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C III] λ 1909 emission as an alternative to Ly α in the reionization era: the dependence of C III] and Ly α at 3 < z < 4 from the VANDELS survey

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ABSTRACT

The velocity offset of Ly α emission from a galaxy's systemic redshift is an excellent tracer of conditions that enable the escape of Ly α photons from the galaxy, and potentially the all-important hydrogen ionizing Lyman continuum photons. However at z \geq 6, Ly α is often heavily attenuated by the neutral intergalactic medium. Here we investigate the utility of C III] $\lambda\lambda$ 1907, 1909 emission, usually the brightest ultraviolet (UV) line after Ly α , as a proxy estimating the Ly α velocity offset ($\Delta v_{\rm Ly}\alpha$). To do so, we use analogues of reionization era galaxies based upon 52 star-forming galaxies with robust C III] detections drawn from the VANDELS survey. Our sample spans a broad UV magnitude range of $-18.5 < M_{\rm UV} < -22.0$, with a sample average value of EW(C III]) = 5.3 Å. We find a slight increase of EW(C III]) with increasing EW(Ly α), but find a large range of EW(C III]) $\sim 1-13$ Å particularly at EW(Ly α) < 10 Å. Using the C III] line peak as the systemic redshift, we calculate $\Delta v_{\rm Ly} \alpha$ and recover the previously reported trend of decreasing $\Delta v_{\text{LV}\alpha}$ with increasing EW(Ly α). Interestingly, we find an anticorrelation between $\Delta v_{\text{LV}\alpha}$ and EW(C III]), which also displays a dependence on the UV absolute magnitude. We derive a multivariate fit to obtain $\Delta v_{\mathrm{Ly}\alpha}$ using both EW(C III]) and M_{UV} , finding that $\Delta v_{\mathrm{Ly}\alpha}$ is more strongly dependent on EW(C III]), with a weaker but nonnegligible dependence on $M_{\rm UV}$. We find that for a fixed EW(C III]), UV-bright Ly α emitting galaxies show smaller values of $\Delta v_{\mathrm{LV}\alpha}$, which suggests that such galaxies may be undergoing more bursty star-formation compared to the UV-fainter ones, akin to a population of extremely UV-bright galaxies identified at z > 10.

Key words: galaxies: evolution – galaxies: high-redshift.

1 INTRODUCTION

A largely uncharted period in the history of the Universe is the Epoch of Reionization (EoR), which began with the emergence of the first stars, black holes, and galaxies after the big bang. It is widely believed that ultraviolet radiation from early sources governed the reionization of the intergalactic medium (IGM), which ended at a redshift $z \sim$ 5.5–6, (Fan et al. 2006; Becker & Bolton 2013; McGreer, Mesinger & D'Odorico 2015; Salvador-Solé et al. 2017; Bosman et al. 2022).

The duration of the EoR can be constrained by, amongst other probes, measures of the scattering and polarization of the cosmic microwave background (Haiman & Knox 1999; Zahn et al. 2012; Bennett et al. 2013; Planck Collaboration X 2016; Choudhury, Mukherjee & Paul 2021), the opacity of hydrogen absorption in the spectra of distant guasars (Erb et al. 2010; Stark et al. 2014; Rigby et al. 2015; Weisz & Boylan-Kolchin 2017; Garaldi, Gnedin & Madau 2019; Bosman et al. 2022), and the visibility or otherwise of Lyman- α (Ly α) emission in star-forming galaxies (SFGs) since this is resonantly scattered by neutral gas (Pentericci et al. 2011; Pritchard & Loeb 2012; Jones et al. 2023; Trapp, Furlanetto & Davies

2023). Constraining the sources responsible for reionization requires knowing their abundances and ionizing capabilities.

With JWST, at least constraining the average production rate of ionizing photons in galaxies at z > 6 is now within reach (Cameron et al. 2023; Curti et al. 2023; Katz et al. 2023; Saxena et al. 2023b; Tang et al. 2023). However, as the neutral fraction of the IGM increases, it becomes impractical to directly observe the escaping ionizing photons from galaxies beyond $z \simeq 4$ (e.g. Inoue et al. 2014). Direct measures of the leakage of Lyman continuum (Ly C) photons is only possible in lower redshift analogues, where such observations can be linked to other observables (e.g. Fletcher et al. 2019; Izotov et al. 2021; Flury et al. 2022; Saxena et al. 2022; Pahl et al. 2023). A further tool in understanding the physics of how Ly C photons may escape is via high-resolution cosmological hydrodynamical simulations (e.g. Maji et al. 2022; Choustikov et al. 2023).

Ly α emission has historically been used to explore how ionizing radiation can escape, since both Ly α and Ly C photons rely on relatively dust-free environments to escape (Verhamme et al. 2017; Izotov et al. 2021; Maji et al. 2022; Naidu et al. 2022; Hayes & Scarlata 2023a). The velocity offset of Ly α with respect to the systemic velocity determined from other nebular lines in particular can trace the geometry of the distribution of neutral gas in an H II region, with Ly α peaks close to the systemic velocity probing relatively low-column densities that facilitate both the escape of Ly α

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as well as Ly C photons from galaxies (Dijkstra, Haiman & Spaans 2006; Dijkstra 2014; Verhamme et al. 2015, 2017).

Fundamentally, however, the increasing neutrality of the IGM renders Ly α emission an unreliable tool in the EoR (e.g. Santos 2004; Dijkstra et al. 2006; Hayes et al. 2010; Behrens & Niemeyer 2013; Inoue et al. 2014) and we must rely on other emission lines to infer the leakage of Ly C photons. Rest-frame ultraviolet spectroscopy has revealed strong emission lines of the semi-forbidden C III] $\lambda\lambda$ 1907, 1909 doublet (referred to as C III] hereafter) over $z\sim2-7$ (Shapley et al. 2003; Erb et al. 2010; Stark et al. 2014, 2015) and various authors have discussed whether studies of C III] emission may be a valuable probe as a substitute for Ly α (Stark et al. 2014, 2015, 2017; Ding et al. 2017). These observations however pre-date *JWST*. With *JWST* results delving deeper into the reionization era with spectroscopy, we are already seeing how *JWST* is transforming identification of these emission lines (e.g Bunker et al. 2023; Saxena et al. 2023a; Tang et al. 2023)

Several studies have previously explored correlations between the equivalent widths (EW) of Ly α and C III] (Shapley et al. 2003; Erb et al. 2010; Stark et al. 2014, 2015; Rigby et al. 2015; Ding et al. 2017; Schaerer et al. 2018; Hutchison et al. 2019; Le Fèvre et al. 2019; Marchi et al. 2019; Cullen et al. 2020; Ravindranath et al. 2020; Llerena et al. 2021) in low-redshift SFGs where Ly α is unattenuated by the neutral IGM. Marchi et al. (2019) found a strong correlation between the two lines, but noted that their sample may be biased as the sample was selected to have both Ly α and C III] emission, questioning whether this relation could change if the entire population of galaxies with CIII] emission were considered. The basic idea behind any possible correlation is that the conditions which support the efficient production and escape of Ly α radiation (i.e. low-dust content, low metallicity, partial coverage of neutral hydrogen) may be similar to those required for the production of collisionally excited emission lines such as CIII]. However, Rigby et al. (2015) pointed out that the correlation is driven primarily by the strongest emitters of Ly α and C III] [EW(Ly α) ~ 50 Å & $EW(C III) \sim 5 \text{ Å}$, whereas for weaker emitters there is no convincing correlation. Moreover, at low metallicities ($Z = 0.1 - 0.2 Z_{\odot}$), the correlation may weaken further (Nakajima et al. 2018b).

Regardless, any correlation between Ly α and C III] may offer the prospect of using C III] emission from galaxies at z>6, now feasible with *JWST*, to constrain the Ly α properties and therefore the nature of the ionizing sources and their environment. Using spectroscopic data from the VANDELS survey for SFGs in the redshift range $z\sim3-4$ where both Ly α and C III] are visible, we examine in further detail the utility of C III] as a proxy for Ly α .

