged as a promising candidate for such applications, offering various advantages over conventional quantum well lasers [6-8]. Their three-dimensional carrier confinement leads to a discrete energy state, resulting in low threshold current density ( $J_{th}$ ), nearly-zero linewidth enhancement factor, ultra-fast gain dynamics, high temperature operation, *etc* [6, 9-14]. Furthermore, the strong carrier localization in QDs provides defect tolerance, making them suitable for monolithic integration on Si substrates, which inevitably suffer from high density of defects during the heteroepitaxy [15-18].

Although InAs/InAlGaAs QD lasers are promising candidates for C- and L-band light sources, their development has been hampered by epitaxial growth challenges. One of the primary obstacles in InAs/InP material system is the anisotropic indium diffusion coefficient along [110] and  $[1\bar{1}0]$  directions on an InP (001) substrate, leading to the formation of elongated quantum dash (Qdash) structures along the  $[1\bar{1}0]$  direction, which disrupts the desired three-dimensional carrier confinement and undermines many benefits of QDs [9, 19-21]. Additionally, the broader dot-size distributions induced by moderate lattice mismatch ( $\sim 3.2\%$ ) between InAs and InP result in significant inhomogeneous spectral broadening, which limits modal gain and hinders the realization of high-performance InAs/InAlGaAs QD lasers [6, 22, 23].

Over the past two decades, a great deal of effort has been made to improve the size/shape uniformity of InAs/InAlGaAs QDs. Various growth strategies have been explored, including the adjustments to growth temperature, V/III ratio, growth rate, growth interruption and use of As<sub>2</sub> instead of As<sub>4</sub> [19, 24-26]. These aim at accelerating the dot formation process and thereby suppressing the anisotropic surface diffusion of indium. Furthermore, modifications to the underlying layer composition, such as InAlAs, InGaAs, AlGaAs, GaAs or

InAlGaAs, have been employed to control the growth front, strain field, and phase separation, thereby optimizing QD nucleation and morphology [19, 27-31].

As an alternative solution, the indium flush technique controls the dot-height uniformity and tunes the emission wavelength [32-37]. This involves partially capping the QDs with a thin first capping layer (FCL), followed by annealing at elevated temperature to evaporate the indium or to induce an As/P exchange process, which reduces the height of large dots. A subsequent second capping layer (SCL) is then grown to complete the structure. Most dots present a similar height as the FCL, leading to a significant blueshift in emission wavelength and a narrower full-width at half-maximum (FWHM) of photoluminescence (PL). For example, Luo *et al.* demonstrated a peak wavelength blueshift from 1690 nm to 1570 nm and a PL FWHM reduction from 124 meV to 87 meV for 5-stacked InAs/InGaAsP QDs, by using an optimal SCL growth temperature that maximizes the indium migration and As/P exchange reaction [34]. Our previous work optimized the FCL thickness and annealing temperature, achieving a narrow PL FWHM of 50.9 meV for a 5-layer InAs/InAlGaAs QD stack [32]. However, the employment of the indium flush technique in the InAs/InAlGaAs/InP QD system remains relatively underexplored, particularly with varied composition of FCL and thickness adjustment for FCL/SCL combinations.

Furthermore, there is a growing demand toward realizing monolithic InAs/InAlGaAs QD lasers on Si as an on-chip light source for Si photonics, because QDs are less sensitive to defects generated from the interface between III-V and Si [15, 38, 39]. While O-band InAs/GaAs QD lasers have shown substantial advances in monolithic integration on Si platforms, the situation for the C-/L-band InAs/InAlGaAs QD lasers is more complex. The large lattice mismatch of  $\sim 8\%$  between InP and Si ( $\sim 4\%$  for GaAs and Si) results in much higher densities of threading dislocation, as well as antiphase boundaries and thermal cracks [40, 41]. These issues make the direct growth of electrically pumped InAs/InAlGaAs QD lasers on Si (001) extremely difficult, and only a few reports have demonstrated such devices to date [42, 43]. For instance, Zhu *et al.* [42] reported pulsed InAs/InAlGaAs QD lasers on Si with  $J_{th}$  of 1.6 kA/cm<sup>2</sup>, total output power of 110 mW, and maximum operating temperature of 80 °C, using InGaAs/InP defect filter layer and V-groove patterned (001) Si.

In this work, we present an epitaxial approach to develop InAs/InAlGaAs QD lasers on both InP and Si substrates. The QD active regions were grown by molecular beam epitaxy (MBE) as its sub-monolayer precision and solid-source evaporation offer superior control over interface abruptness and stoichiometry compared to metal-organic chemical vapor deposition (MOCVD), where gas-phase reactions can introduce variability. We first examine a modified indium flush technique, utilizing a strained InAlGaAs layer as part of the FCL and varying strained/unstrained FCL thickness combination, to improve dot height uniformity and control emission wavelength. This growth strategy is then used to fabricate 7-stack InAs/InAlGaAs QD lasers on (001) InP substrates, demonstrating high-performance lasers with low  $J_{th}$

and high-temperature operation. This technique is further exploited to grow InAs/InAlGaAs QD lasers directly on Si (001) substrates. Lasing operation up to 100 °C is demonstrated for InAs/InAlGaAs QD lasers directly grown on Si with a low  $J_{th}$  of 1.35 kA/cm<sup>2</sup> at room temperature (RT). These results represent significant progress in InAs/InAlGaAs QD laser development and highlight the potential for InP- and Si-based InAs/InAlGaAs QD lasers in optical communication and Si photonics.

## II. EXPERIMENT

All the active regions of the InAs/InAlGaAs QD samples were grown by solid-source MBE system equipped with a valved arsenic cracker source on InP (001) and Si (001) substrates. Prior to growth, the substrates were degassed in the buffer chamber at 400 °C for an hour and then transferred to the growth chamber for a 1-min deoxidation at 500 °C under As<sub>2</sub> overpressure. The detailed growth information regarding InAs/InAlGaAs QDs will be presented in each respective section. Scanning transmission electron microscopy (STEM) was performed to investigate the cross-sectional structural characteristics. In addition, atomic force microscopy (AFM) was employed to examine the surface morphology of QD samples, enabling direct comparison before and after stacking. PL measurements at RT were conducted using a Nanometrics RPM2000 system, incorporating a 635 nm continuous-wave laser at an excitation power density of 430 W/cm<sup>2</sup> and a wavelength-extended InGaAs detector with a cutoff at 2  $\mu$ m.

Fabry-Pérot (FP) lasers were fabricated on both InP and Si substrates. Ridges were formed slightly above the active region by conventional photolithography and wet chemical etch using a mixture of HCl and H<sub>3</sub>PO<sub>4</sub>. For Si-based lasers, an additional wet etch step was employed to form a second mesa structure and expose the n-type InP contact layer for a top-top contact configuration. After the ridge formation, a 400 nm SiO<sub>2</sub> passivation layer was deposited by plasma-enhanced chemical vapor deposition, and contact windows were opened by reactive ion etching. Ti/Au (20/400 nm) was deposited on the exposed ridge using a sputtering system as the p-contact. For InP-based lasers, substrates were thinned to 150  $\mu$ m and subsequently, an n-type Ni/AuGe/Ni/Au (10/100/10/200 nm) contact layer was deposited on the polished surface using a thermal evaporator. For Si-based lasers, the same n-type metallization was first deposited on the exposed n-type InP contact layer, followed by substrate thinning to 150  $\mu$ m. Both were annealed at 380 °C for 1 min to form ohmic contacts. No facet coating was applied.

