



# Search for pairs of muons with small displacements in $pp$ collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

The ATLAS Collaboration\*

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## ABSTRACT

A search for new phenomena giving rise to pairs of opposite electrically charged muons with impact parameters in the millimeter range is presented, using  $139 \text{ fb}^{-1}$  of  $\sqrt{s} = 13$  TeV  $pp$  collision data from the ATLAS detector at the LHC. The search targets the gap in coverage between existing searches targeting final states with leptons with large displacement and prompt leptons. No significant excess over the background expectation is observed and exclusion limits are set on the mass of long-lived scalar supersymmetric muon-partners (smuons) with much lower lifetimes than previously targeted by displaced muon searches. Smuon lifetimes down to 1 ps are excluded for a smuon mass of 100 GeV, and smuon masses up to 520 GeV are excluded for a proper lifetime of 10 ps, at 95% confidence level. Finally, model-independent limits are set on the contribution from new phenomena to the signal-region yields.

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## 1. Introduction

The Standard Model (SM) contains many particles that have a significant lifetime which, when produced at a collider, travel a certain distance before decaying away from the primary proton-proton ( $pp$ ) interaction. Despite this, the majority of beyond the standard model (BSM) searches at the Large Hadron Collider (LHC) focus on prompt decays, and are not optimized for particles that travel a macroscopic distance. Many BSM theories predict particles that have significant lifetimes including  $R$ -parity-conserving supersymmetry (SUSY) [1–7] as well as  $R$ -parity-violating SUSY models [8,9], models like split-SUSY [10,11], exotic scenarios such as universal extra dimensions [12,13] and gauge-mediated SUSY breaking (GMSB) [14–16].

In GMSB SUSY models the lightest SUSY particle (LSP) is a nearly massless gravitino, and the next-to-lightest SUSY particle (NLSP) becomes long-lived due to the small coupling to the LSP. Well-motivated versions of this model have a stau ( $\tilde{\tau}$ ) as the single NLSP, or a selectron ( $\tilde{e}$ ), smuon ( $\tilde{\mu}$ ), and  $\tilde{\tau}$  as a set of degenerate co-NLSPs [17]. In these models, pair-produced sleptons ( $\tilde{\ell}$ ) of the same flavor decay into an invisible gravitino and a charged lepton of the same flavor as the parent  $\tilde{\ell}$ .

This paper presents a search for supersymmetric partners of the muon ( $\tilde{\mu}$ ) with a lifetime of  $\mathcal{O}(1 - 10)$  ps, targeting a gap in coverage between prompt slepton searches, and displaced slepton searches which have optimal sensitivity for lifetimes around  $\mathcal{O}(100 - 1000)$  ps. This regime has been highlighted as a possible

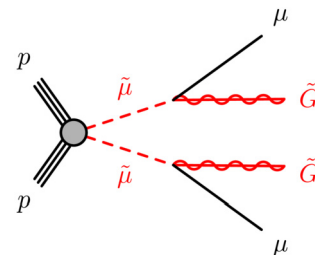


Fig. 1. Decay topology of the simplified model considered where smuons ( $\tilde{\mu}$ ) are pair produced and each smuon decays to a gravitino ( $\tilde{G}$ ) and a muon ( $\mu$ ).

blind spot in BSM searches at the LHC [18]. Fig. 1 shows a diagram of the targeted signal. A combination of results from the LEP experiments excludes the superpartners of the right-handed muons ( $\tilde{\mu}_R$ ) of any lifetime for masses less than 96.3 GeV [19–23]. Previous searches for long-lived sleptons have been performed by the ATLAS [24] and CMS [25] collaborations, excluding smuons up to 700 GeV and 620 GeV respectively, for a lifetime of 100 ps.

## 2. ATLAS detector

The ATLAS detector [26] at the LHC covers nearly the entire solid angle around the collision point.<sup>1</sup> It consists of an inner

<sup>1</sup> ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the  $z$ -axis along the beam pipe. The  $x$ -axis points from the IP to the center of the LHC ring, and the  $y$ -axis points upwards. Cylindrical coordinates  $(r, \phi)$  are used in the transverse plane,  $\phi$  being

\* E-mail address: [atlas.publications@cern.ch](mailto:atlas.publications@cern.ch).

tracking detector surrounded by a thin superconducting solenoid, electromagnetic and hadron calorimeters, and a muon spectrometer incorporating three large superconducting air-core toroidal magnets.

The inner-detector system (ID) is immersed in a 2 T axial magnetic field and provides charged-particle tracking in the range  $|\eta| < 2.5$ . The high-granularity silicon pixel detector covers the vertex region and typically provides four measurements per track, the first hit normally being in the insertable B-layer installed before Run 2 [27,28]. It is followed by the silicon microstrip tracker, which usually provides eight measurements per track. These silicon detectors are complemented by the transition radiation tracker (TRT), which enables radially extended track reconstruction up to  $|\eta| = 2.0$ . The TRT also provides electron identification information based on the fraction of hits (typically 30 in total) above a higher energy-deposit threshold corresponding to transition radiation.

The calorimeter system covers the pseudorapidity range  $|\eta| < 4.9$ . Within the region  $|\eta| < 3.2$ , electromagnetic calorimetry is provided by barrel and endcap high-granularity lead/liquid-argon (LAr) calorimeters, with an additional thin LAr presampler covering  $|\eta| < 1.8$  to correct for energy loss in material upstream of the calorimeters. Hadron calorimetry is provided by the steel/scintillator-tile calorimeter, segmented into three barrel structures within  $|\eta| < 1.7$ , and two copper/LAr hadron endcap calorimeters. The solid angle coverage is completed with forward copper/LAr and tungsten/LAr calorimeter modules optimized for electromagnetic and hadronic energy measurements respectively.

The muon spectrometer (MS) comprises separate trigger and high-precision tracking chambers measuring the deflection of muons in a magnetic field generated by the superconducting air-core toroidal magnets. The field integral of the toroids ranges between 2.0 and 6.0 Tm across most of the detector. Three layers of precision chambers, each consisting of layers of monitored drift tubes, cover the region  $|\eta| < 2.7$ , complemented by cathode-strip chambers in the forward region, where the background is highest. The muon trigger system covers the range  $|\eta| < 2.4$  with resistive-plate chambers in the barrel, and thin-gap chambers in the endcap regions.

Interesting events are selected by the first-level trigger system implemented in custom hardware, followed by selections made by algorithms implemented in software in the high-level trigger [29]. The first-level trigger accepts events from the 40 MHz bunch crossings at a rate below 100 kHz, which the high-level trigger further reduces in order to record events to disk at about 1 kHz.

An extensive software suite [30] is used in the reconstruction and analysis of real and simulated data, in detector operations, and in the trigger and data acquisition systems of the experiment.

### 3. Data and simulated event samples

The data set used in this analysis was collected by the ATLAS detector in  $pp$  collisions provided by the LHC during Run 2 from 2015 to 2018. The beams collided at a center-of-mass energy of  $\sqrt{s} = 13$  TeV and with a minimum separation of 25 ns between consecutive proton bunch crossings. The average number  $\langle \mu \rangle$  of additional  $pp$  interactions per bunch crossing (pile-up) ranged from 14 in 2015 to about 38 in 2017–2018. After data-quality requirements [31], applied to ensure that all parts of the detector were operational during data-taking, the data sample amounts to a total integrated luminosity of  $139 \text{ fb}^{-1}$ . Candidate events were selected by a di-muon trigger [29,32], where the trigger has no

explicit selection cuts on the transverse impact parameter of the muon's track with respect to the beam line,  $d_0$ .

To evaluate signal sensitivity, Monte Carlo (MC) events in a simplified GMSB SUSY model were simulated with up to two additional partons at leading order using MADGRAPH5\_aMC@NLO v2.6.1 [33] with the NNPDF2.310 parton distribution function (PDF) set [34], and interfaced to PYTHIA 8.230 [35] using the A14 set of tuned parameters [36]. The smuon decay was simulated using GEANT4 [37]. The impact of multiple interactions in the same and neighboring bunch crossings (pileup) was modeled by overlaying each hard-scattering event with simulated minimum-bias events generated with PYTHIA 8.210 [35] using the A3 tune [38] and NNPDF2.310 PDF set [34]. Signal cross sections were calculated at next-to-leading order in  $\alpha_s$ , with soft-gluon emission effects added at next-to-leading-logarithm accuracy [39–43]. The nominal cross section and uncertainty were taken from an envelope of predictions using different PDF sets and factorization and renormalization scales [44]. The simplified model used for interpretation gives a cross section of  $5.4^{+0.1}_{-0.2}$  pb for a  $\tilde{\mu}$  with mass 50 GeV and  $0.1 \pm 0.05$  fb for a  $\tilde{\mu}$  with mass 700 GeV, where the cross section assumes degenerate left- and right-handed smuons. Simulated events were generated for  $\tilde{\mu}$  masses ranging from 50–700 GeV and lifetimes of 0.1 ps, 1 ps, 10 ps, and 100 ps.

The dominant SM background for this search originates from semileptonic  $B$ -hadron decays,  $b\bar{b} \rightarrow \mu^+\mu^-$ . A data driven background method is used to estimate the number of background events in the signal region, described in Section 5, and any subdominant background processes are included within this method. SM processes with prompt leptons such as  $Z$ +jets,  $W$ +jets,  $t\bar{t}$ , single top quark, and di-boson were assessed with simulation and found to be negligible.

### 4. Object reconstruction and event selection

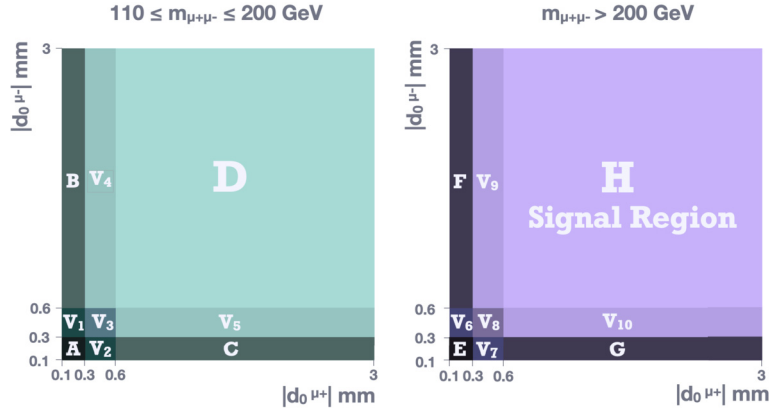
Candidate events are required to have at least one  $pp$  interaction vertex with a minimum of two associated tracks, where each track has  $p_T > 500$  MeV [45]. In events with multiple vertices, the primary vertex is defined as the vertex with the highest scalar sum of the squared transverse momenta of its associated tracks.

