

Periodical oscillation of particle-laden laminar flow within a tubular photocatalytic hydrogen production reactor predicted by DEM simulation

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Abstract

Besides their wide existence in various industrial processes, nanoscale particle suspensions are also the important media for some emerging technologies such as photocatalytic hydrogen production. The circulating flow properties of the nanoparticles in the fluid are of great concern for their practical use. In our study, a modified experimental system was set up based on Malvern laser particle analyzer that can estimate the nanoparticle concentration and size distribution in a laminar nanoparticle circulating flow. We found that the particle concentration and size distribution were periodical oscillation with time in such flow. Understanding the oscillation mechanism is capable of promoting the energy efficiency of photocatalytic hydrogen production. A simulation based on Discrete Element Method (DEM) was conducted to understand this particular oscillation mechanism by counting the single particle movement and trajectory properties in the solid-liquid suspension. The simulation results agree well with the tendency obtained by the experimental results and capable of better understanding the oscillation characteristics. The simulation results reveal that the nanoparticles tend to gather in the middle region (the higher velocity region) of the tube after several cycles. Moreover, gravity is of great significance in the circulating flow of solid-liquid suspension because the particle swarms tend to distribute a little below the axial center line of the straight tube. These obtained results are credible for understanding the nanoscale particle transport phenomenon in many natural or industrial processes. In particular, our results are helpful for the understanding and effective control of the movement

34 and distribution of photocatalyst particles in the tubular photocatalytic reactor, which is believed
35 to significantly affect the incident radiation distribution and finally the energy conversion
36 efficiency of the photocatalytic process.

37 **Key Words:** particle-laden flow, DEM simulation, periodical oscillation, particles trajectory

38

39 **1. Introduction**

40 The suspensions of micro- or nanoscale particles in liquid exist widely in various fields, e.g.
41 advanced material processing, electronic technology, chemical engineering, petroleum, food
42 processing, waste treatment etc. It is also the important media for some emerging technologies
43 such as the enhancement of heat transfer, photocatalytic energy conversion and environmental
44 remediation [1-4]. In all of these processes, the particle transportation and deposition have
45 significant impact on the properties application promotion of such kind of suspension. For
46 example, in photocatalytic process, sedimentation of the photocatalyst which leading to ineffective
47 light absorption must be avoided [5]. For another example, in nanofluids heat transfer
48 enhancement, the deposition of nanoparticles will change the wettability of the heater surface, thus
49 affecting the overall heat transfer performance [6].

50 Photocatalytic hydrogen production from water splitting using solar energy is one of the
51 ultimate reactions to solve energy and environmental issues [7]. In order to realize the industrial
52 application of photocatalytic hydrogen production technology, a large number of studies have been
53 carried out in the past twenty years. And there are two main research fields, one is the exploration
54 for efficient visible light driven photocatalyst [5,8-12], the other one is the development of high
55 efficiency photocatalytic reactors [13-17]. As the photocatalyst has been studied for many years,
56 and the quantum efficiency of the photocatalyst has reached 93% with noble metal loading [18]
57 and reached 62% without noble metal loading[19], whereas the energy conversion efficiency is
58 still not able to exceed 5% in present [20,21]. It is believed that the great difference between the
59 above two efficiencies is due to the inefficient utilization of the light in the photocatalytic reactor.
60 And the suspension state and transport phenomenon of photocatalytic particles in the reactor has
61 significant impact on the absorption of light; as a result, determine the overall energy conversion
62 efficiency of the photocatalytic reaction.

63 For the clearer understanding of the transportation phenomenon in nanoparticle flow, many

64 studies have been conducted by means of experimental and/or numerical method. In experimental
65 study, one mainly concerns about the concentration distribution, velocity distribution and particle
66 size distribution, and the experimental methods mainly include intrusive and non-intrusive method.
67 Sampling method[22] and probe method[23-25] are two most frequently-used methods in the
68 intrusive method. Although this kind of method is easy to use and intuitive, it inevitably disturbs
69 the flow field due to the existence of sampling equipment and probe. To overcome this defect,
70 some non-intrusive methods have been developed, such as the light technique [26], radiation
71 technique[27], ultrasonic attenuation techniques [28], tomographic technique[29] etc. Among
72 these techniques, light technique is the most widely used technique, for it is convenient, safe, and
73 responsive and has simple structure. Laser diffraction (LD) technique is a kind of light technique
74 based on the light diffraction by particles, which has become a popular method to measure the
75 particle size and concentration distribution [30,31], due to its high speed, good reliability and high
76 reproducibility [32].

77 Although many experimental methods have been developed, they still have some limitations,
78 such as low accuracy, high cost, disturbance to the flow field etc. Also, Numerical methods have
79 been employed as a very important supplementary method to give insight into the particle
80 distribution and flow characteristics in the particle suspensions. In general, simulation of the
81 solid-liquid suspension can be conducted by either continuum or discrete methods, depending on
82 the time and length scales employed for solid phase. The former is often based on the two-fluid
83 model (TFM), whereas the latter is usually based on the discrete element method (DEM). Very
84 often, the discrete phase method (DPM) coupled with computational fluid dynamics (CFD) have
85 been used, which gives the DPM-CFD or combined continuum and discrete model (CCDM). TFM
86 is available to study the fluid mixture movement as a whole and the phase distribution. Hu et al.
87 [33] investigated the effects of the slurry flow and catalyst distributions in the reactor on
88 photocatalysis for hydrogen production with an algebraic slip mixture model (ASM). Hatami et al.
89 [34] studied steady and unsteady magneto-hydrodynamic (MHD) Couette flows between two
90 parallel infinite plates with a two-fluid model through numerical Differential Quadrature Method
91 (DQM) and analytical Differential Transformation Method (DTM), respectively. The forced
92 convection of laminar TiO₂-water nanofluid flow in a parallel plate microchannel has been studied
93 by a modified homogeneous flow model and the dispersion model [35]. Although the TFM model

94 is used widely and suitable to most two-phase flow problems, some detailed particle-scale
95 information might be missed, such as the trajectories and forces acting on individual particles,
96 which are very significant to explain the mechanisms governing the complicated flow behavior.
97 CCDM has thus been increasingly used for studying various solid-liquid flow phenomena [36].
98 Morris and Brady [37] studied the pressure-driven flow of a non-neutrally buoyant suspension by
99 the suspension-balance model of NB considering the particles migration which is a DEM model.
100 Peng et al. [38] investigated the influence of primary particle size distribution (PPSD) on
101 aggregation behavior and the resulting effect on yield stress of a concentrated colloidal suspension
102 with DEM model. Chaumeil and Crapper [39] investigated the agglomeration and deposition on a
103 constricted tube collector of colloidal size particles immersed in a liquid.

104 The literature review above indicate that most of the studies by DEM modeling are conducted
105 in a one-way channel, which means the particle suspension will pass the channel for only one time
106 in simulation. Whereas in many photocatalytic reactors, nanoparticle suspensions are in fact
107 circulated in a loop, considering that such flowing reactor is cost-effective, easy for scale-up and
108 for nanoparticles recycling [40-44].

109 In particular, sequencing batch reactor (SBR) is commonly used in photocatalytic technology
110 to decrease the consumption of pump power [16,45]. And particles need to be resuspension after a
111 period of time when the SBR photocatalytic reactor is in operation, so it is very important to know
112 the real process of particles spreading with the suspension flow. Thus, a initial condition that
113 particles added at the beginning was select to simulate the process. And in the circulating flow, the
114 gravity cannot be neglected considering that the particles in solid-liquid suspension will keep
115 moving in the reactor under the action of gravity [46]. And the equilibrium position of the
116 particles after enough times of circulating flow should be of great interest in photocatalytic
117 application, for it is extremely important to know the location and of particles in the photocatalytic
118 reactor which has great impact on the light absorption. Therefore, the equilibrium position of these
119 particles is also discussed in this study. In our previous study, an experimental research is
120 conducted to investigate the changes of particles concentration and size distribution with time
121 during the circulating flow. In this study, a DEM two-way coupling model was employed to obtain
122 the details of particles migration and also the key macroscopic parameters for comparison with
123 experimental results. This model focused on understanding the periodical oscillation of particle

124 concentration and size distribution with time in such circulating flow. The particle volume fraction,
125 the flow field of the particle-laden flow in the circulating flow system, the particle trajectory and
126 the particle size spatial distribution were discussed by this model. This model provides an intuitive
127 and credible insight on the periodical oscillation mechanism of photocatalyst particles tubular
128 photocatalytic reactor, and eventually explains incident radiation distribution and the energy
129 conversion efficiency of the photocatalytic process.