## III. INDIUM FLUSH FOR INAS/INALGAAS QD GROWTH

### A. Mechanism Revealed by STEM Investigations

While indium flush has been widely applied to InAs/GaAs QD growth to manipulate the dot uniformity and hence enhance gain [44-46], its adoption in InAs/InAlGaAs/InP material system remains less explored, especially in the quaternary alloy barrier matrices. Therefore, we studied the mechanism of indium flush in an InAs/InAlGaAs material system. Two single-layer QD samples were grown to study the effect of

indium flush, of which the growth parameters were discussed in our previous work [32]. Reference sample A (without indium flush) consists of 500 nm InAlAs and 100 nm InAlGaAs grown on an n-type InP (001) substrate. 5.5 monolayers (MLs) of InAs were deposited directly onto InAlGaAs at 485 °C with a V/III ratio of 18, followed by a 10-second growth interruption under As overpressure to promote QD ripening [26]. A 100 nm InAlGaAs capping layer was then deposited at the same temperature. Additional 100 nm InAlAs, 100 nm InAlGaAs, and surface QDs were grown sequentially to complete the reference sample for PL measurements. A high indium growth rate of 0.42 ML/s and As<sub>2</sub> were used for the QD growth to suppress the anisotropic indium surface diffusion. All layers except the InAs QD layer are lattice-matched to InP. Sample B (with indium flush) was grown under identical conditions except that after the QD deposition and growth interruption, a 4-nm InAlGaAs FCL was grown at 485 °C. Then, the substrate temperature was increased to 540 °C and held for 1 min under As<sub>2</sub> flux with a beam equivalent pressure of  $7 \times 10^{-6}$  Torr (the indium flush step), followed by the growth of the remaining 96 nm InAlGaAs SCL at 500 °C [47]. The emission wavelength was blue shifted by over 300 nm from 1824 nm for sample A to 1522 nm for sample B, and the PL FWHM improved significantly from 89.2 to 47.9 meV. High-angle annular-dark-field (HAADF) STEM images of sample A and B are presented in Fig. 1 (a) and (b), respectively. For sample A without indium flush, a large QD size dispersion is observed, with a mean QD height of  $5.46 \pm 0.84$  nm. In contrast, the QDs in sample B with indium flush are clearly truncated, with a reduced average height of  $3.54 \pm 0.48$  nm, as the larger dots are affected by the flush while smaller dots are fully buried. This facilitates the size quantization-assisted blue shift and improves the uniformity of the dot ensemble dramatically.

Apart from the QD height modifications, HAADF images of sample B also show dark-contrasted regions connecting adjacent dots and bright-contrasted regions above large, truncated dots. Electron energy loss spectroscopy (EELS) measurements confirm these to be Al-rich and In-rich regions, respectively, as seen in Figure 1 (c) and (d). Based on these distinctive observations, we propose a growth model, as shown in Fig. 1 (e – j). Initially, QDs form via Stranski-Krastanov growth mode. Then, the 4 nm quaternary InAlGaAs FCL is deposited; as the average dot height is  $5.46 \pm 0.84$  nm, only larger dots ( $> 4$  nm) remain exposed while smaller dots ( $< 4$  nm) are buried. The elastically relaxed InAs island apices of exposed QDs create energetically unfavorable nucleation sites for Ga and Al adatoms, thereby limiting the FCL growth atop the QDs [22, 44, 46–48]. Strain from the InAlGaAs FCL promotes indium migration across the FCL surface, forming a partial capping layer to reduce the surface energy [44]. Upon substrate temperature elevation, indium desorbs from both the partially capped dot layer and the FCL, leaving the top of the FCL Al-rich, thereby accounting for the dark-contrasted regions between adjacent dots in atomic-number-sensitive STEM-HAADF images. During the subsequent SCL growth, indium atoms readily incorporate above the truncated dots because of the smaller lattice mismatch, i.e. less compressive

strain, forming In-rich regions above truncated dots as shown in Fig. 1(b).

These findings offer important insight into the complex dynamics occurring during the indium flush process in quaternary material matrices, thereby contributing to the further optimization and broader applicability of this technique.

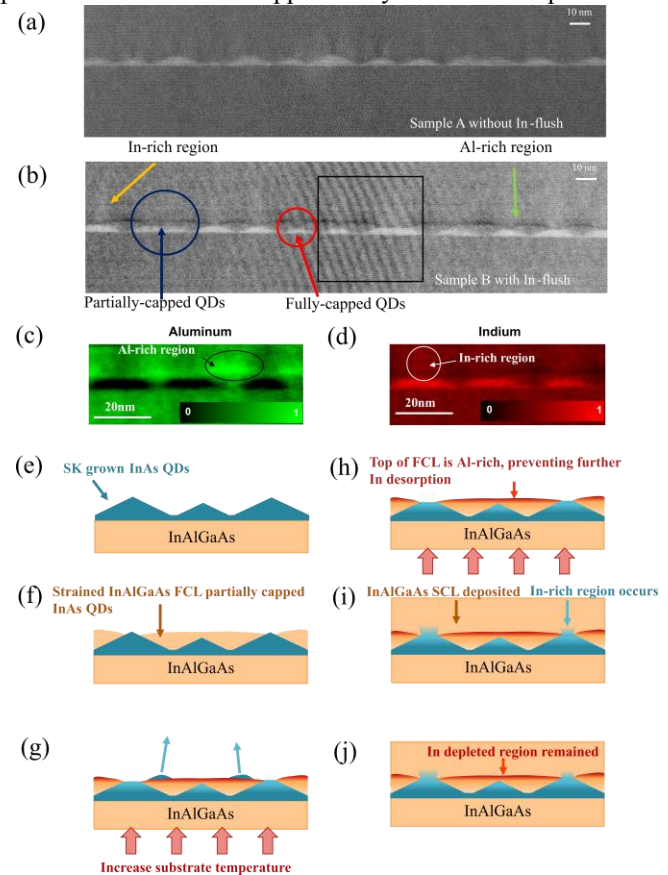


Fig. 1. HAADF images of QD morphologies and adjacent layers of sample A (without indium flush) in (a) and sample B (with indium flush) in (b). (c) and (d) Individual EELS maps of sample B, showing the relative distribution of Al (green intensity channel) and In (red intensity channel). (e – j) The schematic illustration of the proposed impact of the indium flush technique on QDs and adjacent regions. (e) Formation of InAs QDs. (f) The deposition of FCL. (g) The substrate temperature is elevated, and indium atoms desorb from both FCL and QDs. (h) Desorbed indium atoms are flushed away. (i) Indium-rich region and (j) indium-depleted (Al-rich) region formed after the deposition of SCL (Adapted from [47] under CC BY 4.0).