The work explored in this paper complements earlier studies by Marchi et al. (2019) and Llerena et al. (2021) that have used VANDELS data to investigate these emission lines. Llerena et al. (2021) explored the global properties of C III] derived from stacked spectra within the VANDELS data set. The unique aspect of this work is the emphasis on a larger sample of individual emitters compared to Llerena et al. (2021) and, crucially, the offset between C III] and Ly α and its connection with Lyman continuum leakage. Marchi et al. (2019) studied C III] and Ly α including offsets; however, a key difference in this work is that we do not pre-select for Ly α .

The layout of this paper is as follows: In Section 2, we describe the VANDELS data used in this work and the various line measurement methods employed. We present the main results based on this spectroscopic data in Section 3. Finally, we summarize the findings of this paper in Section 4 and place the results into context with those in the literature and discuss the C III] and Ly α relationship along with the implications of the Ly α offsets in the reionization era.

Throughout the paper, we assume the following cosmology: $\Omega_{\rm m} = 0.3$, $\Omega_{\Lambda} = 0.7$, $H_0 = 70\,{\rm km\,s^{-1}\,Mpc^{-1}}$. We use the notation of a positive EW to imply emission, and all logarithms are in base 10 unless otherwise specified.

2 DATA AND MEASUREMENTS

2.1 VANDELS survey

The spectra used in this study all come from the VANDELS survey, which is a Visible Multi-Object Spectrograph (VIMOS) survey of the CANDELS fields using the VIMOS spectrograph on European Southern Observatory (ESO)'s Very Large Telescope (McLure et al. 2018; Pentericci et al. 2018). The survey covers a total area of 0.2 square degrees centred on the CANDELS UKIDSS Ultra Deep Survey (UDS) and Chandra Deep Field South fields (CDFS). It secured high signal-to-noise ratio (SNR) spectra at optical wavelengths (0.48 < λ < 1.0 μ m) for 2100 galaxies in the redshift interval $1.0 \le z \le 7.0$ (Garilli et al. 2021). Around 85 per cent of the galaxies targeted by VANDELS are at $z \sim 3$. The unique aspect of the VANDELS survey is that each source has a significant exposure time, ranging from 20 h to a maximum of 80 h (Pentericci et al. 2018). Full details and description of the survey and target selection can be found in McLure et al. (2018).

2.2 C III] and Ly α emitter selection and line measurement

In this work, we focus on spectroscopically confirmed star-forming galaxies from VANDELS data release 4 (DR4) across both UDS and CDFS fields (Garilli et al. 2021) in the redshift range of $3 \le z \le 4$, only selecting those galaxies that have a redshift probability of being 95 per cent correct (a redshift reliability flag of 3 or 4, see Pentericci et al. 2018 for more details). This redshift range ensures that both the C III] and Ly α emission lines lie within the spectral range. These initial selection criteria resulted in a sample of 773 objects.

To identify C III] and Ly α emission from this parent sample, both the 1D and 2D spectra were visually inspected using the PANDORA software. Based on the robustness of the C III] emission from visual inspection, the sample was split into three different groups based on the 'confidence level' (CL hereafter). CL 3 was assigned to those sources with clear C III] emission in both their 1D and 2D spectra, unaffected by sky features or residual noise. Spectra that showed clear C III] emission in either 1D or 2D spectra (but not both) were assigned a CL of 2, whereas possible line detections likely contaminated by residual noise/sky features were assigned a CL of 1. Out of an initial parent sample of 773 galaxies, C III] emission (CL 1–3) was identified in 280 galaxies, with 139 galaxies assigned the highest confidence level of CL3.

As the main aim of this paper is to study the dependence of C III] emission on Ly α , we then proceeded to identify Ly α emission from all those galaxies that had possible C III] emission. We did not introduce a confidence level criterion for Ly α emission, as the reliability and significance of any Ly α emission was determined using emission line fitting (as discussed later). Across our parent sample, 62 galaxies were found to exhibit both Ly α and C III] (CL3) via visual inspection.

We then measured the strength of the emission lines in these 62 objects. The first step involved determining a 'systemic' redshift using the C III] line, for which we used a single Gaussian fit and the

¹https://www.iasf-milano.inaf.it/software/

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systemic redshift was determined from the peak of this Gaussian (see e.g. Saxena et al. 2020). Using the systemic redshift, each spectrum was de-redshifted and a single Gaussian was fitted to Ly α , C III] and all other rest-ultraviolet (UV) lines (such as He II and C IV) visible in the spectra, with the local continuum value for each line estimated by polynomial fitting to nearby regions free from other emission or absorption features. EWs and integrated fluxes were determined from these Gaussian fits. If the signal-to-noise (S/N) ratio of the continuum was below 2σ , the 2σ noise level was used to determine a lower limit for the line EW.

Thanks to accurate systemic redshifts measured from the C III] line, we could also accurately measure the Ly α velocity offset from the systemic velocity. By comparing the centroid of the observed Ly α emission with that expected for a rest-frame wavelength of 1215.67 Å using the galaxy systemic redshift, the velocity offset of the Ly α peak was determined for all galaxies in our sample. The emission line fluxes, widths, EWs, and the velocity offset of Ly α compared to systemic redshift are listed in Table A1.

2.3 Identification of possible AGN

Prior to further analysis, we considered it important to identify possible active galactic nuclei (AGN) in our sample to remove any biases. As C III] has a high ionization potential (24.4 eV), it is possible that some of the stronger C III] emitters are AGN. To explore this, we searched for another high ionization line, He II λ 1640 (54.4 eV), which is also commonly seen in AGN.

We then attempted to identify possible AGN using rest-UV line ratio diagnostics. Specifically, we used EW(C III]) versus C III] λ 1909/He II λ 1640 proposed by Nakajima et al. (2018b), which is shown in Fig. 1. This resulted in the identification of nine possible AGN out of the 62 CL3 C III] and Ly α emitters, which were removed from our sample. Further, cross-matching with the deep X-ray deep catalogue that exists in CDFS (Luo et al. 2017) revealed one more source identified to be an AGN, which was also removed from our sample. We note that using the C III] and He II based AGN selection, we have reported seven new possible AGN from VANDELS, which were not identified in previous studies (e.g Saxena et al. 2020).

For those galaxies with no clear He II detection, we estimated an upper flux limit by integrating the 1σ continuum-subtracted flux at the expected position (red stars with arrows in Fig. 1). Since these values were consistent for star-forming galaxies, they were retained in the sample.

2.4 Final sample

The final sample of C III] and Ly α emitting SFGs in the redshift range $z\approx 3$ –4 comprises of 52 galaxies. Fig. 2 provides a breakdown of the overall sample and the various selection steps that were implemented to select the sources.

In Fig. 3, we show the parameter space occupied by our C III] emitter sample. In the left panel, we show the distribution of $M_{\rm UV}$ and redshift for our final sample. As is evident from the figure, our sample spans a relatively large range of UV magnitudes (-22 to -18.5), and the flux-limited nature of the survey naturally results in an incompleteness of fainter UV magnitudes at high redshifts. In the middle panel we show EW(C III]) as a function of $M_{\rm UV}$, which shows a slight decrease in EW(C III]) at brighter $M_{\rm UV}$. Finally in the right panel we show the EW(C III]) as a function of redshift, which also highlights the known overdensity in CDFS at $z \sim 3.5$ (Skelton

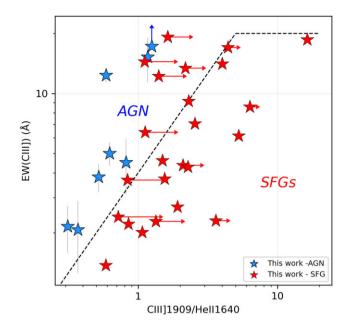


Figure 1. C III]/He II flux ratio versus EW(C III]) for CL3 C III] and Ly α emitting sources in our sample, with the dashed line representing the boundary between SFGs and AGN according to the photoionization models of Nakajima et al. (2018b). Red stars indicate those classified as SFG's according to the boundaries. The red stars with arrows are C III] and Ly α emitters with no He II detection in their emission spectra and illustrate upper limits for those sources, which are consistent with photoionization via star formation. Blue stars are objects with He II detections likely photoionized by AGN. Using this diagnostic (along with the CDFS deep X-ray catalogue, Luo et al. 2017), we identify nine objects from our sample of 139 CL3 C III] and Ly α emitters likely dominated by AGN, and remove them from our sample going forward.

