

1 Probing Mg Intercalation in Tetragonal Tungsten Bronze Framework 2 $V_4Nb_{18}O_{55}$

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14 **ABSTRACT:** While commercial Li-ion batteries offer the highest energy densities of current rechargeable battery technolo-
15 gies, their energy storage limit has almost been achieved. Therefore, there is considerable interest in Mg batteries, which
16 could offer increased energy densities compared to Li-ion batteries if a high-voltage electrode material, such as a transition
17 metal oxide, can be developed. However, there are currently very few oxide materials which have demonstrated reversible
18 and efficient Mg^{2+} insertion and extraction at high voltages; this is thought to be due to poor Mg-ion diffusion kinetics within
19 the oxide structural framework. Herein, the authors provide conclusive evidence of electrochemical insertion of Mg^{2+} into the
20 tetragonal tungsten bronze $V_4Nb_{18}O_{55}$, with a maximum reversible electrochemical capacity of 75 mA h g^{-1} , which corresponds
21 to a magnesiated composition of $Mg_4V_4Nb_{18}O_{55}$. Experimental electrochemical magnesiation/de-magnesiation revealed a
22 large voltage hysteresis with charge/discharge ($1.12 \text{ V vs Mg/Mg}^{2+}$); by limiting magnesiation to a composition of
23 $Mg_2V_4Nb_{18}O_{55}$, this hysteresis can be reduced down to only 0.5 V . Hybrid-exchange Density Functional Theory (DFT) calcula-
24 tions suggest that a limited number of Mg sites are accessible via low-energy diffusion pathways, but that larger kinetic bar-
25 riers need to be overcome to access the entire structure. The reversible Mg-ion intercalation involved concurrent V and Nb
26 redox activity and changes in crystal structure, confirmed by an array of complimentary methods including powder X-ray
27 diffraction, X-ray absorption spectroscopy, and energy-dispersive X-ray spectroscopy. Consequently, it can be concluded that
28 the tetragonal tungsten bronzes show promise as intercalation electrode materials for Mg batteries.

29 Introduction

30 With the growing need for high energy-density batteries,
31 the development of new energy storage technologies is par-
32 amount. While Li-ion batteries have achieved great com-
33 mercial success,¹ and are the most energy-dense recharge-
34 able storage technology for portable applications, the tech-
35 nology already heavily optimized, and future increases in
36 energy density are predicted to be incremental.² Therefore,
37 there is a need to develop new, low cost, and sustainable
38 battery chemistries for applications such as future electric
39 vehicle and grid storage. Mg metal anodes have a lower pro-
40 pensity to form dendrites (that can cause cell failure) from
41 stripping and plating, which means that it could be used di-
42 rectly as a metallic anode,³⁻⁵ which would offer up to a five-
43 fold increase in anodic volumetric energy density (Mg metal
44 = $3833 \text{ mA h cm}^{-3}$) compared to Li-ion graphite anodes
45 ($\sim 800 \text{ mA h cm}^{-3}$), leading to an overall doubling in energy
46 density if a suitably high-energy density electrode material
47 can be found.⁶ In comparison to Li^+ , the divalent ion Mg^{2+}
48 transfers twice the charge ($2 e^-$ per Mg^{2+} ion) when moving

49 between anode and cathode. This phenomenon decouples
50 the effect of electronic changes from the availability of va-
51 cant sites in the crystal structure, creating opportunities to
52 create novel fundamental insights into the limits of interca-
53 lation electrochemistry. Furthermore, the Mg battery offers
54 a step-change in sustainability; Mg is three orders of magni-
55 tude more abundant than Li in the Earth's crust, and can be
56 harvested from a range of commercial sources, including
57 seawater.⁷ Finally, if Mg batteries were developed, it would
58 be possible to repurpose existing Li-ion battery pilot lines
59 for their manufacture without the large expense and tech-
60 nical challenges of a completely new plant design.

Achieving reversible Mg^{2+} insertion within electrode mate-
61 rial candidates at desirable cathodic potentials ($> 2 \text{ V vs}$
62 Mg/Mg^{2+}) has proved extremely challenging; whilst the
63 Mg^{2+} ion possesses a similar ionic radius to Li^+ ($0.72 \text{ vs } 0.76$
64 Å respectively),⁸ the increased charge of the former, results
65 in stronger bonding interactions with anions in both the
66 host electrode structure and the electrolyte. Consequently,
67 both Mg-ion diffusion in the electrode structure, and Mg-ion
68

1 transfer between electrolyte and electrode, are suggested to 57
2 present much higher kinetic barriers to reversible interca- 58
3 lation compared to Li^+ . Therefore, stable materials which re- 59
4 versibly cycle against Mg metal anodes, have largely con- 60
5 sisted of heavy metal chalcogenides such as sulfides, where 61
6 the soft nature of the sulfide anion leads to weaker anion- 62
7 intercalant interactions. Therefore, more facile Mg interca- 63
8 lation kinetics.⁹⁻¹¹ However, the operating voltages of sul- 64
9 fides vs. Mg/Mg^{2+} are too low to achieve higher energy den- 65
10 sity than even first-generation Li-ion cells.^{12,13} Therefore, 66
11 oxide hosts, which are typically lighter and can allow higher 67
12 voltages vs. Mg/Mg^{2+} , are of great interest as electrode ma- 68
13 terials for devices with high energy densities.¹⁴⁻¹⁸ In com- 69
14 parison, there has also been numerous studies of alterna- 70
15 tive electrode chemistries, including Prussian Blue ana- 71
16 logues,¹⁹ selenides,²⁰ phosphates,²¹ silicates,²² titanates,²³ 72
17 and borates.²⁴ 73
18 Recent studies on Mg battery anodes have confirmed that a 74
19 far more comprehensive and holistic approach is required 75
20 to irrefutably verify reversible Mg intercalation in oxide 76
21 hosts, i.e. chemical evidence of changing Mg content, spec- 77
22 troscopic evidence (e.g. XAS) of changing redox states in the 78
23 host material, and evidence of crystal structure changes 79
24 (from XRD and/or PDF analysis).^{17,25} To date, this compre- 80
25 hensive approach has only been conducted on relatively few 81
26 structures including the λ polymorph of MnO_2 , and the α and 82
27 ζ polymorphs of V_2O_5 .^{14,26-28} Many other reports have not 83
28 provided a full, comprehensive proof of Mg-ion intercala- 84
29 tion by combining chemical, diffraction, and spectroscopic 85
30 methods.²⁹ 86
31 The importance of this holistic approach (required to con- 87
32 firm intercalation) is demonstrated in the literature in sev- 88
33 eral instances. For example, α - V_2O_5 has a wide range of re- 89
34 ported capacities, which heavily depends on the water con- 90
35 tent within the structure or electrolyte.^{27,30} Unfortunately, 91
36 the existence of H_2O is precluded in a Mg battery because of 92
37 its severe incompatibility with the metal anode. Further- 93
38 more, Verrelli *et al.* identified that significant proton (H^+) in- 94
39 tercalation can take place instead of Mg^{2+} under many con- 95
40 ditions, arising from residual water present in the electro- 96
41 lyte, or from H^+ generated from electrolyte decomposition.³¹ 97
42 Furthermore, other oxides (such as α - MoO_3) are prone to 98
43 competing conversion reactions, and irreversible MgO for- 99
44 mation, with cycling.¹⁶ 100
45 For the above reasons, there are a dearth of suitable oxid- 101
46 candidates for Mg battery electrodes, and therefore there 102
47 are significant efforts to find a material with the combina- 103
48 tion of capacity, voltage, and kinetics required for an en- 104
49 ergy-dense, energy-efficient Mg battery. The greatest capac- 105
50 ity achieved for an oxide electrode (α - V_2O_5) in a very dry 106
51 electrolyte is 140 mA h g^{-1} (room temperature), and >280 107
52 mA h g^{-1} (110°C).^{28,32} but the efficiency of the Mg-ion inser- 108
53 tion/removal was poor in both cases, (voltage hysteresis of 109
54 0.7 and 0.8 V, respectively). While unconfirmed at present 110
55 it was suggested the source of these high kinetic barriers 111
56 were strong interactions between Mg^{2+} and the host oxide

lattice. Therefore, structures with intrinsically high intercalant-diffusion properties could remove the kinetic bottleneck observed within oxide electrodes. Fundamentally, the number of oxides that can intrinsically intercalate Mg^{2+} remains small at present, which hampers our ability to understand design rules for effective intercalation.