### B. Modified Indium Flush Technique

More flexibility is given to optimizing the growth condition for high density and uniform QDs due to the capability of indium flush to effectively tune the emission wavelength and QD height uniformity. Theoretically, increasing the InAs deposition thickness leads to a shape transition from elongated QD dashes to round-shaped islands, hence increasing QD density [8]. Consequently, to increase QD density, we further increased the InAs thickness based on previous work [32]. High QD densities of  $4.4 \times 10^{10} \text{ cm}^{-2}$  and  $5.2 \times 10^{10} \text{ cm}^{-2}$  were



achieved by an InAs coverage of 6.5 MLs and 6.8 MLs, respectively. However, the emission wavelength also increases with increasing InAs coverage [32]. For the sample with 6.5 ML InAs coverage, even using indium flush technique with an unstrained InAlGaAs FCL of 3.5 nm, the emission wavelength is still as long as 1726 nm.

To further shift the emission wavelength to within the C-/L-band, we developed a modified indium flush technique that utilized a strained InAlGaAs layer as part of the FCL. By changing the thickness of the strained  $\text{In}_{0.359}\text{Al}_{0.323}\text{Ga}_{0.318}\text{As}$  FCL, which has a smaller lattice parameter than InP, significant blue shift in PL peak wavelength can be achieved. Three samples with different FCL were grown to investigate the capability of wavelength tuning. The schematic QD structure is shown in Fig. 2(a). For sample R1, after the 6.5 ML QD growth and interruption for 10 seconds at 485 °C, a 0.5-nm strained  $\text{In}_{0.359}\text{Al}_{0.323}\text{Ga}_{0.318}\text{As}$  layer and a 2-nm  $\text{In}_{0.528}\text{Al}_{0.238}\text{Ga}_{0.234}\text{As}$  layer were subsequently deposited on the QDs as the FCL. Then the thermal annealing was carried out at 515 °C for 3 min and the remaining 97.5-nm  $\text{In}_{0.528}\text{Al}_{0.238}\text{Ga}_{0.234}\text{As}$  SCL growth resumed at 485 °C. Note that the difference in annealing temperature compared with that in Section III-A is due to the variations in the different facility setup. For sample R2, the FCL consisted of 1.5 nm strained  $\text{In}_{0.359}\text{Al}_{0.323}\text{Ga}_{0.318}\text{As}$  and 1 nm  $\text{In}_{0.528}\text{Al}_{0.238}\text{Ga}_{0.234}\text{As}$  and R3 is a reference sample with unstrained FCL. It should be noted that in our experiment, the stressor layer was grown directly on top of QDs rather than in the second part of the FCL because this should have an immediate impact on the strain field within and around the dots, which more readily affects the indium flush procedure and the material exchange at the interface. The PL spectra of the three samples are shown in Fig. 2(b). The inset in Fig. 2(b) shows a  $1 \times 1 \mu\text{m}^2$  AFM image of the uncapped surface dots of the samples with highly uniform and high density QDs. A significant blue shift from 1624.8 nm to 1576.2 nm is obtained just by changing the strained  $\text{In}_{0.359}\text{Al}_{0.323}\text{Ga}_{0.318}\text{As}$  layer thickness by 1 nm from R1 to R2. The effectiveness of using lower composition InAlGaAs FCL in tuning emission wavelength is also reported by Kwoen *et al.* [49], thanks to the band energy modification and reduced indium interdiffusion during subsequent annealing process. In addition, a 27% enhancement in PL intensity is achieved for the sample with 0.5 nm strained  $\text{In}_{0.359}\text{Al}_{0.323}\text{Ga}_{0.318}\text{As}$ . The reduction in strain in the FCL and the spacer layer brought by this strained InAlGaAs is reported to affect the interdiffusion of indium atoms at the InAs/InAlGaAs interface [50], which might account for the observed improvement in PL intensity. Compared to the reference sample, the peak wavelength of the sample with 1.5 nm strained InAlGaAs is blue shifted significantly by 150 nm. The FWHM of the strained InAlGaAs samples shows a slight decrease (58.2 meV for R1 and 58.3 meV for R2) compared to the reference sample (59 meV).

This modified indium flush technique not only effectively shifts the emission wavelength, thus allowing more flexibility for QD growth optimization, but is also especially advantageous for multi-layer QD stack growth due to the strain compensation effect [35, 51]. More QD stacks and thinner

spacer layers can be expected using this technique, which favors the growth of high gain, low defect density QDs. In the following sections, multi-layer InAs/InAlGaAs/InP QD lasers based on this modified indium flush have been developed and demonstrated high performance in terms of high-temperature operation and low  $J_{th}$ .

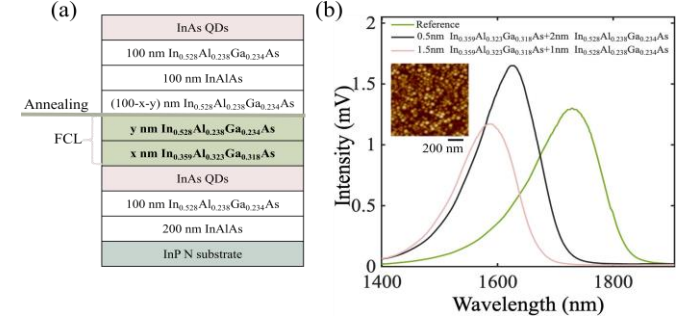


Fig. 2. (a) Schematic structure of single-layer QD sample applying the modified indium flush with different strained FCL. (b) RT PL spectra for single-layer QD samples with different FCL applying modified indium flush. The inset shows  $1 \times 1 \mu\text{m}^2$  AFM of the uncapped surface dots of the samples.

#### IV. INAS/INALGAAS QD LASERS ON (001) INP SUBSTRATE

Compared with InAs/GaAs QD lasers, the performance of InAs/InAlGaAs QD lasers lags, particularly in terms of low  $J_{th}$  and temperature-insensitive operation, mainly due to the severe size/shape inhomogeneity of QDs. To realize high-performance InAs/InAlGaAs QD lasers, achieving uniform QDs with narrow FWHM of PL is critical for enhancing modal gain, low  $J_{th}$ , and high-temperature operation. Here, we employed the modified indium flush technique to effectively control dot height and thus obtain high-density and uniform QDs. This approach resulted in InAs/InAlGaAs QD lasers with low threshold and high-temperature operation.

##### A. 7-stack InAs/InAlGaAs QD laser Growth

The InAs/InAlGaAs QD laser structure was grown on a (001) n-type InP substrate by MBE, and consists of the following layers: Si-doped 200 nm  $\text{In}_{0.524}\text{Al}_{0.476}\text{As}$ , Si-doped 200 nm  $\text{In}_{0.528}\text{Al}_{0.238}\text{Ga}_{0.234}\text{As}$ , a 7-stack InAs QD active region with optimized growth conditions, Be-doped 200 nm  $\text{In}_{0.528}\text{Al}_{0.238}\text{Ga}_{0.234}\text{As}$ , and Be-doped 200 nm  $\text{In}_{0.524}\text{AlAs}$ . A 10 nm Be-doped InGaAs cap was deposited to prevent oxidation during the subsequent transfer to MOCVD, where a Zn-doped 1700 nm InP p-type cladding layer and a Zn-doped 200 nm InGaAs p-contact layer were grown.

