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# A Hydrogen-Initiated Chemical Epitaxial Growth Strategy for In-Plane Heterostructured Photocatalyst

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**ABSTRACT:** Integrating carbon nitride with graphene into a lateral heterojunction would avoid the energy loss within the interlaminar space region on conventional composites. Up to now, its synthesis process is limited to the bottom-up method which lacks the targeting and homogeneity. Herein, we proposed a hydrogen-initiated chemical epitaxial growth strategy at a relatively low temperature for the fabrication of graphene/carbon nitride in-plane heterostructure. Theoretical and experimental analysis proved that methane *via in situ* generation from the hydrogenated decomposition of carbon nitride triggered the graphene growth along the active sites at the edges of confined spaces. With the enhanced electrical field from the deposited graphene (0.5 %), the performances on selective photooxidation and photocatalytic water splitting were promoted by 5.5 and 3.7 times, respectively. Meanwhile, a 7720  $\mu$ mol/h/g(graphene) hydrogen evolution rate was acquired without any co-catalysts. This study provides an up-bottom strategy to synthesize in-plane catalyst for the utilization of solar energy.

**KEYWORDS:** hydrogen-initiated chemical epitaxial growth, carbon nitride/graphene, intralayer heterojunction, photocatalytic water splitting, photooxidation

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Carbon nitride photocatalyst with the advantages of metal-free and visible light response has been widely utilized into energy preparation,<sup>1</sup> pharmaceutical synthesis<sup>2</sup> and environmental governance.<sup>3</sup> However, intrinsic carbon nitride typically suffers from inefficient generation of hot carriers, which results in low photocatalytic performance, as a result of the weak driving force on photo induced excitons. Hybridization with graphene has been regarded as an effective approach for the modification of carbon nitride without extra addition of metal elements.<sup>4</sup> With a high value of charge mobility ( $\mu \approx 10,000 \text{ cm}^2/\text{Vs}$ ), graphene acts as an electron mediator and diverts electrons once excited, realizing the faster separation and transfer of charge carriers.<sup>5,6</sup> Nevertheless, the interlaminar space region formed in the fabricated heterojunction provides rooms for the relaxation of electron-hole pairs.<sup>7</sup> Therefore, spaces on improving the photocatalytic performances of graphene/carbon nitride heterojunctions still exist.

Inosculating carbon nitride and graphene within in-plane interfaces might hold the promise to shorten the transportation path of photo induced charge carriers in comparison of passing through the interlaminar space region in van der Waals heterojunctions.<sup>8-10</sup> Thus this more rational photocatalyst in theory has always been a pursuit in the field of materials science. Based on this, Che *et al.* employed a bottom-up method to obtain (C<sub>ring</sub>)-C<sub>3</sub>N<sub>4</sub> plane heterostructural nanosheets.<sup>11</sup> Despite of enhanced performance on this in-plane heterostructure, the targeting and homogeneity of formed graphene on carbon nitride substrate cannot be guaranteed. Current reasonable strategies for the fabrication of uniform intralayer heterostructures are mainly dependent on the chemical epitaxial growth (CEG) or physical vapour deposition (PVD) strategies.<sup>8,10,12</sup> However, it is rarely reported on the fabrication of graphene/organic semiconductor intralayer heterojunctions. On one side, the requirement of tremendous compatibility on the lattice constants and lattice symmetry from each component of intralayer heterostructures pose difficulties on its construction and on the other side processing conditions such as temperature for both sides' stabilization are extremely uncontrollable.<sup>9,10</sup> More seriously, these requirements can not always be achieved at the same time.

Herein, we proposed a hydrogen-initiated chemical epitaxial growth strategy at a relatively low temperature for the fabrication of graphene/carbon nitride intralayer photocatalyst. Our proof-of-concept demonstration is initially on the basis of the exceptionally similar molecular configuration on carbon nitride and graphene. As demonstrated in Figure S1, both lattices of this 'wistedtwins' have a hexagonal symmetry with similar space group, and the difference in each lattice parameter is only 1.4 %, proving the feasibility of material construction. Moreover, considering the pyrolysis temperature of carbon nitride (starting from 600 °C) is far below the temperature for graphene formation (above 1000 °C using CVD method),<sup>13,14</sup> confined spaces and active sites, which are theoretical and experimental evidenced to be indispensable on graphene deposition at relatively low temperature, were created on carbon nitride substrate. After connecting the carbon rings of graphene with the tri-s-triazine units within carbon nitride, the acquired in-plane heterojunction was found to exhibit enhanced electrical field and a more favourable electronic structure for facilitating carriers' separation and transfer. With the trace amount of deposited graphene (0.5 %), the performances on selective photooxidation and photocatalytic water splitting were promoted to be 5.5 and 3.7 times, respectively. In particular, this in-plane heterostructure exhibited a 7720 µmol h<sup>-1</sup> g<sup>-1</sup>(graphene) hydrogen evolution rate without noble metal as co-catalyst.

## **RESULTS AND DISCUSSION**

The fabrication process of graphene/carbon nitride intralayer heterojunction includes two processes (Figure 1), *i.e.* the creation of confined spaces with active sites at edges on bulk structure and the  $H_2$  initiated CEG procedure for graphene formation. At first, a soft template of sodium bicarbonate was added into a melamine solution followed by heating, another molecule (Figure S2) was synthesized and subsequently heated with unreacted sodium bicarbonate and melamine. During the polymerization process, apparent pore channels were created on 10%-C<sub>3</sub>N<sub>4</sub> and 14%-C<sub>3</sub>N<sub>4</sub> (also named as porous carbon nitride) *via* the released