Muon candidates are reconstructed in the pseudorapidity range  $|\eta| < 2.5$  by matching MS tracks with ID tracks. The muons are required to be matched to the muons that fired the trigger. As requirements on the transverse momentum,  $p_T$ , of the muons are imposed in the online trigger decision, higher  $p_T$  requirements of  $p_T > 20$  GeV are applied to the muons to ensure that trigger efficiencies are constant in the relevant phase-space. The muons must satisfy the Medium identification requirements defined in Ref. [46]. Nominal track reconstruction was used and it was verified by MC simulations that the nominal tracking reconstruction efficiency, including the muon identification efficiency, is flat as a function of  $|d_0|$  in the entire  $|d_0|$  region (0–3) mm that is used in this search. Isolation criteria are applied to each muon in order to suppress contributions of muons from the decays of hadronic sources. The isolation variable named “PflowLoose” defined in Ref. [46] is applied using the working point where the  $\Delta R$  parameter is varied.

To reduce background contributions from prompt SM processes a selection cut of  $|d_0| \geq 0.1$  mm is made. No selection cuts are made on the longitudinal impact parameter,  $z_0$ .

Selected events are required to have at least two muons of opposite electric charge passing all the selection criteria. Although this search probes non-prompt muons, muons from prompt process, such as from the decay of a  $Z$  boson, can be included in the control regions (defined in Section 5) due to the resolution of the transverse impact parameter. In order to reduce contributions of di-muon events from the decay of  $Z$  bosons the invariant mass

the azimuthal angle around the  $z$ -axis. The pseudorapidity is defined in terms of the polar angle  $\theta$  as  $\eta = -\text{Tan}(\theta/2)$ . Angular distance is measured in units of  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ .



**Fig. 2.** A graphic depicting the regions for the extended ABCD background estimation method for Set of Regions 1. Data is split into a 2D plane in  $|d_0|$  with the positive charged muon versus the negative charged muon. For the regions on the left the invariant mass of the two muons must be greater or equal to 110 GeV and less than or equal to 200 GeV, the invariant mass of the two muons must be greater than 200 GeV for the regions on the right. A, B, C, and E are control regions. The regions V1-V10, D, F, and G are validation regions, and H is the signal region.

**Table 1**

The definitions of the three sets of regions that are used to define the CRs, VRs, and SRs, where the columns 2 and 3 refer to the boundaries used to subdivide the planes in  $|d_0|$  as depicted in Fig. 2 for Set of Regions 1.

Set of Regions	Lower displacement region	Higher displacement region	Threshold $m_{\mu^+\mu^-}$	Additional cut
1	$0.1 \leq  d_0  < 0.3$	$0.6 \leq  d_0  < 3$ mm	200 GeV	-
2	$0.1 \leq  d_0  < 0.3$	$0.6 \leq  d_0  < 3$ mm	140 GeV	-
3	$0.1 \leq  d_0  < 0.3$	$0.6 \leq  d_0  < 1.3$ mm	125 GeV	$\Delta R_{\mu^+\mu^-} > 3$ rad.

of the two muons must be  $\geq 110$  GeV. Further selection cuts are made to provide control regions (CRs), validation regions (VRs) and signal regions (SRs) as described in Section 5.

## 5. Background estimation

After selection, the dominant background originates from semi-leptonic  $B$ -hadron decays,  $b\bar{b} \rightarrow \mu^+\mu^-$ . Possible contributions from processes with prompt leptons such as  $Z$ +jets,  $W$ +jets,  $t\bar{t}$ , single top quark, and di-boson were assessed with MC and were found to be negligible.

An extended ABCD method [47] is used to estimate the number of background events in the signal region. The ABCD method requires two uncorrelated variables, chosen here to be the absolute value of the transverse impact parameter,  $|d_0|$ , of the muon track for the positively charged muon and the negatively charged muon, to create a two-dimensional plane in data that is split into four sections. The variable  $|d_0|$  is chosen as it is a good proxy for the displacement of muon tracks. To further separate background dominated regions from signal dominated regions the method is extended with a third variable, the invariant mass of the two muons ( $m_{\mu^+\mu^-}$ ), which is uncorrelated with both impact parameters to define eight regions, A to H, as shown in Fig. 2. The region H, located at high invariant mass and impact parameter values, is defined as the signal region.

Using the ABCD method the number of background events in the signal region H can be estimated by:

$$N_H^{\text{est. bkg.}} = N_A^{\text{data}} \cdot r^{d_0^+} \cdot r^{d_0^-} \cdot r^{m_{\mu^+\mu^-}}, \quad (1)$$

where  $r^{d_0^+}$ ,  $r^{d_0^-}$ , and  $r^{m_{\mu^+\mu^-}}$  are ratios from high to low  $|d_0|$  of the positive and negative charged muons and  $m_{\mu^+\mu^-}$  regions. The ratios can be computed from multiple pairs of regions and so the regions with highest statistics and lowest signal contamination are used to compute the ratios:  $r^{d_0^+} = N_C^{\text{data}}/N_A^{\text{data}}$ ,  $r^{d_0^-} = N_B^{\text{data}}/N_A^{\text{data}}$ , and  $r^{m_{\mu^+\mu^-}} = N_E^{\text{data}}/N_A^{\text{data}}$ , where  $N_i^{\text{data}}$  is the number of observed events in data in region  $i \in \{A, B, C, E\}$ .

The boundaries between regions are defined to be non-adjacent in  $|d_0|$ , such that there are ‘‘gap’’ regions between each of the control regions (CRs) and signal regions (SRs). The gaps are needed to reduce signal contamination in the CRs, and allow for validation regions (VRs) to be defined to confirm the correctness of the background estimation.

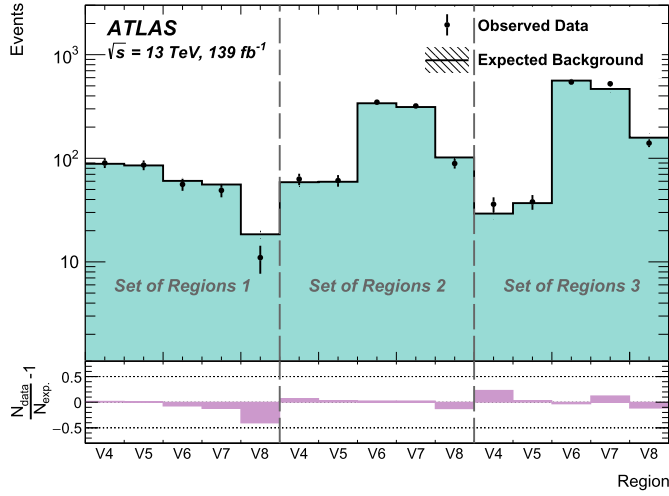
Three overlapping sets of regions are defined in Table 1 and Set of Regions 1 is depicted in Fig. 2. In the Set of Regions 2 and 3 the invariant mass selection threshold of 200 GeV is lowered to 140 and 125 GeV respectively. An additional selection on the angular distance between the two muons,  $\Delta R_{\mu^+\mu^-} > 3$  radians, is applied in Set of Regions 3. The Set of Regions 1 provides the strongest sensitivity to the GMSB model targeted by this search. However, the Set of Regions 2 and 3 provide two additional SRs, and CRs and VRs without significant contamination from signals with masses  $\leq 200$  GeV, where a potentially large signal contamination in the control regions would bias the background estimation and absorb partly an excess of events in the signal region. These sets of regions are also used to set limits on model-independent BSM signal processes. When performing hypothesis tests only one set of regions is used per signal mass and lifetime point as the three regions are overlapping in events.

The extended ABCD method requires the three variables,  $|d_0^{\mu^+}|$ ,  $|d_0^{\mu^-}|$ , and  $m_{\mu^+\mu^-}$ , to be uncorrelated. To quantify potential correlations closure tests are performed in the validation regions V4 to V8 of the three region sets, using regions A, B, C, E, V1, V2 and V3 for the ratios to compute the number of expected events in the validation region. Regions F, G, and V9-V10 are not included in the test due to potential signal contamination. The numbers of estimated and observed events are consistent within statistical uncertainties at the  $1\sigma$  standard deviation level, except for V8 where the standard deviation is found to be  $2\sigma$  in Set of Regions 1. Fig. 3 shows expected and observed number of events in the validation regions V4-V8 for each set of regions and the relative difference for each validation region. Based on these results, a conservative uncertainty of 40% is assigned as non-closure systematic uncertainty in the three signal regions.

**Table 2**

Presented here are the expected and observed number of background events in each SR H defined by the three sets of regions, where combined statistical and systematic uncertainties on expected  $N_H^{\text{bkg}}$  are given. Columns four to six show the 95% CL upper limits on the visible cross section ( $\langle A\epsilon\sigma_{\text{obs}}^{\text{95}} \rangle$ ) and on the number of signal events given the observed and expected number of events ( $S_{\text{obs}}^{\text{95}}$  and  $S_{\text{exp}}^{\text{95}}$ ), where there is a  $\pm 1\sigma$  uncertainty on  $S_{\text{exp}}^{\text{95}}$ . The last two columns indicate the  $CL_B$  value, i.e. the confidence level observed for the background-only hypothesis, and the discovery  $p$ -value ( $p(s=0)$ ), capped at 0.5.

Set of Regions	Expected $N_H^{\text{bkg}}$	Observed $N_H^{\text{data}}$	$\langle A\epsilon\sigma_{\text{obs}}^{\text{95}} \rangle$ [fb]	$S_{\text{obs}}^{\text{95}}$	$S_{\text{exp}}^{\text{95}}$	$CL_B$	$p(s=0)$ (Z)
1	$2.1 \pm 0.8$	1	0.02	3.3	$4.2^{+2.5}_{-1.4}$	0.27	0.50 (0.00)
2	$12.5 \pm 5.2$	7	0.04	5.2	$8.5^{+4.0}_{-2.7}$	0.08	0.50 (0.00)
3	$17.2 \pm 7.4$	14	0.06	8.9	$10.5^{+5.0}_{-3.1}$	0.26	0.50 (0.00)



**Fig. 3.** Expected and observed number of events in the validation regions V4–V8 for each set of regions are shown. The ratio panel shows the relative difference for each validation region.

## 6. Systematic uncertainties

Differences in the efficiencies and resolution between data and MC were corrected for by the application of multiplicative factors. Uncertainties on the factor correcting for differences in trigger efficiency are found to be within 5% for all signal samples and are applied as a systematic uncertainty to the number of expected signal events. The remaining uncertainties that may arise due to the differences between MC and data include muon reconstruction efficiency, the scale and resolution of the momentum of the muon and are found to be negligible,  $< 1\%$ , across all analysis regions and no systematic uncertainties are added in these categories. Statistical uncertainties associated with the simulated MC samples are also accounted for. Theoretical uncertainties include cross section uncertainties of 1.6–3.9%. Modeling uncertainties for the signal process, e.g. effects of varying the factorization and renormalization scales, initial and final state radiation, and underlying effects were found to be negligible and are not included as systematics. In order for the exclusion fit to interpolate between the generated lifetimes and extrapolate above the generated lifetimes a re-weighting procedure is applied to the generated signal MC samples to provide smuon lifetimes between the generated lifetimes. Uncertainties on this procedure were evaluated by observing the percentage change when re-weighting the lifetime of a generated sample to that of a different generated lifetime. From these studies a 12% uncertainty is applied when the lifetime is re-weighted to a higher lifetime. However, when the lifetime is re-weighted to a lower lifetime the agreement between the lifetime of the original and the re-weighted lifetime is well within statistical uncertainty and thus no uncertainty is added in this case.

