# Redshift-dependent RSD bias from intrinsic alignment with DESI Year 1 spectra

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# ABSTRACT

We estimate the redshift-dependent, anisotropic clustering signal in the Dark Energy Spectroscopic Instrument (DESI) Year 1 Survey created by tidal alignments of Luminous Red Galaxies (LRGs) and a selection-induced galaxy orientation bias. To this end, we measured the correlation between LRG shapes and the tidal field with DESI's Year 1 redshifts, as traced by LRGs and Emission-Line Galaxies. We also estimate the galaxy orientation bias of LRGs caused by DESI's aperture-based selection, and find it to increase by a factor of seven between redshifts 0.4-1.1 due to redder, fainter galaxies falling closer to DESI's imaging selection cuts. These effects combine to dampen measurements of the quadrupole of the correlation function ( $\xi_2$ ) caused by structure growth on scales of  $10-80 h^{-1}$  Mpc by about 0.15 per cent for low redshifts (0.4 < z < 0.6) and 0.8 per cent for high (0.8 < z < 1.1), a significant fraction of DESI's error budget. We provide estimates of the  $\xi_2$  signal created by intrinsic alignments that can be used to correct this effect, which is necessary to meet DESI's forecasted precision on measuring the growth rate of structure. While imaging quality varies across DESI's footprint, we find no significant difference in this effect between imaging regions in the Legacy Imaging Survey.

Key words: methods: data analysis - large-scale structure of Universe - dark energy - observations.

# **1 INTRODUCTION**

Measuring the growth of large-scale structure in the Universe informs us about the components that drive it: gravity and dark energy. The main observable used to measure this evolution on large scales is the Kaiser effect (Kaiser 1987). As structure grows, matter falls towards dense regions. This increases the recessional velocity of matter between us and a dense region, while decreasing it for matter falling in from the other side. The result is a 'squashing' effect in redshift space. This is the dominant source of redshift-space distortions (RSD) on scales larger than clusters (around  $10 h^{-1}$  Mpc). Clustering is quantified using the correlation function. This can be expressed as a series of spherical harmonics, of which the quadrupole  $\xi_2$  describes the anisotropic clustering that arises from RSD. On large scales, the growth rate of structure is linearly related to  $\xi_2$ . This makes RSD a powerful test of cosmological parameters and measuring it is one of the two main science goals of the the Dark Energy Spectroscopic Instrument (DESI).

DESI is in the midst of a 5-yr survey, measuring spectra of over 40 million galaxies within 16 000 square degrees of the sky. The instrument is installed on the 4-m Mayall telescope and can gather tens of thousands of extragalactic spectra in one night with its focal plane, which is comprised of 5000 individually controlled robots to position fibres on to galaxies (Levi et al. 2013; DESI Collaboration 2016a,b, 2022).

DESI is forecasting a 0.4–0.7 per cent measurement of the growth rate of structure,  $f\sigma_8$ . This is measured through the comparison of anisotropic and isotropic clustering,  $\xi_2$  and  $\xi_0$ . It is more difficult to obtain high precision on  $\xi_2$  than  $\xi_0$ , so errors on  $f\sigma_8$  are dominated by  $\xi_2$ . To meet DESI's science goals, it is imperative to measure  $\xi_2$ to a precision of at least 0.4–0.7 per cent. A subtle effect that could take up a significant fraction of this error budget is the bias in  $\xi_2$ due to a combination of two effects: intrinsic alignment (IA) and a selection-induced orientation bias (Hirata 2009).

IA refers to physical correlations between galaxy shapes and with galaxy shapes to the underlying density. See Lamman et al.

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**Figure 1.** The parameters used to describe the shape and orientation of the ellipse created by projecting an elliptical galaxy. Here, the ellipticity is measured relative to North. For our measurement, ellipticity is measured relative to the positions of a tracer sample.

(2023a) for a pedagogical guide to IA, Joachimi et al. (2015) and Troxel & Ishak (2015) for detailed reviews. It is historically measured as a contaminant of weak lensing, but IA in upcoming surveys may provide novel constraints on galaxy formation and cosmology (Chisari & Dvorkin 2013; Kurita & Takada 2023; Okumura & Taruya 2023; Xu et al. 2023). For DESI, IA also needs to be understood as a bias in measurements of anisotropic clustering.