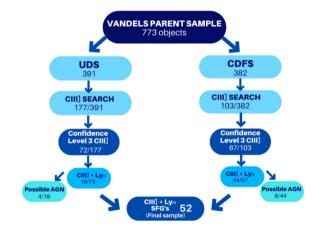


Figure 2. Flowchart showing the various samples and the number of objects retained after each stage of inspection/selection. Our final sample comprises 52 SFGs that are CL3 C III] emitters with Ly α emission and unlikely to host an AGN.

et al. 2014; Straatman et al. 2016; Forrest et al. 2017; Luo et al. 2017; Guaita et al. 2020). Apart from two objects with relatively high EW(C III]) at $z \sim 3.5$, we do not find any strong indication that there is an overall increase in the EW(C III]) in the overdense region compared to the rest of the sample. There is also no appreciable clear trend between EW(C III]) and redshift, although we acknowledge the possibility that C III] detections may be influenced by sample incompleteness.

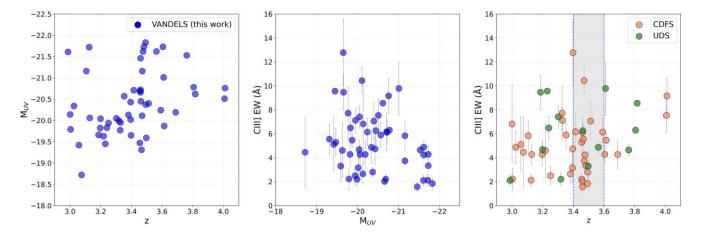


Figure 3. This figure outlines the parameter space which our VANDELS sample occupies. (left): absolute UV magnitudes as a function of redshift for our final sample, demonstrating a wide range of $M_{\rm UV}$ covered by our sample as well as the incompleteness of fainter $M_{\rm UV}$ at high redshift. (middle): rest-frame EW(C III] as a function of $M_{\rm UV}$, showing a mild decrease in the EW(C III] at brighter $M_{\rm UV}$. (right): the distribution of EW(C III] with redshift for sources in the UDS (green) and CDFS (orange) field. The shaded region highlights the known overdensity in the CDFS field at $z \sim 3.5$, where we do find a slight excess of C III] emitters.

Table 1. Average EW(C III]) and the range for C III] emitters that also show Ly α emission in this work and from other studies in the literature at redshifts $z \sim 2$ –5 (Dashes represent no data publicly available).

Data	EW(C III]) (Å)	Range (Å)	N
VANDELS (this work)	5.3	1.6–12.8	52
BRIGHT (this work)	5.0	1.6-9.8	24
FAINT (this work)	5.6	2.2-12.8	28
Nakajima et al. (2018b)	2.0	_	_
Stark et al. (2014)	8.8	1.8-13.5	11
Shapley et al. (2003)	3.6	1.9-5.7	2.0
Rigby et al. (2015)	1.6	0.1 - 4.0	20
Le Fèvre et al. (2019) (stacked)	2.0	_	_
Le Fèvre et al. (2019) (stacked)	2.2	_	_
Llerena et al. (2021) (stacked)	3.9	-	-

3 RESULTS

3.1 Strength of C III] emission

Across our sample of $52 \,\mathrm{C\,III}]$ and Ly α emitting galaxies, we measure an average EW(C III]) of $5.0 \,\text{Å}$. In Fig. 3, we have highlighted the possible UV magnitude incompleteness resulting at higher redshifts because of the flux-limited nature of the sample. To address this, we divided our total sample of 52 galaxies into two sets based on M_{UV} , with the 'Bright' subsample with $M_{\mathrm{UV}} < -20.4$ containing 24 galaxies and the 'Faint' subsample with $M_{\mathrm{UV}} \geq -20.4$ containing 28 galaxies.

The galaxies in the Bright subsample have EW(C III]) in the range 1.6–9.8 Å and an average EW(C III]) of 5 Å. Galaxies in the Faint subsample contains 28 galaxies showing a much larger range of EW(C III]) in the range 2.2–12.8 Å, with an average value of 5.6 Å, which is marginally larger than that of the Bright subsample. These values along with those from other C III]-selected samples in the literature are given in Table 1.

In Table 1, we also compare our sample with similar samples targeting C III] emitters available in the literature (Shapley et al. 2003; Stark et al. 2014; Le Fèvre et al. 2019; Cullen et al. 2020; Llerena et al. 2021). Our VANDELS sample represents a valuable increase in the total number of galaxies with both C III] and Ly α EWs over the full redshift range ($z \sim 0$ –10.6) (as quoted in Ravindranath

et al. 2020) from 59 to 111. Specifically, within the redshift range $z \sim 3-4$, the sample increases from 38 to 80 galaxies with EW(C III])'s ranging from 1.6–12.8 Å. Within our own sample we find that the UV-faint galaxies have a slightly higher average compared to the UV-bright galaxies.

We note that although the Bright subsample is more complete across higher redshifts, the Faint subsample may be more representative of typical star-forming galaxies at z>6 that are now being routinely discovered by various *JWST* surveys (Tacchella et al. 2022; Cameron et al. 2023; Curti et al. 2023; Endsley et al. 2023; Saxena et al. 2023a). Therefore, going forward, we colour code the VANDELS sample using their UV absolute magnitude when studying their Ly α and C III] properties where appropriate.

3.2 Relationship between C III] and Ly α line strengths

In this section, we explore the relationships between C III] and Ly α emission in the VANDELS data from 3 < z < 4 and compare our results to earlier studies at similar redshifts in the literature. We also discuss the dependence of Ly α velocity offsets.

Comparing with our original sample of only C III]-selected galaxies from VANDELS, we find that C III] emitters that show Ly α emission have higher C III] EWs overall compared to only C III] emitters. This can be seen in Fig. 4 where we compare the histogram of the full CL3 C III] emitters in the full VANDELS data set (blue) and C III] emitters with detectable Ly α (green). The median C III] EW of galaxies with both C III] and Ly α is \sim 5 Å(\pm 2.5), compared to a median C III] EW of only 2.9 Å (\pm 4) for those across the full CL.3 sample (see also Cullen et al. 2020). These Ly α emitters having systematically higher EW(C III]) may be explained by the presence of harder ionizing radiation fields in such galaxies, tracing low metallicities and young stellar ages, as was also discussed by Cullen et al. (2020).

To explore any dependence of C III] and Ly α strengths on the absolute UV magnitudes of our final sample, we also investigate the line flux ratios between these two lines in both the Bright and Faint subsamples that we had created above. This reveals an average Ly α /C III] flux ratio of 8.42 for Bright ($M_{\rm UV} < -20.4$) galaxies, and a ratio of 5.99 for Faint ($M_{\rm UV} > -20.4$) galaxies in our final sample. This implies that, while possessing similar Ly α strengths, UV-faint

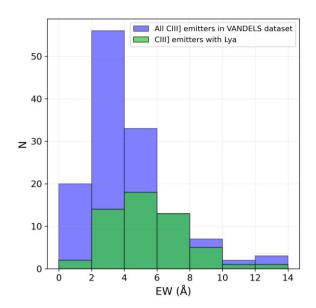


Figure 4. Histogram showing the distribution of EW(C III]) for CL3 C III]-emitters in our analysis. The blue histogram illustrates the full CL3 sample of C III], while the green illustrates the subsample of galaxies that have both C III] and Ly α emission. We find that C III] emitters that also have Ly α emission tend to have higher EW(C III]) when compared to C III] emitters that do not have strong Ly α emission.

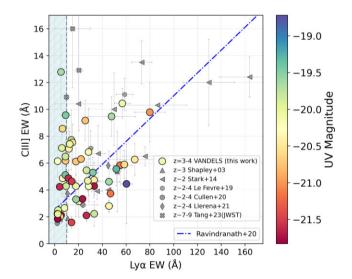


Figure 5. Rest-frame EWs of C III] and Ly α in our VANDELS sample, colour coded by absolute UV magnitude, compared with samples from the literature spanning a redshift range $z \simeq 2$ –4 (Shapley et al. 2003; Stark et al. 2014; Le Fèvre et al. 2019; Cullen et al. 2020; Llerena et al. 2021). We also include in the figure z > 7 measurements from *JWST/NIRSpec* from Tang et al. (2023) ($z \simeq 7$ –9). The blue dashed line shown represents the line of best fit outlined in Ravindranath et al. (2020) while the vertical grey line and shaded region presents a cut off at EW(Ly α) \leq 10 where our sample has a number of strong C III] emitters with weak Ly α emission.

galaxies exhibit a more pronounced C III] line intensity, once again potentially tracing the presence of younger stars and perhaps more 'bursty' star-formation histories.