Fig. 3 presents HAADF STEM images for the seven-stack QD laser along the (a)  $[110]$  and (b)  $[1\bar{1}0]$  directions to confirm the dot morphology. The InAs QDs exhibit a truncated lens shape with an average size of 25.7 nm along  $[110]$  and 47.3 nm along  $[1\bar{1}0]$  direction. Although minor elongation is observed, the nanostructures are distinct from the Qdash that presents continuous-phase InAs layer along the  $[1\bar{1}0]$  direction resembling a QW [26]. The height along  $[110]$  direction is 3.4 nm, slightly higher than 2.9 nm observed along  $[1\bar{1}0]$ , indicating a minor anisotropy. Importantly, the positions of the QDs in each layer are randomly distributed, rather aligned with

the bottom of QDs in layers above or below, evidencing negligible strain coupling effects.

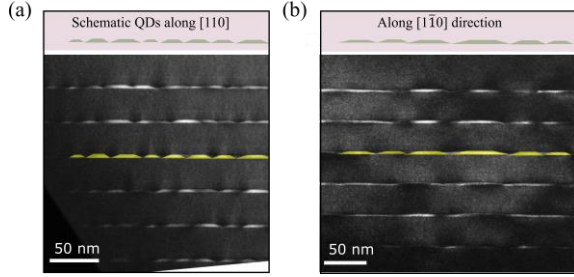


Fig. 3 HAADF images along both (a) [110] and (b)  $[1\bar{1}0]$  directions of the seven-QD-layer laser on InP substrate and their schematic shapes.

### B. Low threshold InAs/InAlGaAs QD lasers

Broad-area FP lasers with a ridge width of 15  $\mu\text{m}$  were first fabricated and tested. The fabricated FP lasers were characterized under pulsed injection (1 % duty cycle, 1  $\mu\text{s}$  pulse width) to minimize the self-heating effect. Fig. 4(a) displays typical power-current ( $LI$ ) curves for varied cavity lengths. The threshold current ( $I_{th}$ ) and  $J_{th}$  for devices with 3000, 2000, 1000, and 500  $\mu\text{m}$  cavity lengths were measured to be 198 mA (440 A/cm<sup>2</sup>), 145 mA (483 A/cm<sup>2</sup>), 92 mA (613 A/cm<sup>2</sup>), and 89 mA (1,189 A/cm<sup>2</sup>), respectively. These correspond to  $J_{th}$  per QD layer of 63, 69, 88, and 170 A/cm<sup>2</sup>. The achieved  $J_{th}$  per QD layer of 63 and 69 A/cm<sup>2</sup> for the 3000 and 2000  $\mu\text{m}$  devices outperformed prior reports for C-/L-band InAs/InAlGaAs QD lasers on (001) InP. Fig. 4(b) presents temperature-dependent  $LI$  characteristics of the 2000  $\mu\text{m}$  device, showing a maximum operating temperature of 130  $^{\circ}\text{C}$ . The inset of Fig. 4(b) depicts  $J_{th}$  versus temperature on a logarithm scale, in which the  $J_{th}$  increases from 483 A/cm<sup>2</sup> at 20  $^{\circ}\text{C}$  to 6,628 A/cm<sup>2</sup> at 130  $^{\circ}\text{C}$ . The characteristic temperature ( $T_0$ ), a measure of temperature sensitivity of  $I_{th}$  or  $J_{th}$  evaluated from  $J_{th}(T) = J_0 \times \exp(T/T_0)$ , was 48.2 K (below 70  $^{\circ}\text{C}$ ) and 44.9 K (above 70  $^{\circ}\text{C}$ ). The RT peak lasing wavelength is 1624 nm and the shift rate is 0.37 nm/K. Note that maximum operating temperatures for the 3000, 1000, and 500  $\mu\text{m}$  devices were 110, 120, and 120  $^{\circ}\text{C}$ , respectively (not shown here).

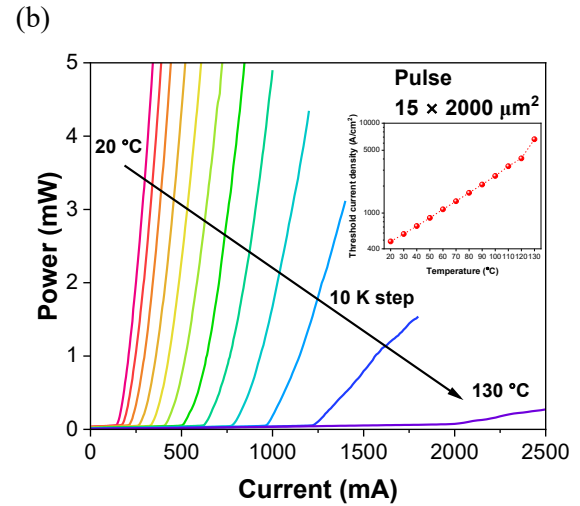
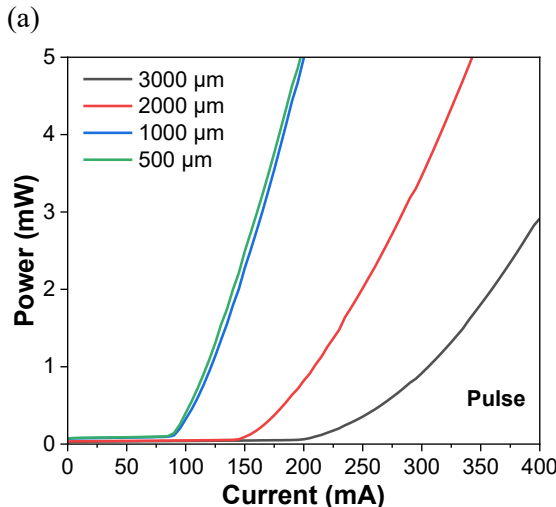


Fig. 4. (a) RT  $LI$  for the InAs/InAlGaAs QD lasers with cavity width of 15  $\mu\text{m}$  and different cavity lengths. (b) Temperature-dependent  $LI$  characteristics for the 15  $\mu\text{m} \times 2000 \mu\text{m}$  device. The inset presents the  $J_{th}$  versus temperature. (Adapted from [52] under CC BY 4.0)

Fig. 5 exhibits continuous-wave (CW) temperature-dependent  $LI$  curves of the 15  $\mu\text{m} \times 2000 \mu\text{m}$  device, confirming a maximum operating temperature of 35  $^{\circ}\text{C}$ . The device yields  $I_{th}$  of 393 mA ( $J_{th}$  of 1.31 kA/cm<sup>2</sup>; 187 A/cm<sup>2</sup> per QD layer), series resistance of 1.6  $\Omega$ , and turn-on voltage of  $\sim 0.7$  V. The reduced  $T_0$  of 32 K was evaluated, highlighting the substantial impact of self-heating under CW operation. As shown in the inset of Fig. 5, the RT peak lasing wavelength at an injection current of  $1.1 \times I_{th}$  and wavelength shift rate are evaluated as 1631 nm and 0.76 nm/K, respectively.