This effect arises from the extent to which galaxy shapes are correlated with the underlying tidal field. The primary axis of Luminous Red Galaxies (LRGs) tend to be aligned along strands of density and point towards denser regions. This creates a clustering bias when combined with DESI's aperture-based target selection. An elliptical galaxy with its primary axis pointed at the observer will have a more concentrated light profile on the sky and a higher fraction of its light will fall within the aperture. This makes DESI more likely to observe galaxies which lie in density filaments that are parallel to the line of sight (LOS). For a visualization of this effect, see Fig. 1 in Lamman et al. (2023b). Studies have explored the effects of orientationdependent selection in Sloan Digital Sky Survey (SDSS) catalogues with differing results (Martens et al. 2018; Obuljen, Percival & Dalal 2020; Singh, Yu & Seljak 2021). We expect this effect to be more pronounced with DESI, which has a smaller fibre aperture of 1.5 arcsec in diameter, as opposed to SDSS' 3 arcsec aperture.

A total-magnitude selection would remove this bias from DESI, but spectroscopic success is highly dependent on the surface brightness of an object. Especially for a survey which prioritizes speed, there will be a surface-brightness dependence on the sample which is easier to impose explicitly as a target cut (Zhou et al. 2022). This selection-induced bias in galaxy orientation likely also affects DESI's Emission-Line Galaxy (ELG) sample. However, ELGs are not the primary tracers DESI uses for measuring large-scale structure and they display weaker alignments. While predicted by simulations, there is currently no observed shape alignment in ELG-like (spiral) galaxies (Johnston et al. 2019; Samuroff et al. 2019, 2023).

Measuring IA for the purpose of predicting an RSD bias has a few differences from IA measured in the context of weak lensing, which requires very precise shape measurements with well controlled systematic effects to measure gravitational shear. We only require shape measurements which are more precise than intrinsic shape variation. This is the case with LRGs in DESI's Legacy Imaging Survey, which are relatively large and bright. Therefore it is more valuable for us to use the full redshift sample available than limit to a region which overlaps with a deeper imaging survey, as will be done with other DESI IA measurements (Lange et al. in prep). Also, since this study uses spectroscopic redshifts, we can sufficiently isolate pairs with low separation along the LOS to the degree that we are unconcerned about contamination from weak lensing or across redshift bins. Lamman et al. (2023b) used photometric redshifts from DESI's imaging catalogue to estimate that this effect could lower DESI's measurement of  $\xi_2$  by about 0.5 per cent for LRGs. In this work, we use DESI's Year one spectra (DESI Collaboration, in prep) to produce estimates which can be used to correct DESI's RSD measurements. We measure the tidal alignment of LRGs as traced by LRGs and ELGs, assess the impacts of imaging on the IA measurement, and estimate the redshift dependence of the selection-induced shape polarization. We report the resulting redshift-dependent bias for DESI's Year one RSD results and discuss sources of systematic uncertainties.

# 2 DESI CATALOGUES

# 2.1 Imaging

DESI's targets are chosen from DR9 of the Legacy Imaging Survey (Dey et al. 2019; Myers et al. 2023). This contains imaging of sources in 14 000\$ square degrees of the extragalactic sky from three different telescopes: Mayall *z*-band Legacy Survey at the Mayall telescope at Kitt Peak (MzLS), the DECam Legacy Survey from the Blanco telescope at Cerro Tololo (DECaLS), and the Beijing-Arizona Sky Survey from the BOK telescope at Kitt Peak (BASS) (Zou et al. 2017). A region in the South Galactic cap also contains imaging from the Dark Energy Survey (DES).

Their shapes are fit using TRACTOR (Lang, Hogg & Mykytyn 2016). After deconvolving with a point-spread function (PSF), shapes are modelled at the pixel level with several light profiles: exponential disc, de Vaucouleurs, Sersic, PSF, and round-exponential. The default shape parameters are chosen based on a modified  $\chi^2$  criteria which avoids over-fitting bright targets as round-exponentials. Measuring intrinsic alignment requires shape orientations, so where galaxies are fit as circles (PSF and round-exponentials), we use shape parameters from the best-fit between exponential disc and de Vaucouleurs. This will not affect our final results as the DESI target selection does not depend on these derived shape parameters.

The parameters used to describe each projected galaxy image are its primary axis, *a*, secondary axis, *b*, and orientation  $\theta$  (Fig. 1). This is used to describe the shape of an ellipse relative to some direction using  $\epsilon_+$ :

$$\epsilon_{+} = \frac{a-b}{a+b}\cos(2\theta) \tag{1}$$

DESI's LRG target selection from this catalogue includes a cut based on the expected flux which falls within a DESI fibre, which corresponds to a 1.5 arcsec diameter aperture. The *z*-band magnitude within the aperture is limited to  $z_{\rm fibre} < 21.61$  in the Northern Galactic Cap and  $z_{\rm fibre} < 21.60$  in the Southern Galactic Cap. For more information DESI's target selection, see Raichoor et al. (2020) and Zhou et al. (2020, 2022).