Bronze-type frameworks such as $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ are composed of corner-sharing polyhedra along the c axis, provide large diffusion planes which can accommodate intercalated ions. The $\text{T-Nb}_2\text{O}_5$ phase, or the related bronze-type and Wadsley-Roth crystallographic shear phases $\text{V}_4\text{Nb}_{18}\text{O}_{55}$, $\text{Nb}_{18}\text{W}_{16}\text{O}_{93}$, and $\text{Nb}_{16}\text{W}_5\text{O}_{55}$, have been shown to have very low barriers to Li^+ diffusion in the bulk phase.^{33,34} For example, the activation energy barrier for Li^+ self-diffusion in $\text{T-Nb}_2\text{O}_5$ was found to be in the range 58 to 91 meV (from NMR studies) in comparison a typical Li-ion intercalation anode (760 meV for spinel-structured $\text{Li}_4\text{Ti}_5\text{O}_{12}$).^{35,36}

The advantages of bronze-type phases can be understood using a comparison with the related structure of ReO_3 , in which octahedral Re^{6+} ions form a corner sharing MO_6 octahedral framework (i.e. a perovskite with vacant A sites). ReO_3 can accommodate Li^+ in the vacant interstices; however, correlated rotations to the MO_6 octahedra upon Li incorporation act to stabilize the Li ions, distorting the structure, limiting ionic mobility and resulting in capacity fade.³⁷ In contrast, bronze structures containing M^{5+} ions, such as $\text{T-Nb}_2\text{O}_5$ and $\text{V}_4\text{Nb}_{18}\text{O}_{55}$, can resist such framework distortions and offer a more rigid network for reversible metal intercalation.^{33,38} The edge-sharing polyhedra present within ab planes provide a critical increase in framework rigidity (in comparison with corner-sharing only frameworks). Upon Li^+ or Mg^{2+} incorporation, the polyhedra are effectively locked in place by shared edges and cannot tilt. Therefore, the inserted ions remain in a 'frustrated' coordination, and can diffuse through the structure with low energy barriers.^{39,40} Bronze frameworks are therefore candidate materials for high Mg^{2+} mobility electrodes.

Herein, we report the electrochemical characterization of $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ (which possesses a distorted $\text{T-Nb}_2\text{O}_5$ structure) in Mg^{2+} electrolytes. The structure was solved in 2010 by Börrnert *et al.*,⁴¹ and first reported as a high-power anode material in Li-ion batteries by Qian *et al.* in 2018.³⁸ $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ shows the intrinsic ability to reversibly intercalate Mg^{2+} with lower hysteresis than known oxide materials (0.5 vs 0.7 V). This was supported by hybrid-exchange DFT investigations, which revealed $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ possessed low-energy Mg diffusion pathways (ca. 0.65 eV activation energy barrier), and suggested design rules for new high Mg-mobility oxide materials based on the tetragonal tungsten bronzes. We confirm that $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ possesses reduced overpotentials and greater energy efficiency than comparable reported oxide materials in dry electrolyte where Mg-ion intercalation is backed-up with a combination of XRD, XAS and elemental evidence (Table 1), albeit with low observed capacity.

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114
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116

1
2**Table 1. A summary of different Mg battery electrode materials and their energy densities.**

Material	Gravimetric Energy Density / Wh kg ⁻¹	Applied Current	Energy efficiency / %	Cycling Temperature / °C	Reference	Degree of characterization of discharged and charged states
Chevrel-Mo ₆ S ₈	77	C/2 (0.3 mA cm ⁻²)	~95	60	11	Fully characterized
Spinel-Ti ₂ S ₄	228	C/20 (10 mA g ⁻¹)	~95	60	13	Fully characterized
CuS nanoparticles	250	C/3 (100 mA g ⁻¹)	~60	50	12	No elemental or structural characterization
TiSe ₂	108	C/22 (5 mA g ⁻¹)	~65	25	20	No redox characterization
orthorhombic-V ₂ O ₅	660	C/20 (15 mA g ⁻¹)	~65	110	28	Fully characterized
Mo _{2.48} VO _{9.93}	~250	C/40 (4 mA g ⁻¹)	~88	25	29	No elemental or redox characterization
MoO ₃	270	C/15 (80 μA cm ⁻²)	~70	25	16	No elemental or redox characterization
FePO ₄ F	99	C/20 (8.5 mA g ⁻¹)	~70	25	21	No elemental or redox characterization
FeSiO ₄	825	C/50 (6.62 mA g ⁻¹)	~75	55	22	No elemental or redox characterization
Cr-doped Li ₄ Ti ₅ O ₁₂	~60	C/10 (17.5 mA g ⁻¹)	~60	25	23	No structural or redox characterization
Ni(CN) ₆ ·nH ₂ O	102	C/5 (10 mA g ⁻¹)	~95	25	19	No elemental/redox characterization, electrolyte contained H ₂ O
V ₄ Nb ₁₈ O ₅₅ (this work) ^a	75	C/3 (15 mA g ⁻¹)	81	110	This work ^a	?????add
V ₄ Nb ₁₈ O ₅₅ (this work) ^b	183	C/3 (15 mA g ⁻¹)	56	110	This work ^b	?????add

^a Where the capacity was limited to 30 mA h g⁻¹. ^b where the cell was cycled to the limits of the electrolyte stability window, corresponding to 75 mA h g⁻¹.

5 Experimental Section

6 Phase-pure V₄Nb₁₈O₅₅ nanoparticles were synthesized *via* a
7 two-step process of Continuous Hydrothermal Flow Syn-
8 thesis (CHFS) of an amorphous Nb/V oxide solid-solution,
9 followed by a 5 h heat-treatment step in air at 600 °C. The
10 details of the CHFS process are described in detail in previ-
11 ous publications, and a diagram of the process is given in
12 Figure S1a in the Supporting Information (SI).^{42,43} Briefly,
13 the CHFS process consisted of three high-pressure dia-
14 phragm pumps, which provided three separate pressurized
15 feeds into a Confined Jet Mixer (CJM, Figure S1b) which is
16 where mixing of reagents and initial precipitation occurred.
17 A V⁴⁺ aqueous precursor was prepared using protocol de-
18 scribed in the SI. The V⁴⁺ precursor was diluted with D.I. wa-
19 ter and combined with ammonium niobium oxalate hydrate
20 (hydrated C₄H₄NNbO₉, 99.99%, Sigma Aldrich, Dorset, UK)
21 to give a combined precursor solution with 36.4 mM and

22 163.6 mM V⁴⁺ and Nb⁵⁺, respectively, i.e. a 4:18 V:Nb molar
23 ratio, with a total metal concentration of 0.2 M. This com-
24 bined precursor solution was pumped (via pump 2, flow
25 rate of 40 mL min⁻¹) to mix with an aqueous flow of room
26 temperature D.I. water (in-flow dilution, delivered via
27 pump 3, flow rate of 40 mL min⁻¹) in a T-piece. The resulting
28 solution was subsequently mixed with a superheated water
29 feed at 450 °C (via pump 1, flow rate of 80 mL min⁻¹) in the
30 CJM. This CJM was designed to eliminate any possible block-
31 ages under typical experimental conditions, and to enable
32 rapid, turbulent mixing [Reynolds number (Re) > 6900] of
33 the incoming hot and cold feeds. The mixing temperature in
34 the CJM was *ca.* 335 °C, and the residence time was *ca.* 7.2 s
35 before the nanoparticle slurry was cooled in-flow via a pipe-
36 in-pipe heat exchanger and continuously collected at *ca.* 40
37 °C from the exit of a backpressure regulator valve [which
38 maintained the pressure of the whole system at 24 MPa
39 (Tescom Model 26-1700, Tescom BPR, Hamilton, UK)]. The

1 resulting particle slurry was cleaned by dialysis until the 60
2 conductivity of the supernatant was $< 60 \mu\text{S}^{-1}$. The concen- 61
3 trated aqueous sludge was then frozen at -40°C and slowly 62
4 heated from -40 to 25°C in a VirTis Genesis 35 XL at *ca.* 13 63
5 Pa for 24 h. The resulting amorphous V and Nb oxide was 64
6 heat-treated at 600°C for 5 hours (1°C min^{-1} ramp rate) in 65
7 a furnace (air atmosphere) to yield the phase-pure 66
8 $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ product. 67
9 The Powder X-Ray Diffraction (PXRD) pattern of the pris- 68
10 tine material was collected using a Stoe StadiP diffractome- 69
11 ter in transmission mode (coupled θ - 2θ geometry), using 70
12 $\text{Mo-K}\alpha$ radiation ($\lambda = 0.709 \text{ \AA}$), with the sample contained in 71
13 a borosilicate capillary. Three patterns were collected in the 72
14 2θ range of 2 to 40° , with a step size of 0.5° in 2θ and a col- 73
15 lection time of 100 s per step, and these three patterns were 74
16 combined into one to generate the final diffraction pattern. 75
17 The PXRD pattern of the demagnesiated (charged) elec- 76
18 trode was collected on Bruker D8 Advance X-ray Diffrac- 77
19 tometer ($\text{Cu-K}\alpha$ radiation, $\lambda = 1.54 \text{ \AA}$) in reflection mode. 78
20 Synchrotron PXRD measurements ($\lambda = 0.413 \text{ \AA}$) were per- 79
21 formed on an electrochemically magnesiated $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ 80
22 sample at 11-BM, Advanced Photon Source, Argonne in a 81
23 Kapton capillary. Rietveld refinement was performed using 82
24 RIETICA software. 83
25 Transmission Electron Microscopy (TEM) and Energy-Dis- 84
26 persive Spectroscopy (EDS) images of the pristine 85
27 $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ were collected using a 200 keV JEOL JEM-2100 86
28 TEM fitted with an Oxford Instruments X-MaxN 80-T Silicon 87
29 Drift Detector (SDD, GATAN) and processed with AZtec® 88
30 software. TEM and EDX on magnesiated and demagnesiated 89
31 electrodes was performed using JEOL-3010 microscope op- 90
32 erated at 300 kV. Samples were prepared for imaging by 91
33 dispersing them in methanol (with the aid of ultrasoni- 92
34 cation) and pipetting a drop of the resulting suspension 93
35 onto carbon-coated Cu TEM grids (Agar Scientific, UK). The 94
36 solvent was then evaporated, leaving a deposition of parti- 95
37 cles on the grid which was then inserted into the TEM mi- 96
38 croscope for analysis. 97
39 X-ray absorption spectroscopy (XAS) measurements in 98
40 transmission mode were taken of the sample at the vana- 99
41 dium K-edge (*ca.* 5465 eV) and Nb K-edge (*ca.* 18986 eV) at 100
42 the B18 beamline at Diamond Light Source (UK). Nb L_{3-} , V 101
43 $L_{2,3-}$ and O K-edge XAS was performed at the beamline 4-ID- 102
44 C, Advanced Photon Source, Argonne National Laboratory, 103
45 USA. At 4-ID-C, spectroscopy was completed simultane- 104
46 ously in both the Total-Electron-Yield (TEY) and Total-Flu- 105
47 orescence Yield (TFY) mode utilizing photocurrent for the 106
48 TEY and a silicon drift diode detector for the TFY, in order 107
49 to make direct surface to bulk comparisons. Data were ob- 108
50 tained at a spectral resolution of $\sim 0.2 \text{ eV}$, with a 2 s dwell 109
51 time. 3 scans were performed on each sample, at each ab- 110
52 sorption edge, and scans were averaged in order to maxim- 111
53 ize the signal to noise ratio. The Nb L_{3-} edge was scanned in 112
54 the range 2360 to 2440 eV and V $L_{2,3-}$ and O K-edges were 113
55 scanned in the range 500 to 560 eV. The V and O energy 114
56 scales were normalized using a SrTiO_3 standard. 115
57 Electrodes were prepared by hand-grinding the active ma- 116
58 terial with conductive carbon SuperP® (MTI Corp., Califor- 117
59 nia, USA) and binder in an active material:carbon:binder 118