130 **2. Experimental set-up and materials**

131 In our previous work, we had established an experimental set-up as has been shown in [Fig 1](#)
132 [47]. It has six parts: main flow channel, centrifugal pump, wet sample injector, laser particle
133 analyzer and connecting tube. The main flow channel working as test section in this study is made
134 from acrylic glass and has an inner diameter of 30 mm and a total length of 0.5 m which is
135 transparent. Connecting tube is the silicone hose which the inner diameter is 10mm, and the total
136 length of it is 1.7 m. The centrifugal pump is used to provide driving power of the circulating flow.
137 A Malvern Spraytec laser particles analyzer was used to monitor particle concentration and
138 particles size distribution simultaneously. Degussa P25 titanium dioxide was employed as model
139 particles and main parameters for this widely used photocatalyst, and the specifications of P25
140 TiO₂ powder can be found in Ref. [48]. Again, it is worth mentioning that photocatalytic energy
141 conversion is a very attractive technology for which one of the key issues is the well-suspension of
142 the particles in the reactor [3,49,50]. A very low flow velocity may result in the sedimentation of
143 particles while a too high flow velocity could lead to a significantly reduced exposure time of the
144 photocatalyst to solar light. One must guarantee that all the useful incoming photons are used and
145 do not escape without having intercepted a particle in the reactor [3,51]. In our design, the average
146 volume fraction of nanoparticles is $\varphi=1\times 10^{-4}$ in the tube. It can thus be considered as a dilute flow
147 and the possible interactions between particles can be safely ignored during theoretical analysis.

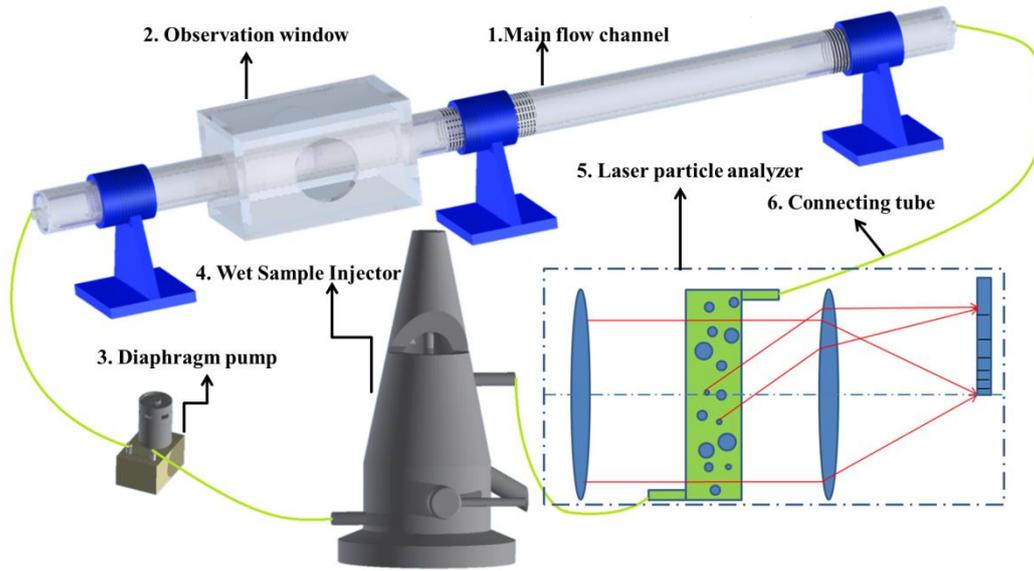


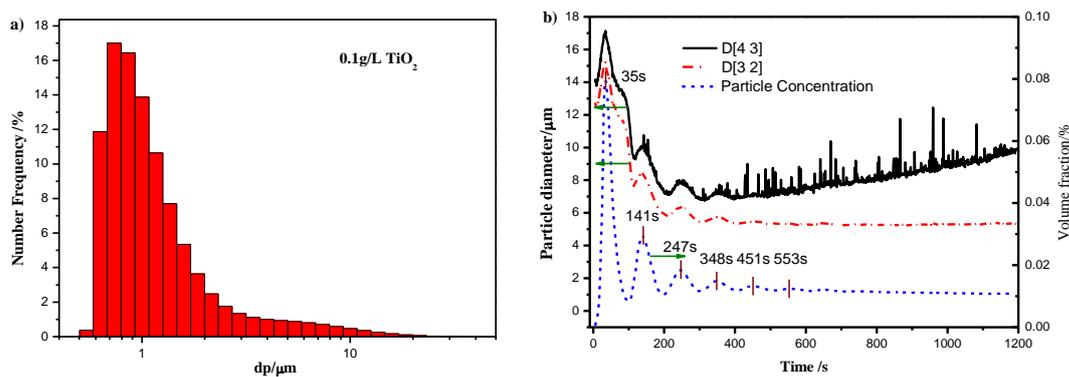
Fig.1. Schematic of experiment set-up[47]

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 149
 150
 151 Firstly, the system was filled with deionized water and circulated by pump with the flowrates
 152 adjusted to 15 L/h. And the P25 particles are dispersed in deionized water in previous, which the
 153 particle size distribution is showed in Fig. 2a. And it can be seen that the range of the particle
 154 diameter is 0.50-36.87 μm . Then the prepared P25-TiO₂ suspension was introduced into the wet
 155 sample injector and the laser diffraction equipment starts to record the particles volume fraction
 156 and size distribution simultaneously. Fig.2b shows the mean particle size and concentration versus
 157 time in the repeatedly circulating flow. Here, the upper two curves corresponds to the variations of
 158 two kinds of mean particle sizes, D₃₂ and D₄₃, respectively and the bottom curve corresponds to
 159 the variation of particles volume fraction against time. Here, D₃₂ is the mean particle size taking
 160 into account both volume and surface area of the particles [52], while the D₄₃ is the average
 161 particle size based on the volume moment [53]. The corresponding expression for the two average
 162 particle sizes can be derived respectively by:

$$D_{32} = \frac{\sum V_i D_i^3}{\sum V_i D_i^2} \quad (1)$$

$$D_{43} = \frac{\sum V_i D_i^4}{\sum V_i D_i^3} \quad (2)$$

165 Obvious periodic oscillation of both mean particle sizes and particles volume fraction with time
 166 can be noted in Fig.2b. And we can define the distance between two adjacent peaks as the period
 167 of the oscillation curve. It can be found that the period of the three curves in the Fig. 2b is nearly
 168 the same. The oscillation of the particle volume fraction curve with time is particularly significant
 169 in the initial 600 seconds, and after that the particle volume fraction of the particles in suspension
 170 tends to be a constant value, which we define this state as quasi-steady state. According to the
 171 definition of mean particle size and D_{43} , D_{43} is more sensitive to the large particles, while smaller
 172 particles contribute more to the D_{32} . It is also noted in the Fig. 2b that the D_{32} and D_{43} curves are
 173 also oscillation curves, but the tendency of them are some different comparing to the particle
 174 volume fraction curve. It can be observed that the particle volume fraction curve fluctuates around
 175 an average value, while the D_{43} and D_{32} curves fluctuate with an obviously downward trend.

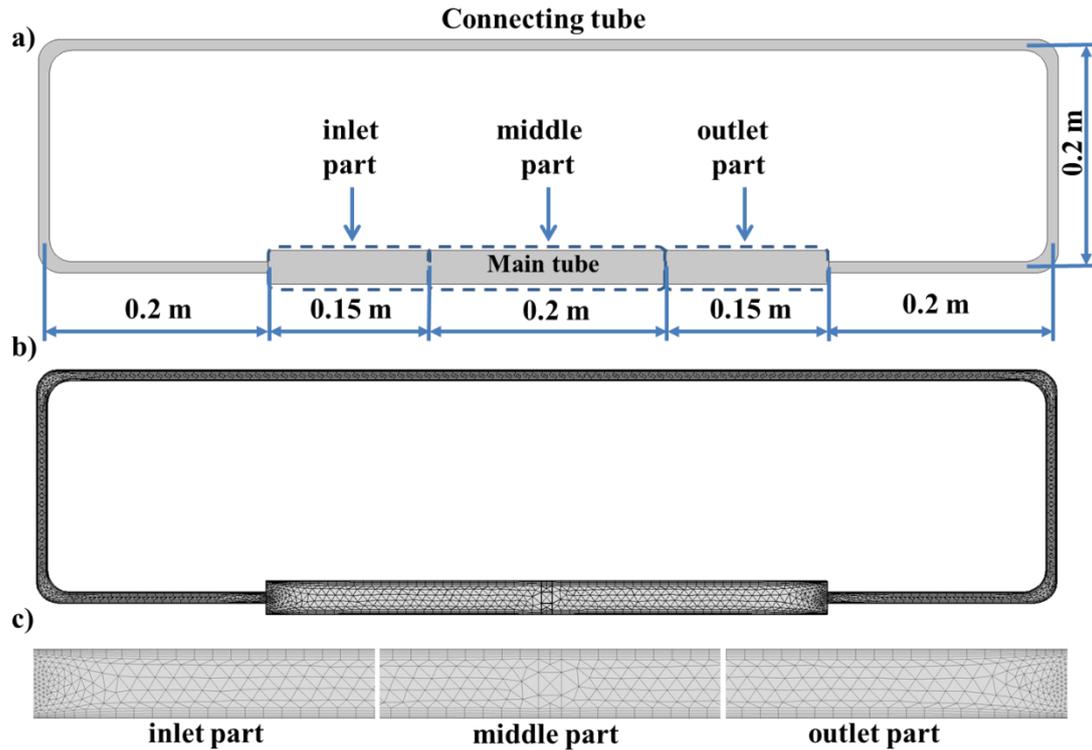


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 177 **Fig 2** a) Particles size distribution of the P25-TiO₂ powder dispersed in water, b) Variation of
 178 various mean particle sizes and particle volume fraction with time.[47]