This shape fitting and target selection is dependent upon imaging quality, which varies across sky regions. To qualify the effect of imaging quality on shape parameters, we separate the LRGs into three sky regions: The MzLS and BASS region, the DECaLS region which does not contain DES imaging, and the DES region. We compare the reported axis ratio, b/a, of the reported galaxy shapes in each region in Fig. 2. The MzLS and BASS region reports more eccentric LRG shapes than the other regions, and the region with highest-quality imaging, DES, reports the roundest shapes. While this may indicate an overcorrecting of the PSF in MzLS and BASS imaging, we measure the IA signal independently in these regions and do not find any significant impact on final results (Section 3.2).



**Figure 2.** The distribution of projected axis ratios for LRGs in three DESI imaging regions: the North Galactic cap, which contains imaging from MzLS and BASS, the portion of the South Galactic cap which contains DECaLS but no DES imaging, and the DES region. Residuals from the mean are plotted below. These reported shapes have been deconvolved with a point spread function to account for imaging conditions. The region with the highest quality imaging, DES, reports the least eccentric shape.

# 2.2 Spectroscopy

We used spectra from DESI's internal data release, Iron, which is comprised of data from commissioning through Year 1 of the survey (DESI Collaboration, in prep). It contains spectra of 2.9 million LRGs and 4.0 million ELGs from observations taken during 2020 December 14, through 2022 June 13 (Zhou et al. 2022; Guy et al. 2023; Raichoor et al. 2023). For approximate redshift distributions of the samples, see Fig. 2 of DESI Collaboration (2023a). We used Iron's large-scale structure catalogueue which includes a veto on targets based on hardware and imaging conditions. The catalogueue contains weights to account for completeness based on the probability that each target was assigned. LRGs also have a weight to account for redshift failures. From this catalogueue, we ran a basic redshiftquality cut to obtain our final samples of 2.5 million LRGs and 3.1 million ELGs.

The projected correlation functions used to calibrate our measurements (Section 3.1) were made with this Year 1 data. Due to internal blinding policies, our estimation of DESI's  $\xi_2$  measurements are calibrated with the smaller, publicly available spectroscopic catalogueue from DESI's Survey Validation (Lan et al. 2023; DESI Collaboration 2023a,b). Our determination of the  $\xi_2$  signal which arises from IA is independent of the RSD  $\xi_2$  signal.

### **3 INTRINSIC ALIGNMENT MEASUREMENT**

#### 3.1 IA formalism

As in Lamman et al. (2023b), we measure the correlation between galaxy shapes and density by averaging the ellipticity of each LRG relative to the separation vector between it and nearby galaxies in the tracer sample.<sup>1</sup>

$$\mathcal{E}(r_{\rm p}) = \langle \epsilon_+(a, b, \theta) \rangle,\tag{2}$$

<sup>1</sup>Code available here: github.com/cmlamman/ellipse\_alignment

For a given galaxy-tracer pair, *a* and *b* are the axis lengths of the galaxy shape and  $\theta$  is the orientation of the galaxy relative to the separation vector between it and the tracer. This is measured as a function of the projected separation between them,  $r_{\rm p}$ .

We limit the separation of pairs along the LOS,  $r_{\Pi}$ , to  $\pm \Pi_{\text{max}} = 30 h^{-1}$  Mpc. This, along with clustering, is taken into account in our model when estimating how far along the LOS the IA measurement is averaged over.

For measuring the IA of our full LRG sample, we divide the tracer catalogueue into 100 sky regions based on right ascension and declination with an equal number of galaxies in each.  $\mathcal{E}(r_p)$  is measured independently in each region using its tracers and the full shape catalogueue, then averaged over every pair. This average included the catalogueue weights described in Section 2.2 for both the shape and tracer samples. The average of these 100 measurements and standard error is our final measurement.

IA is often quantified using a form of correlation functions generalized to include information about galaxy shapes (Mandelbaum et al. 2006). The IA correlation function relating galaxy shapes and density is

$$\xi_{g+}(r_{\rm p},\,\Pi) = \frac{S_{+}D - S_{+}R_{\rm D}}{R_{\rm S}R_{\rm D}}.$$
(3)

 $S_+D$  is the count of data-data pairs weighted by the orientation of shapes,  $S_+$ , relative to a tracer sample *D*. Measuring this as function of projected separation and averaging over each data pair,  $S_+D(r_p)/DD(r_p)$ , is equivalent to  $\mathcal{E}(r_p)$ .  $S_+R$  represents the data shapes relative to a random sample, which has an expectation value of 0.  $R_S R_D$  is the random-random count. Integrating  $\xi_{g+}$  along the LOS direction,  $\Pi$ , produces the projected IA correlation function:

$$w_{g+}(r_{p}) = \int_{-\Pi_{max}}^{\Pi_{max}} d\Pi \,\xi_{g+}(r_{p},\,\Pi).$$
(4)