ratio of 60:20:20 wt%. The binder was polyvinylidene difluoride (Solef® PVDF, Solvay Inc., Illinois, USA) dissolved in N-methylpyrrolidone (NMP) (Sigma-Aldrich, Missouri, USA) as a 6 wt% solution. The electrode slurry was cast on $1/2''$ stainless steel 316 mesh electrodes (325 x 325 wire cloth, Gerard Daniel Worldwide, Pennsylvania, USA) or 30 μm thick Al foil and allowed to dry overnight in air. These electrodes were dried under an IR lamp, and pressed at 8 tons between two pieces of $8 \times 8 \text{ cm}^2$ stainless steel to improve contact between active material and current collector, and further dried at 100°C overnight under vacuum (model 1407 Shel Lab Economy Vacuum Oven). The electrodes possessed active material loadings in the range 2.2 to 4.1 mg cm^{-2} . For the synchrotron XRD electrode, the magnesiated material was extracted from the cell and sealed in a borosilicate glass capillary.

Dried electrodes were assembled into 2032 coin-type cells in an Ar-filled glovebox (water and oxygen $\leq 0.1 \text{ ppm}$). Electrochemical cells contained an activated carbon cloth as a counter electrode (ACC, ACC-5092-20, American Technical Trading Inc., Pleasantville, New York, USA), a glass fiber separator (grade 691, 28297-289, VWR International, USA) and $0.5 \text{ M Mg}[\text{N}(\text{SO}_2)_2(\text{CF}_3)_2]_2 \cdot (\text{C}_9\text{H}_{20}\text{N})(\text{N}(\text{SO}_2)_2(\text{CF}_3)_2)$ ionic liquid electrolyte (abbreviated as $\text{MgTFSI}_2\text{-PY}_{14}\text{TFSI}$) with low H_2O content ($\sim 43 \text{ ppm}$). The electrolyte was made by dissolving MgTFSI_2 (M1208c, Solvionic, France) in $\text{PY}_{14}\text{TFSI}$ (Pyr0408a, Solvionic, France). For these 2-electrode cells, the potential of cathode was calibrated by considering the ACC anode potential, which was originally $2.2 \text{ V vs Mg/Mg}^{2+}$ and linearly proportional to the state-of-charge (SoC).³² The polarization of the ACC electrode was quantified by measuring the voltage change of a symmetric ACC|ACC cell with SoC, using the same electrolyte and cycling temperature used for the ACC| $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ cells. Subtracting the capacitive contributions from the measured potential vs. ACC allowed for the determination of the cathodic voltage vs. carbon and vs. Mg/Mg^{2+} as a function of the SoC.

Electrochemistry was carried out at 110°C in the potential ranges 1.2 to -1.7 V or 1.2 to -2.5 V vs. ACC as indicated in the text. The charge/discharge rate (15 mA g^{-1} which corresponds to a C/3 C-rate assuming a maximum capacity of 75 mA h g^{-1}) was galvanostatically controlled by a Bio-Logic VMP3 potentiostat. After magnesiation and demagnesiation of the $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ samples, the electrodes were recovered, rinsed in acetonitrile five times, and dried at room temperature under vacuum for 1 minute before characterization.

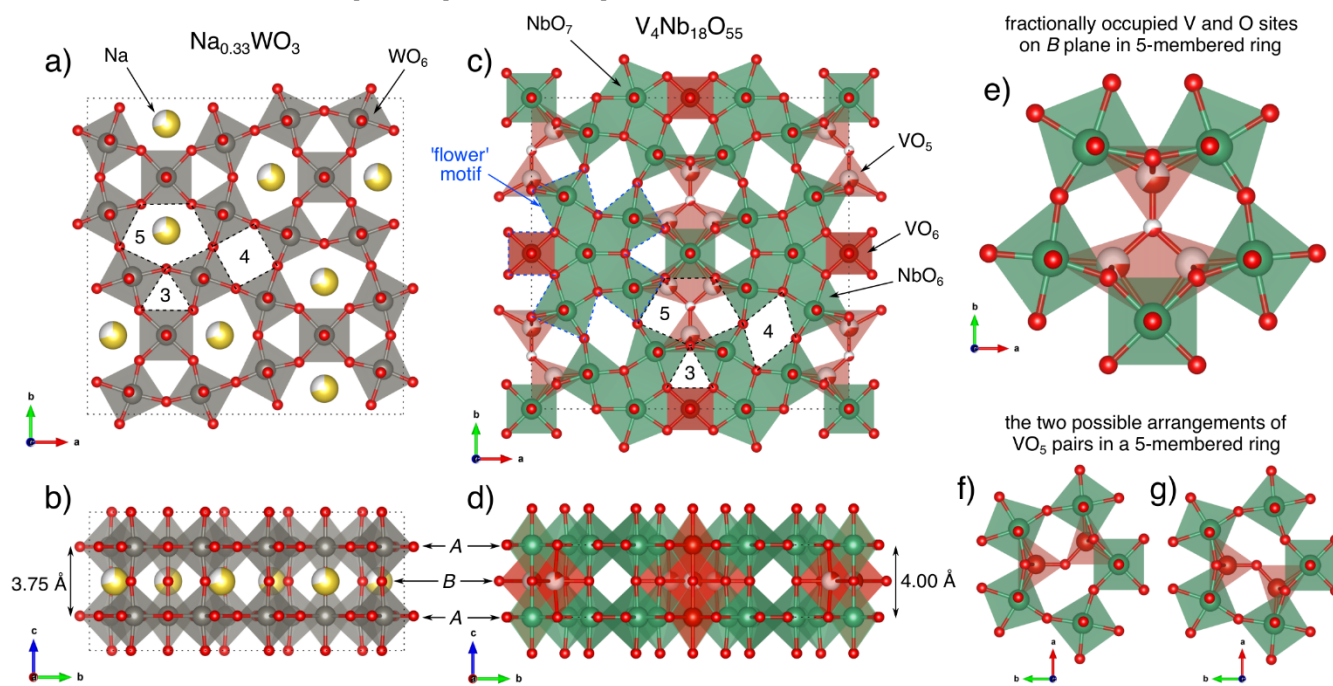
To investigate the discharged state, an electrode was discharged to -1.7 V vs ACC , yielding a capacity of 68 mA h g^{-1} (on a stainless-steel grid electrode substrate, Figure S3a), which corresponds to a theoretical stoichiometry of $\text{Mg}_{3.5}\text{V}_4\text{Nb}_{18}\text{O}_{55}$, with mass loading 4.1 mg cm^{-2} , and will hereafter be labelled as "VNO(Discharged_1)". For the High-Angle Annular Dark-Field Scanning Transmission Electron Microscopy (HAADF-STEM) elemental mapping, an electrode (mass loading 2.89 mg cm^{-2}) was discharged to a capacity of 160 mA h g^{-1} (Figure S3b) and named "VNO(Discharged_2)". For the XRD analysis, an electrode was prepared on Al foil (mass loading 3.8 mg cm^{-2}), as Al is much less attenuating than stainless steel thereby giving a