A background uncertainty of 40% for the non-closure of the extended ABCD method is assigned to the estimated number of background events in the signal region.

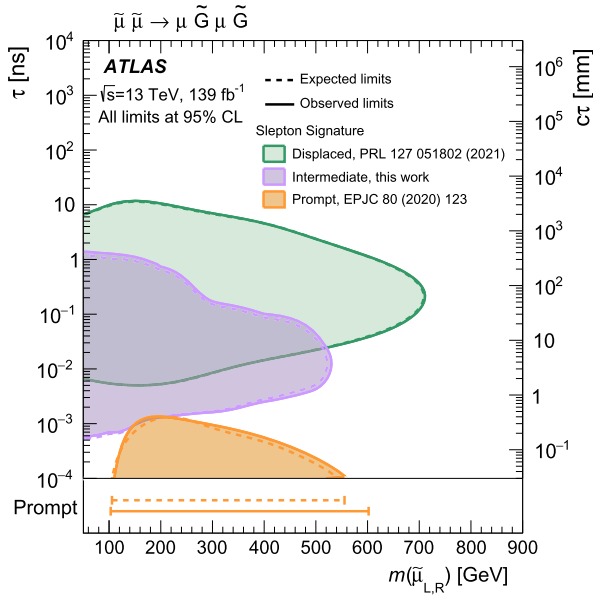
The uncertainty in the combined 2015–2018 integrated luminosity is 1.7% [48], obtained using the LUCID-2 detector [49] for the primary luminosity measurements.

## 7. Results

The observed number of events in all signal regions, shown in Table 2, is found to be compatible with the background expectation within statistical and systematic uncertainties.

The results are used to set model-independent limits on the contribution of generic BSM signals in each of the SRs defined by the three sets of regions, assuming no signal contamination in the CRs. Possible signal leakage to the CRs can produce a bias in the background estimation, leading to conservative limits. For the GMSB model the signal contamination for smuon masses of 300, 400 and 500 GeV in Set of Regions 1 is negligible, for Set of Regions 2 the signal contamination is negligible for smuon mass 200 GeV and for the Set of Region 3 the signal contamination is negligible for smuon masses of 50 and 100 GeV. Table 2 shows the results of a model-independent fit, performed using the HistFitter package [50], where the  $CL_s$  prescription [51] is used to set upper limits at 95% CL on the visible cross-section ( $\langle A\epsilon\sigma_{\text{obs}}^{\text{95}} \rangle$ ), where  $A$  is the acceptance and  $\epsilon$  the efficiency, as well as on the observed ( $S_{\text{obs}}^{\text{95}}$ ) and expected ( $S_{\text{exp}}^{\text{95}}$ ) number of events from potential new physics processes in the SRs. The  $p$ -value and the corresponding significance for the background only hypothesis are also evaluated.

Model-dependent exclusion limits for GMSB SUSY models on the smuon masses and lifetimes are derived at 95% confidence level following the  $CL_s$  prescription [51]. A combined likelihood fit is performed in regions A, B, C, E, and H, including the possible signal contribution in the control regions. The Set of Regions 1 provides the best expected sensitivity across all the plane and therefore is the only set of regions used in the model dependent fit. The HistFitter package [50] is used for statistical interpretation, and all systematic uncertainties are treated as Gaussian nuisance parameters during the fitting procedure. A re-weighting procedure is applied to the generated signal samples to provide signal lifetime points between those that are generated. Interpolation is used to provide smooth results throughout the plane, connecting the discrete mass and lifetime values that were simulated and the re-weighted lifetimes. The results are presented in Fig. 4 where smuon lifetimes down to 1 ps and smuon masses up to 520 GeV are excluded, assuming degenerate left- and right-handed smuons. The results from a previous search for displaced leptons (named here as the Displaced Slepton Signature) with large impact parameter ( $3 \text{ mm} < |d_0| < 300 \text{ mm}$ ) are also shown [24]. A search for direct slepton production with prompt decay [52] is reinterpreted using the RECAST framework [53] to cover lifetimes below the picosecond regime (named here as the Prompt Slepton Signature). This is the first explicit reinterpretation of prompt lepton searches



**Fig. 4.** Expected (dashed) and observed (solid) exclusion contours for  $\tilde{\mu}_L \tilde{\mu}_R$  NISP production as a function of the left- and right-handed smuons,  $\tilde{\mu}_{L,R}$ , mass and lifetime at 95% CLs, for the Displaced Slepton Signature (Phys. Rev. Lett. 127 051802 (2021) [24]), the Intermediate Slepton Signature (the result of this search in this paper) and the Prompt Slepton Signature (Eur. Phys. J. C 80 (2020) 123 [52]) reinterpretation. The lines below the graphs show the expected and observed limits from the prompt search where the smuons are assumed to be prompt.

in the long-lived regime in ATLAS. The search presented in this paper bridges the gap between both previous searches.

## 8. Conclusion

A search has been presented for pairs of opposite electrically charged muons with impact parameters in the millimeter range using  $139 \text{ fb}^{-1}$  of  $\sqrt{s} = 13 \text{ TeV}$   $pp$  collision data from the ATLAS detector. This search addresses a gap in coverage of possible new physics signatures between existing searches for leptons with large displacement and prompt leptons. Results are consistent with the SM background prediction. This search provides unique sensitivity to long-lived scalar supersymmetric muon-partners (smuons) with much lower lifetimes than previously targeted by ATLAS searches. Smuon lifetimes down to 1 ps and smuon masses up to 520 GeV are excluded at 95% confidence level.

As no requirements are made on missing energy, displaced vertices, or jets, this result is model-independent and applicable to any BSM model producing at least two opposite sign, isolated displaced muons with transverse momenta greater than 20 GeV.

## Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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G. Aad<sup>103, </sup>, B. Abbott<sup>121, </sup>, K. Abeling<sup>55, </sup>, N.J. Abicht<sup>49, </sup>, S.H. Abidi<sup>29, </sup>, A. Abouhorma<sup>35e, </sup>, H. Abramowicz<sup>152, </sup>, H. Abreu<sup>151, </sup>, Y. Abulaiti<sup>118, </sup>, A.C. Abusleme Hoffman<sup>138a, </sup>, B.S. Acharya<sup>69a, 69b, </sup>, C. Adam Bourdarios<sup>4, </sup>, L. Adamczyk<sup>86a, </sup>, L. Adamek<sup>156, </sup>, S.V. Addepalli<sup>26, </sup>, M.J. Addison<sup>102, </sup>, J. Adelman<sup>116, </sup>, A. Adiguzel<sup>21c, </sup>, T. Adye<sup>135, </sup>, A.A. Affolder<sup>137, </sup>, Y. Afik<sup>36, </sup>, M.N. Agaras<sup>13, </sup>, J. Agarwala<sup>73a, 73b, </sup>, A. Aggarwal<sup>101, </sup>, C. Agheorghiesei<sup>27c, </sup>, A. Ahmad<sup>36, </sup>, F. Ahmadov<sup>38, </sup>, W.S. Ahmed<sup>105, </sup>, S. Ahuja<sup>96, </sup>, X. Ai<sup>62a, </sup>, G. Aielli<sup>76a, 76b, </sup>, M. Ait Tamliah<sup>35e, </sup>, B. Aitbenkikh<sup>35a, </sup>, I. Aizenberg<sup>170, </sup>, M. Akbiyik<sup>101, </sup>, T.P.A. Åkesson<sup>99, </sup>, A.V. Akimov<sup>37, </sup>, D. Akiyama<sup>169, </sup>, N.N. Akolkar<sup>24, </sup>, K. Al Khoury<sup>41, </sup>, G.L. Alberghi<sup>23b, </sup>, J. Albert<sup>166, </sup>, P. Albicocco<sup>53, </sup>, G.L. Albouy<sup>60, </sup>, S. Alderweireldt<sup>52, </sup>, M. Aleksa<sup>36, </sup>, I.N. Aleksandrov<sup>38, </sup>, C. Alexa<sup>27b, </sup>, T. Alexopoulos<sup>10, </sup>, A. Alfonsi<sup>115, </sup>, F. Alfonsi<sup>23b, </sup>, M. Algren<sup>56, </sup>, M. Alhroob<sup>121, </sup>, B. Ali<sup>133, </sup>, H.M.J. Ali<sup>92, </sup>, S. Ali<sup>149, </sup>, S.W. Alibocus<sup>93, </sup>, M. Aliev<sup>37, </sup>, G. Alimonti<sup>71a, </sup>, W. Alkakh<sup>55, </sup>, C. Allaire<sup>66, </sup>, B.M.M. Allbrooke<sup>147, </sup>, J.F. Allen<sup>52, </sup>, C.A. Allendes Flores<sup>138f, </sup>, P.P. Allport<sup>20, </sup>, A. Aloisio<sup>72a, 72b, </sup>, F. Alonso<sup>91, </sup>, C. Alpigiani<sup>139, </sup>, M. Alvarez Estevez<sup>100, </sup>, A. Alvarez Fernandez<sup>101, </sup>, M.G. Alviggi<sup>72a, 72b, </sup>, M. Aly<sup>102, </sup>, Y. Amaral Coutinho<sup>83b, </sup>,