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CO<sub>2</sub> from the pyrolysis of sodium bicarbonate,<sup>15</sup> which could be confirmed from the disappeared (100) peak from XRD profiles (Figure S3) and the appeared pores from electron and atomic force microscopy images (Figure S4-6). Meanwhile, as the intercalating effect of Na<sup>+</sup> from sodium bicarbonate,<sup>16</sup> a gradual decline can be observed in the (002) peak of porous carbon nitride with increasing proportions of soft template (Figure S3), which is indicative of breaking of van der Waals bonding within the layers of bulk carbon nitride.<sup>17</sup> Since the generated molecules in the precursor contain acylamino groups at edges (Figure S2), which could act as a polymerization inhibitor in the polymerization process of carbon nitride. Therefore, a large number of small and thin carbon nitride fragments with thickness less than 1 nm appeared (Figure S4c and 4g). These fragments can be gradually hydrogenated into methane at 500 °C in H<sub>2</sub> atmosphere, which cannot otherwise be achieved in N<sub>2</sub> (Figure S7). It is known that methane is the carbon gas source in the CEG process for graphene synthesis. Thus the porous carbon nitride was thermally treated in hydrogen at 500 °C. After the hydrogenation process, small carbon nitride fragments disappeared and the thickness of H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> and H<sub>2</sub>-14%-C<sub>3</sub>N<sub>4</sub> (also named as hydrogenated carbon nitride) exhibited an improved tendency compared to that of porous carbon nitride (Figure S4d, S4h and S8). Same phenomenon can also be observed on TEM and SEM images (Figure S9-10), especially on H<sub>2</sub>-14%-C<sub>3</sub>N<sub>4</sub> (Figure S9d), a partial hydrogenated state that some pores were not completely patched was recorded. Therefore, we reckon that the decomposition of the small fragments in hydrogen released methane, which then *in situ* formed graphene along the edge of pores (confined spaces) to produce a unique intralayer graphene/carbon nitride heterostructure. On this occasion, the volumes of three samples, g-C<sub>3</sub>N<sub>4</sub>, 10%-C<sub>3</sub>N<sub>4</sub>, and H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub>, at the same mass were compared (Figure S10f), revealing that the column volume of 10%-C<sub>3</sub>N<sub>4</sub> significantly increased because of the exfoliation and porosity. However, the volume of H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> decreased to almost the same as that of g-C<sub>3</sub>N<sub>4</sub>, indicating the lose of porosity due to the graphene filling. Same conclusion can also be reached from the BET results (Figure S11). By contrast, with less

pore maker of NaHCO<sub>3</sub> in the first step of Figure 1, no obvious damage can be observed on 2%- $C_3N_4$  and 6%- $C_3N_4$ , which still remained the bulk structure (Figure S5-6). While after the calcination under H<sub>2</sub> atmosphere, the (002) peak intensity of both H<sub>2</sub>-2%- $C_3N_4$  and H<sub>2</sub>-6%- $C_3N_4$  (Figure S8) experienced further decrease and their surface areas increased by two and three times, respectively, as compared with that of g- $C_3N_4$  (Figure S11 and Table S1). This confirms the further exfoliation, rather than the formation of graphene, occurred on 2%- $C_3N_4$  and 6%- $C_3N_4$  without confined spaces. Therefore, porous carbon nitride (10%- $C_3N_4$  and 14%- $C_3N_4$ ) and hydrogenated carbon nitride (H<sub>2</sub>-10%- $C_3N_4$  and H<sub>2</sub>-14%- $C_3N_4$ ) will be emphatically discussed in this work.

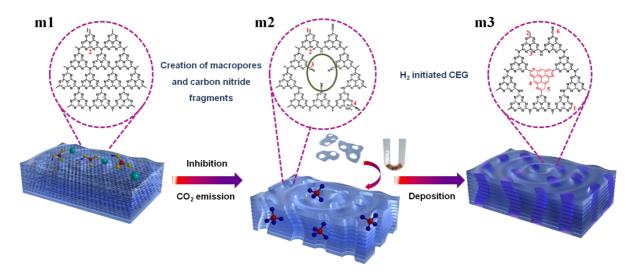


Figure 1. Schematic illustration of the formation processes of intralayer heterojunctions. m1, m2 and m3 are the proposed molecular structures of g-C<sub>3</sub>N<sub>4</sub>, porous carbon nitride and hydrogenated carbon nitride, respectively.

The microscopic changes during this modified CEG process are initially supported by the NMR spectra (Figure 2a): on pristine  $g-C_3N_4$ , two main peaks at 165 and 155 ppm are evident, corresponding to the carbon atoms at position 1 and the position 2 in m1, respectively (Figure 1). In addition to these two main peaks, there is an obvious peak centered at approximately 40 ppm on porous carbon nitride, which is attributed to the CH<sub>2</sub> group (position 4 in m2), indicating the tri-s-triazine unit of the porous carbon nitride is destroyed. Besides, acylamino groups in