To further explore correlations between Ly α and C III] strengths in our sample, in Fig. 5 we show the distribution of EW(C III]) and EW(Ly α), colour-coded by $M_{\rm UV}$. For demonstration purposes, we show the line of best fit as obtained by Ravindranath et al. (2020) for

a combined sample from the literature of Ly α -selected star-forming galaxies at $z\simeq 0$ –7 (depicted as the blue dashed line). From our sample alone, we do not find a clear correlation between EW(C III]) and EW(Ly α), which may be expected given how our sample has been selected. Compared to other literature samples, we have selected our galaxies based on C III] detections alone. This means that there may be a number of Ly α -emitting galaxies with no C III] detection, which will be missed by our selection, rendering our sample relatively incomplete for any statistical analysis.

Interestingly however, at EW(Ly α) < 10 Å, we observe a large scatter in EW(C III]) shown via the grey dashed line and shaded region. The EW(C III]) of galaxies in this regime ranges from \sim 1.6–13.0 Å, which represents the widest range of EW(C III]) for any given EW(Ly α) bin. Large EW(C III]) values often indicate conditions with high ionization parameters and low metallicities, whereas Ly α emission is affected by line-of-sight attenuation from regions of high neutral gas or dust. These effects may explain the low observed EW(Ly α). Another possible explanation is offered by the presence of AGN powering the C III] emission in this regime, which could perhaps have been missed by the UV line ratio diagnostics that we employed earlier.

We note that UV luminous galaxies in our sample (red points in Fig. 5) occupy the lower part of the distribution, with low EW(C III]) as well as low EW(Ly α). These luminous galaxies are perhaps tracing more evolved, less star-forming systems in our sample, which would explain the weak EWs.

We further show C III] detections reported from galaxies at z > 7 from Tang et al. (2023) using JWST/NIRSpec spectroscopy. Interestingly, the galaxies from Tang et al. (2023) also show very high EW(C III]) values with low EW(Ly α). At these redshifts, the low EW(Ly α) can be explained by near complete attenuation of Ly α emission by the intervening neutral IGM. This highlights that at the highest redshifts, we would expect to see strong C III] even in the absence of strong Ly α , as is the case for a handful of our VANDELS galaxies at intermediate redshifts.

3.3 Ly α velocity offset as a function of Ly α equivalent width

In this section, we briefly explore the Ly α properties of our galaxies. Across our sample, we measure Ly α velocity offsets ranging from $166-1051\,\mathrm{km\,s^{-1}}$, with an average value of $\Delta v_{\mathrm{Ly}\alpha} \simeq 533\,\mathrm{km\,s^{-1}}$. In Fig. 6, we show the velocity offset from systematic redshift as a function of EW(Ly α), colour-coded by M_{UV} , finding qualitatively an anticorrelation between the Ly α velocity offset and strength. This anticorrelation has also been previously reported in the literature for star-forming galaxies across redshifts (e.g. Erb et al. 2010; Nakajima et al. 2018a; Prieto-Lyon et al. 2023; Roy et al. 2023; Saxena et al. 2023b; Tang et al. 2023).

We also compare our Ly α measurements with those available for a handful of reionization era Ly α -emitting galaxies at z>6 that likely reside in ionized bubbles. Shown in the figure are Ly α velocity offsets measured from Lyman- α Emitters (LAEs) as compiled by Endsley et al. (2023) using ground-based observations with systemic redshifts measured from other rest-UV lines or far-infrared lines from ALMA.Additionally, we show measurements from LAEs at z>6 from the latest *JWST* surveys, which include the faint LAEs presented in Saxena et al. (2023b) using deep spectroscopy from the *JWST* Advanced Deep Extragalactic Survey (JADES), and relatively brighter LAEs selected from CEERS (Jung et al. 2023; Tang et al. 2023). We additionally show the measurement from GN-z11 at z=10.6 from Bunker et al. (2023).

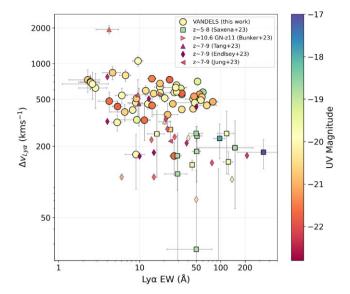


Figure 6. Ly α EW versus Ly α velocity offset for the VANDELS C III] data in the context of *JWST* literature (Bunker et al. 2023; Endsley et al. 2023; Jung et al. 2023; Saxena et al. 2023b; Tang et al. 2023), colour coded using $M_{\rm UV}$. Overall, we find a declining Ly α velocity offset from systemic redshift with increasing Ly α EWs, consistent with earlier work (e.g. Erb et al. 2010; Nakajima et al. 2018a; Prieto-Lyon et al. 2023; Saxena et al. 2023b). We further note that UV-bright galaxies often tend to have weaker Ly α EWs and higher velocity offsets.

Overall, the z>6 LAEs agree with the general trend observed between EW(Ly α) and $\Delta v_{\rm Ly}\alpha$ in our sample, and the scatter we find in our sample is highly consistent with other high-redshift LAE observations (e.g. Erb et al. 2010; Stark et al. 2014; Willott et al. 2015; Nakajima et al. 2018a; Choudhury et al. 2021; Prieto-Lyon et al. 2023; Tang et al. 2023). This demonstrates that the z>6 LAEs must likely reside in patches of relatively ionized IGM, where the Ly α attenuation and profile is largely being determined by scattering through the interstellar medium (ISM).

As has been noted by other studies, the anticorrelation between Ly α velocity offset and EW is likely driven by the density and geometry of neutral gas in the systems, where a higher density leads to increased absorption/scattering of Ly α photons along a line of sight, leading to a reduction in the observed EW and an increase in the velocity offset from systemic, which can be explained using the 'expanding shell model' described in Verhamme, Schaerer & Maselli (2006) and Dijkstra, Gronke & Venkatesan (2016). This explanation is supported by Marchi et al. (2019), who used a similar VANDELS data set and found a correlation between the Ly α velocity offset and the shift in the interstellar absorption lines, where galaxies with high-ISM outflow velocities also showed small Ly α velocity offsets.

Prieto-Lyon et al. (2023) identified a notably reduced mean value of $\Delta v_{\rm Ly\alpha}$ at 205 km s⁻¹ for some of the faintest galaxies (avg. $M_{\rm UV} \sim -17.8$) within a similar redshift range (3 < z < 5). This observation is expected since their sample primarily encompasses galaxies with limited diversity in terms of their stellar masses and consequently lower neutral gas densities, where a higher $\Delta v_{\rm Ly\alpha}$ may be less frequently observed.

Assuming that the observed Ly α properties in our sample of galaxies are regulated by attenuation/scattering by neutral gas in the ISM at redshifts where the neutral IGM is not expected to attenuate Ly α significantly, we can now attempt to understand the physical conditions that regulate both the strength and velocity profiles of

Ly α emission as well as the strength of C III] emission. This can then be used to understand how C III] emission may potentially be used as a diagnostic for Ly α properties in reionization era galaxies that do not reside in ionized bubbles, and therefore do not show Ly α emission

3.4 C III] as a tracer of Ly α velocity offset

Having confirmed that the Ly α properties are indeed sensitive to the neutral gas density that controls the escape of Ly α photons along a line of sight, we now compare the velocity offset of Ly α with the strength of C III] emission to investigate whether the observed EW(C III]) can be used as a diagnostic of neutral gas content, and consequently Ly α and Ly C escape from the ISM of high-redshift galaxies.