1 greater signal to noise ratio, and similarly discharged to 32
 2 -1.7 V vs ACC, yielding a similar capacity to 33
 3 VNO(Discharged_1) of 67 mA h g⁻¹ and this sample was la- 34
 4 belled "VNO(Discharged_3)" (Figure S3c). 35
 5 To investigate the reversibility of the Mg insertion/removal 36
 6 process, two charged electrodes were produced; the first 37
 7 charged electrode (mass loading 2.7 mg cm⁻²) was cycled in 38
 8 conditions mimicking the reversible cycling behavior seen 39
 9 in Figure 4a, i.e. a discharge of 75 mA h g⁻¹, followed by a 40
 10 charge of 75 mA h g⁻¹, and this sample is hereafter named 41
 11 "VNO(Charged_1)" (Figure S3d). The second charged elec- 42
 12 trode (mass loading 2.2 mg⁻¹) was discharged to a greater 43
 13 capacity (137 mA h g⁻¹) and subsequently overcharged (300 44
 14 mA h g⁻¹), to further investigate the irreversibility of the Mg 45
 15 insertion reaction, generating an electrode that was labeled 46
 16 as "VNO(Charged_2)" (Figure S3e). First-principles calcula- 47
 17 tions were performed using the periodic Density Functional 48
 18 Theory (DFT) code CRYSTAL17.⁴⁴ Electronic exchange and 49
 19 correlation were approximated using the hybrid- exchange 50
 20 functional PBE0, which is known to give accurate estimates 51
 21 of the band structure of metal oxide systems with highly 52
 22 correlated electrons,⁴⁵ and is efficiently implemented in the 53
 23 CRYSTAL code. All-electron atom-centered Gaussian basis 54
 24 sets were used for all atoms, available from the CRYSTAL 55
 25 online database (www.crystal.unito.it), indicated by the fol- 56
 26 lowing labels: Nb (Nb_986-31(631d)G_dallolio_1996), V 57
 27 (V_86-411d31G_harrison_1993), O (O_8-57
 28 411d1_cora_2005), Li 58
 29 (Li_5-11(1d)G_baranek_2013_LiNbO₃). The Coulomb and 59
 30 exchange series were truncated with thresholds of 10⁻⁷, 60
 31 10⁻⁷, 10⁻⁷, 10⁻⁷ and 10⁻¹⁴. Reciprocal space was sampled

using a Pack-Monkhorst net,⁴⁶ with a shrinking factor of IS = 4 along each periodic direction. The self-consistent field (SCF) procedure was performed up to a convergence threshold of ΔE = 10⁻⁸ Hartree per unit cell. Full geometry optimizations (lattice parameters and atomic positions) were performed using the default convergence criteria in CRYSTAL17. Grimme's semiclassical D3 scheme to account for the effects of dispersion was used to give an improved structural description.⁴⁷⁻⁴⁹

Results and Discussion

V₄Nb₁₈O₅₅ possesses an orthorhombic crystallographic unit cell and belongs to the *Cmmm* space group (no. 65) with lattice parameters $a = 17.355 \text{ \AA}$, $b = 17.698 \text{ \AA}$ and $c = 3.971 \text{ \AA}$. The structure of V₄Nb₁₈O₅₅ is related to the family of tetragonal tungsten bronzes (M_xWO₃, where M = Na, K, Pb, Sn), in which corner-sharing WO₆ octahedra are linked through their axial O-ions in the *c* direction, and equatorial O-ions in the *ab* plane. This results in a WO₃ framework displaying three, four and five-member ring tunnels extending in the *c* direction, with partial occupancy of the 5-membered rings with the large M cations such as Na⁺ (Figure 1a,b). In V₄Nb₁₈O₅₅, however, some of the 5-member rings are occupied by 7-coordinated pentagonal bipyramidal Nb⁵⁺ ions, which are edge-sharing with neighboring octahedral units in the *ab* plane, forming flower-like topology (Figure 1c). These edge-sharing NbO₇ units provide the structural rigidity to avoid correlated polyhedral rotations or other unfavorable distortions upon ion intercalation, which is missing from the M_xWO₃ framework.



61

62 Figure 1. a) A ($\sqrt{2} \times \sqrt{2} \times 2$) unit cell of a prototypical tetragonal tungsten bronze Na_{0.33}WO₃ viewed along the *c* direction and b) viewed
 63 along the *ab* plane, highlighting the metal-rich 'A' and oxygen-rich 'B' planes. c) The crystallographic unit cell of V₄Nb₁₈O₅₅ viewed
 64 along the *c* direction, and d) viewed along the *ab* plane. The 3, 4 and 5-membered rings are indicated in both the structures by
 65 the black dotted lines. e) 5-membered rings, fractionally occupied by V and O-ions. f) and g) Possible geometric arrangements of VO₅ pairs
 66 in the 5-membered rings.

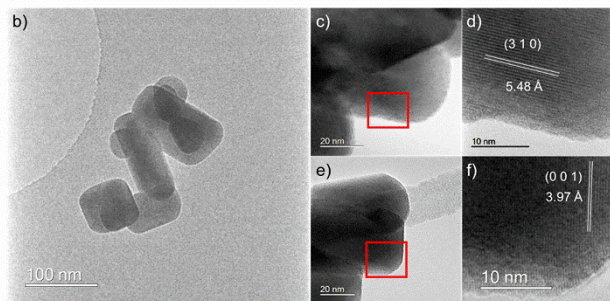
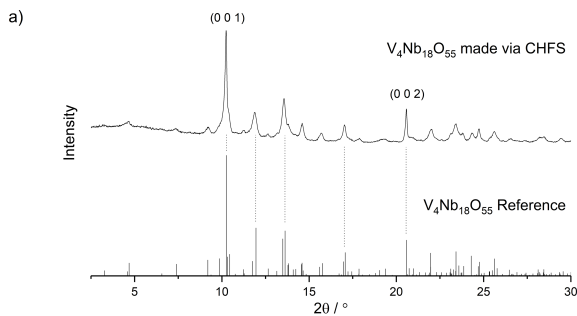


Figure 2. a) The XRD pattern (Mo-K α radiation, $\lambda = 0.709 \text{ \AA}$) of the $V_4Nb_{18}O_{55}$ sample with a $V_4Nb_{18}O_{55}$ reference pattern (ICSD 420781) at the bottom, with the (0 0 1) and (0 0 2) diffraction peaks highlighted (top). Dotted lines have been superimposed on the diagram to highlight peak-matching between the reference pattern and the experimental pattern. TEM images of the $V_4Nb_{18}O_{55}$ sample: b) a representative cluster of $V_4Nb_{18}O_{55}$ nanoparticles, c) a rod identified by TEM analysis, with a selected area for lattice plane analysis shown by a red rectangle, d) lattice plane analysis of the area highlighted in c), e) a separate rod identified by TEM analysis, with a selected area for lattice plane analysis shown by a red rectangle, f) lattice plane analysis of the area highlighted in e).

The octahedra bridging between these NbO_7 units are VO_6 , rather than NbO_6 . A distinctive feature of W and Nb bronze-type frameworks are layered metal-rich 'A' and oxygen-rich 'B' planes when the structure is viewed along ab planar directions (Figures 1b,d). The c -parameter length of $\sim 4 \text{ \AA}$ in $V_4Nb_{18}O_{55}$ is similar to that of $T-Nb_2O_5$, and is typical of structures with axial corner sharing NbO_6 octahedra in this direction. A key structural feature of $V_4Nb_{18}O_{55}$ is the fractional occupation of some of remaining 'empty' (i.e. not occupied by Nb^{5+} ions) 5-membered rings by pairs of trigonal-bipyramidal VO_5 units, which share a bridging O-ion (Figure 1e). These VO_5 pairs occupy $\sim 35\%$ of the 'empty' 5-membered rings, are disordered throughout the lattice, and reside on the 'B' plane, in contrast to the 'structural' Nb and V ions on the 'A' plane. There are two possible arrangements of the VO_5 pairs, indicated in Figures 1f and 1g. Note that the VO_5 pairs are always arranged sharing an O-ion across a corner, and never share an edge. $V_4Nb_{18}O_{55}$ was successfully synthesized in a two-step synthesis process, as described in the Methods section. PXRD data analysis revealed that the material possessed the pure $V_4Nb_{18}O_{55}$ structure ($Cmmm$ space group), with peaks that had an excellent match with the reference pattern (ICSD collect code 420781, Figure 2a). LeBail refinement of the diffraction pattern revealed lattice

parameters of $a = 17.3658(12) \text{ \AA}$, $b = 17.6500(14) \text{ \AA}$, $c = 3.97886(6) \text{ \AA}$, $V = 1219.55(13) \text{ \AA}^3$, which are in concordance with the values reported by Börrnert *et al.* (Figure S2a).⁴¹

The diffraction peaks with hkl indices (0 0 1) and (0 0 2) displayed enhanced intensities with respect to the calculated intensities. This was most visible in the LeBail fit of this material in Figure S2a, where the fit included a (0 0 1) texture term. This suggested that the material formed rod-type particles (with the long axis aligned with the c direction). This was confirmed by TEM analysis, which revealed that the primary particles possessed a rod-like morphology (Figure 2b). The long-axis of these particles was shown to be parallel to the c -axis, (Figure 2b-f), with the (0 0 1) lattice plane normal parallel to, and the (3 1 0) plane normal perpendicular to, the long axis. Particle size distribution analysis (200 particles) revealed an average long-axis dimension of $89 \pm 39 \text{ nm}$ and aspect ratio 1.8 ± 0.7 . Given the small aspect ratio of this sample, the relative sharpness of the (0 0 1) and (0 0 2) diffraction peaks suggested increased disorder within the atomic arrangement within ab planes relative to the stacking order along the c -axis, i.e. the ($h k 0$) diffraction peaks were disproportionately broadened in comparison to the (0 0 l) peaks. As diffusion is expected to occur in the ab plane,³⁶ the shortened diffusion distances resulting from the particle morphology should favor Mg^{2+} intercalation into the $V_4Nb_{18}O_{55}$ lattice. EDS analysis confirmed the expected homogenous distribution of V and Nb ions throughout the $V_4Nb_{18}O_{55}$ particle (Figure S5), and a V:Nb ratio of ca. 4:18 (Figure S6).