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the precursor could convert into cyano groups because of the dehydration effect during the polymerization (130 ppm in Figure 2a).<sup>18</sup> As a result, this damage process forms one kind of intralayer heterostructure as shown in m2. This structure is also reflected by the increased value in Zeta potential of porous carbon nitride as the cyano groups have a strong electronwithdrawing ability (Table S1). Whereas after the hydrogenation process, the signal of the CH<sub>2</sub> group on porous carbon nitride disappeared, with a peak emerging at 171 ppm, assigned to the methine carbon at position 1, structure m3.<sup>19</sup> This suggests that the damaged tri-s-triazine unit is repaired. Noticeably, an peak at 137 ppm associated with sp<sup>2</sup> carbon also emerged, consistent with the formation of graphene during the hydrogen treatment.<sup>20</sup> As a result, another intralayer heteojunctions connected by the in situ generated graphene are produced. Apart from this, a peak at 32.5 ppm was also found on H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub>, which originates from the sp<sup>3</sup> carbon (position 5 in m3). XPS and FTIR studies were also performed to elucidate the changes of chemical states and surface composition of the prepared samples. As observed from C1s spectra (Figure S12), the peak at 285.4 eV assigned to amino groups on  $g-C_3N_4^{21}$  upshift to 286 eV on 14%-C<sub>3</sub>N<sub>4</sub> after the damage process. Meanwhile the peak intensity gradually increases from 2.64 % on g-C<sub>3</sub>N<sub>4</sub> to 6.18 % on 14%-C<sub>3</sub>N<sub>4</sub>, suggesting that the marginal amino groups of g-C<sub>3</sub>N<sub>4</sub> gradually changed to cyano groups with the increasing amount of soft template, since cyano groups have the similar C1s binding energy to that of amino groups.<sup>18</sup> This can also be proven by the FTIR spectra (Figure S13), where the peak intensity of cyano groups at 2173 cm<sup>-1</sup> turns stronger and that of amino groups at 3090 cm<sup>-1</sup> becomes lower with more soft template being added in the damage process.<sup>22</sup> Besides, it was interesting to note that the ratio between C=N and C=C bonds remained unchanged on all of the products after the damage process (Figure S12b), suggesting that no obvious changes occurred on the tri-s-triazine unit of the basal plane of carbon nitride. While on the hydrogenated samples, the ratio of C=C to C=N bonds rose from 0.23 in g-C<sub>3</sub>N<sub>4</sub> to 0.36 in H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub>, verifying the growth of graphene

(Figure 2b and S14). Based on the NMR, XPS and FT-IR results, the possible molecular structures of the prepared samples are proposed in Figure 1.

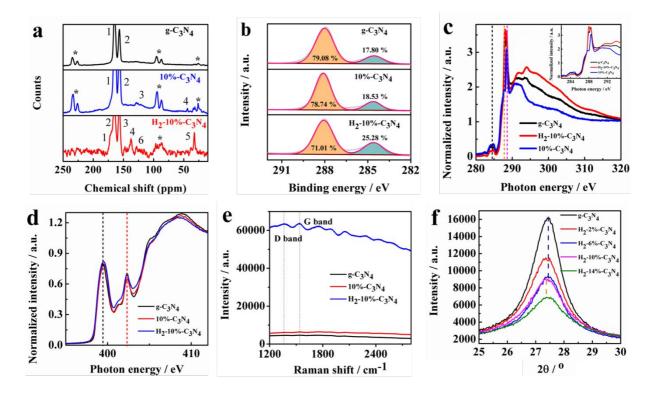


Figure 2. Microscopic evidences of the formation of graphene/carbon nitride intralayer heterojunction. (a) <sup>13</sup>C solid-state NMR spectra. (b) C1s XPS spectra. (c-d) NEXAFS results. (e) Raman spectra and (f) Magnified XRD patterns of as-prepared samples.

The presence of graphene in the hydrogenated carbon nitride is confirmed by synchrotronbased near edge X-ray absorption fine structure (NEXAFS) spectroscopy measurements. All the samples show peaks at 284.5, 288 and 288.4 eV in C K-edge spectra (Figure 2c) which are the characteristic resonances of the defective carbon,  $\pi^*_{C=C}$  and  $\pi^*_{C=N-C}$ , respectively.<sup>4</sup> For the peak at 284.5 eV, the order of intensities from low to high is H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub>, g-C<sub>3</sub>N<sub>4</sub> and 10%-C<sub>3</sub>N<sub>4</sub>, indicating that more carbon defects existed in the porous carbon nitride, and following hydrogenation, some of the defects were repaired. For the peak at 288 eV, small protuberances could be observed in g-C<sub>3</sub>N<sub>4</sub> and 10%-C<sub>3</sub>N<sub>4</sub> whereas H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> exhibited a single distinct peak, which demonstrates the successful deposition of the intraformational carbon ring in the basal plane of carbon nitride. In the N K-edge spectra (Figure 2d), two obvious peaks at 399.4

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and 402.3 eV can be observed which are respectively the  $\pi^*$  resonances of C=N-C and C-N bonding. The reduced peak intensity at 402.3 eV on H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> also confirmed that the N-3C bonding were substituted by the C=C bonding from graphene, which is consistent with the C K-edge results. The same conclusion can also be drawn from Raman spectroscopy. As shown in Figure 2e, distinct peaks of D (1370 cm<sup>-1</sup>) and G (1540 cm<sup>-1</sup>) bands can be seen on H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> which are not evident in the graphitic carbon nitride and 10%-C<sub>3</sub>N<sub>4</sub> samples, indicating the existence of graphitic carbon and defective carbon (mainly N doped carbon).<sup>23</sup> Besides, in XRD profiles, although the intensity of the (002) peaks in hydrogenated carbon nitride exhibited increased tendency compared with that of porous carbon nitride (Figure S8), the (100) peaks also fail to re-emerge, confirming the absence of the CN polymer in the filled macropores of the porous carbon nitride. More importantly, the (002) peaks of H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> and H<sub>2</sub>-14%-C<sub>3</sub>N<sub>4</sub> shift slightly from 27.4° in g-C<sub>3</sub>N<sub>4</sub> to 27.3° (Figure 2f), indicating the existence of carbon in the inplanar of carbon nitride.<sup>24</sup> These combined results clearly show the derivation of g-C<sub>3</sub>N<sub>4</sub> structure in the modified CEG processes, in which graphene grows along the edge plane of porous carbon nitride from methane *via* hydrogenation of carbon nitride fragments.