In Fig. 7, we show the distribution of Ly α velocity offset as a function of EW(C III]), once again colour-coded by $M_{\rm UV}$ for galaxies in our sample. Qualitatively speaking, there appears to be a tentative anticorrelation between the observed Ly α offset and EW(C III]), with UV-faint galaxies showing less scatter and UV-bright galaxies showing a larger scatter. Therefore, any possible correlation between Ly α velocity offset and EW(C III]) must take into account the UV absolute magnitudes which seem to affect the relation.

In light of this, we derive a correlation between C III] and $\Delta v_{\rm Ly\,\alpha}$, taking into account $M_{\rm UV}$. To do this, we employ a multivariate fit to quantify $\Delta v_{\rm Ly\,\alpha}$ as a function of both EW(C III]) and $M_{\rm UV}$, standardizing the independent variables (x_i) , which takes the form:

$$\frac{\Delta v_{\rm Ly\alpha} - \bar{\Delta v}_{\rm Ly\alpha}}{\sigma_{\Delta v}} = \sum_{i=1}^{2} A_i \frac{x_i - \bar{x_i}}{\sigma_i},\tag{1}$$

where A_i are the coefficients for EW(C III]) and $M_{\rm UV}$, $\bar{x_i}$ and σ_i are the mean and standard deviation, respectively, of the independent variables. Performing a linear regression on this equation gives us the coefficients -0.388 for EW(C III]) and 0.16 for $M_{\rm UV}$. The values of the coefficients as well as the mean and standard deviations of the independent variables are given in Table 2.

The dashed blue line in Fig. 7 shows the relationship that we have derived between Ly α velocity offset and EW(C III]) only, which shows a more promising anticorrelation between the two quantities with a Spearman's rank correlation coefficient = -0.36 and p-value = 0.0008. The shaded region represents the 1σ uncertainty on this relation. The influence of $M_{\rm UV}$ on this relationship is evident in the Figure, with a notable split observed between the majority of UV luminous points falling below the trend line.

Interestingly, the z>7 Ly α and C III] emitting galaxies from Tang et al. (2023) appear to line up well within the 1σ uncertainty on the relation that we have derived, as shown in Fig. 7, suggesting the presence of large ionized bubbles around them. GN-z11 on the other hand seems to lie just outside the 1σ band. This may be consistent with the very low Ly α escape fraction reported in Bunker et al. (2023), which in combination with a relatively large velocity offset of $\sim 500 \, \mathrm{km \, s^{-1}}$ suggests that GN-z11 is likely not situated within an ionized bubble (see also Hayes & Scarlata 2023a), and the emergent Ly α has been heavily attenuated by the neutral IGM surrounding it. Based on the relations that we have derived, the Ly α velocity offset emerging from the ISM will be predicted to be much lower (up to $300 \, \mathrm{km \, s^{-1}}$ lower) based on our multivariate regression.

We find that for a fixed $M_{\rm UV}$, an increase in EW(C III]) results in a decreasing $\Delta v_{\rm Ly\,\alpha}$. However, interestingly, at a fixed EW(C III]) emission, UV-bright galaxies tend to exhibit a lower Ly α velocity offset compared to UV-faint galaxies. Notably, as mentioned previously,

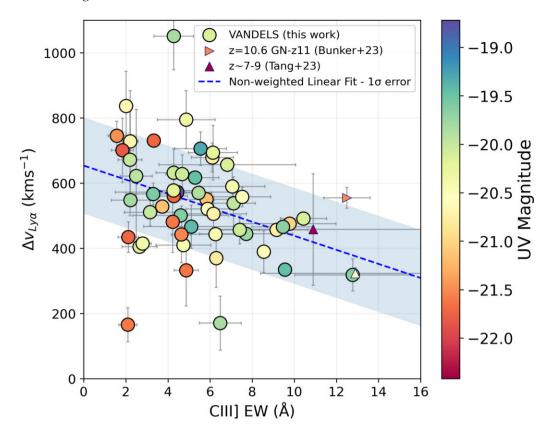


Figure 7. C III] EW versus Ly α velocity offset for the VANDELS C III] data, colour coded using $M_{\rm UV}$. The blue dashed line represents a non-weighted linear regression fit, and the shaded area represents the 1σ error for the VANDELS C III] data. The bright UV galaxies have a larger scatter due to the increased gas compared to the faint UV galaxies.

Table 2. Best-fitting coefficients to predict Ly α velocity offset from the systemic redshift using EW(C III]) and $M_{\rm UV}$.

Variable	A_i	$ar{x_i}$	σ_i	
$\Delta v_{\mathrm{Ly}\alpha}$	_	539.5	158.6	
$\Delta v_{\text{Ly}\alpha}$ EW(C III])	-0.388	5.3	2.5	
$M_{ m UV}$	+0.16	-20.4	0.7	

the UV-bright galaxies tend to lie below the best-fitting line we show in Fig. 7. This behaviour suggests that for the same EW(C Π]), UV-bright galaxies may be in a more 'bursty' phase of star-formation, potentially facilitating a more efficient clearing of channels through which Ly α can escape. It has indeed been shown that galaxies in the early Universe may be expected to follow a more bursty star-formation that boosts their observed UV magnitudes (e.g. Endsley et al. 2023; Simmonds et al. 2024). Stronger outflows driven by a stronger starburst aligns with expectations from Ly α escape in an expanding shell model (e.g. Verhamme et al. 2006; Dijkstra et al. 2016), where neutral gas is expelled, allowing the Ly α emission to escape close its systemic velocity.

Overall, our results indicate that the C III] emission is likely originating from an ionizing radiation field likely driven by young stars. In such a scenario, the UV magnitude serves as a measure of the number of stars undergoing formation. To validate this observation comprehensively, additional UV-bright galaxies with strong C III] and Ly α signals are required. Erb et al. (2014) also noted velocity offset to be correlated with UV luminosity, however found that their faint LAE selected subsample of galaxies had smaller Ly α velocity

offsets and did not see the correlations present in their UV-brighter comparison sample possibly due to lack of dynamical range in mass and luminosity.

In summary, our findings indicate that for a fixed $M_{\rm UV}$, a higher EW(C III]) is associated with a lower velocity offset. Conversely, for a fixed EW(C III]), a brighter $M_{\rm UV}$ results in a lower velocity offset. We propose that this effect may be contributing to the observed patterns. The observed inverse correlation between EW(C III]), Ly α velocity offset and its dependence on galaxy UV magnitude, therefore, promotes the utility of C III] as a proxy for Ly α from galaxies in the epoch of reionization. With increased C III] line detections across samples of z>6 galaxies with NIRSpec, our derived correlations may lead to an additional albeit indirect estimate of the leakage of Ly α and from that potentially the Ly C photons from star-forming galaxies in the early Universe.

4 DISCUSSION AND SUMMARY

By employing alternative indicators in the absence of direct Ly α observations, we can still make significant strides in studying the processes that regulate the escape of Ly α as well as Ly C photons, essential to understand the key drivers of reionization. While Ly α remains a crucial component for understanding reionization, its unavailability limits our capacity to obtain velocity offsets, which are crucial in estimating parameters such as Ly α /Lyman continuum escape fractions and bubble sizes (Jones et al. 2023; Prieto-Lyon et al. 2023; Saxena et al. 2023b; Witstok et al. 2023).

Latest results constraining the Ly α fractions and EW distributions at z > 6 from *JWST* have found the Ly α fractions to be low, which is

expected given the increasing neutral fraction of the IGM (e.g. Jones et al. 2023). This means that at the highest redshifts, unless star-forming galaxies are surrounded by ionized regions, observing Ly α emission will become increasingly challenging. Particularly beyond z > 9.5 as the prominent O and H lines shift beyond the detectable range of the *JWST* NIRSpec, C III] may emerge as a next best feature to both confirm redshifts and deduce additional properties. Therefore, we propose a potential solution to use the relatively bright C III] emission line in rest-UV, coupled with $M_{\rm UV}$, as a tracer for Ly α velocity offsets from systemic through the equation provided in equation (1).

With an increase in the number of detections of rest-UV lines at z>8 (Bunker et al. 2023; Fujimoto et al. 2023; Hsiao et al. 2023; Tang et al. 2023; Topping et al. 2024), including the detection of C III] emission at z=12.5 reported by D'Eugenio et al. (2023), it is becoming essential to maximize the utility of these emission lines in inferring all possible properties of galaxies and their surroundings. Our proposal to employ C III] as an indirect tracer of Ly α velocity offsets (Fig. 7) represents an approach to utilize insights from this strong rest-UV line to better understand the escape of Ly α photons from the ISM of galaxies in the reionization era, possibly evaluating their contribution towards reionizing their local bubble.