For predicting the RSD bias that arises from IA,  $\mathcal{E}(r_p)$  is the most direct relevant observable. Unlike  $w_{g+}$ ,  $\mathcal{E}(r_p)$  is normalized by data pairs, not randoms. DD can be expressed as

$$DD(r_{p}) = RR \int_{-\Pi_{max}}^{\Pi_{max}} d\Pi \frac{DD(r_{p}, \Pi)}{RR}$$
$$= RR \int_{-\Pi_{max}}^{\Pi_{max}} d\Pi (1 + \xi(r_{p}, \Pi)) = RR(2\Pi_{max} + w_{p}(r_{p})).$$
(5)

Here,  $\xi$  and  $w_p$  are the typical correlation function and projected correlation function, as opposed to those weighted by shape alignments.  $w_{g+}$  can be expressed as

$$w_{g+}(r_{p}) = \frac{1}{RR} \int_{-\Pi_{max}}^{\Pi_{max}} d\Pi S_{+} D(r_{p}, \Pi) = \frac{S_{+} D(r_{p})}{RR}.$$
 (6)

Therefore, a given  $w_{g+}$  and  $\mathcal{E}$  made with the same  $\Pi_{max}$  and same clustering  $w_p$  are related as

$$w_{g+}(r_p) = (2\Pi + w_p(r_p))\mathcal{E}(r_p) = L\mathcal{E}(r_p).$$
 (7)

Here we have introduced *L*, which can be understood as the effective LOS distance that  $\mathcal{E}$  is measured over, adjusted to account for clustering which decreases the average LOS-separation of pairs. *L* is included in our final model of the RSD bias (Section 5).<sup>2</sup> While  $L\mathcal{E}$  is functionally equivalent to  $w_{g+}$ , we notate *L* and  $\mathcal{E}$  separately to be explicit about how the quantity was estimated.

 $^{2}L$  here is equivalent to  $L_{\text{eff}}$  in Lamman et al. (2023b).



**Figure 3.** The IA signal of LRGs over the entire redshift range, 0.4 < z < 1.1, compared to the estimate made in (Lamman et al. 2023b) with photometric distances. The photometric estimate was made with 17.5 million galaxies, compared to 2.5 million LRGs for the spectroscopic sample, but necessarily averaged over a larger radial distance. This is adjusted for here, which shows the 'relative' IA signal that has been calibrated by the effective radial depth *L*.

We can compare our spectroscopic IA measurement with a similar one made with photometric data (Lamman et al. 2023b) by scaling by L, as shown in Fig. 3. Although made with seven times fewer galaxies, the spectroscopic sample from Iron can be measured in smaller LOS bins and provide us with similar level of precision.

Spectroscopic data also allows us to better isolate the sample in radial bins and explore redshift dependence. To compare our IA signal between samples of different target classes and redshift distributions,  $\mathcal{E}(r_p)$  needs to be calibrated by *L* as well as the galaxy clustering bias, *b*. For bias-independent comparisons, we scale by a relative bias. The bias of a sample 2 relative to sample 1 is

$$b_{\rm rel} = \frac{b_2}{b_1} = \frac{D(z_1)}{D(z_2)} \left(\frac{w_{\rm p2}}{w_{\rm p1}}\right)^{1/2},\tag{8}$$

where D(z) is the linear growth function.

Therefore, when comparing IA measurements across samples we use the value  $(L/b_{rel})\mathcal{E}(r_p)$ . While L is taken into account when estimating the final RSD bias, b does not affect the final result. This is because the amplitude of the power spectrum quadrupole effect arises from the correlation of the galaxy density field and the selection-induced shape polarization, the latter of which is independent of bias.

When calculating distances and the growth factor, we assume a flat  $\Lambda$ CDM cosmology with  $\Omega_m = 0.286$ ,  $\Omega_{\Lambda} = 0.714$ , and  $H_0 = 69.6 \text{ km s}^{-1} \text{ Mpc}^{-1}$ .

#### 3.2 Dependence on imaging region

The amplitude of IA can strongly depend upon imaging quality and the methods used to estimate shapes. This is in part due to difficulties in accurately modelling imaging processes, and in part due to isophotal twisting (Fasano & Bonoli 1989), which causes the outer regions of galaxies to have a stronger alignment signal than the inner regions. This has been measured in SDSS-III's Baryon Oscillation Spectroscopic Survey (BOSS) LOWZ, DES, and Large Synoptic Survey Telescope (LSST) (Singh, Mandelbaum & More 2015; Leonard, Mandelbaum & Collaboration 2018; Zuntz et al. 2018; Georgiou et al. 2019; MacMahon & Leonard 2023). This is



Figure 4. Measurement of the tidal alignment of LRG shapes made independently in areas from three regions of DESI's Legacy Imaging Survey described in Section 2.2. DES has the highest quality imaging, but there is no significant effect on our total averaged IA signal.

a concern for our imaging catalogue, as imaging quality and shapes vary across region (Fig. 2). A lower signal due to poor imaging will result in an underestimate of our final estimate of the RSD bias.