Furthermore, different from the porous sample (Figure S15), uneven distributions of C and N elements on  $H_2$ -10%- $C_3N_4$  are evident, and carbon-rich areas are easily discernable (Figure 3). Correspondingly, in EDX spectra some areas of  $H_2$ -10%- $C_3N_4$  obey the same C/N ratio as that in the porous carbon nitride and some areas contain less N. The high-resolution TEM images (Figure 3c and Figure S16) confirm an interlayer crystal lattice in hydrogenated carbon nitride which locates within the position of macro-pore in porous carbon nitride. The

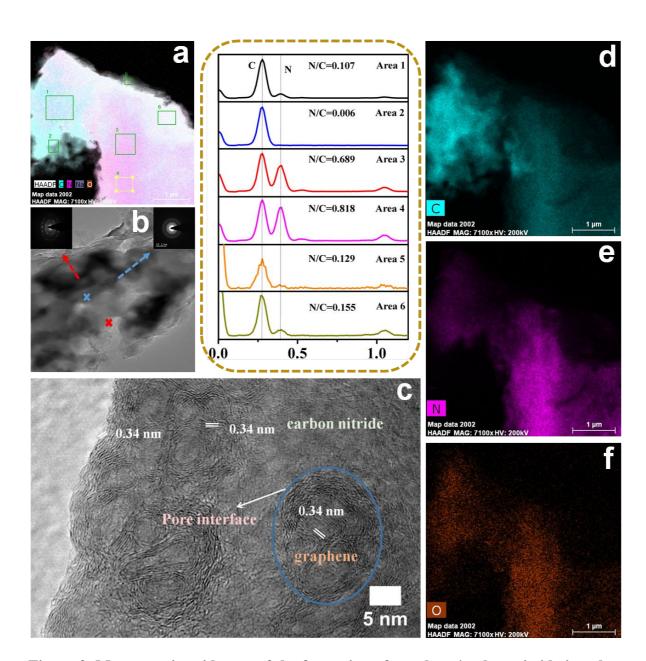


Figure 3. Macroscopic evidences of the formation of graphene/carbon nitride intralayer heterojunction. (a) High angle annular dark field scanning TEM images (HAADF-STEM) and corresponding EDX spectra of six areas of hydrogenated carbon nitride, different C/N ratio in EDX spectra indicated the filling in macropores of porous carbon nitride was graphene. (b-c) HRTEM image of reconstructed carbon nitride and selected-area electron diffraction images of two different areas. (d-f) Energy-dispersive X-ray spectroscopy (EDX) elemental mapping images of hydrogenated carbon nitride.

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lattice spacing is 3.40 Å<sup>12</sup> and obvious boundaries (blue cycle in Figure 3c) can be found between carbon nitride and graphene, indicating that the carbon grows along the cyclic edge of macropores on the porous carbon nitride. In addition, the selected-area electron diffraction image (inset of Figure 3b) proves the presence of the in-planar order of carbon and amorphous carbon nitride. These results further confirm that the uniform graphene/carbon nitride intralayer heterojunction was obtained.

Graphitic carbon nitride (bulk nanostructure), 2%-C<sub>3</sub>N<sub>4</sub> and 6%-C<sub>3</sub>N<sub>4</sub> (with bulk structure and cyano groups at edge) and urea derived carbon nitride (porous without cyano groups) were selected as contrast samples but failed to obtain graphene (Figure S17). Therefore, nano-pores (confined spaces) and cyano groups at edge (active site) together are of crucial factors for the graphene deposition at relatively low temperatures. The reaction rates on a planar solid surface in open space, and that on the solid walls within a nano-pores (Figure 4a) were compared. Based on the collision theory shown in equation ,<sup>25</sup> the reaction rate of graphene growth is directly proportional to the collision frequency.

$$r(T) = Z_{\rho} exp\left(\frac{-E_a}{RT}\right) \qquad \qquad \text{Eq. 1}$$

where Z is the collision frequency, Ea represents activation energy, R and T are constant and temperature, respectively. In our case the reaction is between a diluted gaseous phase and the reaction sites containing reactant B embedded in a solid surface. Considering a diluted gaseous reactant A moved within the vicinity (one mean free path distance to) of a solid surface where reactant B sites were embedded within a solid planar wall in open space, with all A molecules moving perpendicularly towards the surface. A collision between the A molecule and the wall will be likely to happen when the molecule is within the range of a mean free path to the wall, which will result in a reaction when a type B molecule is simultaneously within a sphere of diameter  $r_{cs} = r_1 + r_2$  along the trajectory of the A molecule, where the radius of molecule A and B are respectively  $r_1$  and  $r_2$  and  $r_{cs}$  represents the radius of the collision cross section

(Figure 4a). Within the time interval  $\Delta t$  the distance travelled by the molecule A is  $v_{mean}\Delta t$ , while the total number of A molecules enclosed

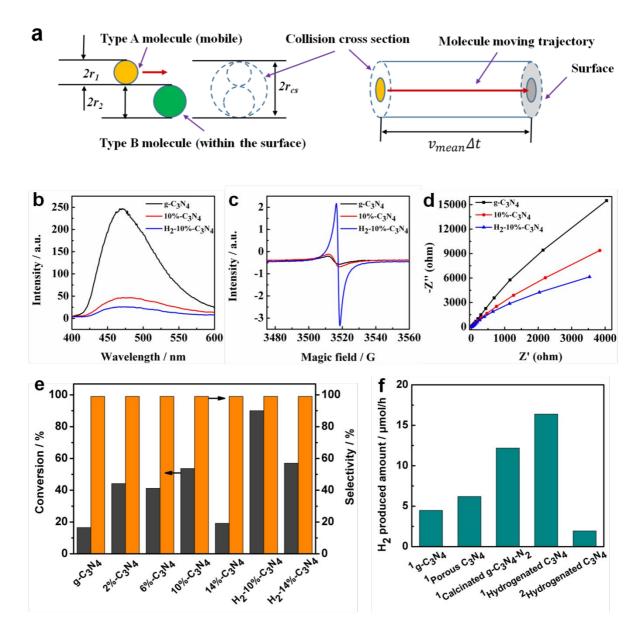


Figure 4. (a) A schematic graph of the collision process between a type A and a type B molecule, with collision cross section and path way distance demonstrated. (b) Fluorescence spectrum (PL). (c) Solid ESR spectra. (d) EIS spectra. (e) Selective photooxidation performances and (f) Photocatalytic performances on water splitting using the prepared samples. 1. Photocatalytic reaction conducted with > 420 nm filter and Pt as co-catalyst, and 2. without the filter and co-catalyst.