Evaluation of the electrochemical properties of $V_4Nb_{18}O_{55}$ was conducted at 110°C in a cell with a Mg^{2+} ionic liquid electrolyte against an activated carbon cloth (ACC) counter-electrode (Figure 3). The electrolyte, 0.5 M $Mg(TFSI)_2$ in $PY_{14}TFSI$, was chosen as it has the greatest stability at cathodic potentials of all known Mg-containing electrolytes (with minimal degradation observed in the range -1.7 V and 1.2 V vs carbon which corresponds to 0.5 V and 3.4 V vs Mg/Mg^{2+}),²⁸ and therefore, the observed cathodic electrochemistry was easier to attribute to cathodic activity as opposed to electrolyte decomposition. A cycling temperature of 110°C was chosen as oxides typically only intercalate significant Mg at elevated temperatures, i.e. 50°C or 110°C .^{14,50}

The discharge behavior to the limit of the voltage stability window (represented by a blue line in Figure 3a and in Figure S3c) revealed two distinct regions of activity; a sloping plateau in the range 2.3 to 1.6 V vs Mg/Mg^{2+} , followed by a region of increasing capacity with further reduction in potential to 0.7 V vs Mg/Mg^{2+} (Figure 3a, S3c). dQ/dV analysis (Figure 3b) revealed that at least three distinct chemical processes occurred during charge; a higher potential redox couple centered at $\sim 2.6 \text{ V}$ vs Mg/Mg^{2+} , and two overlapping redox couples centered at $\sim 1.4 \text{ V}$ vs Mg/Mg^{2+} . The presence of the latter dual redox couple was revealed by a double peak in the charge cycle (found in the range 2.0 to 2.5 V vs Mg/Mg^{2+}), which is not evident in the discharge peak below 1.5 V vs Mg/Mg^{2+} due to overlapping peak locations and possible concurrent electrolyte decomposition. Cycling in this range proved reasonably reversible, with 67% capacity retention over 50 cycles (Figure 3c). The large

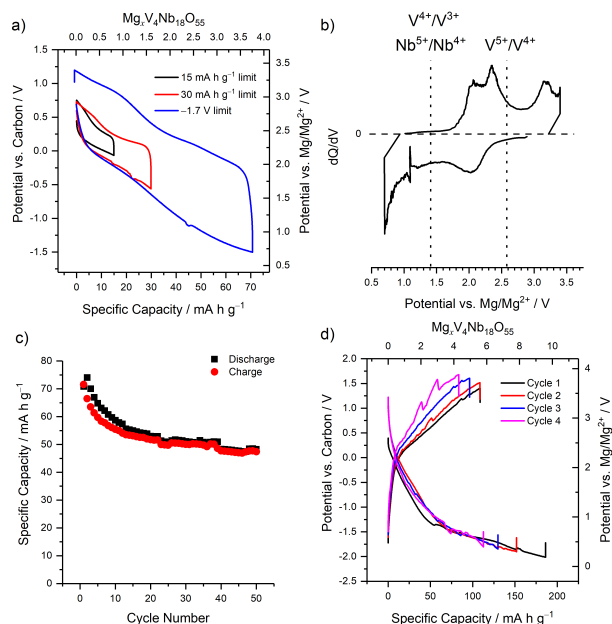


Figure 3. a) The discharge behavior for the first cycle of each depth of discharge, highlighting the increase in overpotential with depth of discharge (mass loading of electrode was 3.0 mg cm^{-2}) b) dQ/dV plot of the first full cycle with a -1.7 V voltage limit vs. ACC on the discharge. c) Capacity of discharge and charge vs cycle number for the -1.7 V vs. ACC discharge limit, using the same electrode as in a). d) Voltage vs Capacity for four cycles to a lower voltage limit of -2.5 V vs. ACC (mass loading of electrode was 2.8 mg cm^{-2}).

overpotentials observed with cycling across the electrolyte stability window (ca. 1.12 V) could be reduced significantly by limiting capacity to 15 and 30 mA h g^{-1} , (giving overpotentials of 0.33 and 0.47 V respectively, see Figure S3a-b), which compares extremely favorably with other oxide materials ($>0.8 \text{ V}$ or higher).²⁸ This gave a round-trip energy efficiency of 81% , which is the highest observed for oxide materials in Mg batteries (Table 1) where Mg intercalation is proven by structural, redox, and elemental analysis. When a discharge limit beyond the electrolyte stability window was used (-2.5 V vs. ACC, Figure 3d), the measured capacity increased significantly (initial discharge capacity of 186 mA h g^{-1}). However, this capacity proved much less reversible and decreased to 113 mA h g^{-1} by the 4th cycle. The discharge capacity was consistently higher than the charge capacity (108 mA h g^{-1} on the first cycle) for these deep discharge tests, which suggested irreversible electrolyte decomposition was responsible for the capacity observed at these low potentials ($<0.7 \text{ V}$ vs Mg/Mg^{2+}). Comparison with literature reports revealed the

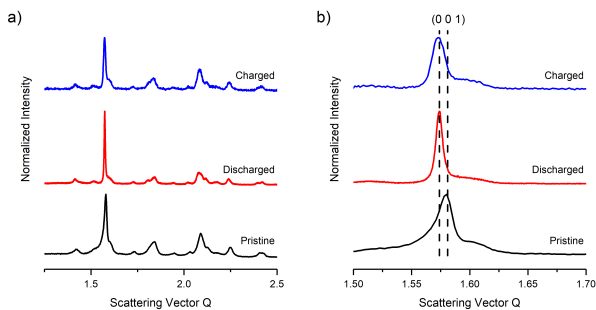
Mg-ion behavior of $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ corresponded well to its Li-ion behavior; in Li-ion, the redox processes occurred at $\sim 3.0 \text{ V}$ and $\sim 1.75 \text{ V}$ vs Li/Li^+ , respectively, thereby possessing a similar redox potential separation ($\sim 1.2 \text{ V}$).³⁸ These processes are attributed to the $\text{V}^{5+}/\text{V}^{4+}$ couple and a combination of the $\text{V}^{4+}/\text{V}^{3+}$ and $\text{Nb}^{5+}/\text{Nb}^{4+}$ couples, respectively, as indicated by three peak maxima within the charge process (only two were visible on the discharge as the lowest discharge maxima is not reached within the cycling window, and the $\text{V}^{4+}/\text{V}^{3+}$ and $\text{Nb}^{5+}/\text{Nb}^{4+}$ couples overlap. These redox changes are corroborated with XAS analysis below.

Five electrodes of the material were recovered following electrochemical discharge (Mg insertion), or electrochemical discharge and charge (Mg insertion and removal) as described in the Experimental section. These were labelled "VNO(Discharged_x)" ($x = 1, 2$ or 3) and "VNO(Charged_y)" ($y = 1$ or 2) respectively, and their electrochemistries are shown in Figure S4. While voltage plateaus were occasionally observed consistent with electrolyte degradation, they provided useful insights into both structural and redox changes in $\text{V}_4\text{Nb}_{18}\text{O}_{55}$ with cycling when analyzed by a combination of EDX, XRD and XAS analysis. The observed Mg:Nb ion ratios within VNO(Discharged_1) from EDX were $4:18$, which corresponded well to the predicted stoichiometry ($\text{Mg}_{3.5}\text{V}_4\text{Nb}_{18}\text{O}_{55}$) from the observed discharged capacity (68 mA h g^{-1}), and given the low H_2O content of the electrolyte, this suggested Mg intercalation (rather than proton intercalation) had indeed taken place. HAADF-STEM and EDX mapping of VNO(Discharged_2) confirmed a uniform dispersion of Mg (ions) found within the particles, consistent with predominant insertion vis-à-vis surface reactivity with the electrolyte (Figure S10). Conversely for VNO(Charged_1) and VNO(Charged_2), the observed Mg:Nb ratios were $0.81:18$ and $0.13:18$, respectively. This observation suggested a limited proportion of Mg was irreversibly trapped in the active material of the respective electrodes, although this could be somewhat reduced when the electrode was overcharged in the case of VNO(Charged_2).

A comparison of the XRD patterns of the pristine, discharged and charged materials provided further strong evidence for Mg^{2+} intercalation behavior (Figure 4). While the exact nature of the shifts were not obvious on inspection of the $(0\ 0\ 1)$ peak (Figure 4b) due to zero errors, peak broadening and anisotropy effects, LeBail fitting (which incorporates all diffraction peaks and accounts for these additional effects) revealed that VNO(Discharged_2) displayed an anisotropic expansion of the lattice parameters, with the largest increase in the b parameter, and a $\sim 1\%$ increase in unit cell volume overall, as indicated by the shift to lower scattering vector in the discharged electrode (Table 2 and

Table 2. The lattice parameters and their standard errors calculated from LeBail refinement for the pristine sample, and the electrodes VNO(Discharged_2), VNO(Charged_1) and VNO(Charged_2).