In this study, we have used the rest-UV spectra for a sample of star-forming galaxies in the redshift range $\approx 3-4$ from the VANDELS survey with confirmed detections of Ly α and C III] lines, to explore the utility of using the C III] λ 1909 line, often the second brightest emission line after Ly α in the rest-UV, to infer the intrinsic Ly α properties of galaxies in the reionization epoch, where Ly α is significantly attenuated by the neutral IGM. Starting with the detection of C III] emission from 391 objects in the UDS and CDFS fields from the VANDELS survey, we find that 36 per cent of objects show C III] emission, with an average EW \sim 3 Å across the total C III] population. Using the C III] line as a prerequisite, we have then searched for Ly α emission in the sample, recovering a final sample of 52 star-forming galaxies with robust detections of both Ly α and C III] in their spectra.

We find that the equivalent widths of Ly α and C III] are not significantly correlated, finding an increased scatter in EW(C III]) for galaxies with particularly low EW(Ly α) < 10 Å, in agreement with previous studies. We show that galaxies with strong Ly α emission have a higher EW(C III]) on average (\sim 5 Å) compared to those without Ly α , highlighting that Ly α emitting galaxies that also show C III] emission are likely tracing galaxies undergoing rapid star-formation. In this work, we report some of the highest EW(C III]) values (ranging from 1.6–12.8 Å) when compared with those published in the literature at comparable redshifts with a range of $z \sim 2$ –4.

Using the peak of the C III] emission as an indicator of the systemic redshift of galaxies in our sample, we calculate Ly α velocity offsets, $\Delta v_{{\rm Ly}\alpha}$, in the range 166–1051 km s^-1, with an average of 533 km s^-1. We find a weak anticorrelation between EW(Ly α) and $\Delta v_{{\rm Ly}\alpha}$, as has been previously reported in the literature, attributed to the role of scattering of Ly α photons by the neutral gas in the ISM.

Interestingly, we report an anticorrelation between EW(C III]) and $\Delta v_{{\rm Ly}\alpha}$, which appears to also depend on the absolute UV magnitude of the galaxy. This dependence of EW(C III]) on both $\Delta v_{{\rm Ly}\alpha}$ and $M_{{\rm UV}}$ is captured using a multi-variate equation, as shown in equation (1) and Table 2, where we find that $\Delta v_{{\rm Ly}\alpha}$ anticorrelates more strongly with EW(C III]) (coefficient of -0.388), with a comparatively weaker but non-negligible correlation with $M_{{\rm IIV}}$ (coefficient of +0.16).

From our multivariate equation, we find that for a fixed $M_{\rm UV}$, an increase in EW(C III]) leads to a decrease in $\Delta v_{\rm Ly}\alpha$. At a fixed EW(C III]), on the other hand, UV-bright galaxies show lower Ly α

velocity offsets compared to UV-faint galaxies. In our sample, we find that UV-bright galaxies tend to lie below the best-fitting relationship between only EW(CIII) and $\Delta v_{{\rm Ly}\alpha}$, suggesting that UV-bright galaxies with strong Ly α and CIII] emission may be undergoing a more 'bursty' phase of star formation, aiding the efficient clearing of channels through which Ly α can escape, consistent with the 'expanding shell model'. Although this is promising within the context of a population of extremely UV-bright galaxies that are now being discovered $z\gtrsim 10$, to validate this trend, additional spectroscopy for galaxies at z<6 with comparably bright absolute UV magnitudes is essential.

Thanks to JWST spectroscopy, it is now possible to detect C III] emission out to $z\gtrsim 10$, and therefore the relationship between $\Delta v_{\rm Ly\alpha}$ on EW(C III]) and M_{UV} that we have presented in this work may provide insights into the nature of the intrinsic Ly α emission that may be escaping from the ISM of these galaxies, before it encounters a relatively neutral IGM and gets attenuated along the line of sight. Obtaining a better handle on the intrinsic production and escape of Ly α photons from galaxies in the reionization era can provide unprecedented insights into the dominant contributors of ionizing photon towards cosmic reionization.

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DATA AVAILABILITY

The data underlying this paper are publicly available from the VANDELS data release page http://vandels.inaf.it/dr4.html and from the ESO archive http://archive.eso.org/cms.html. The PYTHON code used to analyze the data will be made available upon request.

REFERENCES

Astropy Collaboration, 2013, A&A, 558, A33

Bacon R., Piqueras L., Conseil S., Richard J., Shepherd M., 2016, Astrophysics Source Code Library, record ascl:1611.003

Becker G. D., Bolton J. S., 2013, MNRAS, 436, 1023

Behrens C., Niemeyer J., 2013, A&A, 556, A5

Bennett C. L. et al., 2013, ApJS, 208, 20

Bosman S. E. I. et al., 2022, MNRAS, 514, 55

Bunker A. J. et al., 2023, A&A, 677, A88

Cameron A. J. et al., 2023, A&A, 677, A115

Choudhury T. R., Mukherjee S., Paul S., 2021, MNRAS, 501, L7

Choustikov N. et al., 2023, MNRAS 529 4

Cullen F. et al., 2020, MNRAS, 495, 1501

Curti M. et al., 2023, preprint (arXiv:2304.08516)

D'Eugenio F. et al., 2023, preprint (arXiv:2311.09908)

Dijkstra M., 2014, Publ. Astron. Soc. Aust., 31, e040

1600 M. H. Cunningham et al.

Dijkstra M., Haiman Z., Spaans M., 2006, ApJ, 649, 14 Dijkstra M., Gronke M., Venkatesan A., 2016, ApJ, 828, 71

Ding J. et al., 2017, ApJ, 838, L22

Endsley R., et al., 2023, preprint (arXiv:2306.05295)

Erb D. K., Pettini M., Shapley A. E., Steidel C. C., Law D. R., Reddy N. A., 2010, ApJ, 719, 1168

Erb D. K. et al., 2014, ApJ, 795, 33

Fan X. et al., 2006, AJ, 132, 117

Fletcher T. J., Tang M., Robertson B. E., Nakajima K., Ellis R. S., Stark D. P., Inoue A., 2019, ApJ, 878, 87

Flury S. R. et al., 2022, ApJ, 930, 126

Forrest B. et al., 2017, ApJ, 838, L12

Fujimoto S. et al., 2023, preprint (arXiv:2308.11609)

Garaldi E., Gnedin N. Y., Madau P., 2019, ApJ, 876, 31

Garilli B. et al., 2021, A&A, 647, A150

Guaita L. et al., 2020, A&A, 640, A107

Haiman Z., Knox L., 1999, in de Oliveira-Costa A., Tegmark M., eds, ASP Conf. Ser. Vol. 181, Microwave Foregrounds. Astron. Soc. Pac., San Francisco, p. 227

Hayes M. J., Scarlata C., 2023a, ApJ, 954, L14

Hayes M. et al., 2010, Nature, 464, 562

Hsiao T. Y.-Y. et al., 2023, preprint (arXiv:2305.03042)

Hunter J. D., 2007, Comput. Sci. Eng., 9, 90

Hutchison T. A. et al., 2019, ApJ, 879, 70

Inoue A. K., Shimizu I., Iwata I., Tanaka M., 2014, MNRAS, 442, 1805

Izotov Y. I., Worseck G., Schaerer D., Guseva N. G., Chisholm J., Thuan T. X., Fricke K. J., Verhamme A., 2021, MNRAS, 503, 1734

Jones G. C. et al., 2023, A&A, 683, A238

Jung I. et al., 2023, preprint (arXiv:2304.05385)

Katz H. et al., 2023, MNRAS, 518, 592

Le Fèvre O. et al., 2019, A&A, 625, A51

Llerena M. et al., 2021, A&A, 659, A16

Luo B. et al., 2017, ApJS, 228, 2

Maji M. et al., 2022, A&A, 663, A66

Marchi F. et al., 2019, A&A, 631, A19

McGreer I. D., Mesinger A., D'Odorico V., 2015, MNRAS, 447, 499

McLure R. J. et al., 2018, MNRAS, 479, 25

Naidu R. P. et al., 2022, MNRAS, 510, 4582

Nakajima K., Fletcher T., Ellis R. S., Robertson B. E., Iwata I., 2018a, MNRAS, 477, 2098