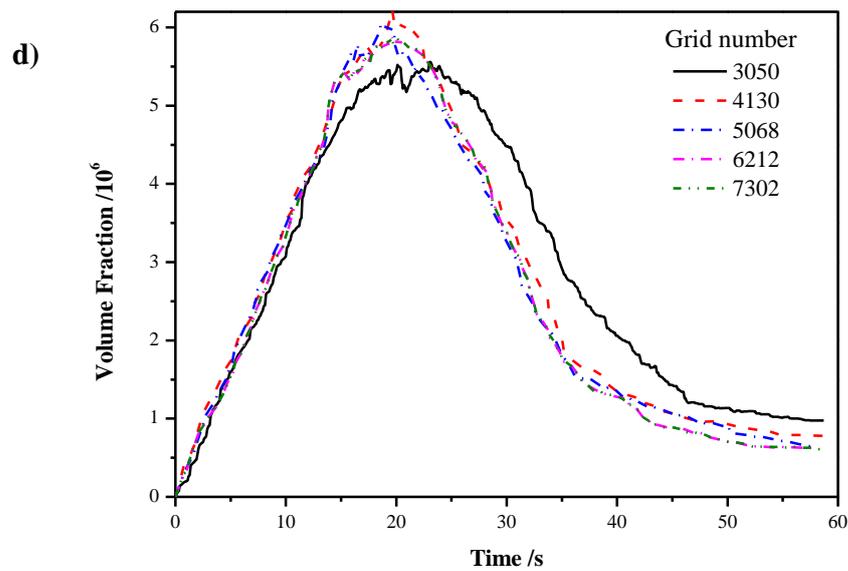
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 180 **3. Numerical simulation of the nanoparticle-laden flow**
 181 **3.1 Physical model for the tubular reactor**

182 As has been mentioned above, the particle volume fraction has been measured by LD method,
 183 but only volume fraction cannot help us to know the mechanism of the particle transport
 184 phenomenon in such a circulating flow system. To explain the transport phenomenon of the
 185 particles in laminar circular flow, especially the details of the movement and variation of the
 186 particles with different size during the transport process, a simulation based on the Euler-Lagrange
 187 model was conducted. The geometry model for the simulation is showed in Fig. 3a, which is
 188 established according to the experimental set-up, but the particle size analyze and pump are
 189 neglected for the sake of simplicity. Here the main tube has a diameter of 30 mm with length of

190 0.5 m, and the connecting pipe has a diameter of 10 mm with length of 1.7m, which is the same
 191 with experimental set-up. The mesh used in this simulation is showed in Fig 3b, which is obtained
 192 from unstructured triangular grid. For the boundary layer, two layers of quadrilateral grids are
 193 adopted near the boundary. The mesh dependency analysis based on particle volume fraction was
 194 then conducted and the results are showed in Fig.3d. It turned out that the mesh with a grid
 195 number of 6212 is accurate enough for this simulation and also time-efficient.



196



197

198 Fig.3 a) The Geometry model, b) computing mesh of the simulation, c) Local enlarged figure of

199 the mesh, **d)** mesh dependency analysis based on particle volume fraction

200 On the other hand, as the particles movement will be described under the Lagrange viewpoint,

201 the particles need to be traced one by one. However, it will be an extremely large amount of

202 particles in the tube in practice, for example, the total volume of the channel is about 0.423 L,

203 supposing the particles size is 3 μm and volume fraction is 0.01%, the amount of the particles will

204 be 2.994×10^9 , which is a too huge to cost an impossible computational resource. To solve this

205 problem, only 10,000 particles are taken into account in the simulation, and these particles can

206 present the migration principle of the particles with different size. For the convenience of

207 analyzing, we also divided the main tube into three sections, i.e. inlet, middle and outlet sections,

208 respectively. Furthermore, the properties of particles and the base fluid have significant influence

209 on the interactions between particles and liquid and/or between particles and particles. The

210 properties of particles mainly include density, shape, size, specific surface area and surface charge

211 et al. The density of the P25 TiO_2 particles was defined as 4000 kg/m^3 according to Ref.[48], and

212 the size of particle was defined as the particle size distribution as Fig. 2a shows which was

213 obtained from the experimental data. The shape of the particles also has a significant impact on the

214 drag force between liquid and solid phase, whereas the shape is often defined as spherical in

215 numerical study to simply the computational model [54]. To solve this problem, some drag

216 correlations have been developed to modify the influence of particle shape on drag force between

217 particles and fluid [55,56]. In this study, as the particle was set as a common photocatalyst P25

218 TiO_2 , and the SEM of this kind of particles shows that the morphology generally exhibits spherical

219 [57], the particles in this study are all considered as spherical particles. And it needs to be further

220 pointed out that if the current computational model is needed to simulate the other non-spherical

221 particles, one merely need to add some corresponding drag correlations mentioned above. On the

222 other hand, particle aggregation is a common phenomenon in suspension especially for tiny

223 particles such as nanoparticles, and there are many influence factors of the aggregation

224 phenomenon, such as the temperature of suspension, the pH of the suspension, ions in the base

225 solution, particle size, shape, method of preparation and solid concentration [1,58]. Since the

226 particle aggregation phenomenon has noticeable impact on many applications, there have been

227 many theoretical [59-61] and experimental [62-64] studies focus on the problem. In this study,

228 aggregation was also taken into account by introducing a real particle size distribution (PSD)

229 which had been obtained from experimental data [47]. In other words, the particles aggregation
 230 was assumed to be equilibrium state and the particles in this study are aggregated secondary
 231 particles, thus the aggregation effect is taken into consideration in reality. Table 1 lists the
 232 parameters of particles and base liquid used in the computational model of this study.

233 **Table 1** Some parameters of particles and base liquid used in the computational model

Parameters	Description
Particles	P25 TiO ₂
Base fluid	Deionized water
Density of particles	4000 kg/m ³
Particle size	Using the PSD shown in Fig. 2a
Particle quantity in the model	10000
Density of the base fluid	1000 kg/m ³
Viscosity of the base fluid	1 mPa • s

234

235 3.2 Mathematical model

236 Considering that the particles concentration is very low in our study, to simplify the
 237 mathematic model, we assume that the particles hardly influence the fluid field. Therefore, the
 238 simulation can be done in two steps. Firstly, a single phase flow in the geometry was solved and
 239 then the trajectory of particles was computed. In particular, in our case it is worth noting that the
 240 flow region is a closed circular channel, which only has wall boundary but not has inlet or outlet
 241 boundary. To solve this problem, the whole circular channel will be divided into two parts, the
 242 main tube and the connecting tube. The flow direction of the circular is set to be anticlockwise, so
 243 the right head of the connecting tube is set as the inlet boundary of the connecting tube which is
 244 also set as the outlet boundary of the main flow tube and the left end of the connecting tube is set
 245 as the outlet boundary which is also set as the inlet boundary of the main flow tube. The circular
 246 flowrate is set as 15 L/h which is the same as the experimental condition.

247

248 As has been noted, in our study the trajectory of particles will be computed by solving
 249 ordinary differential equations using Newton's law of motion. In this method, the particles are
 250 treated as point masses. The specification of particle diameter is mostly used for size-dependent
 251 forces, such as the drag and Brownian motion forces. As the particles concentration is very low in
 252 the liquid phase, it can be supposed that the liquid phase affects the motion of the particles but not
 253 vice-versa and one particle will not have obvious interactions with other particles. Thus, the

254 particle tracing equation according to Newton's second law can be expressed as [Eq. \(3\)](#):

$$255 \quad \frac{d}{dt} (m_p \mathbf{v}) = \mathbf{F}_D + \mathbf{F}_g + \mathbf{F}_B \quad (3)$$

256 Where the m_p is the mass of the particle, the \mathbf{v} is the velocity vector of particle; \mathbf{F}_D , \mathbf{F}_g and \mathbf{F}_B are
 257 three kinds of forces that can affect particles motion, and their calculation methods will be
 258 introduced in detail. \mathbf{F}_D is the drag force of the fluid which is defined as:

$$259 \quad \mathbf{F}_D = \left(\frac{1}{\tau_p} \right) m_p (\mathbf{u} - \mathbf{v}) \quad (4)$$

260 where the τ_p is the particle velocity response time, \mathbf{u} is the velocity of the fluid, \mathbf{v} is the velocity of
 261 the particle. There are a large number of expressions for the particle response time. It depends on
 262 the drag law, and selecting the appropriate drag law needs a consideration of the relative Reynolds
 263 number Re_r of particles in the flow, which is given by:

$$264 \quad Re_r = \frac{\rho \|\mathbf{u} - \mathbf{v}\| d_p}{\mu} \quad (5)$$

265 where the ρ and μ are the density and viscosity of the fluid, respectively, and d_p is the diameter of
 266 the particles. With the dimensionless number Re_r , τ_p can be given by:

$$267 \quad \tau_p = \frac{4\rho_p d_p^2}{3\mu C_D Re_r} \quad (6)$$

268 where the C_D is the drag coefficient. Here, the parameter d_p is given by the particles size
 269 distribution shown in [Fig. 2a](#). According to this particles size distribution, the range of the Re_r can
 270 be evaluated. As the drag force drives the particles to follow the fluids movement, the term

271 $\|\mathbf{u} - \mathbf{v}\|$ that describe the difference between particle velocity and fluid velocity will be closed to
 272 0 during the transport process, therefore the minimum value of the Re_r will be very clear to zero.