A. Ambler <sup>105, [id](#)</sup>, C. Amelung <sup>36</sup>, M. Amerl <sup>102, [id](#)</sup>, C.G. Ames <sup>110, [id](#)</sup>, D. Amidei <sup>107, [id](#)</sup>,  
 S.P. Amor Dos Santos <sup>131a, [id](#)</sup>, K.R. Amos <sup>164, [id](#)</sup>, V. Ananiev <sup>126, [id](#)</sup>, C. Anastopoulos <sup>140, [id](#)</sup>, T. Andeen <sup>11, [id](#)</sup>,  
 J.K. Anders <sup>36, [id](#)</sup>, S.Y. Andrean <sup>47a,47b, [id](#)</sup>, A. Andreazza <sup>71a,71b, [id](#)</sup>, S. Angelidakis <sup>9, [id](#)</sup>, A. Angerami <sup>41, [id](#), [ag](#)</sup>,  
 A.V. Anisenkov <sup>37, [id](#)</sup>, A. Annovi <sup>74a, [id](#)</sup>, C. Antel <sup>56, [id](#)</sup>, M.T. Anthony <sup>140, [id](#)</sup>, E. Antipov <sup>146, [id](#)</sup>,  
 M. Antonelli <sup>53, [id](#)</sup>, D.J.A. Antrim <sup>17a, [id](#)</sup>, F. Anulli <sup>75a, [id](#)</sup>, M. Aoki <sup>84, [id](#)</sup>, T. Aoki <sup>154, [id](#)</sup>, J.A. Aparisi Pozo <sup>164, [id](#)</sup>,  
 M.A. Aparo <sup>147, [id](#)</sup>, L. Aperio Bella <sup>48, [id](#)</sup>, C. Appelt <sup>18, [id](#)</sup>, A. Apyan <sup>26, [id](#)</sup>, N. Aranzabal <sup>36, [id](#)</sup>,  
 C. Arcangeletti <sup>53, [id](#)</sup>, A.T.H. Arce <sup>51, [id](#)</sup>, E. Arena <sup>93, [id](#)</sup>, J-F. Arguin <sup>109, [id](#)</sup>, S. Argyropoulos <sup>54, [id](#)</sup>,  
 J.-H. Arling <sup>48, [id](#)</sup>, A.J. Armbruster <sup>36, [id](#)</sup>, O. Arnaez <sup>4, [id](#)</sup>, H. Arnold <sup>115, [id](#)</sup>, Z.P. Arrubarrena Tame <sup>110</sup>,  
 G. Artoni <sup>75a,75b, [id](#)</sup>, H. Asada <sup>112, [id](#)</sup>, K. Asai <sup>119, [id](#)</sup>, S. Asai <sup>154, [id](#)</sup>, N.A. Asbah <sup>61, [id](#)</sup>, J. Assahsah <sup>35d, [id](#)</sup>,  
 K. Assamagan <sup>29, [id](#)</sup>, R. Astalos <sup>28a, [id](#)</sup>, S. Atashi <sup>161, [id](#)</sup>, R.J. Atkin <sup>33a, [id](#)</sup>, M. Atkinson <sup>163</sup>, N.B. Atlay <sup>18, [id](#)</sup>,  
 H. Atmani <sup>62b</sup>, P.A. Atmasiddha <sup>107, [id](#)</sup>, K. Augsten <sup>133, [id](#)</sup>, S. Auricchio <sup>72a,72b, [id](#)</sup>, A.D. Auriol <sup>20, [id](#)</sup>,  
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 P. Balek <sup>86a, [id](#)</sup>, E. Ballabene <sup>23b,23a, [id](#)</sup>, F. Balli <sup>136, [id](#)</sup>, L.M. Baltes <sup>63a, [id](#)</sup>, W.K. Balunas <sup>32, [id](#)</sup>, J. Balz <sup>101, [id](#)</sup>,  
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 M. Begalli <sup>83d, [id](#)</sup>, M. Begel <sup>29, [id](#)</sup>, A. Behera <sup>146, [id](#)</sup>, J.K. Behr <sup>48, [id](#)</sup>, J.F. Beirer <sup>55, [id](#)</sup>, F. Beisiegel <sup>24, [id](#)</sup>,  
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 G. Bianco <sup>23b,23a, [id](#)</sup>, O. Biebel <sup>110, [id](#)</sup>, R. Bielski <sup>124, [id](#)</sup>, M. Biglietti <sup>77a, [id](#)</sup>, T.R.V. Billoud <sup>133, [id](#)</sup>, M. Bindi <sup>55, [id](#)</sup>,  
 A. Bingul <sup>21b, [id](#)</sup>, C. Bini <sup>75a,75b, [id](#)</sup>, A. Biondini <sup>93, [id](#)</sup>, C.J. Birch-sykes <sup>102, [id](#)</sup>, G.A. Bird <sup>20,135, [id](#)</sup>,  
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 K. Bjørke <sup>126, [id](#)</sup>, I. Bloch <sup>48, [id](#)</sup>, C. Blocker <sup>26, [id](#)</sup>, A. Blue <sup>59, [id](#)</sup>, U. Blumenschein <sup>95, [id](#)</sup>, J. Blumenthal <sup>101, [id](#)</sup>,  
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 M. Bona <sup>95, [id](#)</sup>, M. Boonekamp <sup>136, [id](#)</sup>, C.D. Booth <sup>96, [id](#)</sup>, A.G. Borbély <sup>59, [id](#)</sup>, I.S. Bordulev <sup>37, [id](#)</sup>,