To qualify the impact of imaging quality, we compare our IA signal across the three different imaging regions used in the Legacy Imaging Survey: DES, DECaLS, and MZLS + BASS. Each region has varying survey completeness, so to avoid edge effects we made these measurements in a limited area with the most completeness in each region. The result and size of each sample is shown in Fig. 4. We do not find a significant difference in the signals, which is in part due to measurement noise. Additionally, although the BASS + MZLS region may be overcorrecting for the PSF and producing more eccentric shapes, systematic imaging effects are uncorrelated with the tidal field. A small change in ellipticity does not propagate as an order-unity error on this signal, which is a very small response to the tidal shear. This may still be an issue for higher signal-to-noise detections beyond DESI Year 1.

As a null test, we reproduced this measurement using the crosscomponent of shape,  $\epsilon_{\times} = \frac{a-b}{a+b}\sin(2\theta)$  instead of  $\epsilon_{+}$  in equation (2). This was consistent with 0 on all scales.

#### 3.3 Dependence on redshift and tracer sample

The redshift dependence of IA is unclear (Samuroff, Mandelbaum & Blazek 2021; Samuroff et al. 2023; Zhou et al. 2023), and cannot be directly observed without accounting for luminosity differences across redshift bins. DESI's LRG sample is designed to have a constant comoving volume with redshift, which results in more luminous, and therefore more aligned, galaxies in high-redshift samples. However, since we are only inferring a systematic bias and not any physical trends, we only require the IA of each sample. The IA RSD bias is proportional to the amplitude of this signal, so if not properly accounted for, it could manifest in DESI's results as a false evolution of the growth rate as measured by the quadrupole of the correlation function. Therefore, we separate our LRG tracer sample into three subsamples based on redshift and measure the correlation of LRG shapes in each.

The samples are plotted in Fig. 5 and displayed in Table 1. To compare the strength of tidal alignment between redshifts, the signal is adjusted based on the clustering in each sample, as described in



**Figure 5.** Comparison of the intrinsic alignment of LRGs between spectroscopic redshift bins. The *y*-axis is scaled by the effective depth of the measurement *L* and the galaxy bias  $b_{rel}$ , which here is defined as  $b_{rel}(z = 0.7) = 1$ . These were calculated using the projected correlation function from DESI's Year one data. Errors here only include the statistical difference of the signal between sky regions and not from *b* or *L*. Nearby galaxies broadly display a weaker alignment, though here we have not accounted for luminosity differences across samples.

Section 3.1. As expected, we find the weakest signal for nearby galaxies (0.4 < z < 0.6).

We also measured the alignment of LRGs to the tidal field as traced by ELGs, as opposed to the same LRG sample (Fig. 6). In the overlapping redshift range of the LRG and ELG samples, 0.8 < z < 1.1, we find a similar IA signal once both samples are adjusted for clustering. Although some regions of DESI's Year one footprint are less complete for ELGs, this is accounted for in the catalogue completeness weights described in Section 2 and we find no impact of this on our IA measurement.

A jackknife covariance matrix for the spectroscopic LRG measurement made over all redshift ranges is shown in Fig. 7 and demonstrates no correlation between the bins of projected separation that the IA measurement  $\mathcal{E}$  was made in.

# 4 SELECTION-INDUCED SHAPE POLARIZATION

The final RSD bias is proportional to both the degree to which galaxies are aligned along the tidal field (IA) and the degree to which galaxies are aligned along the LOS due to target selection, or 'shape polarization'. For DESI, the latter plays a strong role in the redshift dependence of the RSD bias. Redder, fainter galaxies fall closer to the aperture-magnitude cut that is used to select DESI targets. Therefore their orientation will have a stronger impact on whether or not they are selected; an elongated galaxy aligned with the LOS will have more light concentrated within a sky aperture.

Lamman et al. (2023b) estimated the shape polarization of DESI LRGs from a parent sample without the aperture-magnitude cut. This was done by generating many 3D light profiles for each galaxy based on the expected distribution of triaxial shapes from Padilla & Strauss (2008). These light profiles were assigned random orientations then put through an aperture-magnitude cut. The average ellipticity of the

selected shapes relative to the LOS,  $\epsilon_{LOS}$ , is the selection-induced shape polarization.