273 On the other hand, if both the term $\|\mathbf{u} - \mathbf{v}\|$ and d_p reach the maximum value, the maximum of

274 Re_r can be obtained. In this simulation, the maximum value of the $\|\mathbf{u} - \mathbf{v}\|$ most possibly

275 happens at the initial state which the particles are static. And the maximum velocity of the fluid in

276 the main tube can be calculated according to the velocity distribution of the laminar flow, and the

277 value is 0.0118 m/s, so the maximum value of $\|\mathbf{u} - \mathbf{v}\|$ will be 0.0118 m/s. For the maximum

278 particle size is 36.87 μm , the maximum value of the Re_r can be evaluated according to the [Eq. 3](#), is

279 0.4203, that is, the range of the Re_r is 0-0.4203. According to the previous researches, the most
 280 common used drag law is the Stokes drag law, but it only can be used in the case that the $Re_r \ll 1$.
 281 Thus, another famous drag law called Oseen correction is used [65], which the drag coefficient is
 282 given by:

$$C_D = \frac{24}{Re_r} \left(1 + \frac{3}{16} Re_r \right) \quad (7)$$

283
 284
 285 \mathbf{F}_g is the gravity force which is given by:

$$\mathbf{F}_g = m_p \mathbf{g} \frac{(\rho_p - \rho)}{\rho_p} \quad (8)$$

287 \mathbf{F}_B is the Brownian Force which can describe the Brownian motion of the particles. Actually,
 288 the Brownian motion of the particles leads to spreading of particles from regions of high particle
 289 concentration to low concentration. The expression of the Brownian Force is:

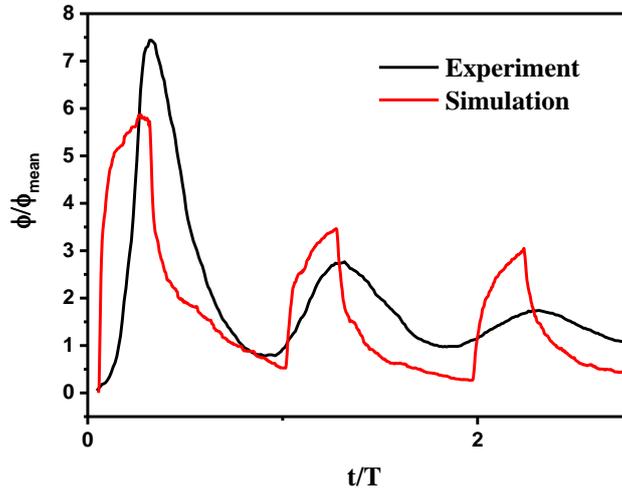
$$\mathbf{F}_B = \zeta \sqrt{\frac{6\pi k_B \mu T d_p}{\Delta t}} \quad (9)$$

291 where the Δt is the time step taken by the solver, T is the absolute fluid temperature, k_B is the
 292 Boltzmann constant and the value is 1.3806×10^{-23} J/K, and ζ is a normally distributed random
 293 number with a mean of zero and unit standard deviation which is created for each particle, at each
 294 time step.

295 3.3 Model validation

296 To validate the model, the result of the variation of the particle volume with time by this
 297 simulation is compared with the experimental result, which is showed in **Fig 4**. As has been
 298 mentioned above, the particle volume fraction in the simulation is much smaller than the
 299 experimental condition for the huge computational source consumption, the particle volume
 300 fraction in **Fig 4** is normalized by dividing the mean particle volume fraction in the quasi-steady
 301 state. Also, for it has some simplified assumptions in the physical model, so the times of the two
 302 results are not exactly equivalent, therefore a normalized time is used by dividing the period of the
 303 oscillation curve, T . From the **Fig 4**, one can find that the tendencies of two curves are very similar
 304 in the first 3 periods, but it can be also seen that the continuity of the simulation curve is poor
 305 compared with the experimental results, this may due to the particle amount is much smaller than

306 the experimental condition. However, in this study, we only want to reveal the particle movement
 307 mechanism and how this movement results in an oscillation of particle volume fraction and mean
 308 particle size. From this viewpoint, the result of this simulation result also exhibit a very obvious
 309 oscillation tendency, so it can help us to understand the mechanism.



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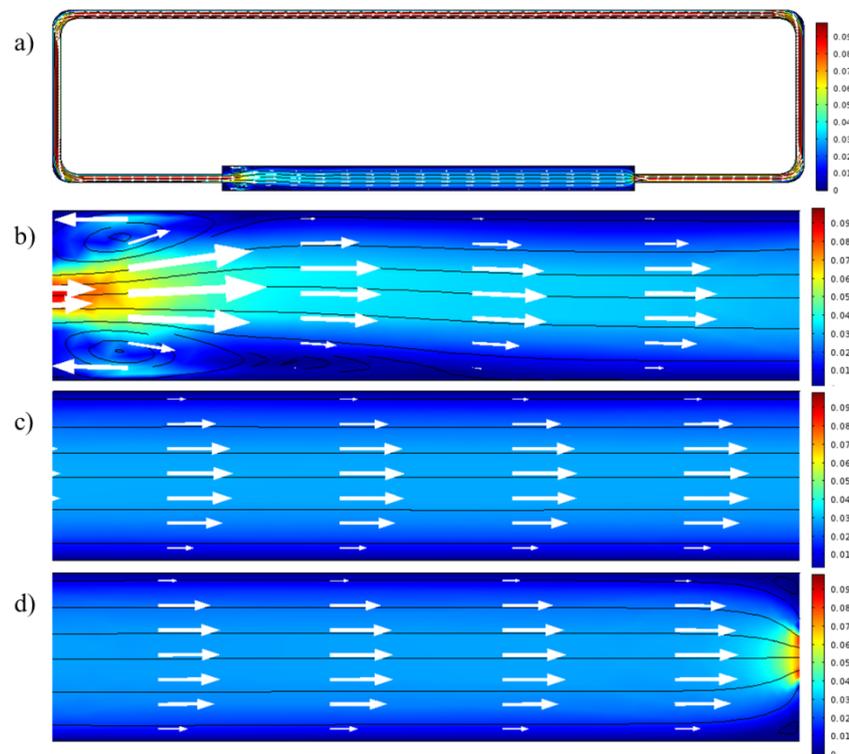
311 **Fig. 4.** Variation of particle volume fraction with time in experiment and simulation

312 4. Results and discussion

313 4.1 The flow field of the particle-laden flow in the circulating flow system

314 As the particles moving in the fluids are significantly affected by the drag force of fluid, the
 315 flow field of the flow is believed to be extremely important for the latter analysis of particles
 316 movement. The velocity distribution obtained by the simulation is showed in [Fig. 5](#), which is a
 317 typical laminar tube flow. [Fig. 5b](#) shows the velocity distribution in the inlet part, and it can be
 318 seen that the fluid flow from the thinner connecting tube into the thicker main tube. As the fluid
 319 velocity is larger in the thinner connecting tube, there is a high velocity region in the front-end of
 320 the inlet part, and there also appears inverse flow with two symmetrical vortexes in the top and the
 321 bottom of this region. Then the fluid flow in the main tube gradually passes into the fully
 322 developed flow. [Fig. 5c](#) shows the velocity distribution in the middle part, which shows the
 323 character of fully developed flow. And the velocity distribution exhibits a parabolic distribution,
 324 which is a typical laminar tube flow. This kind of flow has a character that the velocity in the
 325 center of tube reaches the maximum value, while the velocity near the wall is low and the velocity
 326 at wall is zero. That is to say, the velocity distribution shows an obvious difference in radial
 327 direction. And [Fig. 5d](#) shows the velocity distribution in the outlet part, and it can be seen that the
 328 velocity distribution in the front of this part is very similar to the middle part. However, it appears

329 a high velocity region in the last end of this part, for the fluid flow from the thicker main tube to
330 the thinner connecting tube.