H.M. Borecka-Bielska <sup>109, [id](#)</sup>, L.S. Borgna <sup>97, [id](#)</sup>, G. Borissov <sup>92, [id](#)</sup>, D. Bortoletto <sup>127, [id](#)</sup>, D. Boscherini <sup>23b, [id](#)</sup>,  
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D. Boye <sup>29, [id](#)</sup>, I.R. Boyko <sup>38, [id](#)</sup>, J. Bracinik <sup>20, [id](#)</sup>, N. Brahimi <sup>62d, [id](#)</sup>, G. Brandt <sup>172, [id](#)</sup>, O. Brandt <sup>32, [id](#)</sup>,  
F. Braren <sup>48, [id](#)</sup>, B. Brau <sup>104, [id](#)</sup>, J.E. Brau <sup>124, [id](#)</sup>, R. Brenner <sup>170, [id](#)</sup>, L. Brenner <sup>115, [id](#)</sup>, R. Brenner <sup>162, [id](#)</sup>,  
S. Bressler <sup>170, [id](#)</sup>, D. Britton <sup>59, [id](#)</sup>, D. Britzger <sup>111, [id](#)</sup>, I. Brock <sup>24, [id](#)</sup>, G. Brooijmans <sup>41, [id](#)</sup>, W.K. Brooks <sup>138f, [id](#)</sup>,  
E. Brost <sup>29, [id](#)</sup>, L.M. Brown <sup>166, [id](#), [m](#)</sup>, L.E. Bruce <sup>61, [id](#)</sup>, T.L. Bruckler <sup>127, [id](#)</sup>, P.A. Bruckman de Renstrom <sup>87, [id](#)</sup>,  
B. Brüers <sup>48, [id](#)</sup>, D. Bruncko <sup>28b, [id](#), [\\*](#)</sup>, A. Bruni <sup>23b, [id](#)</sup>, G. Bruni <sup>23b, [id](#)</sup>, M. Bruschi <sup>23b, [id](#)</sup>, N. Brusino <sup>75a, 75b, [id](#)</sup>,  
T. Buanes <sup>16, [id](#)</sup>, Q. Buat <sup>139, [id](#)</sup>, D. Buchin <sup>111, [id](#)</sup>, A.G. Buckley <sup>59, [id](#)</sup>, M.K. Bugge <sup>126, [id](#)</sup>, O. Bulekov <sup>37, [id](#)</sup>,  
B.A. Bullard <sup>144, [id](#)</sup>, S. Burdin <sup>93, [id](#)</sup>, C.D. Burgard <sup>49, [id](#)</sup>, A.M. Burger <sup>40, [id](#)</sup>, B. Burghgrave <sup>8, [id](#)</sup>,  
O. Burlayenko <sup>54, [id](#)</sup>, J.T.P. Burr <sup>32, [id](#)</sup>, C.D. Burton <sup>11, [id](#)</sup>, J.C. Burzynski <sup>143, [id](#)</sup>, E.L. Busch <sup>41, [id](#)</sup>,  
V. Büscher <sup>101, [id](#)</sup>, P.J. Bussey <sup>59, [id](#)</sup>, J.M. Butler <sup>25, [id](#)</sup>, C.M. Buttar <sup>59, [id](#)</sup>, J.M. Butterworth <sup>97, [id](#)</sup>,  
W. Buttinger <sup>135, [id](#)</sup>, C.J. Buxo Vazquez <sup>108, [id](#)</sup>, A.R. Buzykaev <sup>37, [id](#)</sup>, G. Cabras <sup>23b, [id](#)</sup>, S. Cabrera Urbán <sup>164, [id](#)</sup>,  
L. Cadamuro <sup>66, [id](#)</sup>, D. Caforio <sup>58, [id](#)</sup>, H. Cai <sup>130, [id](#)</sup>, Y. Cai <sup>14a, 14e, [id](#)</sup>, V.M.M. Cairo <sup>36, [id](#)</sup>, O. Cakir <sup>3a, [id](#)</sup>,  
N. Calace <sup>36, [id](#)</sup>, P. Calafiura <sup>17a, [id](#)</sup>, G. Calderini <sup>128, [id](#)</sup>, P. Calfayan <sup>68, [id](#)</sup>, G. Callea <sup>59, [id](#)</sup>, L.P. Caloba <sup>83b, [id](#)</sup>,  
D. Calvet <sup>40, [id](#)</sup>, S. Calvet <sup>40, [id](#)</sup>, T.P. Calvet <sup>103, [id](#)</sup>, M. Calvetti <sup>74a, 74b, [id](#)</sup>, R. Camacho Toro <sup>128, [id](#)</sup>,  
S. Camarda <sup>36, [id](#)</sup>, D. Camarero Munoz <sup>26, [id](#)</sup>, P. Camarri <sup>76a, 76b, [id](#)</sup>, M.T. Camerlingo <sup>72a, 72b, [id](#)</sup>,  
D. Cameron <sup>126, [id](#)</sup>, C. Camincher <sup>166, [id](#)</sup>, M. Campanelli <sup>97, [id](#)</sup>, A. Camplani <sup>42, [id](#)</sup>, V. Canale <sup>72a, 72b, [id](#)</sup>,  
A. Canesse <sup>105, [id](#)</sup>, M. Cano Bret <sup>80, [id](#)</sup>, J. Cantero <sup>164, [id](#)</sup>, Y. Cao <sup>163, [id](#)</sup>, F. Capocasa <sup>26, [id](#)</sup>, M. Capua <sup>43b, 43a, [id](#)</sup>,  
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J.I. Carlotto <sup>13, [id](#)</sup>, B.T. Carlson <sup>130, [id](#), [u](#)</sup>, E.M. Carlson <sup>166, 157a, [id](#)</sup>, L. Carminati <sup>71a, 71b, [id](#)</sup>, A. Carnelli <sup>136, [id](#)</sup>,  
M. Carnesale <sup>75a, 75b, [id](#)</sup>, S. Caron <sup>114, [id](#)</sup>, E. Carquin <sup>138f, [id](#)</sup>, S. Carrá <sup>71a, 71b, [id](#)</sup>, G. Carratta <sup>23b, 23a, [id](#)</sup>,  
F. Carrio Argos <sup>33g, [id](#)</sup>, J.W.S. Carter <sup>156, [id](#)</sup>, T.M. Carter <sup>52, [id](#)</sup>, M.P. Casado <sup>13, [id](#), [j](#)</sup>, M. Caspar <sup>48, [id](#)</sup>,  
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V. Cavasinni <sup>74a, 74b, [id](#)</sup>, Y.C. Cekmecelioglu <sup>48, [id](#)</sup>, E. Celebi <sup>21a, [id](#)</sup>, F. Celli <sup>127, [id](#)</sup>, M.S. Centonze <sup>70a, 70b, [id](#)</sup>,  
K. Cerny <sup>123, [id](#)</sup>, A.S. Cerqueira <sup>83a, [id](#)</sup>, A. Cerri <sup>147, [id](#)</sup>, L. Cerrito <sup>76a, 76b, [id](#)</sup>, F. Cerutti <sup>17a, [id](#)</sup>, B. Cervato <sup>142, [id](#)</sup>,  
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J. Chan <sup>171, [id](#)</sup>, W.Y. Chan <sup>154, [id](#)</sup>, J.D. Chapman <sup>32, [id](#)</sup>, E. Chapon <sup>136, [id](#)</sup>, B. Chargeishvili <sup>150b, [id](#)</sup>,  
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L. Chevalier <sup>136, [id](#)</sup>, V. Chiarella <sup>53, [id](#)</sup>, G. Chiarelli <sup>74a, [id](#)</sup>, N. Chiedde <sup>103, [id](#)</sup>, G. Chiodini <sup>70a, [id](#)</sup>,  
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 S.H. Connell <sup>33c, [id](#)</sup>, I.A. Connelly <sup>59, [id](#)</sup>, E.I. Conroy <sup>127, [id](#)</sup>, F. Conventi <sup>72a, [id](#), [al](#)</sup>, H.G. Cooke <sup>20, [id](#)</sup>,  
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 S. Crépé-Renaudin <sup>60, [id](#)</sup>, F. Crescioli <sup>128, [id](#)</sup>, M. Cristinziani <sup>142, [id](#)</sup>, M. Cristoforetti <sup>78a,78b, [id](#)</sup>, V. Croft <sup>115, [id](#)</sup>,  
 J.E. Crosby <sup>122, [id](#)</sup>, G. Crosetti <sup>43b,43a, [id](#)</sup>, A. Cueto <sup>100, [id](#)</sup>, T. Cuhadar Donszelmann <sup>161, [id](#)</sup>, H. Cui <sup>14a,14e, [id](#)</sup>,  
 Z. Cui <sup>7, [id](#)</sup>, W.R. Cunningham <sup>59, [id](#)</sup>, F. Curcio <sup>43b,43a, [id](#)</sup>, P. Czodrowski <sup>36, [id](#)</sup>, M.M. Czurylo <sup>63b, [id](#)</sup>,  
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 V. D'Amico <sup>110, [id](#)</sup>, J. Damp <sup>101, [id](#)</sup>, J.R. Dandoy <sup>129, [id](#)</sup>, M.F. Daneri <sup>30, [id](#)</sup>, M. Danninger <sup>143, [id](#)</sup>, V. Dao <sup>36, [id](#)</sup>,  
 G. Darbo <sup>57b, [id](#)</sup>, S. Darmora <sup>6, [id](#)</sup>, S.J. Das <sup>29, [id](#), [am](#)</sup>, S. D'Auria <sup>71a,71b, [id](#)</sup>, C. David <sup>157b, [id](#)</sup>, T. Davidek <sup>134, [id](#)</sup>,  
 B. Davis-Purcell <sup>34, [id](#)</sup>, I. Dawson <sup>95, [id](#)</sup>, H.A. Day-hall <sup>133, [id](#)</sup>, K. De <sup>8, [id](#)</sup>, R. De Asmundis <sup>72a, [id](#)</sup>,  
 N. De Biase <sup>48, [id](#)</sup>, S. De Castro <sup>23b,23a, [id](#)</sup>, N. De Groot <sup>114, [id](#)</sup>, P. de Jong <sup>115, [id](#)</sup>, H. De la Torre <sup>108, [id](#)</sup>,  
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 J. Del Peso <sup>100, [id](#)</sup>, F. Del Rio <sup>63a, [id](#)</sup>, F. Deliot <sup>136, [id](#)</sup>, C.M. Delitzsch <sup>49, [id](#)</sup>, M. Della Pietra <sup>72a,72b, [id](#)</sup>,  
 D. Della Volpe <sup>56, [id](#)</sup>, A. Dell'Acqua <sup>36, [id](#)</sup>, L. Dell'Asta <sup>71a,71b, [id](#)</sup>, M. Delmastro <sup>4, [id](#)</sup>, P.A. Delsart <sup>60, [id](#)</sup>,  
 S. Demers <sup>173, [id](#)</sup>, M. Demichev <sup>38, [id](#)</sup>, S.P. Denisov <sup>37, [id](#)</sup>, L. D'Eramo <sup>40, [id](#)</sup>, D. Derendarz <sup>87, [id](#)</sup>, F. Derue <sup>128, [id](#)</sup>,  
 P. Dervan <sup>93, [id](#)</sup>, K. Desch <sup>24, [id](#)</sup>, C. Deutsch <sup>24, [id](#)</sup>, F.A. Di Bello <sup>57b,57a, [id](#)</sup>, A. Di Ciaccio <sup>76a,76b, [id](#)</sup>,  
 L. Di Ciaccio <sup>4, [id](#)</sup>, A. Di Domenico <sup>75a,75b, [id](#)</sup>, C. Di Donato <sup>72a,72b, [id](#)</sup>, A. Di Girolamo <sup>36, [id](#)</sup>,  
 G. Di Gregorio <sup>5, [id](#)</sup>, A. Di Luca <sup>78a,78b, [id](#)</sup>, B. Di Micco <sup>77a,77b, [id](#)</sup>, R. Di Nardo <sup>77a,77b, [id](#)</sup>, C. Diaconu <sup>103, [id](#)</sup>,  
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 M. Donadelli <sup>83c, [id](#)</sup>, B. Dong <sup>108, [id](#)</sup>, J. Donini <sup>40, [id](#)</sup>, A. D'Onofrio <sup>77a,77b, [id](#)</sup>, M. D'Onofrio <sup>93, [id](#)</sup>, J. Dopke <sup>135, [id](#)</sup>,  
 A. Doria <sup>72a, [id](#)</sup>, N. Dos Santos Fernandes <sup>131a, [id](#)</sup>, M.T. Dova <sup>91, [id](#)</sup>, A.T. Doyle <sup>59, [id](#)</sup>, M.A. Draguet <sup>127, [id](#)</sup>,  
 E. Dreyer <sup>170, [id](#)</sup>, I. Drivas-koulouris <sup>10, [id](#)</sup>, A.S. Drobac <sup>159, [id](#)</sup>, M. Drozdova <sup>56, [id](#)</sup>, D. Du <sup>62a, [id](#)</sup>,  
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 O.A. Ducu <sup>27b, [id](#)</sup>, D. Duda <sup>52, [id](#)</sup>, A. Dudarev <sup>36, [id](#)</sup>, E.R. Duden <sup>26, [id](#)</sup>, M. D'uffizi <sup>102, [id](#)</sup>, L. Duflot <sup>66, [id](#)</sup>,  
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 L. Vacavant <sup>103, [id](#)</sup>, V. Vacek <sup>133, [id](#)</sup>, B. Vachon <sup>105, [id](#)</sup>, K.O.H. Vadla <sup>126, [id](#)</sup>, T. Vafeiadis <sup>36, [id](#)</sup>, A. Vaitkus <sup>97, [id](#)</sup>,  
 C. Valderanis <sup>110, [id](#)</sup>, E. Valdes Santurio <sup>47a, 47b, [id](#)</sup>, M. Valente <sup>157a, [id](#)</sup>, S. Valentinetti <sup>23b, 23a, [id](#)</sup>,  
 A. Valero <sup>164, [id](#)</sup>, E. Valiente Moreno <sup>164, [id](#)</sup>, A. Vallier <sup>103, [id](#), [af](#)</sup>, J.A. Valls Ferrer <sup>164, [id](#)</sup>,  
 D.R. Van Arneman <sup>115, [id](#)</sup>, T.R. Van Daalen <sup>139, [id](#)</sup>, A. Van Der Graaf <sup>49, [id](#)</sup>, P. Van Gemmeren <sup>6, [id](#)</sup>,  
 M. Van Rijnbach <sup>126, 36, [id](#)</sup>, S. Van Stroud <sup>97, [id](#)</sup>, I. Van Vulpen <sup>115, [id](#)</sup>, M. Vanadia <sup>76a, 76b, [id](#)</sup>, W. Vandelli <sup>36, [id](#)</sup>,  
 M. Vandenbroucke <sup>136, [id](#)</sup>, E.R. Vandewall <sup>122, [id](#)</sup>, D. Vannicola <sup>152, [id](#)</sup>, L. Vannoli <sup>57b, 57a, [id](#)</sup>, R. Vari <sup>75a, [id](#)</sup>,  
 E.W. Varnes <sup>7, [id](#)</sup>, C. Varni <sup>17a, [id](#)</sup>, T. Varol <sup>149, [id](#)</sup>, D. Varouchas <sup>66, [id](#)</sup>, L. Varriale <sup>164, [id](#)</sup>, K.E. Varvell <sup>148, [id](#)</sup>,  
 M.E. Vasile <sup>27b, [id](#)</sup>, L. Vaslin <sup>40</sup>, G.A. Vasquez <sup>166, [id](#)</sup>, F. Vazeille <sup>40, [id](#)</sup>, T. Vazquez Schroeder <sup>36, [id](#)</sup>,  
 J. Veatch <sup>31, [id](#)</sup>, V. Vecchio <sup>102, [id](#)</sup>, M.J. Veen <sup>104, [id](#)</sup>, I. Veliscek <sup>127, [id](#)</sup>, L.M. Veloce <sup>156, [id](#)</sup>, F. Veloso <sup>131a, 131c, [id](#)</sup>,  
 S. Veneziano <sup>75a, [id](#)</sup>, A. Ventura <sup>70a, 70b, [id](#)</sup>, A. Verbytskyi <sup>111, [id](#)</sup>, M. Verducci <sup>74a, 74b, [id](#)</sup>, C. Vergis <sup>24, [id](#)</sup>,  
 M. Verissimo De Araujo <sup>83b, [id](#)</sup>, W. Verkerke <sup>115, [id](#)</sup>, J.C. Vermeulen <sup>115, [id](#)</sup>, C. Vernieri <sup>144, [id](#)</sup>,  
 P.J. Verschuuren <sup>96, [id](#)</sup>, M. Vessella <sup>104, [id](#)</sup>, M.C. Vetterli <sup>143, [id](#), [ak](#)</sup>, A. Vgenopoulos <sup>153, [id](#), [f](#)</sup>,  
 N. Viaux Maira <sup>138f, [id](#)</sup>, T. Vickey <sup>140, [id](#)</sup>, O.E. Vickey Boeriu <sup>140, [id](#)</sup>, G.H.A. Viehhauser <sup>127, [id](#)</sup>, L. Vigani <sup>63b, [id](#)</sup>,  
 M. Villa <sup>23b, 23a, [id](#)</sup>, M. Villaplana Perez <sup>164, [id](#)</sup>, E.M. Villhauer <sup>52</sup>, E. Vilucchi <sup>53, [id](#)</sup>, M.G. Vincter <sup>34, [id](#)</sup>,  
 G.S. Virdee <sup>20, [id](#)</sup>, A. Vishwakarma <sup>52, [id](#)</sup>, A. Visibile <sup>115</sup>, C. Vittori <sup>36, [id](#)</sup>, I. Vivarelli <sup>147, [id](#)</sup>, V. Vladimirov <sup>168</sup>,  
 E. Voevodina <sup>111, [id](#)</sup>, F. Vogel <sup>110, [id](#)</sup>, P. Vokac <sup>133, [id](#)</sup>, J. Von Ahnen <sup>48, [id](#)</sup>, E. Von Toerne <sup>24, [id](#)</sup>,  
 B. Vormwald <sup>36, [id](#)</sup>, V. Vorobel <sup>134, [id](#)</sup>, K. Vorobev <sup>37, [id](#)</sup>, M. Vos <sup>164, [id](#)</sup>, K. Voss <sup>142, [id](#)</sup>, J.H. Vosseveld <sup>93, [id](#)</sup>,  
 M. Vozak <sup>115, [id](#)</sup>, L. Vozdecky <sup>95, [id](#)</sup>, N. Vranjes <sup>15, [id](#)</sup>, M. Vranjes Milosavljevic <sup>15, [id](#)</sup>, M. Vreeswijk <sup>115, [id](#)</sup>,  
 R. Vuillermet <sup>36, [id](#)</sup>, O. Vujanovic <sup>101, [id](#)</sup>, I. Vukotic <sup>39, [id](#)</sup>, S. Wada <sup>158, [id](#)</sup>, C. Wagner <sup>104</sup>, J.M. Wagner <sup>17a, [id](#)</sup>,  
 W. Wagner <sup>172, [id](#)</sup>, S. Wahdan <sup>172, [id](#)</sup>, H. Wahlberg <sup>91, [id](#)</sup>, R. Wakasa <sup>158, [id](#)</sup>, M. Wakida <sup>112, [id](#)</sup>, J. Walder <sup>135, [id](#)</sup>,  
 R. Walker <sup>110, [id](#)</sup>, W. Walkowiak <sup>142, [id](#)</sup>, A. Wall <sup>129, [id](#)</sup>, T. Wamorkar <sup>6, [id](#)</sup>, A.Z. Wang <sup>171, [id](#)</sup>, C. Wang <sup>101, [id](#)</sup>,  
 C. Wang <sup>62c, [id](#)</sup>, H. Wang <sup>17a, [id](#)</sup>, J. Wang <sup>64a, [id](#)</sup>, R.-J. Wang <sup>101, [id](#)</sup>, R. Wang <sup>61, [id](#)</sup>, R. Wang <sup>6, [id](#)</sup>,  
 S.M. Wang <sup>149, [id](#)</sup>, S. Wang <sup>62b, [id](#)</sup>, T. Wang <sup>62a, [id](#)</sup>, W.T. Wang <sup>80, [id](#)</sup>, W. Wang <sup>14a, [id](#)</sup>, X. Wang <sup>14c, [id](#)</sup>,  
 X. Wang <sup>163, [id](#)</sup>, X. Wang <sup>62c, [id](#)</sup>, Y. Wang <sup>62d, [id](#)</sup>, Y. Wang <sup>14c, [id](#)</sup>, Z. Wang <sup>107, [id](#)</sup>, Z. Wang <sup>62d, 51, 62c, [id](#)</sup>,  
 Z. Wang <sup>107, [id](#)</sup>, A. Warburton <sup>105, [id](#)</sup>, R.J. Ward <sup>20, [id](#)</sup>, N. Warrack <sup>59, [id](#)</sup>, A.T. Watson <sup>20, [id](#)</sup>, H. Watson <sup>59, [id](#)</sup>,  
 M.F. Watson <sup>20, [id](#)</sup>, E. Watton <sup>59, 135, [id](#)</sup>, G. Watts <sup>139, [id](#)</sup>, B.M. Waugh <sup>97, [id](#)</sup>, C. Weber <sup>29, [id](#)</sup>, H.A. Weber <sup>18, [id](#)</sup>,  
 M.S. Weber <sup>19, [id](#)</sup>, S.M. Weber <sup>63a, [id](#)</sup>, C. Wei <sup>62a</sup>, Y. Wei <sup>127, [id](#)</sup>, A.R. Weidberg <sup>127, [id](#)</sup>, E.J. Weik <sup>118, [id](#)</sup>,  
 J. Weingarten <sup>49, [id](#)</sup>, M. Weirich <sup>101, [id](#)</sup>, C. Weiser <sup>54, [id](#)</sup>, C.J. Wells <sup>48, [id](#)</sup>, T. Wenaus <sup>29, [id](#)</sup>, B. Wendland <sup>49, [id](#)</sup>,  
 T. Wengler <sup>36, [id](#)</sup>, N.S. Wenke <sup>111</sup>, N. Wermes <sup>24, [id](#)</sup>, M. Wessels <sup>63a, [id](#)</sup>, K. Whalen <sup>124, [id](#)</sup>, A.M. Wharton <sup>92, [id](#)</sup>,  
 A.S. White <sup>61, [id](#)</sup>, A. White <sup>8, [id](#)</sup>, M.J. White <sup>1, [id](#)</sup>, D. Whiteson <sup>161, [id](#)</sup>, L. Wickremasinghe <sup>125, [id](#)</sup>,  
 W. Wiedenmann <sup>171, [id](#)</sup>, C. Wiel <sup>50, [id](#)</sup>, M. Wielers <sup>135, [id](#)</sup>, C. Wiglesworth <sup>42, [id](#)</sup>, D.J. Wilbern <sup>121</sup>,  
 H.G. Wilkens <sup>36, [id](#)</sup>, D.M. Williams <sup>41, [id](#)</sup>, H.H. Williams <sup>129</sup>, S. Williams <sup>32, [id](#)</sup>, S. Willocq <sup>104, [id](#)</sup>,  
 B.J. Wilson <sup>102, [id](#)</sup>, P.J. Windischhofer <sup>39, [id](#)</sup>, F.I. Winkel <sup>30, [id](#)</sup>, F. Winklmeier <sup>124, [id](#)</sup>, B.T. Winter <sup>54, [id](#)</sup>,  
 J.K. Winter <sup>102, [id](#)</sup>, M. Wittgen <sup>144</sup>, M. Wobisch <sup>98, [id](#)</sup>, Z. Wolffs <sup>115, [id](#)</sup>, R. Wölker <sup>127, [id](#)</sup>, J. Wollrath <sup>161</sup>,  
 M.W. Wolter <sup>87, [id](#)</sup>, H. Wolters <sup>131a, 131c, [id](#)</sup>, A.F. Wongel <sup>48, [id](#)</sup>, S.D. Worm <sup>48, [id](#)</sup>, B.K. Wosiek <sup>87, [id](#)</sup>,