Nakajima K. et al., 2018b, A&A, 612, A94

Pahl A. J., Shapley A., Steidel C. C., Reddy N. A., Chen Y., Rudie G. C., Strom A. L., 2023, MNRAS, 521, 3247 Pentericci L. et al., 2011, ApJ, 743, 132

Pentericci L. et al., 2018, A&A, 616, A174

Planck Collaboration X, 2016, A&A, 594, A10

Prieto-Lyon G. et al., 2023, ApJ, 956, 136

Pritchard J. R., Loeb A., 2012, Rep. Prog. Phys., 75, 086901

Ravindranath S., Monroe T., Jaskot A., Ferguson H. C., Tumlinson J., 2020, ApJ, 896, 170

Rigby J. R., Bayliss M. B., Gladders M. D., Sharon K., Wuyts E., Dahle H., Johnson T., Peña-Guerrero M., 2015, ApJ, 814, L6

Roy N. et al., 2023, ApJ, 952, L14

Salvador-Solé E., Manrique A., Guzman R., Rodríguez Espinosa J. M., Gallego J., Herrero A., Mas-Hesse J. M., Marín Franch A., 2017, ApJ, 834, 49

Santos M. R., 2004, MNRAS, 349, 1137

Saxena A. et al., 2020, A&A, 636, A47

Saxena A. et al., 2022, MNRAS, 511, 120

Saxena A. et al., 2023a, A&A, 678, A68

Saxena A. et al., 2023b, preprint (arXiv:2306.04536)

Schaerer D., Izotov Y. I., Nakajima K., Worseck G., Chisholm J., Verhamme A., Thuan T. X., de Barros S., 2018, A&A, 616, L14

Shapley A. E., Steidel C. C., Pettini M., Adelberger K. L., 2003, ApJ, 588, 65

Simmonds C. et al., 2024, MNRAS, 527, 6139

Skelton R. E. et al., 2014, ApJS, 214, 24

Stark D. P. et al., 2014, MNRAS, 445, 3200

Stark D. P. et al., 2015, MNRAS, 450, 1846

Stark D. P. et al., 2017, MNRAS, 464, 469

Straatman C. M. S. et al., 2016, ApJ, 830, 51

Tacchella S. et al., 2022, ApJ, 927, 170

Tang M. et al., 2023, MNRAS, 526, 1657

Taylor M. B., 2005, 347, 29

The Pandas Development Team, 2020, https://doi.org/10.5281/zenodo.35091 34

Topping M. W. et al., 2024, MNRAS, 529, 3301

Trapp A. C., Furlanetto S. R., Davies F. B., 2023, MNRAS, 524, 5891

Verhamme A., Schaerer D., Maselli A., 2006, A&A, 460, 397

Verhamme A., Orlitová I., Schaerer D., Hayes M., 2015, A&A, 578, A7

Verhamme A., Orlitová I., Schaerer D., Izotov Y., Worseck G., Thuan T. X.,

Guseva N., 2017, A&A, 597, A13

Weisz D. R., Boylan-Kolchin M., 2017, MNRAS, 469, L83

Willott C. J., Carilli C. L., Wagg J., Wang R., 2015, ApJ, 807, 180

Witstok J. et al., 2023, A&A, 682, A40,

Zahn O. et al., 2012, ApJ, 756, 65

APPENDIX A: DATA TABLE

Table A1. Line measurements for confidence level 3 (CL3) C III] and Ly α emitters in the VANDELS sample.