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within the pathway of the cross section is  $n_{\pm x} = Nv_{mean}\Delta t\pi r_{cs}^2/6$ , where *N* is the number of molecule A per volume, and  $v_{mean}$  represents the mean of the Maxwell-Boltzmann velocity distribution of type A molecules. The reason for a denominator of 6 being included is that, the molecules movement in a 3D space may be decomposed with respect to a Cartesian coordinate system along 6 directions, *e.g.*,  $\pm x$ ,  $\pm y$ , *and*  $\pm z$ , and that for a medium in equilibrium the number of molecules moving along the  $\pm x$  direction will be equal.

The collision frequency on a planar wall in open space can now be calculated as the ratio of the total number of the type A molecules moving towards the wall over the length of the pass way, *i.e.*,

$$Z_{wall} = \frac{n_{+x}}{v_{mean}\Delta t} = \frac{N v_{mean}\Delta t \pi r_{cs}^2/6}{v_{mean}\Delta t} = \frac{N}{6}\pi r_{cs}^2 \qquad \text{Eq. 2}$$

In contrast, with a nano-pore with a diameter, D, of the same order of magnitude with the mean free path  $\lambda$  of the type A molecules, the situation is dramatically different even assuming other conditions are unchanged. Here the geometrical constrain plays a key role, which affects the molecules moving along 4 out of the 6 directions ( $\pm x$  and  $\pm y$ ). In this case a fraction of the molecules currently moving in these 4 directions will strike into the wall of the nano-pore within a single cycle along the path way before the next collision occurs. A comparative illustration between a large pore (or a planar surface in open space) and a nano-pore is given in Figure S18a and 18b. The quantitative expression of the collision frequency with a nano-pore can then be given below,

$$Z_{pore} = \frac{n_{+x} + n_{-x} + n_{+y} + n_{-y}}{v_{mean}\Delta t} = \frac{4Nv_{mean}\Delta t\pi r_{cs}^2/6}{v_{mean}\Delta t} = \frac{2N}{3}\pi r_{cs}^2 \quad \text{Eq. 3}$$

Herein, a four-fold increase in the collision frequency, hence also in the reaction rate, would be expected. In the case of this study, the typical nano-pore diameter is  $\sim$ 50 nm, which is compared to the molecule mean free path in the air under standard pressure of 58 nm.<sup>26</sup> Then, a significant increase in both collision frequency and the reaction rate inside the nano-pore, compared to that on a planar surface in open space, is expected. Besides, it has been reported that the dissociation

energy for a H-N bond is 386 kJ/mol,<sup>27</sup> while the energy required for break one of the three bond is 355.3 kJ/mol,<sup>28</sup> with the latter being about 10 % lower than the former. Therefore, the cyano groups at the edge of nano-pores would also provide deposition sites for carbon free radical. Totally, combined the advantages from confined spaces and cyano groups, intralayered graphene/carbon nitride heterojunction *via* CEG process at low temperatures could be successfully obtained.

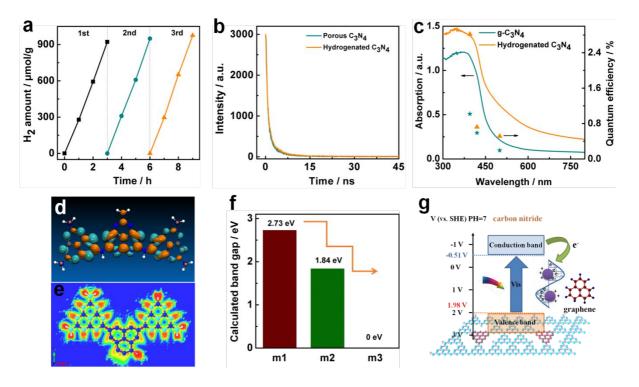


Figure 5. (a) Stability of hydrogenated carbon nitride in photocatalytic water splitting reaction with three cycles. (b) Time-resolved decay spectra of samples before and after the graphene deposition. (c) Wavelength dependent of quantum efficiency of  $g-C_3N_4$  and hydrogenated  $C_3N_4$ . (d) DFT calculations of modules  $g-C_3N_4$ . (e) Electron localization function of hydrogenated carbon nitride, blue, white and cyan atoms represent carbon, hydrogen and nitrogen elements, respectively. (f) Calculated band gap of corresponding modellings of  $g-C_3N_4$ , porous carbon nitride as well as hydrogenated carbon nitride and (g) Diagram of electron transportation path in graphene/carbon nitride intralayered heterojunction.

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After inserting nano-graphene into the lateral plane of carbon nitride, the light response ability was improved. As shown in UV-Vis spectra (Figure S19a), g-C<sub>3</sub>N<sub>4</sub> and 10%-C<sub>3</sub>N<sub>4</sub> are impeded by the weak light harvesting with almost no light response when the wavelength is beyond 500 nm, while on H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> obvious improvement could be observed in the full visible region, even in NIR region this intralayed heterojunction still has response. This enhancement is precisely attributed to the plasmonic effect of incorporated nano-graphene.<sup>29</sup> On this occasion, a photo-to-thermal experiment under inert atmosphere was conducted to prove the effect of added plasmonic nano-graphene (Figure S19b). The temperature of g-C<sub>3</sub>N<sub>4</sub> can increase to 97 °C within 30 min light irradiation mostly due to the heat dissipation from the recombination of photoinduced charge carriers. While after adding trace amount (0.5 %) of nano-graphene into the basal plane of carbon nitride (Table S2), the temperature is sharply increased to 118 °C, which is comparable to that of graphene with same mass. Excluding the heat generated from the the recombination of photo-induced electron-hole pairs on carbon nitride, the extra elevated temperature on the hydrogenated carbon nitride (more than 20 °C) confirms the enhanced electrical field on the graphene/carbon nitride heterojunction.<sup>30,31</sup> Moreover, the enhancement after the addition of graphene into the confined spaces of carbon nitride was also verified using the finite element method. Embedding graphene nano-columns into the confined spaces of porous carbon nitride has a dramatic impact over the original homogenous electrical field (E) distribution (Figure S19 and Figure S20). High field concentration is observed around the edge of the graphene nano-column. The electrical field is plotted along a horizontal line, drawn around the top right corner of the graphene columns where the local field maximum is found. A comparison is made between local intensity at the maximum on the heterostructure and that on the plain g-C<sub>3</sub>N<sub>4</sub> with the ratio  $I/I_0 = (E/E_0)^2$  being plotted in Figure S19c. The enhancement ripples the full spectrum and is wavelength dependent. This observation is similar to the situation when metal nano-particals (NPs) are setting on semiconductor substrate.32