K.W. Woźniak<sup>87, [id](#)</sup>, S. Wozniwski<sup>55, [id](#)</sup>, K. Wraight<sup>59, [id](#)</sup>, C. Wu<sup>20, [id](#)</sup>, J. Wu<sup>14a, 14e, [id](#)</sup>, M. Wu<sup>64a, [id](#)</sup>,  
M. Wu<sup>114, [id](#)</sup>, S.L. Wu<sup>171, [id](#)</sup>, X. Wu<sup>56, [id](#)</sup>, Y. Wu<sup>62a, [id](#)</sup>, Z. Wu<sup>136, [id](#)</sup>, J. Wuerzinger<sup>111, [id](#)</sup>, T.R. Wyatt<sup>102, [id](#)</sup>,  
B.M. Wynne<sup>52, [id](#)</sup>, S. Xella<sup>42, [id](#)</sup>, L. Xia<sup>14c, [id](#)</sup>, M. Xia<sup>14b, [id](#)</sup>, J. Xiang<sup>64c, [id](#)</sup>, X. Xiao<sup>107, [id](#)</sup>, M. Xie<sup>62a, [id](#)</sup>,  
X. Xie<sup>62a, [id](#)</sup>, S. Xin<sup>14a, 14e, [id](#)</sup>, J. Xiong<sup>17a, [id](#)</sup>, D. Xu<sup>14a, [id](#)</sup>, H. Xu<sup>62a, [id](#)</sup>, L. Xu<sup>62a, [id](#)</sup>, R. Xu<sup>129, [id](#)</sup>, T. Xu<sup>107, [id](#)</sup>,  
Y. Xu<sup>14b, [id](#)</sup>, Z. Xu<sup>52, [id](#)</sup>, Z. Xu<sup>14a, [id](#)</sup>, B. Yabsley<sup>148, [id](#)</sup>, S. Yacoob<sup>33a, [id](#)</sup>, N. Yamaguchi<sup>90, [id](#)</sup>,  
Y. Yamaguchi<sup>155, [id](#)</sup>, E. Yamashita<sup>154, [id](#)</sup>, H. Yamauchi<sup>158, [id](#)</sup>, T. Yamazaki<sup>17a, [id](#)</sup>, Y. Yamazaki<sup>85, [id](#)</sup>, J. Yan<sup>62c, [id](#)</sup>,  
S. Yan<sup>127, [id](#)</sup>, Z. Yan<sup>25, [id](#)</sup>, H.J. Yang<sup>62c, 62d, [id](#)</sup>, H.T. Yang<sup>62a, [id](#)</sup>, S. Yang<sup>62a, [id](#)</sup>, T. Yang<sup>64c, [id](#)</sup>, X. Yang<sup>62a, [id](#)</sup>,  
X. Yang<sup>14a, [id](#)</sup>, Y. Yang<sup>44, [id](#)</sup>, Y. Yang<sup>62a, [id](#)</sup>, Z. Yang<sup>62a, [id](#)</sup>, W.-M. Yao<sup>17a, [id](#)</sup>, Y.C. Yap<sup>48, [id](#)</sup>, H. Ye<sup>14c, [id](#)</sup>,  
H. Ye<sup>55, [id](#)</sup>, J. Ye<sup>44, [id](#)</sup>, S. Ye<sup>29, [id](#)</sup>, X. Ye<sup>62a, [id](#)</sup>, Y. Yeh<sup>97, [id](#)</sup>, I. Yeletsikh<sup>38, [id](#)</sup>, B.K. Yeo<sup>17a, [id](#)</sup>,  
M.R. Yexley<sup>97, [id](#)</sup>, P. Yin<sup>41, [id](#)</sup>, K. Yorita<sup>169, [id](#)</sup>, S. Younas<sup>27b, [id](#)</sup>, C.J.S. Young<sup>54, [id](#)</sup>, C. Young<sup>144, [id](#)</sup>, Y. Yu<sup>62a, [id](#)</sup>,  
M. Yuan<sup>107, [id](#)</sup>, R. Yuan<sup>62b, [id](#)</sup>, L. Yue<sup>97, [id](#)</sup>, M. Zaazoua<sup>62a, [id](#)</sup>, B. Zabinski<sup>87, [id](#)</sup>, E. Zaid<sup>52, [id](#)</sup>,  
T. Zakareishvili<sup>150b, [id](#)</sup>, N. Zakharchuk<sup>34, [id](#)</sup>, S. Zambito<sup>56, [id](#)</sup>, J.A. Zamora Saa<sup>138d, 138b, [id](#)</sup>, J. Zang<sup>154, [id](#)</sup>,  
D. Zanzi<sup>54, [id](#)</sup>, O. Zaplatilek<sup>133, [id](#)</sup>, C. Zeitnitz<sup>172, [id](#)</sup>, H. Zeng<sup>14a, [id](#)</sup>, J.C. Zeng<sup>163, [id](#)</sup>, D.T. Zenger Jr<sup>26, [id](#)</sup>,  
O. Zenin<sup>37, [id](#)</sup>, T. Ženiš<sup>28a, [id](#)</sup>, S. Zenz<sup>95, [id](#)</sup>, S. Zerradi<sup>35a, [id](#)</sup>, D. Zerwas<sup>66, [id](#)</sup>, M. Zhai<sup>14a, 14e, [id](#)</sup>,  
B. Zhang<sup>14c, [id](#)</sup>, D.F. Zhang<sup>140, [id](#)</sup>, J. Zhang<sup>62b, [id](#)</sup>, J. Zhang<sup>6, [id](#)</sup>, K. Zhang<sup>14a, 14e, [id](#)</sup>, L. Zhang<sup>14c, [id](#)</sup>,  
P. Zhang<sup>14a, 14e, [id](#)</sup>, R. Zhang<sup>171, [id](#)</sup>, S. Zhang<sup>107, [id](#)</sup>, T. Zhang<sup>154, [id](#)</sup>, X. Zhang<sup>62c, [id](#)</sup>, X. Zhang<sup>62b, [id](#)</sup>,  
Y. Zhang<sup>62c, 5, [id](#)</sup>, Y. Zhang<sup>97, [id](#)</sup>, Z. Zhang<sup>17a, [id](#)</sup>, Z. Zhang<sup>66, [id](#)</sup>, H. Zhao<sup>139, [id](#)</sup>, P. Zhao<sup>51, [id](#)</sup>, T. Zhao<sup>62b, [id](#)</sup>,  
Y. Zhao<sup>137, [id](#)</sup>, Z. Zhao<sup>62a, [id](#)</sup>, A. Zhemchugov<sup>38, [id](#)</sup>, K. Zheng<sup>163, [id](#)</sup>, X. Zheng<sup>62a, [id](#)</sup>, Z. Zheng<sup>144, [id](#)</sup>,  
D. Zhong<sup>163, [id](#)</sup>, B. Zhou<sup>107, [id](#)</sup>, H. Zhou<sup>7, [id](#)</sup>, N. Zhou<sup>62c, [id](#)</sup>, Y. Zhou<sup>7, [id](#)</sup>, C.G. Zhu<sup>62b, [id](#)</sup>, J. Zhu<sup>107, [id](#)</sup>,  
Y. Zhu<sup>62c, [id](#)</sup>, Y. Zhu<sup>62a, [id](#)</sup>, X. Zhuang<sup>14a, [id](#)</sup>, K. Zhukov<sup>37, [id](#)</sup>, V. Zhulanov<sup>37, [id](#)</sup>, N.I. Zimine<sup>38, [id](#)</sup>,  
J. Zinsser<sup>63b, [id](#)</sup>, M. Ziolkowski<sup>142, [id](#)</sup>, L. Živković<sup>15, [id](#)</sup>, A. Zoccoli<sup>23b, 23a, [id](#)</sup>, K. Zoch<sup>56, [id](#)</sup>, T.G. Zorbas<sup>140, [id](#)</sup>,  
O. Zormpa<sup>46, [id](#)</sup>, W. Zou<sup>41, [id](#)</sup>, L. Zwalinski<sup>36, [id](#)</sup>,