Object ID	$z_{\rm spec}$	EW_0	C III] λ1909 FWHM	Flux	EW_0	FWHM	Lyα λ1216 Flux	A 21-	М.
		(Å)	$(km s^{-1})$	$\times 10^{-18}$ (erg s ⁻¹ cm ⁻²)	(Å)	$(km s^{-1})$	$\times 10^{-18}$ (erg s ⁻¹ cm ⁻²)	$\Delta v_{\text{Ly}\alpha}$ (km s^{-1})	$M_{ m UV}$
CDFS				(1.8 2.11)			(8"		
CDFS226606	4.0102	9.1 ± 1.5	678	1.28 ± 0.07	25.7 ± 1.7	894	5.69 ± 0.22	456 ± 17	-20.70
CDFS231741	4.0068	7.5 ± 1.4	671	0.71 ± 0.08	14.9 ± 1.5	1039	1.98 ± 0.13	558 ± 30	-20.5
CDFS226868	3.6884	4.3 ± 1.2	490	0.36 ± 0.09	28.6 ± 1.2	913	3.57 ± 0.05	577 ± 10	-20.1
CDFS000303	3.6137	5.5 ± 0.6	583	0.22 ± 0.02	26.8 ± 3.1	832	1.78 ± 0.01	570 ± 7	-19.8
CDFS017345	3.6052	4.28 ± 0.6	863	1.28 ± 0.2	33.4 ± 1.8	815	22.1 ± 0.85	560 ± 12	-21.7
CDFS233973	3.5913	6.1 ± 1.2	983	0.89 ± 0.12	2.5 ± 0.7	484	0.53 ± 0.15	693 ± 84	-20.2
CDFS015347	3.5124	7.1 ± 1.5	927	0.81 ± 0.1	11.9 ± 1.3	1257	1.85 ± 0.18	590 ± 59	-20.4
CDFS020583	3.4932	2.8 ± 0.4	552	0.30 ± 0.04	45.7 ± 3.5	787	8.49 ± 0.16	414 ± 7	-20.3
CDFS122764	3.492	1.8 ± 0.5	625	0.86 ± 0.05	2.5 ± 0.6	553	1.87 ± 0.43	700 ± 99	-21.8
CDFS128455	3.4813	3.3 ± 0.3	734	1.00 ± 0.06	23.4 ± 1.8	837	17.9 ± 0.88	730 ± 20	-21.7
CDFS004717	3.4756	4.2 ± 0.3	737	1.35 ± 0.07	4.2 ± 1.1	619	2.21 ± 0.56	481 ± 94	-21.6
CDFS140458	3.4708	3.7 ± 0.3	659	0.9 ± 0.03	46.9 ± 3.2	993	16.2 ± 0.33	528 ± 10	-21.1
CDFS009692	3.4698	10.4 ± 1.1	885	0.85 ± 0.02	56.8 ± 6.6	890	9.27 ± 0.19	491 ± 10	-20.1
CDFS012637	3.4655	5.6 ± 1.2	611	0.22 ± 0.03	51.8 ± 4.7	1609	3.10 ± 0.22	706 ± 52	-19.3
CDFS105360	3.4598	1.6 ± 0.4	416	0.54 ± 0.14	14.2 ± 1.7	756	8.79 ± 0.74	745 ± 44	-21.4
CDFS006417	3.4597	2.0 ± 0.2	706	0.33 ± 0.04	4.6 ± 1.2	667	1.03 ± 0.25	837 ± 106	-20.6
CDFS012448	3.4596	6.1 ± 1.0	464	1.04 ± 0.09	17.8 ± 2.1	1293	4.88 ± 0.30	679 ± 44	-20.7
CDFS103010	3.455	2.2 ± 0.4	468	0.40 ± 0.05	2.3 ± 0.9	524	0.49 ± 0.18	728 ± 155	-20.6
CDFS208906	3.4548	5.3 ± 1.1	708	0.54 ± 0.05	32.1 ± 4.1	835	5.68 ± 0.12	617 ± 7	-19.4
CDFS246390	3.4149	6.2 ± 1.6	596	0.97 ± 0.07 0.97 ± 0.07	9.1 ± 1.1	884	2.64 ± 0.12	506 ± 42	-20.7
CDFS240390 CDFS225147	3.3976		564			597	0.58 ± 0.06		-20.7
CDFS223147 CDFS212043	3.3976	12.8 ± 2.8	708	0.61 ± 0.05 0.79 ± 0.11	5.3 ± 0.7	885		318 ± 49 410 ± 7	-19.0 -20.4
		4.7 ± 1.3			33.1 ± 2.8		9.06 ± 0.15		
CDFS215423	3.395	3.2 ± 1.2	313	0.51 ± 0.13	20.4 ± 2.8	1181	3.63 ± 0.26	511 ± 42	-20.0
CDFS206968	3.3808	2.7 ± 0.8	369	0.24 ± 0.07	8.2 ± 0.5	813	1.45 ± 0.01	405 ± 15	-20.1
CDFS024919	3.3528	5.9 ± 1.0	486	0.85 ± 0.09	58.6 ± 3.3	857	15.9 ± 0.22	520 ± 5	-20.5
CDFS229681	3.3291	7.1 ± 1.0	596	0.78 ± 0.04	6.0 ± 1.0	907	1.15 ± 0.16	538 ± 81	-19.9
CDFS217955	3.327	7.7 ± 2.2	752	0.72 ± 0.11	14.0 ± 1.3	908	2.04 ± 0.07	444 ± 15	-19.7
CDFS209992	3.251	2.5 ± 0.8	426	0.27 ± 0.04	2.9 ± 1.1	612	0.41 ± 0.14	622 ± 205	-19.9
CDFS221849	3.2177	4.6 ± 2.3	630	0.42 ± 0.03	45.6 ± 9.0	965	5.69 ± 0.11	501 ± 10	-19.6
CDFS000598	3.1965	4.3 ± 0.9	576	0.44 ± 0.04	9.6 ± 1.4	1002	0.89 ± 0.10	1051 ± 104	-19.8
CDFS007847	3.1284	4.3 ± 0.8	480	0.58 ± 0.05	17.8 ± 2.5	1237	2.72 ± 0.32	632 ± 74	-20.0
CDFS113988	3.1246	2.1 ± 0.5	628	1.25 ± 0.28	5.3 ± 0.79	783	6.12 ± 0.70	434 ± 47	-21.7
CDFS023527	3.106	5.8 ± 1.3	682	1.95 ± 0.08	54.9 ± 5.2	1124	30.5 ± 1.23	550 ± 20	-21.1
CDFS032490	3.0747	4.4 ± 3.0	715	0.14 ± 0.01	60.4 ± 4.9	947	2.92 ± 0.09	573 ± 15	-18.7
CDFS206035	3.0578	5.1 ± 2.3	809	0.66 ± 0.10	9.0 ± 1.8	664	1.19 ± 0.22	466 ± 67	-19.4
CDFS003496	3.0253	4.9 ± 1.4	713	0.90 ± 0.10	6.7 ± 1.1	842	1.7 ± 0.26	795 ± 89	-20.3
CDFS015428	3.0037	2.2 ± 0.9	878	0.29 ± 0.03	33.2 ± 3.2	797	5.24 ± 0.13	548 ± 10	-19.7
CDFS022563	3.0004	6.8 ± 3.2	694	0.85 ± 0.06	28.5 ± 2.7	903	8.05 ± 0.27	656 ± 19	-20.1
UDS									
UDS026657 ^a	3.8161	<8.5	980	0.88 ± 0.09	6.6 ± 2.1	531	1.51 ± 0.17	389 ± 64	-20.6
UDS003783ª	3.8056	<6.3	836	0.72 ± 0.15	22.1 ± 7.4	768 725	4.39 ± 0.43	369 ± 89	-20.7
UDS295074	3.7606	4.6 ± 0.5	1020	1.02 ± 0.08	14.2 ± 2.3	735	6.67 ± 0.66	443 ± 87	-21.5
UDS015507	3.61	9.8 ± 2.2	1224	1.65 ± 0.10	79.9 ± 12.8	841	24.9 ± 0.50	475 ± 19	-21.0
UDS151133	3.5628	4.9 ± 0.6	908	1.5 ± 0.13	11.5 ± 2.1	756	7.50 ± 0.91	332 ± 109	-21.6
UDS026015	3.497	3.3 ± 1.3	371	0.19 ± 0.06	17.9 ± 5.8	827	1.47 ± 0.05	566 ± 33	-19.5
UDS003724	3.4616	6.3 ± 1.3	793	0.74 ± 0.10	67.3 ± 21.8	1242	13.7 ± 0.17	443 ± 12	-20.4
UDS376176	3.3188	2.2 ± 0.6	404	0.25 ± 0.06	2.6 ± 1.5	514	0.36 ± 0.08	672 ± 172	-19.9
UDS020394	3.3018	7.4 ± 0.9	698	0.67 ± 0.06	12.6 ± 2.5	823	2.19 ± 0.09	456 ± 42	-20.0
UDS000166	3.24	6.5 ± 0.9	489	0.51 ± 0.05	9.0 ± 2.4	614	0.67 ± 0.08	170 ± 83	-19.8
UDS380039 ^a	3.23	<9.6	649	0.78 ± 0.18	9.2 ± 3.9	744	2.26 ± 0.06	334 ± 20	-19.4
UDS026615	3.199	4.7 ± 0.8	739	0.50 ± 0.07	27.8 ± 10.2	915	6.0 ± 0.33	628 ± 58	-20.0
UDS019280	3.186	9.5 ± 1.5	841	1.03 ± 0.13	47.8 ± 19.9	945	7.06 ± 0.10	466 ± 16	-19.6
UDS197598	3.000	2.1 ± 0.4	827	1.34 ± 0.24	26.7 ± 3.1	516	40.7 ± 1.48	166 ± 52	-21.6
Average (1σ)	_	5.3 (±2.5)	680 (±187)	$0.76 (\pm 0.39)$	17.8 (±19.6)	836 (±226)	3.34 (±8.15)	533 (±158)	_
AGN	2.6005	65115	740	0.05 0.16	167 1 10	2007	2.01 0.21	440 + 10	
CDFS006905b	3.6885	6.5 ± 1.5	749	0.85 ± 0.16	16.7 ± 1.9	2206	2.91 ± 0.31	442 ± 12	-

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Table A1 - continued

Object ID	Zspec	EW ₀ (Å)	C III] λ 1909 FWHM (km s ⁻¹)	Flux $\times 10^{-18}$ (erg s ⁻¹ cm ⁻²)	EW ₀ (Å)	FWHM (km s ⁻¹)	$\begin{aligned} \text{Ly}\alpha \lambda 1216 \\ \text{Flux} \\ \times 10^{-18} \\ (\text{erg s}^{-1} \text{cm}^{-2}) \end{aligned}$	$\Delta v_{\mathrm{Ly}\alpha}$ $(\mathrm{km}\mathrm{s}^{-1})$	$M_{ m UV}$
CDFS006327 ^b	3.4937	4.5 ± 1.4	678	0.19 ± 0.02	26.2 ± 2.4	976	1.39 ± 0.39	606 ± 12	_
CDFS025897b	3.4706	5.0 ± 0.7	929	0.84 ± 0.08	13.6 ± 1.5	732	3.80 ± 0.22	522 ± 12	_
CDFS217560	3.4535	2.1 ± 0.8	534	0.19 ± 0.07	56.8 ± 6.9	827	8.52 ± 0.12	409 ± 12	_
CDFS247279b	3.2959	14.1 ± 3.6	650	1.1 ± 0.05	7.3 ± 1.6	1161	1.49 ± 0.31	916 ± 96	_
CDFS020988b	3.2638	3.8 ± 0.6	486	0.09 ± 0.01	51.7 ± 10.5	921	2.66 ± 0.05	498 ± 12	_
UDS300247 ^b	3.794	15.3 ± 3.9	923	1.03 ± 0.14	22.2 ± 4.0	778	5.24 ± 0.19	452 ± 30	_
UDS006692b	3.7663	2.1 ± 0.5	471	0.34 ± 0.08	7.3 ± 1.2	748	2.03 ± 0.28	455 ± 132	_
UDS382631ab	3.2176	<17.2	1178	2.45 ± 0.17	60.1 ± 6.9	826	29.4 ± 0.40	453 ± 12	_
UDS145830 ^b	3.2094	12.4 ± 1.0	1118	1.16 ± 0.06	81.2 ± 13.3	983	7.51 ± 0.99	66 ± 15	_

^adenotes those objects with a lower limit for C III]

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^bdenotes an AGN candidate that was removed in the final sample and not included in the quoted average values.