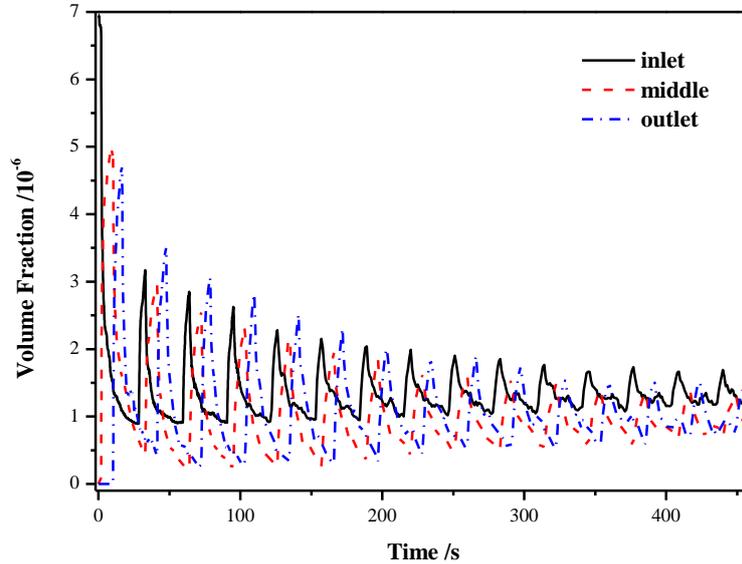
331
332 **Fig.5** Contour of the velocity magnitude of the circular flow with streamlines **a)**the whole
333 channel, **b)** the inlet part **c)**the middle part **d)**the outlet part, of the main tube

334

335 4.2 Particle volume fraction

336 **Fig. 6** shows the variation of the average particle concentration with time during the transport
337 process in three parts of the tube. As has been mentioned in section 3.3, the simulation result is
338 very similar to the experimental results. The most obvious difference between the two results is
339 that the curve of the experiment reaches a constant value in the end, whereas the curve of the
340 simulation just holds a value in a small range. This deviation may be due to that too fewer
341 particles are taken into account in the simulation and the disturbance from pump or other
342 equipment used in experiment has also not been taken into account. However, as the tendency
343 obtained from the simulation is similar to the experiment results, it is therefore reasonable to use
344 the results of simulation to find the mechanism of the observed experimental phenomenon. From
345 Fig 6, it can be found that the highest peak appears at the first peak of the oscillation curve in all
346 the three part of main tube, while the height of the three peaks are different and the peak in the
347 inlet part is the highest. It can also be seen that the peaks of the oscillation curves in different

348 parts also don't coincide in time. These phenomena may mean that the oscillation of particle
349 volume fraction is caused by the particle movement in the circulating flow. And to verify this
350 assumption, we will discuss the particle distribution and trajectory in the next part.



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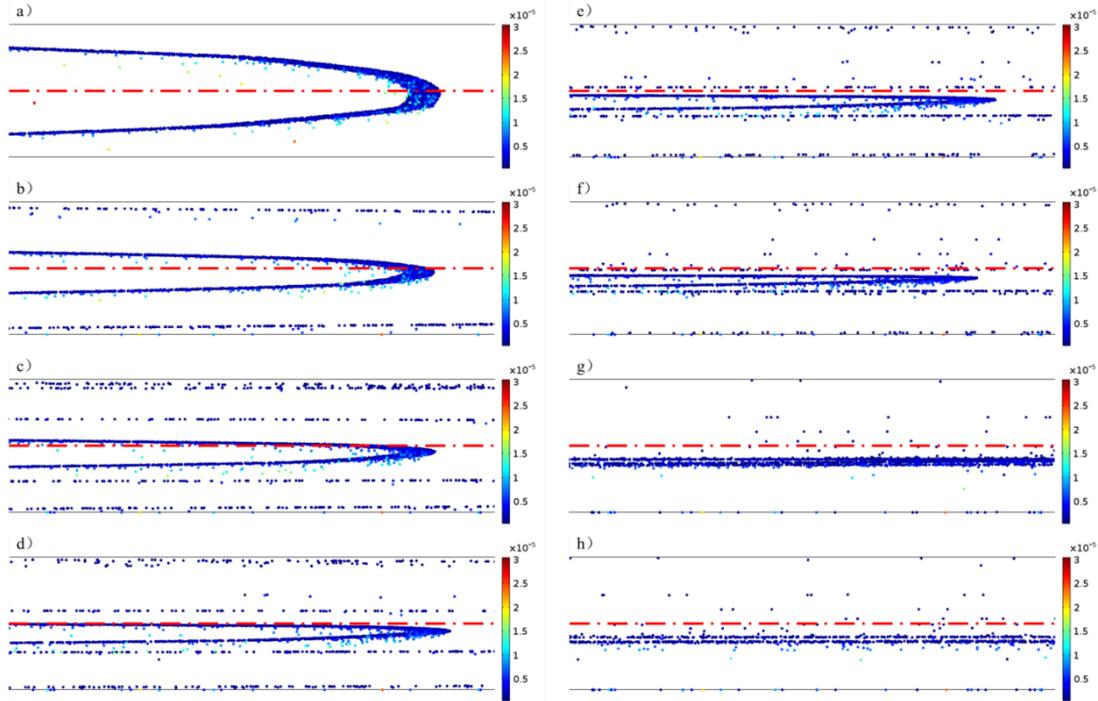
352 **Fig. 6** Variation of volume fraction of particles with time derived from the simulation

353

354 4.3 The particle trajectory

355 Firstly, the trajectories of particles in the times which correspond to the position of the peaks
356 of the particle volume fraction curve are showed in [Fig 7](#). It can be found that the particles
357 distribution exhibits a parabolic like state in the initial time which is very similar to the velocity
358 distribution shown in the [Fig 5](#). And it can be easily found from [Fig 7a-f](#) that the particles in these
359 peak times tend to arrange in a parabolic curve. Combined with the velocity distribution
360 mentioned above, it can be inferred that this phenomenon is caused by the parabolic distribution of
361 radial velocity, that is, the particles near the center will have higher velocity, while the particles
362 near the wall will move slowly. Thus, the particles with higher velocity will pass through the
363 whole circulating system much faster than other particles, and when they reach the counting
364 region, there will be a relative higher particle volume fraction. This process is believed to be a
365 good explanation for the oscillation phenomenon. Furthermore, it can be found that the parabolic
366 shape of particles gets narrower gradually until it almost turns into a line at the quasi-steady state,
367 and this is due to the distance between the particles with different velocity will be longer as time
368 goes on. The narrower parabolic particles arrangement means fewer particles, which can explain
369 the lower and lower peaks high in the oscillation curve. Also, it can be found that the position of

370 the overall parabolic shape in the tube becomes lower and lower during the transport process,
 371 which may due to the gravity force of the particles. After several circles, particles moving in the
 372 main tube will gather in the positions that a little lower than the axis of the main tube.



373
 374 **Fig. 7** schematic of the particles position in the peak times and the quasi-steady state **a)** $t=14.6s$ **b)**
 375 $t=45.6s$ **c)** $t=76.7s$ **d)** $t=108s$ **e)** $t=139s$ **f)** $t=170s$ **g)** $t=360s$ **h)** $t=400s$
 376