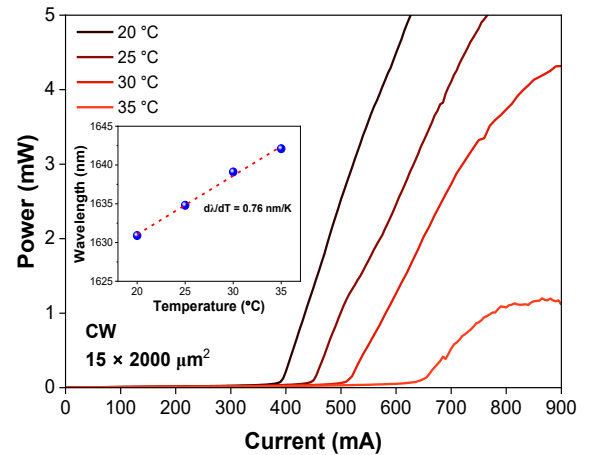


Fig. 5. CW temperature-dependent  $LI$  characteristics of the 15  $\mu\text{m} \times 2000 \mu\text{m}$  device. The inset shows peak lasing wavelength versus temperature.

Low threshold L-band InAs/InAlGaAs QD lasers were demonstrated using the modified indium flush technique. The fabricated seven-stack as-cleaved lasers with a cavity length of 15  $\mu\text{m}$  and cavity lengths of 2000 and 3000  $\mu\text{m}$  under pulsed injection achieved very low  $J_{th}$  per QD layer of 63 and 69 A/cm<sup>2</sup>

and maximum operating temperature of 110 and 130 °C, respectively. While the pulsed performance is promising, the performance under CW—a  $J_{th}$  per QD layer of 187 A/cm<sup>2</sup> and a maximum operating temperature of 35 °C for the 15 × 2000 μm<sup>2</sup> device—is limited, requiring further optimization to enhance CW performance.

### C. High-temperature and ultra-short cavity operation of InAs/InAlGaAs QD lasers

To investigate the impact of a reduced ridge width on the temperature stability, narrow-ridge devices with a cavity width of 5 μm were fabricated. Fig. 6 displays pulsed temperature-dependent  $LI$  curves for the 2000 μm device, yielding  $I_{th}$  of 64.3 mA ( $J_{th}$  of 643 A/cm<sup>2</sup>) and a maximum operating temperature of 140 °C, which is 10 °C higher than the 15 μm × 2000 μm device.  $T_0$  values were 58 K (20 – 70 °C) and 36 K (70 – 140 °C). The enhanced maximum operating temperature for the narrow-ridge device with moderate cavity length of 2000 μm can be attributed primarily to the lower  $I_{th}$  compared to the 15 μm ridge device with the same cavity length—namely, less Joule heating. The inset of Fig. 6 shows a lasing wavelength shift with increasing temperature at an injection current of  $1.1 \times I_{th}$ . The peak lasing wavelength at RT is centered at 1625 nm and redshifts with a shift rate of 0.39 nm/K. Note that the reduced redshift is observed near the maximum operating temperature, which has also been found in literature [12, 53]. This is believed to be due to high current injection-induced variation in effective refractive index [54]. While the redshift driven by bandgap shrinkage dominates over the carrier-induced blueshift as the temperature increases, an excessively increased  $J_{th}$  and reduced quantum efficiency near the maximum operating temperature generate a large amount of unclamped excess carriers, which would mitigate the redshift [55].

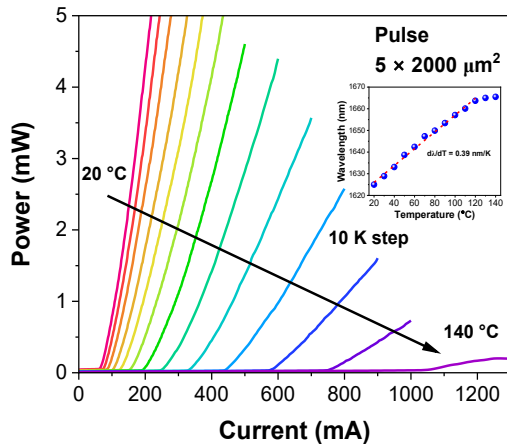


Fig 6. Temperature-dependent  $LI$  curves under pulsed injection for the 5 μm × 2000 μm device. Temperature-induced lasing wavelength shift is shown in the inset.

The 5 μm × 2000 μm device was further characterized under CW operation. Fig. 7 shows CW temperature-dependent  $LI$  curves, with a maximum operating temperature of 60 °C, 25 °C

higher than the 15 μm × 2000 μm device. The  $I_{th}$  ( $J_{th}$ ) at RT was measured as 137 mA (1.37 kA/cm<sup>2</sup>), with a series resistance of 1.6 Ω and turn-on voltage of ~ 0.7 V. The improved  $T_0$  of 40 K was obtained, compared to the 15 μm × 2000 μm device (32.0 K). The inset of Fig. 7 displays a peak lasing wavelength shift at an injection current of  $1.1 \times I_{th}$  as a function of temperature. The RT peak lasing wavelength is 1627 nm, and the shift rate is 0.56 nm/K, lower than 0.76 nm/K (15 μm × 2000 μm).

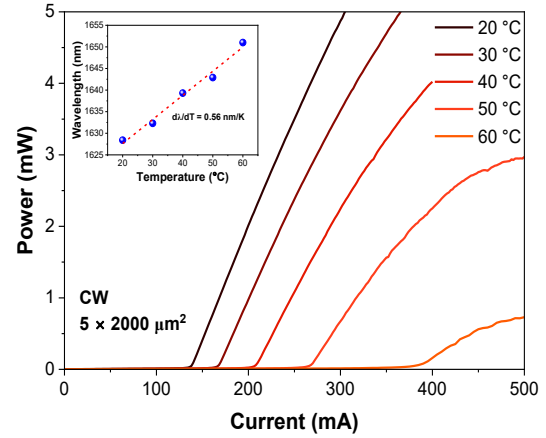


Fig. 7. CW temperature-dependent  $LI$  characteristics for the 5 × 2000 μm<sup>2</sup> device. Inset presents peak lasing wavelength shift as a function of temperature.

Table I summarizes key performance metrics for the devices with different cavity lengths and widths, in terms of  $I_{th}$ ,  $J_{th}$ , single-facet power, maximum operating temperature, and  $T_0$ . For the device with cavity width of 15 μm under pulsed injection, shorter cavity length devices have shown higher maximum operating temperature than longest cavity length (3000 μm) predominantly due to reduction in  $I_{th}$  and Joule heating. However, the 1000 and 500 μm devices degraded by 10 °C, compared with the 2000 μm device. This is because the excessive mirror loss in short cavity length significantly increases the  $J_{th}$ , thereby resulting in increased non-radiative recombination and thermal carrier escape rate at high temperature regime [52]. In other words, excessively high carrier density in short cavity length diminishes the benefit of lower  $I_{th}$  at high temperatures. For the temperature dependence of  $J_{th}$ , shorter cavity devices exhibit higher  $T_0$  at lower temperatures but more pronounced degradation in  $T_0$  at higher temperatures. Similar to the operating temperature, this can be attributed to the higher mirror loss and resultant higher carrier density, further enhancing loss mechanism at high temperature range [52]. Compared with the 15 × 2000 μm<sup>2</sup> device, accordingly, the much lower  $I_{th}$  contributed to the improved thermal stability of the 5 × 2000 μm<sup>2</sup> device.