As this selection is done on aperture magnitudes from an image deconvolved with a point-spread function, the shape selection bias is relatively independent of imaging quality. It matters more to model this effect with imaging that most closely reflects the intrinsic galaxy shape. Therefore the polarization estimate for the entire sample was made using the portion of DESI's footprint with the highest quality imaging, the DES region. Although this results in a noisier measurement, only the average polarization of a sample affects the final RSD bias.

To estimate the polarization of the LRG redshift samples, we averaged the polarization estimates from the parent sample in bins of colour and magnitude. LRGs in the redshift sample were assigned a polarization based on the average in their corresponding bin, and their total average is the polarization estimate of that sample. This was not done for the ELG sample, which were only used as tracers. A demonstration of this mapping can be seen in Fig. 8 and the results are also displayed in Table 1. It is important to note that the polarization varies more across redshift bins than the IA signal, meaning that the redshift dependence of the final RSD bias is more dependent on survey selection than physical alignments.

#### **5 FALSE RSD SIGNATURE IN DESI**

To estimate the RSD bias created by the combination of IA and the selection-induced polarization, we use a non-linear tidal model adopted from Lamman et al. (2023b). The full details can be found in this paper; we give only the results here. We have made minor notation changes for clarity.

The IA signal  $\mathcal{E}$  is combined with the effective LOS-distance *L*, described in Section 3.1, and the non-linear power spectrum *P* as  $\tau$ :

$$r = \frac{2L(r_p)\mathcal{E}(r_p)}{r_p \frac{d}{dr_p} \left[\frac{1}{r_p}\Psi\right]},\tag{9}$$

$$\Psi(R) = \int \frac{K}{2\pi} \frac{dK}{R} \frac{P(K)}{K} J_1(KR)$$
(10)

Here, *K* is 2D Fourier Space and  $J_1$  is the first Bessel function.  $\tau$  is measured independently in each bin of transverse separation,  $r_p$ . The final variable used in our result,  $\bar{\tau}$ , is the average of these determinations with standard error. The transverse bins we used for determining  $\tau$  were linear bins between 5 and 20  $h^{-1}$  Mpc. Since these are relatively large scales, the change from a linear to nonlinear power spectrum had minimal effects on our final result, though it produced more consistent values of  $\tau$  across the transverse bins.

The 'false' signature this produces in the quadrupole of the correlation function,  $\xi_2$ , is

$$\xi_{2, gI}(s) = \epsilon_{\text{LOS}} \frac{\bar{\tau}}{2\sigma_{E1}^2} \int \frac{q^2 \mathrm{d}q}{2\pi^2} P(q) j_2(qs).$$
(11)

Here,  $\epsilon_{\text{LOS}}$  is the selection-induced shape polarization,  $\sigma_{\text{E1}}^2$  is the variance of the shape parameter  $\epsilon_+$  described in Section 2,  $j_2$  is the second spherical Bessel function, and *s* is 3D separation. The relations most relevant for this study can be summarized as

$$\xi_{2,\,\mathrm{gI}} \propto \epsilon_{\mathrm{LRG}} \frac{\bar{\tau}}{\sigma_{\mathrm{EI}}^2} \propto \epsilon_{\mathrm{LRG}} \frac{L\mathcal{E}}{\sigma_{\mathrm{EI}}^2}.$$
(12)

Note that, through cancellation, this result is independent of the amplitude of the power spectrum and galaxy bias, *b*. This is because  $\xi_{2,gI}$  arises from the correlation of the galaxy density field and the selection-induced shape polarization, the latter of which is

**Table 1.** Samples and values used to estimate the RSD bias for three LRG redshift bins and the LRG x ELG cross-correlation.  $r_p$  and s are given in units of  $h^{-1}$  Mpc. The tracer samples used in the top three rows were also used as the shape sample. The last row uses ELG tracers with LRG shapes. The table shows the redshift range and number N of tracers used, and properties of the shape sample: the variance of the real component of ellipticities  $\sigma_{E1}^2$  and the estimated selection-induced polarization of shapes along the LOS,  $\epsilon_{LOS}$ . We do not include uncertainties for these columns as they have negligible statistical errors. The IA signal  $\mathcal{E}(r_p)$  is measured as the ellipticity of shapes relative to the tracer sample.  $L(r_p)$  is the effective LOS-distance that  $\mathcal{E}(r_p)$  is averaged over.  $\tau(r_p)$  is defined in equation (9) and is a combination of  $\mathcal{E}(r_p)$ ,  $L(r_p)$ , and the power spectrum. These are functions of transverse separation,  $r_p$  and are shown in this table as averages over the marked scales. The final column shows the average amplitude of the anisotropic clustering created by IA; the quadrupole of the correlation function without RSD effects. The full estimate of this final result along with the statistical error is shown in Fig. 9.