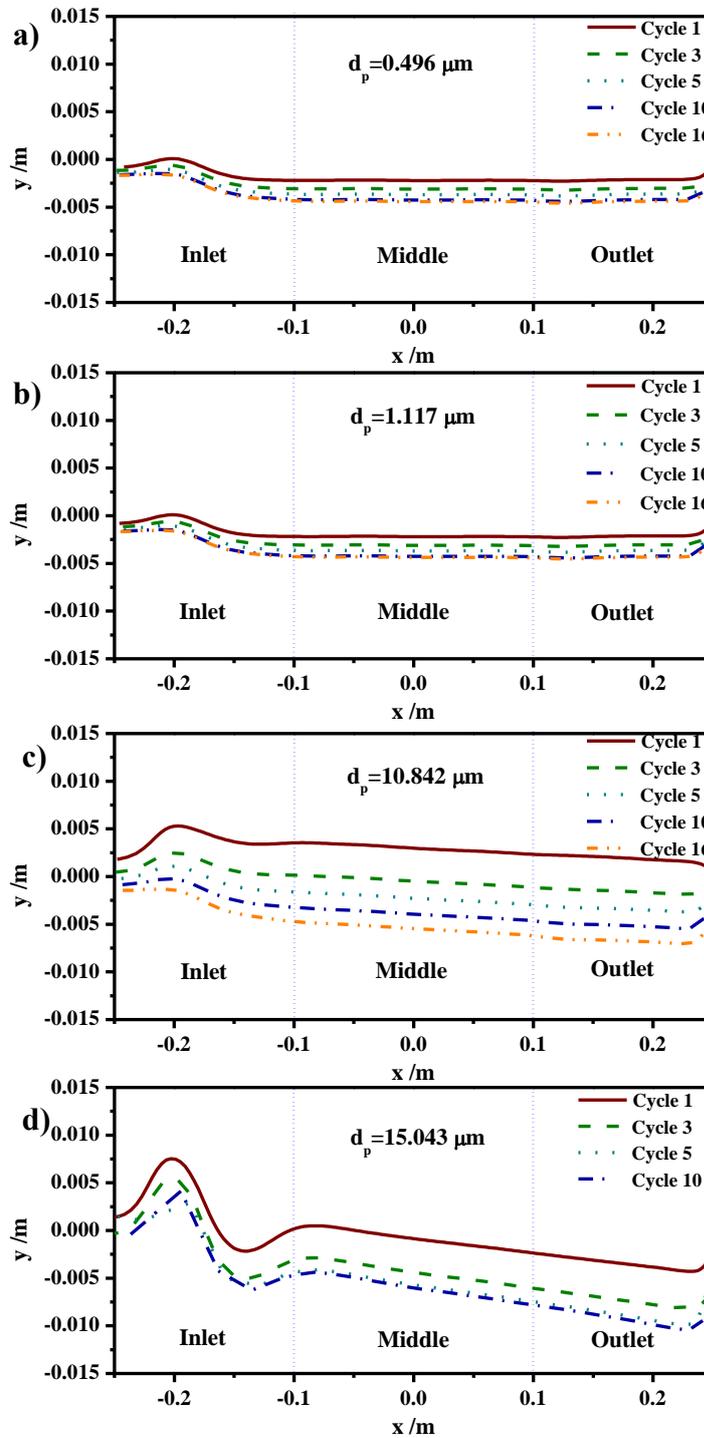
377 **Fig 8** shows the trajectory and evolution process of a single particle counted from the
 378 simulation result. For there are so many particles in the tube, we just choose four particles which
 379 we think are very representative, and the particles size are $0.496\ \mu\text{m}$, $1.117\ \mu\text{m}$, $10.842\ \mu\text{m}$ and
 380 $15.043\ \mu\text{m}$, respectively. We compute the locations of particles in the main tube at any time. As
 381 the particles moves in the circulating system, we also distinguished the different number of cycles
 382 of the particle trajectory. The particle with size of $0.496\ \mu\text{m}$ is selected to represent the little
 383 particle and it can be found that its trajectory in the middle part is almost a straight line. At the
 384 same time, we also found that in the first 10 cycles, the particle trajectory showed a gradual
 385 downward trend, but after the 10th cycle, the particle trajectory remained almost unchanged. In
 386 addition, we also selected the particles with a diameter of $1.117\ \mu\text{m}$ and this part of particles is
 387 treated as median-sized particles. From the **Fig 8 b)**, we can see that the trajectory of this kind of
 388 particles is very similar to that of $0.496\ \mu\text{m}$ particle, which looks straight in the middle part, and
 389 will basically stabilize in one trajectory after the 10th cycle. Furthermore, we can conclude that the

390 motion character of the last mentioned two kinds of particles is that they follow the fluid and keep
391 moving on a straight line in the main tube, that is to say, the settling velocity generated by gravity
392 is very small. After the discussion of small particles, it is necessary to discuss the large particle.
393 Therefore, particle with size of 10.842 μm is also selected. It can be found that this kind of particle
394 is seriously affected by gravity, or the trajectory of this particle in the middle part of the main tube
395 is in the downward inclination. And with the increase of the number of cycles, its trajectory in the
396 main tube shows a very obvious downward trend even after 10th cycles. Finally, we chose the
397 particle size of 15.043 μm , because this particle is the largest particle which is still in motion at
398 400 s in this simulation and can represent the coarse particle. It can be seen that the effect of
399 gravity on the particle is more obvious, and the downward inclination of its trajectory in the main
400 tube is more obvious. But for this particle, we can only take 10 cycles, because the settlement
401 occurred before the 16th cycle, which shows that such particles in the pipeline cannot maintain a
402 continuous circulating movement, and eventually will settle down.

403 **4.4 Particle size distribution**

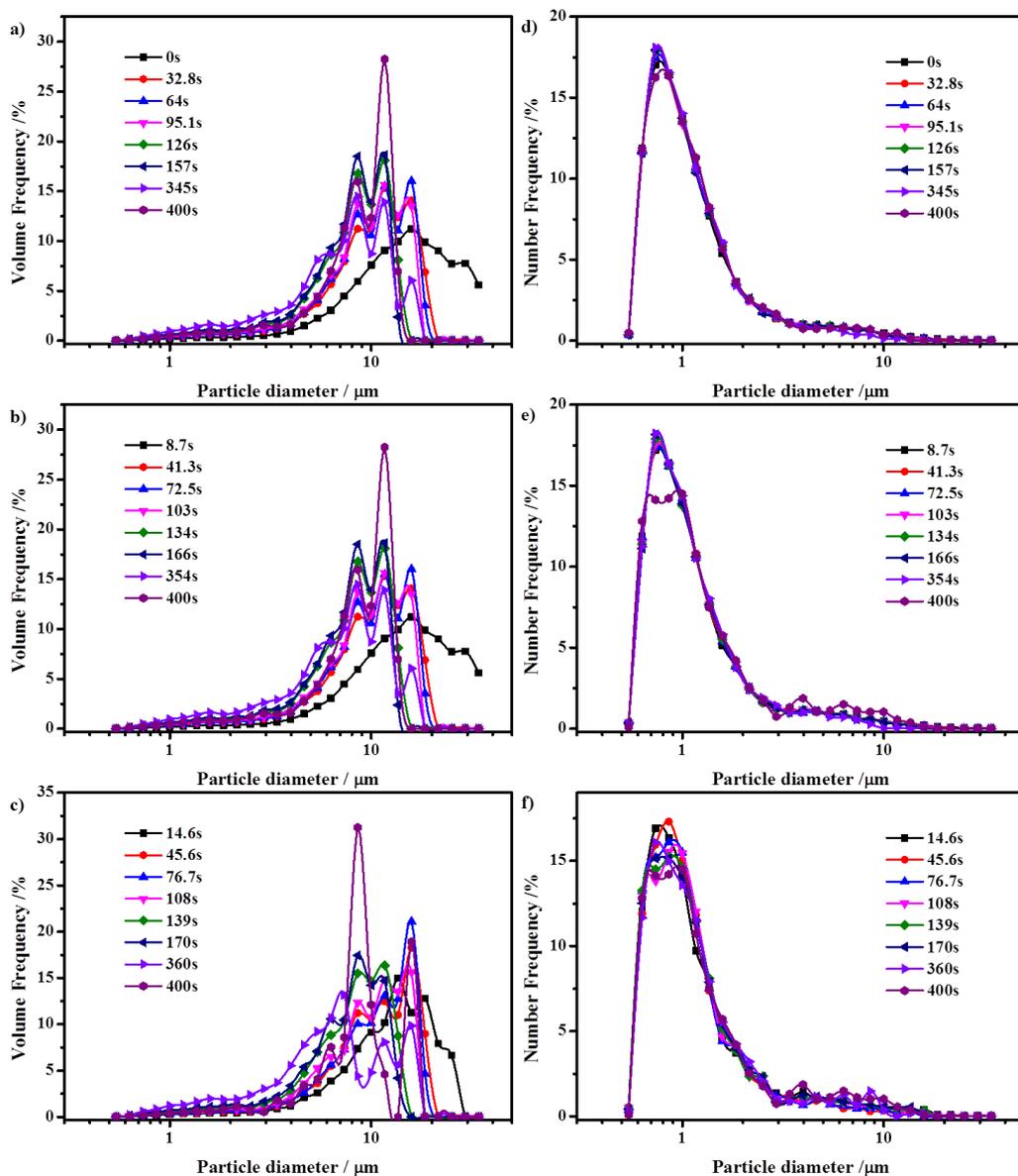
404 The particle size distribution in the different parts of the main tube counted from the simulation
405 results is showed in [Fig. 9](#). Both of the number and volume particle size distributions are taken
406 into account, also the three parts of the tube are computed respectively. For the convenience of
407 description, according to the results in the section 4.3, we can call particles of 1-10 μm as small
408 particles, and particles of more than 10 μm as large particles. It can be found that the volume
409 particle size distribution is a unimodal distribution at the first peak. As the time goes on, the
410 particle size distribution curve gradually becomes an obvious bimodal distribution. If one look at
411 the position and corresponding size of the two peaks separately, one can see that the peak of large
412 particles moves to the left with time, and the height of it increases obviously with time. This trend
413 also proves that the particles with a diameter more than 15 μm are greatly affected by gravity and
414 easy to settle down, which results in some loss of the particles. Therefore, the particles with a
415 diameter slightly smaller than 15 μm occupy a higher volume fraction in the large particles, and it
416 form the peak of large particles in the particle size distribution curve. In addition, it can be seen
417 that the position of the peak of small particles gradually move to left with time. This phenomenon
418 shows that the small particles almost have no loss in the transport process; as a result, volume
419 frequency of this part of particles has obvious advantages especially when the large particles lost.

420 And it can be also seen that the particles size distribution based on number seems not have much
 421 change on the different times of the peaks as the experiment shows, associated with Fig 7, we
 422 further confirm that the group of particles which have the higher velocity will repeat to appear on
 423 the region for counting, so the number particles size distribution is almost no change.