However, a notable performance degradation was observed under CW operation compared with pulsed mode, primarily attributed to self-heating, further exacerbated by non-optimal shallow, wet-etched waveguide ridge geometry and defects in the regrown p-InP cladding layer. Adopting a deep, dry-etched ridge waveguide can enhance optical mode confinement and reduce  $I_{th}$  [56]. Using monolithic growth to avoid regrowth defects [57] and applying high-reflection facet coatings will

TABLE I

SUMMARY OF DEVICE PERFORMANCE FOR INAS/INAlGaAs QD LASERS WITH VARIED CAVITY LENGTHS AND CAVITY WIDTH OF 15  $\mu\text{M}$  AND 5  $\mu\text{M}$

Current injection	Cavity width	Cavity length	$I_{th}$ (mA)	$J_{th}$ (A/cm <sup>2</sup> )	Single facet power (mW) at 0.5 / 1.0 A	Max. temp. (°C)	$T_0$ (< 70 °C)	$T_0$ (> 70 °C)
Pulse	15 $\mu\text{m}$	3000 $\mu\text{m}$	198	440	5.6 / 30	110 °C	44.3 K	41.3 K
		2000 $\mu\text{m}$	145	483	12 / 48	130 °C	48.2 K	44.9 K
		1000 $\mu\text{m}$	92	613	22 / 69	120 °C	50.9 K	39.0 K
		500 $\mu\text{m}$	89	1189	24 / 60	120 °C	55.3 K	37.2 K
	5 $\mu\text{m}$	2000 $\mu\text{m}$	64	643	17 / 43	140 °C	58.0 K	36.0 K
CW	15 $\mu\text{m}$	2000 $\mu\text{m}$	393	1310	1.5 / 4.7 (at 0.8 A)	35 °C	32 K	-
	5 $\mu\text{m}$		137	1370	8.4 / 9.3 (at 0.65 A)	60 °C	40 K	-

further improve CW performance. In addition, epi-side-down mounting on high-thermal-conductivity submounts can greatly improve heat dissipation [58]. These strategies can narrow the gap between CW and pulsed operation.

In addition to demonstrating high-temperature operation of InAs/InAlGaAs QD lasers, we investigated the lasing characteristics of short cavity-length devices to explore their potential for high-speed applications. InAs/InAlGaAs QD lasers with cavity lengths less than 500  $\mu\text{m}$  were characterized at RT under pulsed injection.

As shown in Fig. 8(a), the  $LI$  curves for devices with cavity lengths of 375, 300, 250, 200, and 180  $\mu\text{m}$  reveals  $I_{th}$  of 46, 52, 61, 82, and 95 mA, corresponding to  $J_{th}$  of 2.4, 3.5, 4.8, 8.2, and 10.6 kA/cm<sup>2</sup>, respectively. Notably, even the ultra-short cavity device with a length of 180  $\mu\text{m}$  demonstrated lasing without any facet coating, highlighting the capability of our QD laser structure to support considerably high gain. To the best of our knowledge, this represents the record-short cavity length achieved for C-/L-band InAs/InAlGaAs QD lasers on (001) InP. Fig. 8(b) presents the optical spectra at an injection current of  $1.1 \times I_{th}$ . The peak lasing wavelength blue-shifted from 1589 nm (L-band) to 1540 nm (C-band). This shift is primarily ascribed to the higher threshold gain requirement induced by higher mirror losses in shorter cavities, leading to lasing at shorter wavelengths within the QD ensemble where the modal gain is higher [59, 60].

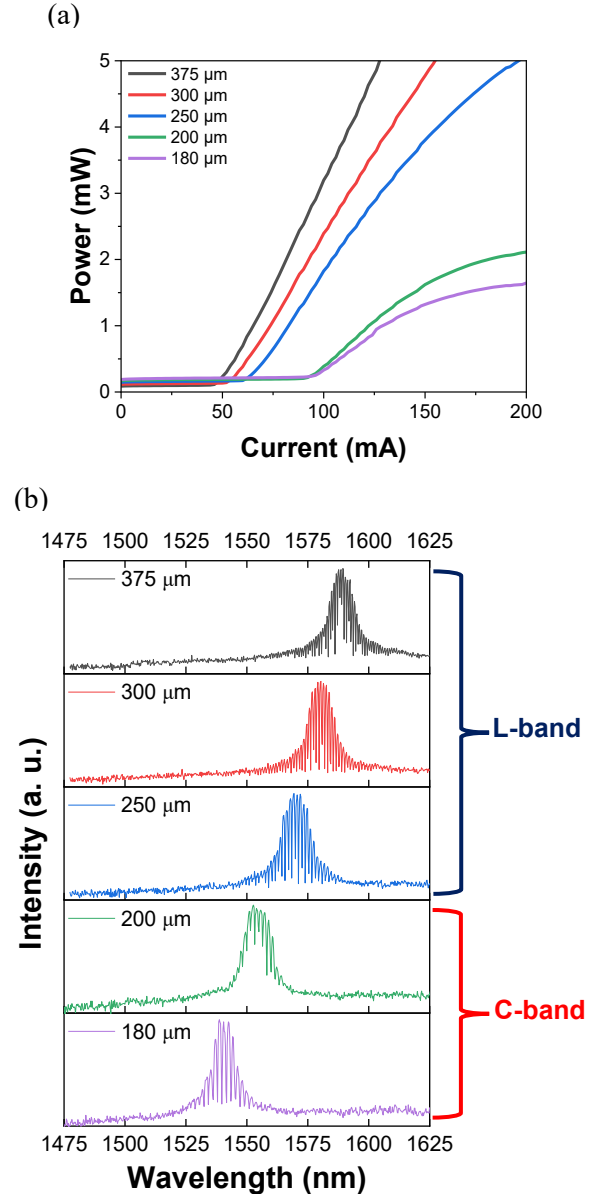


Fig. 8. (a) RT  $LI$  characteristics for devices with cavity length



less than 500  $\mu\text{m}$ . (b) Optical spectra at an injection current of  $1.1 \times I_{th}$  for the short cavity length devices.