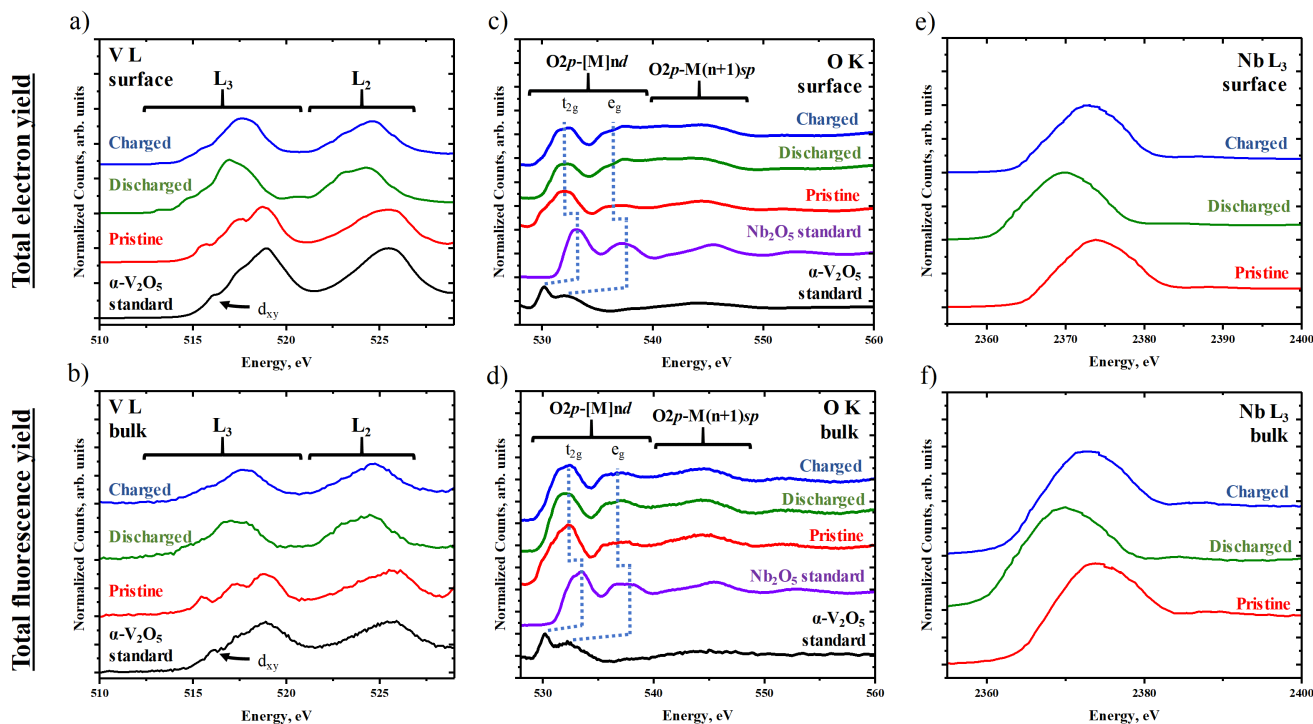
	$a / \text{\AA}$	$\Delta a / \%$	$b / \text{\AA}$	$\Delta b / \%$	$c / \text{\AA}$	$\Delta c / \%$	$V / \text{\AA}^3$	$\Delta V / \%$	R_{wp}
Pristine	17.3658(12)	-	17.6500(14)	-	3.97886(6)	-	1219.55(13)	-	7.80
VNO(Discharged_2)	17.4045(9)	0.22	17.7385(9)	0.50	3.99242(4)	0.34	1232.58(10)	1.07	8.47
VNO(Charged_1)	17.404(2)	0.22	17.674(2)	0.14	3.9859(3)	0.18	1226.0(2)	0.53	9.44
VNO(Charged_2)	17.403(4)	0.21	17.653(5)	0.02	3.9851(4)	0.16	1224.3(5)	0.39	3.25



1
2 Figure 4. a) XRD patterns of the pristine, discharged and 31
3 charged materials, VNO(Discharged_2) and VNO(Charged_2) 32
4 respectively. b) a comparison of the (0 0 1) diffraction peak of 33
5 the pristine and charged materials, revealing a small shift in dif- 34
6 fraction peak location. 35
7 Figure 4). Partial removal of Mg^{2+} was suggested by the 36
8 cell unit cell volume almost exactly half-way in between 37
9 VNO(Discharged_2) and the pristine sample. 38
10 VNO(Charged_2) demonstrated greatest similarity with the 39
11 pristine sample, implying further (but not complete) Mg re- 40
12 moval with overcharging (the XRD pattern of 41
13 VNO(Charged_1) is shown in the SI). Interestingly, the b pa- 42
14 rameter was returned to its original value with over-charg- 43
15 ing, whereas the changes in a and c were largely irreversible, 44
16 respectively. This suggested that Mg^{2+} was primarily 45
17 trapped in crystallographic sites which distorted these spec- 46
18 ific parameters. Rietveld refinement of the discharged and 47
19 charged electrodes was attempted using the predicted 48
49

20 magnesianated structures from the DFT analysis (discussed
21 later); however, the differences in diffraction pattern peak
22 intensities with Mg inclusion were so minimal, Mg location
23 (and any potential defect formation) could not be extracted
24 with any confidence. Moreover, the peak broadening evi-
25 dent in the discharged and charged states (Figure 4a) rela-
26 tive to the pristine sample suggested some loss of crystal-
27 linity, or crystallographic strain, with Mg intercalation.
28 Therefore, local structure methodologies, such as pair distri-
29 bution function (PDF), are suggested in further studies to
30 accurately determine Mg site occupancies in this material.

31 To determine the effect of Mg insertion on elemental oxida-
32 tion states, XAS analysis was conducted on pristine
33 $V_4Nb_{18}O_{55}$, and electrodes VNO(Discharged_1) and
34 VNO(Charged_2) on the V $L_{2,3}$ -, O K, and Nb L_3 -edges. Addi-
35 tionally, V and Nb K-edge data were collected for the pristin-
36 e compounds, and is discussed in the SI. XAS data was col-
37 lected using both total electron and fluorescence yield de-
38 tectors. Signals from electron yields correspond to the
39 chemical state of the surface layer of the material, whereas
40 fluorescence yields correspond to approximately 100 nm
41 into the electrode, thus having a notable contribution from
42 the bulk crystal structure, especially considering the partic-
43 le size of the materials. Comparison between electron yield
44 and fluorescence yields therefore allowed for the distinc-
45 tion between bulk and surface reactions with Mg inser-
46 tion/removal. The V $L_{2,3}$ -edge spectra of V_2O_5 and pristine
47 $V_4Nb_{18}O_{55}$ were similar in shape and position in both elec-
48 tron and fluorescence excited electron transitions from the
49 vanadium core $2p_{3/2}(L_3)$ and $2p_{1/2}(L_2)$ levels to unoccupied



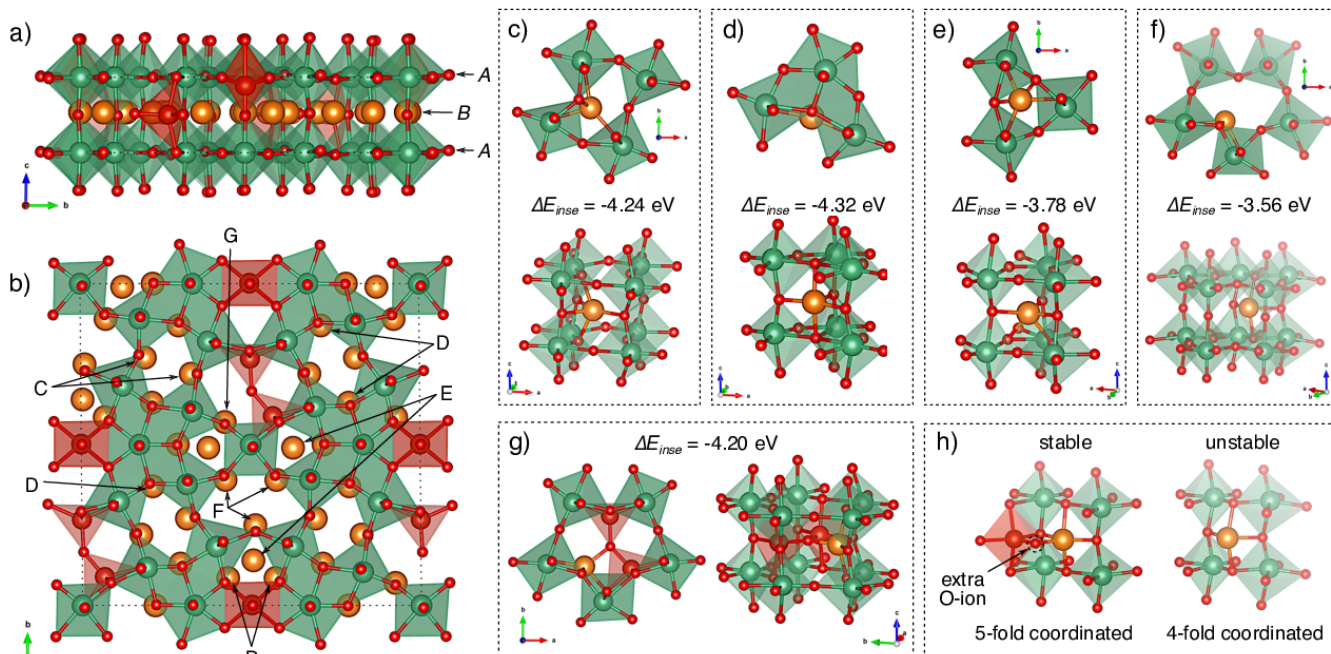
50
51 Figure 5. Integrated X-ray absorption spectra of the V $L_{2,3}$ (a,b), O K (c,d), and Nb L_3 (e,f) edges of pristine $V_4Nb_{18}O_{55}$, discharged to
52 -1.7 V and charged VNO. Standards of orthorhombic V_2O_5 and Nb_2O_5 are shown. Electron yields, corresponding to surface chemistry,
53 of each edge are on top (a,c,e). Fluorescence yields, corresponding to bulk properties, of each edge are on the bottom (b,d,f).