424
 425 **Fig 8** particles motion trace in the main tube a) $d_p = 0.496 \mu\text{m}$, b) $d_p = 1.117 \mu\text{m}$ c) $d_p = 10.842 \mu\text{m}$,
 426 d) $d_p = 15.043 \mu\text{m}$

427 From the prospective of photocatalytic reaction, amount of the particles especially the smaller
 428 particles are more benefit to the reaction efficiency. The result derived from the simulation
 429 inspires that in the circulating system, effects of gravity is only remarkable for the large particles
 430 which the size are more than 10 μm , and a very low circulating flowrate such as 15 L/h is enough
 431 to make sure the small particles keep suspending. And for the amount of the small particles are
 432 quite larger than that of large particles, adopting a lower flowrate in the photocatalytic application
 433 is reasonable and also energy efficient.



434
 435 **Fig.9** Particles size distribution based on volume obtained from a) inlet part, b) middle part, c)
 436 outlet part, and Particles size distribution based on number obtained from d) inlet part, e) middle
 437 part, f) outlet part

438

439 **5 Conclusions**

440 The suspension of micro- or nanoscale particles in liquid is used widely and the circulating flow
441 is also regularly chosen as the flow form in the application of this kind of two-phase fluid. In this
442 study, a DEM method is used in a simulation of the circulating laminar flow system to explain the
443 transport phenomenon of particles. It can be found the particle concentration and size distribution
444 varies with time and appears periodical oscillation, which is related to the non-uniform velocity
445 distribution of particles in the radial direction of tube. In this process, the particle size distribution
446 changes from unimodal distribution to bimodal distribution, for the particles with a size less than
447 10 μm almost have no loss in the transport process. Furthermore, particles tend to gather in the
448 middle region of the tube after several cycles, which is also the region with the higher velocity in
449 the shear flow. Also, with the effects of gravity force, the integral particles group tends to
450 distribute in the region a little lower than the center.

451 As a case study, this study mainly focus on the particle transport phenomenon in photocatalytic
452 reactor, so the kind of particles and base liquid, operating mode, initial condition and other
453 parameters were specified to photocatalytic reactor. However, this simulation can also be extended
454 to other applications, such as heat transfer enhancement by nanofluid, general solid-liquid reaction
455 et al. And it can be implemented by simply replacing some parameters or introducing some
456 correlations in the existing model, which is believed to be useful for understanding the particle
457 transport phenomenon in different applications.

458 **Acknowledgements**

459 The authors gratefully acknowledge the financial support of the National Natural Science
460 Foundation of China (No. 51776165, 51888103) and the financial support from Royal
461 Society-Newton Advanced Fellowship grant (NAF\R1\191163).

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464 **References:**

465 [1] Jing DW, Song DX. Optical properties of nanofluids considering particle size distribution:
466 Experimental and theoretical investigations. RENEW SUST ENERG REV. 2017; 78:452-65.

467 [2] Kudo A, Miseki Y. Heterogeneous photocatalyst materials for water splitting. CHEM SOC REV.
468 2009; 38:253-78.

469 [3] Jing DW, Guo LJ, Zhao L, Zhang XM, Liu H, Li MT, et al. Efficient solar hydrogen production
470 by photocatalytic water splitting: From fundamental study to pilot demonstration. INT J HYDROGEN

471 ENERG. 2010; 35:7087-97.

472 [4] Vajjha RS, Das DK. A review and analysis on influence of temperature and concentration of
473 nanofluids on thermophysical properties, heat transfer and pumping power. INT J HEAT MASS TRAN.
474 2012; 55:4063-78.

475 [5] Fajrina N, Tahir M. A critical review in strategies to improve photocatalytic water splitting
476 towards hydrogen production. INT J HYDROGEN ENERG. 2019; 44:540-77.

477 [6] Kim SJ, Bang IC, Buongiorno J, Hu LW. Effects of nanoparticle deposition on surface wettability
478 influencing boiling heat transfer in nanofluids. APPL PHYS LETT. 2006; 89:153107.

479 [7] KUDO A. Development of photocatalyst materials for water splitting. INT J HYDROGEN
480 ENERG. 2006; 31:197-202.

481 [8] Luo B, Song R, Jing D. ZnCr LDH nanosheets modified graphitic carbon nitride for enhanced
482 photocatalytic hydrogen production. INT J HYDROGEN ENERG. 2017; 42:23427-36.

483 [9] Pulido Melián E, González Díaz O, Ortega Méndez A, López CR, Nereida Suárez M, Doña
484 Rodríguez JM, et al. Efficient and affordable hydrogen production by water photo-splitting using
485 TiO₂-based photocatalysts. INT J HYDROGEN ENERG. 2013; 38:2144-55.

486 [10] Ma Z, Cui Z, Lv Y, Sa R, Wu K, Li Q. Three-in-One: Opened Charge-transfer channel, positively
487 shifted oxidation potential, and enhanced visible light response of g-C₃N₄ photocatalyst through K and
488 S Co-doping. INT J HYDROGEN ENERG. 2020; 45:4534-44.

489 [11] Fang W, Shangguan W. A review on bismuth-based composite oxides for photocatalytic
490 hydrogen generation. INT J HYDROGEN ENERG. 2019; 44:895-912.

491 [12] Li X, Yu J, Low J, Fang Y, Xiao J, Chen X. Engineering heterogeneous semiconductors for solar
492 water splitting. J MATER CHEM A. 2015; 3:2485-534.

493 [13] Ren YX, Zhao L, Jing DW, Guo LJ. Investigation and modeling of CPC based tubular
494 photocatalytic reactor for scaled-up hydrogen production. INT J HYDROGEN ENERG. 2016;
495 41:16019-31.

496 [14] Cao F, Liu H, Wei Q, Zhao L, Guo L. Experimental study of direct solar photocatalytic water
497 splitting for hydrogen production under natural circulation conditions. INT J HYDROGEN ENERG.
498 2018; 43:13727-37.

499 [15] Yang Y, Wei Q, Hou J, Liu H, Zhao L. Solar concentrator with uniform irradiance for particulate
500 photocatalytic hydrogen production system. INT J HYDROGEN ENERG. 2016; 41:16040-7.

501 [16] Wei Q, Yang Y, Liu H, Hou J, Liu M, Cao F, et al. Experimental study on direct solar
502 photocatalytic water splitting for hydrogen production using surface uniform concentrators. INT J
503 HYDROGEN ENERG. 2018; 43:13745-53.

504 [17] Li L, Chen R, Liao Q, Zhu X, Wang G, Wang D. High surface area optofluidic microreactor for
505 redox mediated photocatalytic water splitting. INT J HYDROGEN ENERG. 2014; 39:19270-6.

506 [18] Yan H, Yang J, Ma G, Wu G, Zong X, Lei Z, et al. Visible-light-driven hydrogen production with
507 extremely high quantum efficiency on Pt - PdS/CdS photocatalyst. J CATAL. 2009; 266:165-8.

508 [19] Liu M, Wang L, Max Lu G, Yao X, Guo L. Twins in Cd_{1-x}Zn_xS solid solution: Highly efficient
509 photocatalyst for hydrogen generation from water. ENERG ENVIRON SCI. 2011; 4:1372.

510 [20] Dong Q, Fang Y, Shao Y, Mulligan P, Qiu J, Cao L, et al. Metal-free efficient photocatalyst for
511 stable visible water splitting via a two-electron pathway. . SCIENCE. 2015; 347:967-70.

512 [21] Wang Q, Hisatomi T, Jia Q, Tokudome H, Zhong M, Wang C, et al. Scalable water splitting on
513 particulate photocatalyst sheets with a solar-to-hydrogen energy conversion efficiency exceeding 1%.
514 NAT MATER. 2016; 15:611-5.

515 [22] Barresi A, Baldi G. Solid dispersion in an agitated vessel effect of particle shape and density.
516 CHEM ENG SCI. 1987; 12:2969-72.

517 [23] Huang JK, Lu YJ, Wang H. A new quantitative measurement method for mixing and segregation
518 of binary-mixture fluidized bed by capacitance probe. CHEM ENG J. 2017; 326:99-108.

519 [24] Wang K, Liu G, Liu ZG, Wu J, Yi LT, Zhang JL, et al. Acoustic sensor approaches for sand
520 detection in sand – water two-phase flows. POWDER TECHNOL. 2017; 320:739-47.

521 [25] Felder S, Chanson H. Phase-detection probe measurements in high-velocity free-surface flows
522 including a discussion of key sampling parameters. EXP THERM FLUID SCI. 2015; 61:66-78.

523 [26] Nocentini M, Pinelli D, Magelli F. Dispersion coefficient and settling velocity of the solids in
524 agitated slurry reactors stirred with multiple rushton turbines. CHEM ENG SCI. 2002; 57:1877-84.

525 [27] Roshani GH, Nazemi E, Feghhi SAH, Setayeshi S. Flow regime identification and void fraction
526 prediction in two-phase flows based on gamma ray attenuation. MEASUREMENT. 2015; 62:25-32.