Tracer	Zmin	z <sub>max</sub>	Ν	$\sigma_{\rm E1}^2$	$\epsilon_{ m LOS}$	$\bar{\mathcal{E}}(0 < r_{\rm p} < 20)$	$\bar{L}(0 < r_{\rm p} < 20)$	$\bar{\tau}(5 < r_{\rm p} < 18)$	$\bar{\xi}_{2, gI}(5 < s < 80)$
LRG	0.4	0.6	529 852	0.046	$2.3 \times 10^{-3}$	$1.8 \times 10^{-3}$	95.1 $h^{-1}$ Mpc	$5.9 \pm 0.5  imes 10^{-3}$	0.044
LRG	0.6	0.8	805 181	0.033	$6.7 \times 10^{-3}$	$2.1 \times 10^{-3}$	94.9 $h^{-1}$ Mpc	$7.0 \pm 0.2 \times 10^{-2}$	0.22
LRG	0.8	1.1	896150	0.026	$12.8 \times 10^{-3}$	$1.8 \times 10^{-2}$	$92.9 h^{-1} Mpc$	$5.6 \pm 0.2 \times 10^{-2}$	0.41
ELG	0.8	1.1	591 687	0.026	$12.8 \times 10^{-3}$	$1.9 \times 10^{-3}$	$73.2 \ h^{-1} \ Mpc$	$4.3\pm0.3\times10^{-2}$	0.34



Figure 6. Correlation between LRG shapes and the underlying galaxy density, as traced by both LRGs and ELGs. These samples are both in the redshift range 0.8 < z < 1.1. For comparison, this IA signal is scaled by the samples' clustering, as described in 3.1.



**Figure 7.** The reduced covariance matrix of  $\mathcal{E}$  between bins of transverse separation for our IA measurement with LRG tracers across the full redshift range. The identity matrix has been subtracted from this plot. This demonstrates that there is no evidence for significant correlations between the measurements of  $\mathcal{E}$  in each bin of projected separation.

independent of bias. It does depend on the projected correlation function  $w_p$  through *L*. Also, since the IA signal only affects  $\xi_{2,gI}$  through  $\bar{\tau}$ , which can be determined in transverse bins independently, we can forecast  $\xi_{2,gI}$  beyond the projected scales used to measure  $\mathcal{E}$ .

The variables measured for this estimate are listed in Table 1 and the final quadrupole signature for all our samples is shown in Fig. 9. To provide context for this signal, we estimate the total quadrupole signatures  $\xi_2$  expected for these galaxy samples. They are based on HOD catalogues made with the ABACUSSUMMIT simulations (Hadzhiyska et al. 2021; Maksimova et al. 2021; Yuan et al. prep), and scaled with measurements from DESI's Survey Validation (DESI Collaboration 2023b). Fig. 9 shows 1 per cent of these estimates, which is well above DESI's total error budget for measuring  $\xi_2$ . Since the  $\xi_2$  signal created by the growth of structure is opposite in sign to that created by IA, we have also multiplied this 1 per cent line by -1 for a clear comparison. On the scales used to measure  $f\sigma_8$  (10 <s < 80  $h^{-1}$  Mpc),  $\xi_2$  for LRGs will be dampened by around 0.15 per cent between redshifts 0.4 and 0.6, 0.53 per cent between redshifts 0.6 and 0.8, and 0.80 per cent between redshifts 0.8 and 1.1.  $\xi_2$  as measured by LRG x ELG cross-correlations will be biased by around 0.83 per cent between redshifts of 0.8-1.1. These are a significant fraction of DESI's forecasted error on  $f\sigma_8$  (0.4–0.7 per cent), which is dominated by the error  $\xi_2$ .

We used a Non-linear Alignment model, which has shown to be valid down to  $6h^{-1}$  Mpc for LRGs (Singh, Mandelbaum & More 2015). In principle our estimate can be extended down to the scales of the Fingers of God effect, where peculiar velocities of galaxies create a 'smearing' along the LOS, as opposed to the 'squashing' created by structure growth (Jackson 1972). Here, the sign of  $\xi_2$  switches and this bias will result in an enhancement of the signal. However, as non-linear effects become more apparent here and this effect is less relevant for DESI's main science goals, it is most valuable to interpret the bias on large scales.

#### **6** CONCLUSION

We measure the tidal alignment of LRGs with DESI Year 1 redshifts, using both LRG and ELG tracers. We also estimate a redshiftdependent polarization in LRG orientations relative to the LOS which arises from an aperture-based target selection. Using a nonlinear tidal model, we calculate the signal this will create in DESI's measurements of the quadrupole of the correlation function. It ranges from 0.2–1.1 per cent of the quadrupole signal created by RSD. This is a significant fraction of DESI's full-survey error budget of around 0.4–0.7 per cent for measuring the growth rate with LRGs and ELGs combined.