1
2
3 V 3d states. These results provide evidence that V was in the 60
4 5+ oxidation state throughout the $V_4Nb_{18}O_{55}$ particles, i.e. 61
5 both on the surface and in the bulk.^{14,37,50-55} This conclusion 62
6 was supported by the measurements at the V K-edge in 63
7 transmission mode, and detailed discussion of these spectra 64
8 is included in the SI.^{50,52,53,56} Upon discharge to -1.7 V vs 65
9 ACC, the L_3 and L_2 features were generally red-shifted and 66
10 changed shape, and spectra collected in electron and fluo- 67
11 rescent yields were similar. The maximum intensity in the 68
12 L_3 region was at 517 eV, with additional, less intense fea- 69
13 tures at 513 eV, 518 eV and 521 eV. 70
14 In turn, the L_2 spectrum broadened through the appearance 71
15 of two maxima at 523.5 and 524.2 eV. Overall, the spectra 72
16 for discharged $V_4Nb_{18}O_{55}$ matched theoretical and experi- 73
17 mental reports on compounds containing V^{3+} in distorted 74
18 $[VO_6]$ octahedra, both at the surface and within the bulk of 75
19 the electrode.^{25,50,52,53} This observation was consistent with 76
20 Mg^{2+} intercalation into $V_4Nb_{18}O_{55}$, in agreement with elec- 77
21 trochemical, microscopic and structural analysis. Upon 78
22 charging, the region of maximum intensity of the L_3 -edge 79
23 broadened, and the maximum blue-shifted to ca. 518 eV in 80
24 electron and fluorescence yields. Also, the broad doublet 81
25 was lost in the L_2 region, leading to a single maximum 82
26 slightly above 524 eV, thereby blue-shifting the center of 83
27 gravity. These changes were indicative of electrochemical 84
28 oxidation of V to 4+.^{14,25} This means that reduction of V^{5+} to 85
29 V^{3+} upon discharge was not completely reversible, and fur- 86
30 ther confirmed incomplete removal of Mg upon charge. 87
31 Figure 5c,d depicts integrated O K-edge X-ray absorption 88
32 spectra acquired for pristine and electrochemically tested 89
33 $V_4Nb_{18}O_{55}$ samples, the spectra for V_2O_5 and Nb_2O_5 are de- 90
34 picted alongside for comparison. The O K-edge features re- 91
35 sulted from excited electron transitions from the core O 1s 92
36 levels to unoccupied O 2p states that hybridized with metal 93
37 states.⁵⁷ Within the O K-edge of pristine $V_4Nb_{18}O_{55}$, two 94
38 broad spectral events occurred at ca. 532 and 537 eV. The 95
39 absorption events below 540 eV corresponded to O 2p 96
40 states that hybridized with V 3d and Nb 4d states. The first 97
41 event (532 eV) was very broad with a band dispersion 98
42 (FWHM) of 4 eV and appears to result from contributions 99
43 from t_{2g} and e_g states from $[VO_n]$ moieties and t_{2g} of $[NbO_6]$ 100
44 octahedra, based on comparison with the respective stand- 101
45 ards.⁵⁷ The second event (537 eV) matched with the spec- 102
46 trum of Nb_2O_5 , suggesting that it arose solely from the e_g 103
47 states of $[NbO_6]$ in the structure. Above 540 eV, O 2p states 104
48 hybridized with V 4sp and Nb 5sp states. The center of grav- 105
49 ity of the t_{2g} event of pristine $V_4Nb_{18}O_{55}$ was greater than the 106
50 t_{2g} event of V_2O_5 and less than the t_{2g} event of Nb_2O_5 . This is 107
51 expected because the O 1s absorption threshold blue-shifts 108
52 with increasing atomic number.⁵⁷ Additionally, the energy 109
53 separation of t_{2g} and e_g states increased with atomic num- 110
54 ber. In both electron and fluorescence yields of figure 5c,d 111
55 one feature was centered at 545 eV in the O 2p-M (n+1)sp 112
56 region for pristine $V_4Nb_{18}O_{55}$. This was again in good corre- 113
57 spondence with the addition of features from both V_2O_5 and 114
58 Nb_2O_5 . The similarity between sample and the standards 115
59 was consistent with a similar metal-oxygen coordination

and covalent interactions in these compounds. Upon discharge of the cell to -1.7 V vs. ACC, the dispersion of the features at ~ 532 eV red-shifted by 2 eV and the center of gravity shifted to higher energy for both electron and fluorescence yields. This observation corresponded to occupation of the first states just above the Fermi level through reduction of the metal-oxygen bonds. Correspondingly, the O 2p-M (n+1)sp red-shifted by ca. 1 eV upon discharge, consistent with a decrease of the binding energy of the core electrons when the compound was reduced. Upon charge, the t_{2g} spectral event changed shape as the higher energy shoulder raised slightly for both electron and fluorescence yields, but did not change position (Figure S12). The rest of the spectra changed only subtly. The fact that the spectrum did not return to the pristine state indicated chemical irreversibility, consistent with V $L_{2,3}$ -edge observations. Figure 5e-f depicts the Nb L_3 -edge in electron and fluorescence yields. The L_3 edge of pristine $V_4Nb_{18}O_{55}$ displayed a prominent absorption event with a maximum at 2373.8 eV, followed by a weak event at ca. 2388 eV in both electron and fluorescence yields, similar to Nb^{5+} standards in references 51 and 52. The strong peak is associated with electron transitions from Nb core $2p_{3/2}$ to unoccupied 4d states and the weak peak is associated with Nb $2p_{3/2}$ to 5s transitions.^{58,59} These results provide evidence that Nb was in the 5+ oxidation state on the surface and in the bulk of pristine $V_4Nb_{18}O_{55}$. This conclusion was supported by the measurements at the Nb K-edge in transmission mode, showing that the oxidation state of Nb in pristine $V_4Nb_{18}O_{55}$ was 5+ throughout the particle (Figure S11). Further discussion of Nb K-edge spectra of pristine $V_4Nb_{18}O_{55}$ is in the SI.⁵⁸⁻⁶¹ Upon discharge, the Nb L_3 -edge did not change shape but red-shifted by -3.7 eV in both electron and fluorescence yields, consistent with Nb reduction.^{58,69} Upon charging, the Nb L_3 -edge blue-shifted by $+3$ eV in both electron and fluorescence yields. The change confirms that Nb re-oxidized, but did not completely return to the pristine state, indicating that the electrochemistry is not completely reversible. No evidence existed of chemical gradients (or different oxidation states) between surface and bulk at this edge, based on comparison of data from the two detectors employed, consistent with bulk (rather than surface) reactivity.