Based on the device data with varying cavity lengths and widths, the peak modal gain at threshold was evaluated. The mirror losses for each cavity length were calculated assuming an as-cleaved facet reflectivity of  $R = 0.32$ . Internal loss ( $\alpha_i$ ) and internal quantum efficiency ( $\eta_i$ ) were extracted from the dependence of the external differential quantum efficiency on cavity length, using devices with a 5  $\mu\text{m}$  and varied cavity lengths ranging from 500 to 2000  $\mu\text{m}$ . A linear fit to this data yielded  $\eta_i$  of 20 % and  $\alpha_i$  of  $\sim 12 \text{ cm}^{-1}$ . Note that the extracted low  $\eta_i$  is likely limited by non-radiative recombination at defects introduced during p-InP regrowth, while  $\alpha_i$  is influenced by heterointerface and defect-induced scattering. Future optimization will focus on monolithic MBE growth to reduce these losses. Fig. 9 exhibits the peak modal gain at threshold as a function of current density, calculated from the mirror losses for each cavity length and the evaluated  $\alpha_i$  of  $\sim 12 \text{ cm}^{-1}$ . In addition to the 5  $\mu\text{m}$  cavity width devices, data from devices with 15  $\mu\text{m}$  and 50  $\mu\text{m}$  cavity widths and varying cavity lengths were included in the analysis. A high modal gain of  $76 \text{ cm}^{-1}$  was obtained.

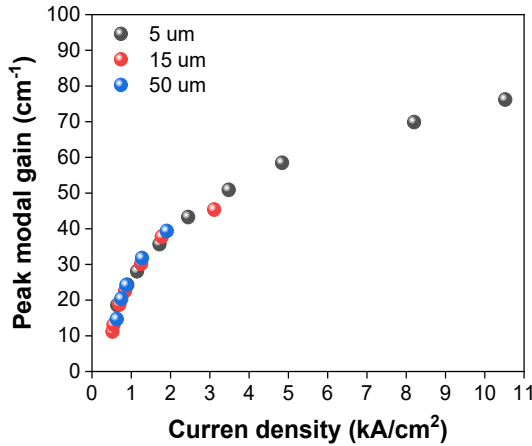


Fig. 9. Peak modal gain versus current density.

## V. INAS/INALGAAS QD LASERS DIRECTLY GROWN ON (001) SI SUBSTRATE

Despite the growing demand for C- and L-band Si-based on-chip light sources, progress in developing InAs/InAlGaAs QD lasers directly grown on (001) Si substrates is stagnant. In contrast to the substantial advances achieved with O-band InAs/GaAs QD lasers on Si, the realization of high-performance InAs/InAlGaAs QD lasers remains more challenging. First, the lattice mismatch between InP and Si ( $\sim 8\%$ ) is larger than GaAs and Si ( $\sim 4\%$ ), generating a much higher density of threading dislocations. Second, the severe size and shape inhomogeneity of InAs QDs grown on InP/Si limits sufficient gain, resulting in higher  $J_{th}$  and limited high-temperature operation. Here, Si substrate pre-treatment [61] and InAsP dislocation filters [62] were employed to suppress the formation of antiphase boundaries and reduce threading dislocation density. On a MOCVD-grown InP/Si

template, 7-stack InAs/InAlGaAs QD lasers were grown using the modified indium flush technique.

### A. 7-stack InAs/InAlGaAs QD lasers growth on Si substrate

Based on the InAs/InAlGaAs QD laser results, an optimized seven-stack InAs/InAlGaAs/InP QD laser structure was further exploited on on-axis Si (001) substrate. First, an InP/GaAs/Si template with a total thickness of 570 nm was grown. The use of a plasma-based surface deoxidization process in conjunction with high-temperature Si annealing result in antiphase boundary-free III-V epitaxial layers [61]. The threading dislocation density of the thin buffer template was measured to be  $1.7 \times 10^9 \text{ cm}^{-2}$ . This dislocation density was further reduced to  $6.6 \times 10^7 \text{ cm}^{-2}$  via two sets of InAsP dislocation filters [62]. The InP/GaAs buffer developed here is considerably thinner than previous work in the literature [42, 63]. Unlike InP-based QD lasers where an n-type InP substrate serves as the n-type contact layer, an undoped (230 nm) and an n-type (770 nm) InP layers were subsequently grown by MOCVD for this role. The InAs/InAlGaAs QDs region was repeated following the previous optimized conditions using the modified indium flush technique in MBE. Finally, the p-type InP cladding and p-InGaAs contact layers were grown by MOCVD to complete the laser structure.

### B. Electrically pumped InAs/InAlGaAs QD lasers on Si

7-stack InAs/InAlGaAs QD FP lasers directly grown on (001) Si substrate with ridge widths of 50  $\mu\text{m}$  and 5  $\mu\text{m}$  were fabricated. The FP lasers were characterized under pulsed injection (1 % duty cycle, 1  $\mu\text{s}$  pulse width) to suppress self-heating effect. Typical  $LI$  characteristics at RT for devices with varied cavity lengths and ridge widths of 50 and 5  $\mu\text{m}$  were measured to evaluate  $J_{th}$  trend (not shown). For the broad-area lasers (50  $\mu\text{m}$  ridge width) with cavity lengths of 2000, 1000, 750, 500, and 350  $\mu\text{m}$ , the measured  $I_{th}$  were 1353, 694, 666, 576, and 524 mA, corresponding to  $J_{th}$  values of 1.35, 1.39, 1.78, 2.30, and 2.79  $\text{kA/cm}^2$ , respectively. The narrow-ridge devices (5  $\mu\text{m}$ ) with cavity lengths of 2000, 1500, 750, 500, and 375  $\mu\text{m}$  exhibit  $I_{th}$  ( $J_{th}$ ) of 323 (3.23), 262 (3.49), 154 (4.10), 119 (4.76), and 104 mA (5.57  $\text{kA/cm}^2$ ), respectively. Fig. 10 plots  $J_{th}$  versus inverse cavity length for both cavity widths. The extracted transparency current density ( $J_{tr}$ ) was calculated to be 0.84 and 2.83  $\text{kA/cm}^2$  for the 50 and 5  $\mu\text{m}$  width devices, respectively.



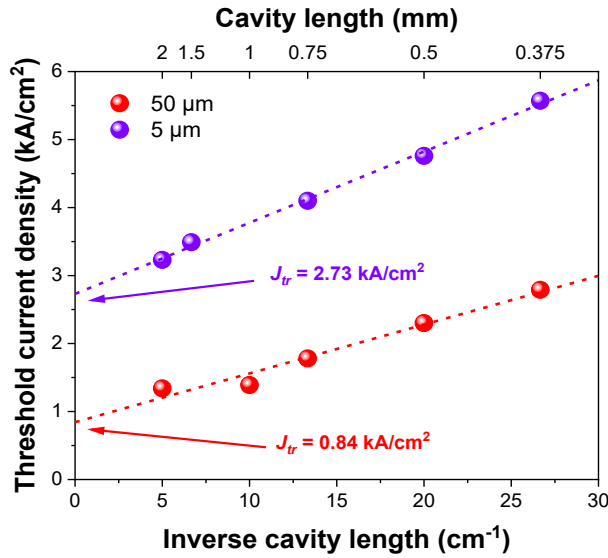


Fig. 10.  $J_{th}$  versus inverse cavity length (cavity length) for the 7-stack InAs/InAlGaAs QD lasers on (001) Si with cavity widths of 50 and 5  $\mu\text{m}$  and varied cavity lengths.