The RSD bias is over five times larger in the highest redshift sample, 0.8 < z < 1.1, than the lowest, 0.4 < z < 0.6. This is partially due to a stronger alignment signal, but mostly due to the selection effect. Galaxies at higher redshifts are redder



Figure 8. Estimates of shape polarization for each of the three redshift bins, or the tendency for these samples to have shapes aligned with the LOS due to an aperture-magnitude cut. These are based on the model in (Lamman et al. 2023b), shown in the grey gridded regions, which estimates the polarization in bins of colour and magnitude for a sample of LRGs without the aperture-magnitude cut. The total polarization estimate for each redshift bin is the average of each galaxy's polarization corresponding to its respective colour bin. The highest redshift sample contains redder, fainter galaxies which are closer to the survey cuts and therefore more likely to have biased galaxy orientations.



**Figure 9.** The anisotropic clustering signal arising from tidal alignment and a selection bias,  $\xi_{2,gI}$ . Statistical errors are shown in the shaded bands, although the total errors are dominated by systematic effects (Section 6). The vertical shaded region falls below the approximate validity of the non-linear IA model. For context, we have also plotted 1 per cent of the expected  $\xi_2$  signal from RSD. This is well above DESI's error budget for measuring the growth rate of structure, which is 0.4–0.7 per cent for LRGs and ELGs combined. Since the  $\xi_2$  signal created by the growth of structure is opposite in sign to that created by IA, we have multiplied the RSD  $\xi_2$  by -1 for an easier comparison. This plot demonstrates that IA will dampen DESI's RSD measurements to a degree that is a significant fraction of the error budget, particularly at higher redshifts. Incorporating these estimates into the  $\xi_2$  measurement will mitigate this effect.

and fainter, falling closer to the target selection cuts. Therefore, their orientation has a stronger influence on whether or not they pass the aperture-magnitude cut and the sample has a stronger orientation polarization. If uncorrected for, this redshift-dependent bias will suppress measurements of the growth rate at all scales, but especially higher redshifts. Therefore it will also bias determinations of how the growth rate evolves, a critical estimator for constraining cosmological models (Kazantzidis & Perivolaropoulos 2021). Our estimates of the quadrupole signature created by IA presented here can be used to adjust initial estimates of  $xi_2$  and mitigate this bias.

These results agree with the previous work in Lamman et al. (2023a), which measured the IA signal using photometric redshifts and estimated that it will produce around a 0.5 per cent decrease on measurements of  $\xi_2$  with the full LRG sample. While large, upcoming photometric surveys can provide constraints on IA, for

this effect it is most important to understand the IA of our particular sample. Additionally, redshifts are necessary to make clean distant cuts to explore redshift dependence in both IA and the shape polarization.

The largest uncertainty in our final results comes from systematic effects in the estimate of the selection-induced shape polarization. This is sensitive to assumptions in the light profiles used for mock selection and the underlying triaxial distribution of shapes, which is based on SDSS imaging (Padilla & Strauss 2008) and not in clear agreement with comparable galaxies in hydrodynamic simulations (Bassett & Foster 2019). This could be significantly improved with a large imaging survey such as the Dark Energy Survey or the upcoming Legacy Survey of Space and Time (Ivezić et al. 2019; Gatti et al. 2021). The methods of imaging and shape fits have a known impact on the inferred intrinsic shapes of galaxies (Georgiou et al. 2019; MacMahon & Leonard 2023).

While we do not expect a significant RSD bias from the shape alignment of ELGs, their spins are known to correlate with the tidal field (Lee 2011). Their shapes could also be biased in DESI's sample, as spectroscopic quality depends upon disc orientation (Hirata 2009). This bias could be explored in DESI through correlations between the quadrupole of the correlation function and fundamental plane residuals (Singh, Yu & Seljak 2021).

The remaining four years of DESI's main survey will produce millions more spectroscopic redshifts, allowing us to refine IA measurements and their redshift dependence. These will also produce higher precision RSD measurements, necessitating the need to incorporate the anisotropic clustering effect caused by IA.

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## DATA AVAILABILITY

The DESI Legacy Imaging Survey is publicly available at legacysu rvey.org and DESI's Early Data Release is publicly available at data .desi.lbl.gov/doc/releases/edr/. Iron covers the DESI Year 1 sample and will be publicly released as part of DESI Data Release 1 (DR1). ABACUSSUMMIT simulations are publicly available at abacusnbody. org. Code for projecting ellipsoids and generating light profiles can be found at github.com/cmlamman/ellipse\_alignment.

Data plotted in this paper are available at zenodo.org/uploads/101 62040.

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