The separation rate of electron-hole pairs of carbon nitride is also boosted as a result of surrounded electrical field from plasmonic nano-graphene. Indeed, the intensity of the PL peak at 470 nm on porous carbon nitride reduced (Figure 4b and Figure S21) and the intensity of the ESR characteristic peak of the unpaired electrons on the carbon nitride units rose (Figure 4c), both indicating the recombination of photo-induced charge carriers was effectively prohibited after the addition of nano-graphene.<sup>33</sup> Furthermore, the heterojunction exhibited a less interfacial electron resistance due to the high electroconductibility of graphene (Figure 4d).

With suitable band structures (Figure S22-24 and Table S1) and enhanced electrical field, the as-obtained carbon nitride sample was then evaluated for photoinduced selective oxidation 1,4-dihydro-2,6-dimethylpyridine-3,5-dicarboxylate (1,4-DHP) (Figure S25-27).<sup>22,34</sup> of Pristine carbon nitride exerted a low conversion of 1,4-DHP (Figure 4e and Table S3) due to the weak light response ability, and constraint of the Coulomb force for charge carriers. Once the bulk carbon nitride was damaged to create the intralayer heterojunction at the edge plane of macropores, the photooxidation rate of 1,4-DHP increased by almost three-fold on 10%-C<sub>3</sub>N<sub>4</sub> with a conversion rate of 53.7 %. The conversion rate for photooxidation of 1,4-DHP experienced a significant enhancement after introducing only 0.5 % plasmonic nano-graphene onto carbon nitride intralayer, with 5.5 times improvement on H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> compared to g- $C_3N_4$ . The prepared samples were also used for photocatalytic hydrogen evolution and 3.7 times activity was acquired on hydrogenated carbon nitride than that of g-C<sub>3</sub>N<sub>4</sub> (Figure 4f). Meanwhile, hydrogenated carbon nitride showed a satisfied stability without any decrease on photocatalytic hydrogen evolution ability after three cycle (Figure 5a). In particular, the metalfree photocatalyst still exhibited photocatalytic ability on water splitting in the absence of cocatalyst, with hydrogen evolution rate of 7720 µmol/h/ggraphene under light irradiation without >420 nm filter, which is higher than reported carbon/carbon nitride composites (Table S4). As re-sintering of carbon nitride resulted in the interlayer stripping, the influence of heating under N<sub>2</sub> rather than hydrogen atmosphere on the carbon nitride was also studied. The

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photocatalytic hydrogen evolution rate of calcinated g-C<sub>3</sub>N<sub>4</sub>-N<sub>2</sub> was 12.2  $\mu$ mol/h, which was lower than that of hydrogenated carbon nitride (16.4  $\mu$ mol/h) (Figure 4f), suggesting that the self-assembly of nano-graphene into carbon nitride is more efficient for the catalyst modification. Furthermore, the time-resolved decay spectra of samples before and after the graphene deposition were obtained for confirming the enhanced performance. As shown in Figure 5b and Table S5, the average lifetime of hydrogenated carbon nitride was 2.14 ns, which was longer than the sample before the graphene deposition (1.99 ns). Meanwhile, a higher visible light irradiated photocurrent was obtained in H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub> than g-C<sub>3</sub>N<sub>4</sub> and 10%-C<sub>3</sub>N<sub>4</sub>, further demonstrating the higher efficiency in the separation of photo-induced charge carriers on the advanced metal-free photocatalyst (Figure S28). As a result, the quantum efficiency of hydrogenated carbon nitride at 395, 420 and 500 nm were all higher than that of g-C<sub>3</sub>N<sub>4</sub> (Figure 5c).

The improvement on light absorption and photocatalytic performances after the graphene formation was proven by DFT calculations (Figure 5d-f), in which part of m1 (Figure 5d), m2 (Figure S29b) and m3 (Figure 5e) molecules are employed as they represent the characteristics of the corresponding samples and are able to affect the optical and electrochemical properties of the relevant intact structures. The porous carbon nitride and hydrogenated carbon nitride presented lowered band gap energies by comparison to g-C<sub>3</sub>N<sub>4</sub>. In particular, the m3 structure exhibited partial metallicity with band gap 0 eV, demonstrating the enhanced optical property of H<sub>2</sub>-10%-C<sub>3</sub>N<sub>4</sub>. To unveil the effect of the fabricated heterostructure on the electronic structure of carbon nitride, frontier molecular orbitals were examined (Figure 5 and Figure S29). With the g-C<sub>3</sub>N<sub>4</sub> system, the LUMO is mainly consisted of C–N bonding orbitals and HOMO is derived from the heterocyclic nitrogen P<sub>z</sub> orbitals.<sup>34</sup> On the porous carbon nitride system, the LUMO shifts from the melem to the damaged unit, while the HOMO mainly locates on the melem moiety, indicative of promoted separation of charge carriers. After incorporation of the carbon ring into hydrogenated carbon nitride, the electrons between some of the C=C bonding

of carbon rings were not well localized and exhibited partial metallization. As the partial metallization of the joint on intralayer heterojunction enabled more free electrons assembled, enhanced electrical field could be obtained on this area. These results are consistent with that calculated by FE method. As a result, combined DFT and FE calculations, conclusions could be reached that on one side, the heterostructure of obtained photocatalyst enhances the electrical field which compels the charge carriers out of the constraint of the Coulomb force. One the other hand, the formed graphene acts as the transferring channel of photoinduced electron, both effectively prohibiting the recombination of electron/hole pairs for redox reactions. (Figure 5g).