To assess the temperature stability of the InAs/InAlGaAs QD lasers on Si, the 5  $\mu\text{m} \times 2000 \mu\text{m}$  device was selected for further characterization. Fig. 11(a) presents the temperature-dependent  $LI$  curves, showing a maximum operating temperature of 100  $^{\circ}\text{C}$ . Note that the RT single-facet power at an injection current of 1 A was measured as 16 mW without showing roll-over. The logarithmic  $J_{th}$  versus temperature plot in the inset of Fig. 11(a) exhibits a typical linear relation across the whole temperature range. The characteristic temperatures  $T_0$  and  $T_1$ , indicating the temperature dependence of  $J_{th}$  and slope efficiency, respectively, were also evaluated, as shown in Fig. 11 (b). The  $T_0$  was extracted as 62.2 K (20 – 100  $^{\circ}\text{C}$ ), slightly better than InP-based QD lasers with the same cavity size (58 K below 70  $^{\circ}\text{C}$  and 36 K above 70  $^{\circ}\text{C}$ ), due to higher thermal conductivity of Si [64]. On the other hand, the  $T_1$  values were calculated to be 116 K (< 60  $^{\circ}\text{C}$ ) and 44 K (> 60  $^{\circ}\text{C}$ ). Since the  $T_1$  is sensitive to the conversion efficiency of carriers into photons above threshold, a notable degradation of slope efficiency at higher temperature range is mainly due to the increased thermal carrier escape/leakage and enhanced non-radiative recombination [52].

Fig. 11 (c) displays optical spectra at an injection current of  $1.1 \times I_{th}$  as the temperature increases. The ground-state lasing at 1585 nm is observed at RT, and the peak lasing wavelength redshifts to 1620 nm at 100  $^{\circ}\text{C}$ , without switching to the excited-state lasing. The corresponding redshift rate was determined to be 0.45 nm/K, as shown in the inset of Fig. 10 (d).

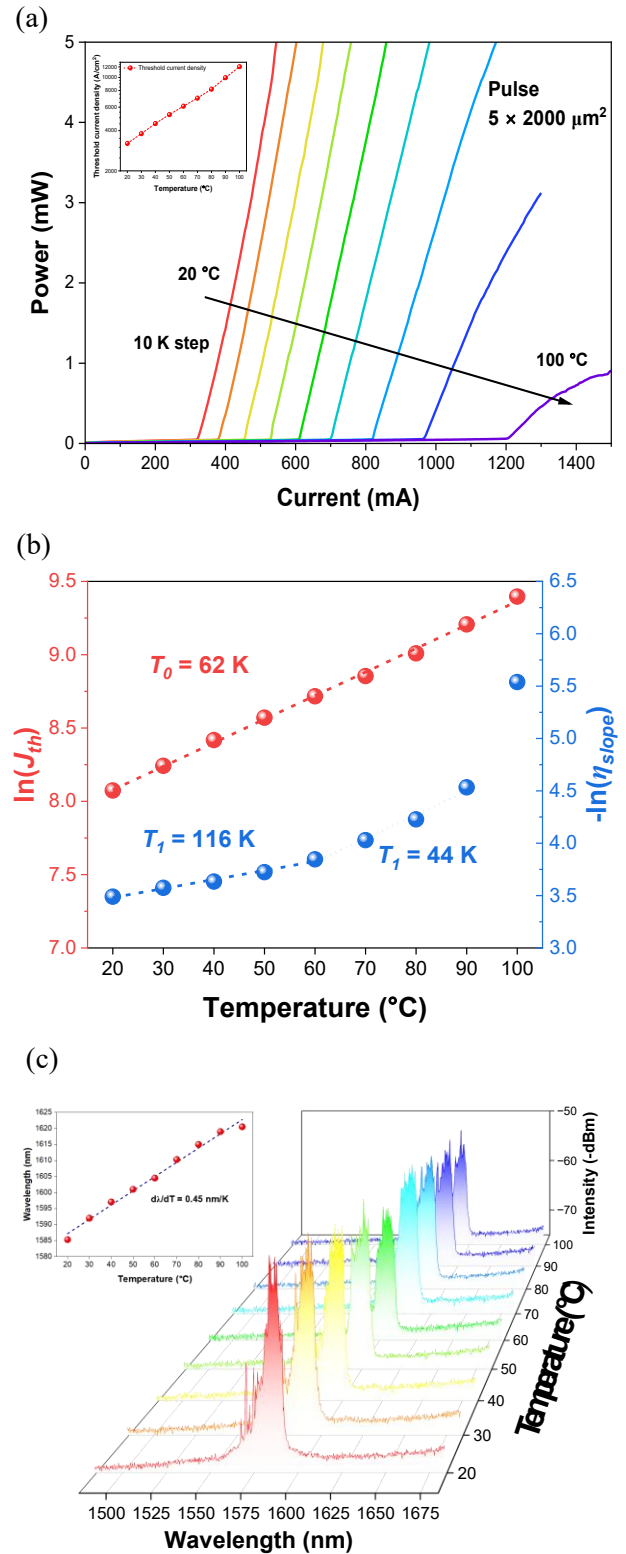


Fig. 11. (a) Temperature-dependent  $LI$  curves for the 5  $\mu\text{m} \times 2000 \mu\text{m}$  device. The temperature-dependent  $J_{th}$  is shown in the inset. (b) Characteristic temperatures  $T_0$  and  $T_1$  (c) Optical spectra at an injection current of  $1.1 \times I_{th}$  in a temperature range of 20 – 100  $^{\circ}\text{C}$ . The inset displays the temperature-induced peak lasing wavelength shift.

## VI. CONCLUSION

This work presents a comprehensive advancement of InAs/InAlGaAs QD lasers on both InP and Si substrates, enabled by the development of a new capping technique to QDs. By introducing a strained partial capping layer during QD growth, we achieved high-density and uniform QDs with reduced height dispersion and emission wavelength tuning, supporting multi-stack InAs/InAlGaAs/InP QD active region with suppressed strain coupling. By employing this approach, 7-stack InAs/InAlGaAs/InP QD lasers on InP exhibit a record-low  $J_{th}$  of 63 A/cm<sup>2</sup> per QD layer ( $15 \times 3000 \mu\text{m}^2$ ) and a maximum operating temperature of 140 °C ( $5 \times 2000 \mu\text{m}^2$ ) under pulsed injection. Under CW operation, the  $5 \times 2000 \mu\text{m}^2$  device achieved a maximum operating temperature of 60 °C, underscoring the need for further thermal and structural optimization. Furthermore, we also achieved electrically injected InAs/InAlGaAs QD lasers directly grown on CMOS-compatible Si (100) substrate, with a low  $J_{th}$  of 1.35 kA/cm<sup>2</sup> ( $50 \times 2000 \mu\text{m}^2$ ) and a maximum operating temperature of 100 °C ( $5 \times 2000 \mu\text{m}^2$ ). These results highlight the potential of the InAs/InAlGaAs QD laser as off-chip or on-chip light sources for optical communication and Si photonics, and mark significant progress toward monolithic integration of long-wavelength QD lasers.

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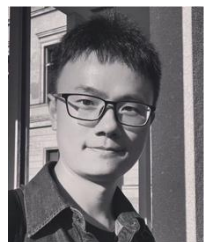


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