527 [28] Xu YQ, Xu CB, Guan ZC, Liu YW, Tian Y, Sheng YN, et al. Numerical simulation method of
528 ultrasonic wave propagation in gas-liquid two-phase flow of deepwater riser. MECH SYST SIGNAL
529 PR. 2019; 118:78-92.

530 [29] Liu L, Fang ZY, Wu YP, Lai XP, Wang P, Song K. Experimental investigation of solid-liquid
531 two-phase flow in cemented rock-tailings backfill using Electrical Resistance Tomography. CONSTR
532 BUILD MATER. 2018; 175:267-76.

533 [30] Black DL, McQuay MQ, Bonin MP. Laser-based techniques for particle-size measurement: A
534 review of sizing methods and their industrial applications. PROG ENERG COMBUST. 1996;
535 22:267-306.

536 [31] Levoguer C. Using laser diffraction to measure particle size and distribution. MET POWDER
537 REP. 2013; 68:15-8.

538 [32] Ma ZH, Merkus HG, de Smet J, Heffels C, Scarlett B. New developments in particle
539 characterization by laser diffraction: size and shape. POWDER TECHNOL. 2000; 111:66-78.

540 [33] Hu XW, Guo LJ. Numerical investigations of catalyst – liquid slurry flow in the photocatalytic
541 reactor for hydrogen production based on algebraic slip model. INT J HYDROGEN ENERG. 2010;
542 35:7065-72.

543 [34] Hatami M, Hosseinzadeh K, Domairry G, Behnamfar MT. Numerical study of MHD two-phase
544 Couette flow analysis for fluid-particle suspension between moving parallel plates. J TAIWAN INST
545 CHEM E. 2014; 45:2238-45.

546 [35] Hedayati F, Domairry G. Nanoparticle migration effects on fully developed forced convection of
547 TiO₂ – water nanofluid in a parallel plate microchannel. PARTICUOLOGY. 2016; 24:96-107.

548 [36] Feng YQ, Yu AB. Assessment of Model Formulations in the Discrete Particle Simulation of Gas–
549 Solid Flow. IND ENG CHEM RES. 2004; 43:8378-90.

550 [37] Morris JF, Brady JF. Pressure-driven flow of a suspension: Buoyancy effects. INT J
551 MULTIPHAS FLOW. 1998; 24:105-30.

552 [38] Peng ZB, Doroodchi E, Evans G. DEM simulation of aggregation of suspended nanoparticles.
553 POWDER TECHNOL. 2010; 204:91-102.

554 [39] Chaumeil F, Crapper M. Using the DEM-CFD method to predict Brownian particle deposition in
555 a constricted tube. PARTICUOLOGY. 2014; 15:94-106.

556 [40] Jing DW, Liu H, Zhang XH, Zhao L, Guo LJ. Photocatalytic hydrogen production under direct
557 solar light in a CPC based solar reactor: Reactor design and preliminary results. ENERG CONVERS
558 MANAGE. 2009; 50:2919-26.

559 [41] Li D, Xiong K, Li W, Yang ZH, Liu C, Feng X, et al. Comparative Study in Liquid-Phase
560 Heterogeneous Photocatalysis: Model for Photoreactor Scale-Up. *IND ENG CHEM RES.* 2010;
561 49:8397-405.

562 [42] Malato S, Blanco J, Vidal A, Richter C. Photocatalysis with solar energy at a pilot-plant scale: an
563 overview. *APPL CATAL B-ENVIRON.* 2002; 37:1-15.

564 [43] Ochoa-Gutiérrez KS, Tabares-Aguilar E, Mueses MÁ, Machuca-Martínez F, Li Puma G. A Novel
565 Prototype Offset Multi Tubular Photoreactor (OMTP) for solar photocatalytic degradation of water
566 contaminants. *CHEM ENG J.* 2018; 341:628-38.

567 [44] Xing Z, Zong X, Pan J, Wang LZ. On the engineering part of solar hydrogen production from
568 water splitting: Photoreactor design. *CHEM ENG SCI.* 2013; 104:125-46.

569 [45] Cao F, Liu H, Wei Q, Zhao L, Guo L. Experimental study of direct solar photocatalytic water
570 splitting for hydrogen production under natural circulation conditions. *INT J HYDROGEN ENERG.*
571 2018; 43:13727-37.

572 [46] Segré G, Silberberg A. Behaviour of macroscopic rigid spheres in Poiseuille flow Part 2.
573 Experimental results and interpretation. *J FLUID MECH.* 1962; 14:136.

574 [47] Geng J, Tang J, Wang Y, Huang Z, Jing D, Guo L. Attenuated Periodical Oscillation
575 Characteristics in a Nanoscale Particle-Laden Laminar Flow. *IND ENG CHEM RES.* 2020;
576 59:8018-27.

577 [48] Kalantary RR, Shahamat YD, Farzadkia M, Esrafil A, Asgharnia H. Photocatalytic degradation
578 and mineralization of diazinon in aqueous solution using nano-TiO₂(Degussa, P25): kinetic and
579 statistical analysis. *DESALIN WATER TREAT.* 2015; 55:555-63.

580 [49] Cao F, Wei QY, Liu H, Lu N, Zhao L, Guo LJ. Development of the direct solar photocatalytic
581 water splitting system for hydrogen production in Northwest China: Design and evaluation of
582 photoreactor. *RENEW ENERG.* 2018; 121:153-63.

583 [50] Wei QY, Yang Y, Hou JY, Liu H, Cao F, Zhao L. Direct solar photocatalytic hydrogen generation
584 with CPC photoreactors: System development. *SOL ENERGY.* 2017; 153:215-23.

585 [51] Jing DW, Liu H, Zhang XH, Zhao L, Guo LJ. Photocatalytic hydrogen production under direct
586 solar light in a CPC based solar reactor: Reactor design and preliminary results. *ENERG CONVERS
587 MANAGE.* 2009; 50:2919-26.

588 [52] Kowalczyk PB, Drzymala J. Physical meaning of the Sauter mean diameter of spherical
589 particulate matter. *PARTICUL SCI TECHNOL.* 2016; 34:645-7.

590 [53] Wang L, Fang NF, Yue ZJ, Shi ZH, Hua L. Raindrop Size and Flow Depth Control Sediment
591 Sorting in Shallow Flows on Steep Slopes. *WATER RESOUR RES.* 2018; 54:9978-95.

592 [54] Song D, Yang Y, Jing D. Insight into the contribution of rotating Brownian motion of
593 nonspherical particle to the thermal conductivity enhancement of nanofluid. *INT J HEAT MASS
594 TRAN.* 2017; 112:61-71.

595 [55] Cao Z, Tafti DK, Shahnam M. Development of drag correlation for suspensions of ellipsoidal
596 particles. *POWDER TECHNOL.* 2020; 369:298-310.

597 [56] Liu X, Gan J, Zhong W, Yu A. Particle shape effects on dynamic behaviors in a spouted bed:
598 CFD-DEM study. *POWDER TECHNOL.* 2020; 361:349-62.

599 [57] Jin J, Li X, Geng J, Jing D. Insights into the complex interaction between hydrophilic
600 nanoparticles and ionic surfactants at the liquid/air interface. *PHYS CHEM CHEM PHYS.* 2018;
601 20:15223-35.

602 [58] French RA, Jacobson AR, Kim B, Isley SL, Penn RL, Baveye PC. Influence of Ionic Strength, pH,

603 and Cation Valence on Aggregation Kinetics of Titanium Dioxide Nanoparticles. ENVIRON SCI
604 TECHNOL. 2009; 43:1354-9.

605 [59] Song D, Hatami M, Wang Y, Jing D, Yang Y. Prediction of hydrodynamic and optical properties
606 of TiO₂/water suspension considering particle size distribution. INT J HEAT MASS TRAN. 2016;
607 92:864-76.

608 [60] Jing D, Hu S, Zhang Y, Luo J. A modified diffusion-limited cluster aggregation model for
609 accurate prediction of the coagulation and fragmentation process in nanoparticle suspension. J PHYS D
610 APPL PHYS. 2019; 52:455305.

611 [61] Kumar S, Ramkrishna D. On the solution of population balance equations by discretization—I. A
612 fixed pivot technique. CHEM ENG SCI. 1996; 51:1311-32.

613 [62] Chowdhury I, Duch MC, Mansukhani ND, Hersam MC, Bouchard D. Colloidal Properties and
614 Stability of Graphene Oxide Nanomaterials in the Aquatic Environment. ENVIRON SCI TECHNOL.
615 2013; 47:6288-96.

616 [63] Luo B, Song R, Jing D. Particle aggregation behavior during photocatalytic ethanol reforming
617 reaction and its correlation with the activity of H₂ production. COLLOID SURFACE A. 2017;
618 535:114-20.

619 [64] Ebini RH, Sorensen CM. Light scattering studies of the sol-to-gel transition in particulate systems.
620 J COLLOID INTERF SCI. 2019; 556:577-83.

621 [65] Yang M, Li S, Marshall JS. Effects of long-range particle - particle hydrodynamic interaction on
622 the settling of aerosol particle clouds. J AEROSOL SCI. 2015; 90:154-60.

623

624