## CONCLUSIONS

In summary, we proposed a hydrogen-initiated chemical epitaxial growth strategy to fabricate intralayered polymer based heterojunction. Compared with the traditional CVD method, this strategy can be realized at relatively low temperatures as a result of the confined space and active site. Combined the experimental and computational methods, the acquired in-plane heterojunction exhibits enhanced electrical field and a more efficient electronic structure to facilitate carriers' separation and transfer. As a result, this photocatalyst was shown to possess better performances on artificial synthesis. This work is expected to shed light in further improvement of CEG process for constructing metal-free photocatalysts in order to achieve feasible photosynthesis.

## METHODS

**Damage of bulk structured carbon nitride.** Porous carbon nitride was prepared by a soft template method. Typically, 5 g melamine and sodium bicarbonate (the mass varied from 0.1, 0.3, 0.5 to 0.7 g to respectively synthesize 2%-C<sub>3</sub>N<sub>4</sub>, 6%-C<sub>3</sub>N<sub>4</sub>, 10%-C<sub>3</sub>N<sub>4</sub> and 14%-C<sub>3</sub>N<sub>4</sub>) were both dissolved in ultrapure water and subsequently heated at 100 °C for water evaporation. After that, the obtained precursor was fully ground with an agate mortar and put into a furnace.

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Heat treatment of the precursor was done at 823 K for 4 h with a heating rate of 2.3 K min<sup>-1</sup>. After the calcination, the samples were thoroughly washed by hot water and ethanol for several times and subsequently dried at 333 K overnight to obtain the porous carbon nitride. For comparison, a sample of bulk g-C<sub>3</sub>N<sub>4</sub> was synthetized by the same synthesis procedure without using the soft template of sodium bicarbonate.

Modified chemical epitaxial growth strategy for graphene/carbon nitride intralayered heterojunctions. Initially, a certain amount of the prepared porous carbon nitride was put into a quartz U-tube which was then put into a temperature programmed heating jacket. The U tube was degassed at 523 K with N<sub>2</sub> for 2 h. Then the gas was changed to  $H_2/N_2$  mixture with the concentration of  $H_2$  5% and the degassed sample was heated to 773 K at a heating rate of 5 K/min. After the temperature arrived at 773 K, it was naturally cooled to room temperature to obtain the hydrogenated carbon nitride. The influence of heating under inert gas rather than hydrogen atmosphere on the carbon nitride was also studied with the same reaction condition but the hydrogen atmosphere was replaced by the ultra-pure nitrogen. The obtained sample was named as calcinated g-C<sub>3</sub>N<sub>4</sub>-N<sub>2</sub>.

Characterizations. The structures of the prepared photocatalysts were characterized through the X-ray diffraction (X'Pert PRO MPD, Holland). N<sub>2</sub> adsorption-desorption isotherms were obtained from a Micromeritics Tristar 3000 at 77 K. Scanning electron microscope images were captured on a FEI Verios XHR 460. Transmission electron microscope and HRTEM images were recorded on a JEOL 2100. HAADF-STEM and EDX elemental mapping were performed by a FEI Titan G2 80-200 TEM/STEM. XPS measurements were performed using a Kratos Axis Ultra DLD spectrometer with an Al K $\alpha$  (1486 eV) X-ray source operating at 225W. NEXAFS tests were carried out with the Soft X-Ray beamline from the Australian Synchrotron<sup>35</sup> and conducted at room temperature under the conditions of ultra-high vacuum (UHV). All spectra were collected in the partial electron yield (PEY) mode using an electron

flood gun to adequately neutralise the sample. The normalizing of all the NEXAFS spectra was processed with the QANT software program developed by the Australian Synchrotron.<sup>36</sup> AFM images were collected by a Bruker icon equipment. Solid-state <sup>13</sup>C Direct Polarization (DP) and Cross-Polarisation (CP) MAS NMR spectra were obtained on a Varian VNMRS WB spectrometer (<sup>1</sup>H 399.868 MHz, <sup>13</sup>C 100.554 MHz) equipped with a 4 mm MAS probe. XPS measurements were performed using a Kratos Axis Ultra DLD spectrometer with an Al Ka (1486 eV) X-ray source operating at 225W. Survey spectra were obtained employing a pass energy of 160 eV, whilst all high resolution spectra was collected via a pass energy of 40 eV. The binding energy scale of all XPS spectra were calibrated by locking the C1s peak to 284.5 eV. Solid-state <sup>13</sup>C Direct Polarization (DP) and Cross-Polarisation (CP) MAS NMR spectra were obtained on a Varian VNMRS WB spectrometer (<sup>1</sup>H 399.868 MHz, <sup>13</sup>C 100.554 MHz) equipped with a 4 mm MAS probe. All samples were spun at 7 kHz in 4 mm zircon rotors with kel-F caps. A 90° pulse of 5 us, contact time of 2 ms (CP-MAS) and recycle times of 2 and 60 seconds were used for the CP-MAS and DP-MAS experiments, respectively. For all experiments, 1K data points were collected, and zero-filled to 8K points, with a line broadening of 100 Hz applied. Chemical shifts were determined against an external adamantine standard (38.5 ppm) and are reported relative to tetramethylsilane (0 ppm). Fourier transform infrared (FT-IR) spectra were collected on a PerkinElmer instrument. Diffused reflectance spectra and photoluminescence spectra (PL) were recorded on a Cary 100 UV-vis Spectrophotometer and a Cary Eclipse Fluorescence Spectrophotometer (Agilent, US), respectively. The time-resolved decay spectra was performed on a FLS1000 with the excitation and detection wavelengths were 375 and 460 nm, respectively.