<sup>1</sup> Department of Physics, University of Adelaide, Adelaide; Australia

<sup>2</sup> Department of Physics, University of Alberta, Edmonton AB; Canada

<sup>3</sup> (a) Department of Physics, Ankara University, Ankara; (b) Division of Physics, TOBB University of Economics and Technology, Ankara; Türkiye

<sup>4</sup> LAPP, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy; France

<sup>5</sup> APC, Université Paris Cité, CNRS/IN2P3, Paris; France

<sup>6</sup> High Energy Physics Division, Argonne National Laboratory, Argonne IL; United States of America

<sup>7</sup> Department of Physics, University of Arizona, Tucson AZ; United States of America

<sup>8</sup> Department of Physics, University of Texas at Arlington, Arlington TX; United States of America

<sup>9</sup> Physics Department, National and Kapodistrian University of Athens, Athens; Greece

<sup>10</sup> Physics Department, National Technical University of Athens, Zografou; Greece

<sup>11</sup> Department of Physics, University of Texas at Austin, Austin TX; United States of America

<sup>12</sup> Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan

<sup>13</sup> Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain

<sup>14</sup> (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Physics Department, Tsinghua University, Beijing; (c) Department of Physics, Nanjing University, Nanjing;

(d) School of Science, Shenzhen Campus of Sun Yat-sen University; (e) University of Chinese Academy of Science (UCAS), Beijing; China

<sup>15</sup> Institute of Physics, University of Belgrade, Belgrade; Serbia

<sup>16</sup> Department for Physics and Technology, University of Bergen, Bergen; Norway

<sup>17</sup> (a) Physics Division, Lawrence Berkeley National Laboratory, Berkeley CA; (b) University of California, Berkeley CA; United States of America

<sup>18</sup> Institut für Physik, Humboldt Universität zu Berlin, Berlin; Germany

<sup>19</sup> Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern; Switzerland

<sup>20</sup> School of Physics and Astronomy, University of Birmingham, Birmingham; United Kingdom

<sup>21</sup> (a) Department of Physics, Bogazici University, Istanbul; (b) Department of Physics Engineering, Gaziantep University, Gaziantep; (c) Department of Physics, Istanbul University, Istanbul; Türkiye

<sup>22</sup> (a) Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá; (b) Departamento de Física, Universidad Nacional de Colombia, Bogotá; (c) Pontificia Universidad Javeriana, Bogotá; Colombia

<sup>23</sup> (a) Dipartimento di Fisica e Astronomia A. Righi, Università di Bologna, Bologna; (b) INFN Sezione di Bologna; Italy

<sup>24</sup> Physikalisches Institut, Universität Bonn, Bonn; Germany

<sup>25</sup> Department of Physics, Boston University, Boston MA; United States of America

<sup>26</sup> Department of Physics, Brandeis University, Waltham MA; United States of America

<sup>27</sup> (a) Transilvania University of Brasov, Brasov; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; (c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; (d) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; (e) University Politehnica Bucharest, Bucharest; (f) West University in Timisoara, Timisoara; (g) Faculty of Physics, University of Bucharest, Bucharest; Romania

<sup>28</sup> (a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice; Slovak Republic

<sup>29</sup> Physics Department, Brookhaven National Laboratory, Upton NY; United States of America

<sup>30</sup> Universidad de Buenos Aires, Facultad de Ciencias Exactas y Naturales, Departamento de Física, y CONICET, Instituto de Física de Buenos Aires (IFIBA), Buenos Aires; Argentina

<sup>31</sup> California State University, CA; United States of America

- 32 Cavendish Laboratory, University of Cambridge, Cambridge; United Kingdom
- 33 <sup>(a)</sup> Department of Physics, University of Cape Town, Cape Town; <sup>(b)</sup> iThemba Labs, Western Cape; <sup>(c)</sup> Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; <sup>(d)</sup> National Institute of Physics, University of the Philippines Diliman (Philippines); <sup>(e)</sup> University of South Africa, Department of Physics, Pretoria; <sup>(f)</sup> University of Zululand, KwaDlangezwa; <sup>(g)</sup> School of Physics, University of the Witwatersrand, Johannesburg; South Africa
- 34 Department of Physics, Carleton University, Ottawa ON; Canada
- 35 <sup>(a)</sup> Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; <sup>(b)</sup> Faculté des Sciences, Université Ibn-Tofail, Kénitra; <sup>(c)</sup> Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; <sup>(d)</sup> LPMR, Faculté des Sciences, Université Mohamed Premier, Oujda; <sup>(e)</sup> Faculté des sciences, Université Mohammed V, Rabat; <sup>(f)</sup> Institute of Applied Physics, Mohammed VI Polytechnic University, Ben Guerir; Morocco
- 36 CERN, Geneva; Switzerland
- 37 Affiliated with an institute covered by a cooperation agreement with CERN
- 38 Affiliated with an international laboratory covered by a cooperation agreement with CERN
- 39 Enrico Fermi Institute, University of Chicago, Chicago IL; United States of America
- 40 LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand; France
- 41 Nevis Laboratory, Columbia University, Irvington NY; United States of America
- 42 Niels Bohr Institute, University of Copenhagen, Copenhagen; Denmark
- 43 <sup>(a)</sup> Dipartimento di Fisica, Università della Calabria, Rende; <sup>(b)</sup> INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; Italy
- 44 Physics Department, Southern Methodist University, Dallas TX; United States of America
- 45 Physics Department, University of Texas at Dallas, Richardson TX; United States of America
- 46 National Centre for Scientific Research "Demokritos", Agia Paraskevi; Greece
- 47 <sup>(a)</sup> Department of Physics, Stockholm University; <sup>(b)</sup> Oskar Klein Centre, Stockholm; Sweden
- 48 Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen; Germany
- 49 Fakultät Physik, Technische Universität Dortmund, Dortmund; Germany
- 50 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden; Germany
- 51 Department of Physics, Duke University, Durham NC; United States of America
- 52 SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh; United Kingdom
- 53 INFN e Laboratori Nazionali di Frascati, Frascati; Italy
- 54 Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany
- 55 II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany
- 56 Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland
- 57 <sup>(a)</sup> Dipartimento di Fisica, Università di Genova, Genova; <sup>(b)</sup> INFN Sezione di Genova; Italy
- 58 II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen; Germany
- 59 SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow; United Kingdom
- 60 LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble; France
- 61 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA; United States of America
- 62 <sup>(a)</sup> Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; <sup>(b)</sup> Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; <sup>(c)</sup> School of Physics and Astronomy, Shanghai Jiao Tong University, Key Laboratory for Particle Astrophysics and Cosmology (MOE), SKLPPC, Shanghai; <sup>(d)</sup> Tsung-Dao Lee Institute, Shanghai; China
- 63 <sup>(a)</sup> Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; <sup>(b)</sup> Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; Germany
- 64 <sup>(a)</sup> Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; <sup>(b)</sup> Department of Physics, University of Hong Kong, Hong Kong; <sup>(c)</sup> Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong; China
- 65 Department of Physics, National Tsing Hua University, Hsinchu; Taiwan
- 66 IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay; France
- 67 Centro Nacional de Microelectrónica (IMB-CNM-CSIC), Barcelona; Spain
- 68 Department of Physics, Indiana University, Bloomington IN; United States of America
- 69 <sup>(a)</sup> INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; <sup>(b)</sup> ICTP, Trieste; <sup>(c)</sup> Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine; Italy
- 70 <sup>(a)</sup> INFN Sezione di Lecce; <sup>(b)</sup> Dipartimento di Matematica e Fisica, Università del Salento, Lecce; Italy
- 71 <sup>(a)</sup> INFN Sezione di Milano; <sup>(b)</sup> Dipartimento di Fisica, Università di Milano, Milano; Italy
- 72 <sup>(a)</sup> INFN Sezione di Napoli; <sup>(b)</sup> Dipartimento di Fisica, Università di Napoli, Napoli; Italy
- 73 <sup>(a)</sup> INFN Sezione di Pavia; <sup>(b)</sup> Dipartimento di Fisica, Università di Pavia, Pavia; Italy
- 74 <sup>(a)</sup> INFN Sezione di Pisa; <sup>(b)</sup> Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy
- 75 <sup>(a)</sup> INFN Sezione di Roma; <sup>(b)</sup> Dipartimento di Fisica, Sapienza Università di Roma, Roma; Italy
- 76 <sup>(a)</sup> INFN Sezione di Roma Tor Vergata; <sup>(b)</sup> Dipartimento di Fisica, Università di Roma Tor Vergata, Roma; Italy
- 77 <sup>(a)</sup> INFN Sezione di Roma Tre; <sup>(b)</sup> Dipartimento di Matematica e Fisica, Università Roma Tre, Roma; Italy
- 78 <sup>(a)</sup> INFN-TIFPA; <sup>(b)</sup> Università degli Studi di Trento, Trento; Italy
- 79 Universität Innsbruck, Department of Astro and Particle Physics, Innsbruck; Austria
- 80 University of Iowa, Iowa City IA; United States of America
- 81 Department of Physics and Astronomy, Iowa State University, Ames IA; United States of America
- 82 Istinye University, Sariyer, Istanbul; Türkiye
- 83 <sup>(a)</sup> Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; <sup>(b)</sup> Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; <sup>(c)</sup> Instituto de Física, Universidade de São Paulo, São Paulo; <sup>(d)</sup> Rio de Janeiro State University, Rio de Janeiro; Brazil
- 84 KEK, High Energy Accelerator Research Organization, Tsukuba; Japan
- 85 Graduate School of Science, Kobe University, Kobe; Japan
- 86 <sup>(a)</sup> AGH University of Krakow, Faculty of Physics and Applied Computer Science, Krakow; <sup>(b)</sup> Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow; Poland
- 87 Institute of Nuclear Physics Polish Academy of Sciences, Krakow; Poland
- 88 Faculty of Science, Kyoto University, Kyoto; Japan
- 89 Kyoto University of Education, Kyoto; Japan
- 90 Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka; Japan
- 91 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata; Argentina
- 92 Physics Department, Lancaster University, Lancaster; United Kingdom
- 93 Oliver Lodge Laboratory, University of Liverpool, Liverpool; United Kingdom
- 94 Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana; Slovenia
- 95 School of Physics and Astronomy, Queen Mary University of London, London; United Kingdom
- 96 Department of Physics, Royal Holloway University of London, Egham; United Kingdom
- 97 Department of Physics and Astronomy, University College London, London; United Kingdom
- 98 Louisiana Tech University, Ruston LA; United States of America
- 99 Fysiska institutionen, Lunds universitet, Lund; Sweden
- 100 Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid; Spain
- 101 Institut für Physik, Universität Mainz, Mainz; Germany
- 102 School of Physics and Astronomy, University of Manchester, Manchester; United Kingdom
- 103 CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France