*In situ* detection of effluent gases under  $H_2$  or  $N_2$  calcination. Porous carbon nitride was placed in a sealed tube furnace, which was connected with an online gas chromatograph (GC, Agilent 7820A).  $H_2$  or  $N_2$  was purged into the reactor for fully replacing the air inside. Then

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the furnace was heated and the effluent gases were analyzed by gas chromatograph equipped with Porapak N and Molsieve 5A columns.

**Photoinduced selective oxidation of 1,4-DHP.** The photooxidation tests were conducted in a customized reactor. The aqueous solution (50 mL) dispersing 10<sup>-4</sup> mol/L of 1,4-DHP and samples (30.0 mg) was irradiated by a LED solar simulator (Newport). After 1 h, the conversion of 1,4-DHP was monitored by a Cary 100 UV-vis Spectrophotometer.

**Photocatalytic water splitting.** The H<sub>2</sub> evolution tests were performed by using a LabSolar-III AG reaction system (Beijing PerfectLight Co.). Photocatalyst (50 mg) was dispersed into a 50 mL of aqueous solution containing 5 mL of TEOA, and 0.8 mL H<sub>2</sub>PtCl<sub>6</sub>·(H<sub>2</sub>O)<sub>6</sub> (3.0 wt%). Before the irradiation, the reaction cell was sealed to thoroughly evacuate for 40 min to remove air. A Xe lamp (300 W) (CEL-HXF300, Beijing Cel Sci-tech Co., Ltd.) equipped with a 420 nm cut-off filter was employed as the visible light source. The reaction system was kept at 6 °C using the cooling water. The generated hydrogen was detected with an online gas chromatograph (GC7920) equipped with a 5A molecular sieve column and a thermal conductive detector (TCD). N<sub>2</sub> was employed as the carrier gas. Same reaction without H<sub>2</sub>PtCl<sub>6</sub>·(H<sub>2</sub>O)<sub>6</sub> as co-catalyst and without cut-off filter was also performed. Besides, the stability of hydrogenated carbon nitride was studied in water splitting reaction with three cycles. In addition, the quantum efficiency of photocatalytic water splitting reaction was tested under the irradiations of monochromatic light including 395, 420 and 500 nm.

**Photo-to-thermal tests.** 50 mg sample was dispersed uniformly on a quartz microfiber filter which was sealed in a quartz reactor. A thermocouple was contacted with the surface of filter for monitoring the temperature of sample. Before reaction, the reactor was replaced by Ar for 2 h to make sure no air inside. Then light with 1 W/cm<sup>2</sup> was irradiated on the sample and simultaneously the temperature of sample was recorded at certain time intervals.

**Simulation details.** Vienna Ab initio Simulations Package (VASP) was used to perform the DFT simulations.<sup>37,38</sup> The projected augmented wave (PAW) method and generalized gradient approximation (GGA) in the Perdew, Burke and Ernzerhof (PBE)<sup>39</sup> parameterization were used. The molecule fragments were modelled in a  $35 \times 18 \times 15$  Å<sup>3</sup> supercell, which allows at least a distance of 10 Å between images in any dimension. The first Brillouin zone was sampled by using a Gamma-centred **k**-point set of  $1 \times 2 \times 2$ . The plane-wave cutoff energy was set to be 400 eV. The coordinates were fully relaxed, and the energy and force were converged within  $10^{-4}$  eV/cell and  $10^{-2}$  eV/Å, respectively.

The finite element modelling (FEM) is on the basis of classical electro-magnetic field theory employed to obtain the electrical field by solving Maxwell equations. COMSOL Multiphysics software<sup>40</sup> was used for modelling the plasmonic effect within the carbon nitride/graphene structure. The size of the simulation domain was of the same order of magnitude with the optical wavelength, so that the effect of the wavelength variation can be accurately evaluated. A FEM model is constructed with the upper part of the domain filled with water while the lower part is filled with the g-C<sub>3</sub>N<sub>4</sub> materials, separated by a p horizontal boundary. A graphene column with a diameter of 50 nm is embedded at the centre of g-C<sub>3</sub>N<sub>4</sub> material. For comparison purpose the second model of plain g-C<sub>3</sub>N<sub>4</sub> block was built. Perfect matching layers (PMLs) were defined both at the top and at the bottom of the model, which were designed to completely absorb all the reflected and transmitted waves from the top and bottom. For all of the FEM modelling results presented in this work, an incident radiation is set at the top, transmitting downwards with an intensity of 10<sup>4</sup> W/m<sup>2</sup>.

## ASSOCIATED CONTENT

## **Supporting Information**

The supporting Information is available free of charge on the ACS Publications website at DOI: Details of experimental section, textural properties, chemical compositions, TEM, SEM and AFM images, FT-IR, nitrogen isotherms, XPS-VB, photooxidation performances of 1,4-DHP, photocurrent curves, FEM and DFT simulations; additional supplementary figures, tables, and text referenced in the article. (PDF)

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## **Author Contributions**

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

## Notes

The authors declare no competing financial interest.

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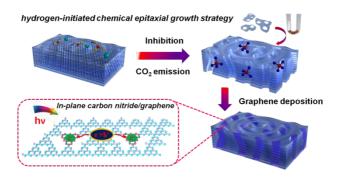
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## **Table of Contents Entry**



We *in situ* deposited nano-graphene along confined spaces with active site of porous carbon nitride. The performances in both selective photooxidation and photocatalytic water splitting were improved owing to the enhanced electrical field and a more efficient electronic structure from the addition of trace amount of plasmonic nano-graphene.