- 104 Department of Physics, University of Massachusetts, Amherst MA; United States of America  
 105 Department of Physics, McGill University, Montreal QC; Canada  
 106 School of Physics, University of Melbourne, Victoria; Australia  
 107 Department of Physics, University of Michigan, Ann Arbor MI; United States of America  
 108 Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America  
 109 Group of Particle Physics, University of Montreal, Montreal QC; Canada  
 110 Fakultät für Physik, Ludwig-Maximilians-Universität München, München; Germany  
 111 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München; Germany  
 112 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya; Japan  
 113 Department of Physics and Astronomy, University of New Mexico, Albuquerque NM; United States of America  
 114 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University/Nikhef, Nijmegen; Netherlands  
 115 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam; Netherlands  
 116 Department of Physics, Northern Illinois University, DeKalb IL; United States of America  
 117 <sup>(a)</sup> New York University Abu Dhabi, Abu Dhabi; <sup>(b)</sup> University of Sharjah, Sharjah; United Arab Emirates  
 118 Department of Physics, New York University, New York NY; United States of America  
 119 Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo; Japan  
 120 Ohio State University, Columbus OH; United States of America  
 121 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK; United States of America  
 122 Department of Physics, Oklahoma State University, Stillwater OK; United States of America  
 123 Palacký University, Joint Laboratory of Optics, Olomouc; Czech Republic  
 124 Institute for Fundamental Science, University of Oregon, Eugene, OR; United States of America  
 125 Graduate School of Science, Osaka University, Osaka; Japan  
 126 Department of Physics, University of Oslo, Oslo; Norway  
 127 Department of Physics, Oxford University, Oxford; United Kingdom  
 128 LPNHE, Sorbonne Université, Université Paris Cité, CNRS/IN2P3, Paris; France  
 129 Department of Physics, University of Pennsylvania, Philadelphia PA; United States of America  
 130 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA; United States of America  
 131 <sup>(a)</sup> Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa; <sup>(b)</sup> Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; <sup>(c)</sup> Departamento de Física, Universidade de Coimbra, Coimbra; <sup>(d)</sup> Centro de Física Nuclear da Universidade de Lisboa, Lisboa; <sup>(e)</sup> Departamento de Física, Universidade do Minho, Braga; <sup>(f)</sup> Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); <sup>(g)</sup> Departamento de Física, Instituto Superior Técnico, Universidade de Lisboa, Lisboa; Portugal  
 132 Institute of Physics of the Czech Academy of Sciences, Prague; Czech Republic  
 133 Czech Technical University in Prague, Prague; Czech Republic  
 134 Charles University, Faculty of Mathematics and Physics, Prague; Czech Republic  
 135 Particle Physics Department, Rutherford Appleton Laboratory, Didcot; United Kingdom  
 136 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette; France  
 137 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA; United States of America  
 138 <sup>(a)</sup> Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; <sup>(b)</sup> Millennium Institute for Subatomic physics at high energy frontier (SAPHIR), Santiago; <sup>(c)</sup> Instituto de Investigación Multidisciplinario en Ciencia y Tecnología, y Departamento de Física, Universidad de La Serena; <sup>(d)</sup> Universidad Andres Bello, Department of Physics, Santiago; <sup>(e)</sup> Instituto de Alta Investigación, Universidad de Tarapacá, Arica; <sup>(f)</sup> Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso; Chile  
 139 Department of Physics, University of Washington, Seattle WA; United States of America  
 140 Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom  
 141 Department of Physics, Shinshu University, Nagano; Japan  
 142 Department Physik, Universität Siegen, Siegen; Germany  
 143 Department of Physics, Simon Fraser University, Burnaby BC; Canada  
 144 SLAC National Accelerator Laboratory, Stanford CA; United States of America  
 145 Department of Physics, Royal Institute of Technology, Stockholm; Sweden  
 146 Departments of Physics and Astronomy, Stony Brook University, Stony Brook NY; United States of America  
 147 Department of Physics and Astronomy, University of Sussex, Brighton; United Kingdom  
 148 School of Physics, University of Sydney, Sydney; Australia  
 149 Institute of Physics, Academia Sinica, Taipei; Taiwan  
 150 <sup>(a)</sup> E. Andronikashvili Institute of Physics, Iv. Javakishvili Tbilisi State University, Tbilisi; <sup>(b)</sup> High Energy Physics Institute, Tbilisi State University, Tbilisi; <sup>(c)</sup> University of Georgia, Tbilisi; Georgia  
 151 Department of Physics, Technion, Israel Institute of Technology, Haifa; Israel  
 152 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv; Israel  
 153 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki; Greece  
 154 International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo; Japan  
 155 Department of Physics, Tokyo Institute of Technology, Tokyo; Japan  
 156 Department of Physics, University of Toronto, Toronto ON; Canada  
 157 <sup>(a)</sup> TRIUMF, Vancouver BC; <sup>(b)</sup> Department of Physics and Astronomy, York University, Toronto ON; Canada  
 158 Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba; Japan  
 159 Department of Physics and Astronomy, Tufts University, Medford MA; United States of America  
 160 United Arab Emirates University, Al Ain; United Arab Emirates  
 161 Department of Physics and Astronomy, University of California Irvine, Irvine CA; United States of America  
 162 Department of Physics and Astronomy, University of Uppsala, Uppsala; Sweden  
 163 Department of Physics, University of Illinois, Urbana IL; United States of America  
 164 Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Valencia; Spain  
 165 Department of Physics, University of British Columbia, Vancouver BC; Canada  
 166 Department of Physics and Astronomy, University of Victoria, Victoria BC; Canada  
 167 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg; Germany  
 168 Department of Physics, University of Warwick, Coventry; United Kingdom  
 169 Waseda University, Tokyo; Japan  
 170 Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot; Israel  
 171 Department of Physics, University of Wisconsin, Madison WI; United States of America  
 172 Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal; Germany  
 173 Department of Physics, Yale University, New Haven CT; United States of America

<sup>a</sup> Also Affiliated with an institute covered by a cooperation agreement with CERN.

<sup>b</sup> Also at An-Najah National University, Nablus; Palestine.

<sup>c</sup> Also at APC, Université Paris Cité, CNRS/IN2P3, Paris; France.

- <sup>d</sup> Also at Borough of Manhattan Community College, City University of New York, New York NY; United States of America.
- <sup>e</sup> Also at Center for High Energy Physics, Peking University; China.
- <sup>f</sup> Also at Center for Interdisciplinary Research and Innovation (CIRI-AUTH), Thessaloniki; Greece.
- <sup>g</sup> Also at Centro Studi e Ricerche Enrico Fermi; Italy.
- <sup>h</sup> Also at CERN, Geneva; Switzerland.
- <sup>i</sup> Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
- <sup>j</sup> Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.
- <sup>k</sup> Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.
- <sup>l</sup> Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America.
- <sup>m</sup> Also at Department of Physics and Astronomy, University of Victoria, Victoria BC; Canada.
- <sup>n</sup> Also at Department of Physics, Ben Gurion University of the Negev, Beer Sheva; Israel.
- <sup>o</sup> Also at Department of Physics, California State University, Sacramento; United States of America.
- <sup>p</sup> Also at Department of Physics, King's College London, London; United Kingdom.
- <sup>q</sup> Also at Department of Physics, Royal Holloway University of London, Egham; United Kingdom.
- <sup>r</sup> Also at Department of Physics, Stanford University, Stanford CA; United States of America.
- <sup>s</sup> Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.
- <sup>t</sup> Also at Department of Physics, University of Thessaly; Greece.
- <sup>u</sup> Also at Department of Physics, Westmont College, Santa Barbara; United States of America.
- <sup>v</sup> Also at Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal; Germany.
- <sup>w</sup> Also at Hellenic Open University, Patras; Greece.
- <sup>x</sup> Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain.
- <sup>y</sup> Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.
- <sup>z</sup> Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia; Bulgaria.
- <sup>aa</sup> Also at Institute of Applied Physics, Mohammed VI Polytechnic University, Ben Guerir; Morocco.
- <sup>ab</sup> Also at Institute of Particle Physics (IPP); Canada.
- <sup>ac</sup> Also at Institute of Physics and Technology, Ulaanbaatar; Mongolia.
- <sup>ad</sup> Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
- <sup>ae</sup> Also at Institute of Theoretical Physics, Ilija State University, Tbilisi; Georgia.
- <sup>af</sup> Also at L2IT, Université de Toulouse, CNRS/IN2P3, UPS, Toulouse; France.
- <sup>ag</sup> Also at Lawrence Livermore National Laboratory, Livermore; United States of America.
- <sup>ah</sup> Also at National Institute of Physics, University of the Philippines Diliman (Philippines); Philippines.
- <sup>ai</sup> Also at Technical University of Munich, Munich; Germany.
- <sup>aj</sup> Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China.
- <sup>ak</sup> Also at TRIUMF, Vancouver BC; Canada.
- <sup>al</sup> Also at Università di Napoli Parthenope, Napoli; Italy.
- <sup>am</sup> Also at University of Colorado Boulder, Department of Physics, Colorado; United States of America.
- <sup>an</sup> Also at Washington College, Chestertown, MD; United States of America.
- <sup>ao</sup> Also at Yeditepe University, Physics Department, Istanbul; Türkiye.
- \* Deceased.