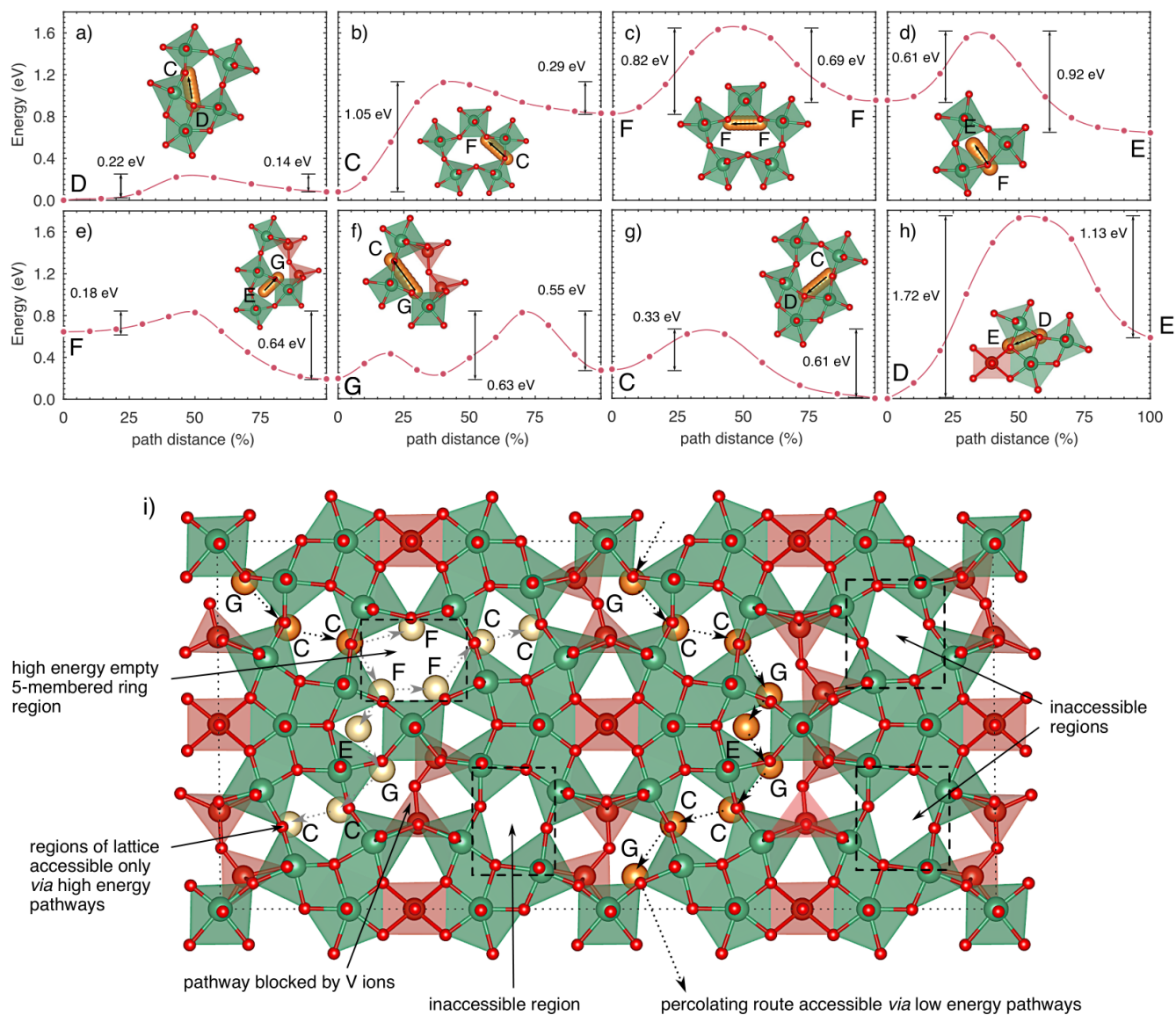
Hybrid-exchange DFT calculations were performed to investigate the electronic and structural properties of $V_4Nb_{18}O_{55}$ (full details in the SI). As previously discussed, the structure contains partially occupied V sites on the oxygen-rich 'B' plane. To represent this key structural feature, but keep the unit cell size within computational limits, two models were constructed based upon the crystallographic (C) and primitive (P) unit cells, giving stoichiometries of $V_4Nb_{22}O_{65}$ and $V_3Nb_{11}O_{34}$ (details in the SI).

These models gave a representation of a closed-shell system (C, containing only d^0 cations) and an O-deficient system (P, containing unpaired electrons) respectively. The latter model was useful to consider, since reduced metal centers are known to occur in similar bronze-type frameworks containing V and Mo.³⁷ The calculated lattice parameters of the



1
2 Figure 6. Structure of $V_4Nb_{18}O_{55}$ (shown as the P cell model, with stoichiometry $V_3Nb_{11}O_{33}$) viewed along a) the a axis, highlighting
3 the metal rich A and oxygen rich B planes and b) viewed along the ‘tunnel’ c axis. Gold spheres indicate stable sites for Mg, grouped
4 according to their similar local geometry, labeled C-G. Optimized geometry of Mg ions in sites C-G are indicated in parts c) – g),
5 viewed along both the c axis, and in the ab plane perspective. The insertion energy of the sites relative to Mg metal is shown. h) The
6 relationship between Mg ion site stability and the coordination number it can achieve.

7 two DFT models are reported in Table 3 and showed good 37
8 agreement with the literature values, despite the deviation 38
9 from experimental stoichiometry in each case. 39
10 A range of insertion sites for Mg-ions were investigated in 40
11 the primitive cell model. The calculations revealed that Mg 41
12 always occupied sites on the oxygen-rich B plane (Figure 42
13 6a). Mg-ions introduced onto the A plane always found sta- 43
14 ble locations on the B plane, minimizing their Coulombic re- 44
15 pulsion from structural Nb and V cations. Figure 6b shows 45
16 insertion sites for Mg-ions in the P cell, grouped together ac- 46
17 cording to their local geometry, and labelled C-G. Figures 6c- 47
18 g show representative optimized structures for the differ- 48
19 ent sites, indicating their local geometry and calculated in- 49
20 sertion energy (ΔE_{inse}) vs. Mg metal. The most stable sites 50
21 for Mg-ions were C and D, located in the 4-membered rings 51
22 (Figures 6c,d). In these sites, Mg-ions could achieve coordi- 52
23 nation to 5 neighboring oxyanions. There were two possible 53
24 local environments for Mg-ions occupying D sites; those 54
25 with only Nb neighbors, which are more stable, and those 55
26 with Nb and V neighbors, which were destabilized by ~ 0.3 56
27 eV. Mg-ions could occupy E sites in the 3-membered rings 57
28 (Figure 6f), which were slightly destabilized with respect to 58
29 C and D sites by ~ 0.4 eV. Mg-ions occupying F sites in the 5- 59
30 membered rings (Figure 6f) were the least stable locations, 60
31 destabilized by ~ 0.75 eV compared to C and D sites. The fi- 61
32 nal possible insertion site was G, indicated in Figure 9g, 62
33 where an Mg-ion occupied a site in the 5-membered ring, 63
34 which was also occupied by interlayer V-ions. G sites were 64
35 stabilized by 0.65 eV compared to F sites. In G sites, the Mg- 65
36 ion could form a bond with the O-ion bridging between the 66
V ions (Figure 6h), resulting in a stable 5-fold coordination, 37
unlike F sites, where the Mg-ion could achieve only an unfavor- 38
able 4-fold coordination. The ability of Mg to coordinate 39
to either 5 or 4 O-ions was therefore a key factor determin- 40
ing the stability in different insertion sites. 41
The diffusion pathways between the sites identified in Fig- 42
ure 6 were evaluated with DFT. Due to the complex structure 43
of the material, and the number of local minima that 44
 Mg^{2+} may occupy, investigating all possible pathways for Mg 45
migration in a unit cell was impractical. Therefore, to gain 46
an understanding of Mg mobility in $V_4Nb_{18}O_{55}$, a selection of 47
pathways was chosen to give a representative behavior in 48
proximity to different structural motifs, and provide a per- 49
colating route demonstrating that migration can occur 50
throughout the structure. Figure 7 a) – h) shows the calcul- 51
ated activation barriers for the selected routes, starting 52
from a stable D site. Mg migration could occur between low 53
energy C and D sites in the 4-membered rings with low ac- 54
tivation barriers in the range 0.14 to 0.61 eV (Figures 7a,g). 55
However, for Mg-ions to migrate into the less stable F sites 56
(in the 5-membered rings), higher barriers existed in the 57
range 0.92 to 1.05 eV (Figures 7b,d). Similarly, moving be- 58
tween F sites had high barriers in the range 0.69 to 0.82 eV. 59
Mg mobility through the empty 5-membered rings, into and 60
between F sites, was therefore kinetically hindered. How- 61
ever, when the 5-membered rings were occupied with inter- 62
layer V ions, Mg could migrate into G sites, in which the 63
Mg-ions were stabilized by the additional neighboring oxy- 64
anion to achieve a more favorable 5-fold coordination (Fig- 65
ure 7g,h). This stabilization provided lower energy barriers 66



1 Figure 7. a) - h) Calculated activation barriers for Mg diffusion through the $V_4Nb_{18}O_{55}$ lattice. Energies are plotted in a consistent
 2 reference frame due to the percolating nature of the pathways investigated, with activation barriers for forward and backward mi-
 3 gration indicated on each plot. The local geometry of the corresponding migration pathways are shown inset. i) Migration routes
 4 throughout the structure using low (black dotted arrows) and high energy (grey dotted arrows) pathways. Gold spheres indicate Mg
 5 sites accessible via low energy pathways a), e), f) and g). White spheres indicate Mg sites accessible only via high energy pathways
 6 b) c) and d) by crossing empty 5-membered ring regions. Green and large red spheres indicate Nb and V, respectively, whereas small
 7 red spheres correspond to O.

8 in the range 0.55 to 0.63 eV for Mg migration through G sites 20
 9 (Figure 7f), compared to when the 5-membered rings were 21
 10 unoccupied (F sites). As the interlayer V ions in these frac- 22
 11 tionally occupied sites are disordered through the lattice, 23
 12 and can be arranged in one of two configurations (Figure 24
 13 6f,g), they could either facilitate migration of Mg, or block 25
 14 migration pathways. It was only when neighboring 5-mem- 26
 15 bered ring were occupied by V ions in a suitable configura- 27
 16 tion that Mg-ions could migrate through the structure along 28
 17 low energy percolating pathways with a maximum activa-
 18 tion barrier of ~ 0.65 eV, using a combination of routes indi-
 19 cated in Figures 7 a), e), f) and g). The arrangement of the

interlayer V ions may be such that pathways were blocked,
 and regions of the lattice inaccessible to migrating Mg-ions.
 An illustration of these two situations is provided in Figure
 7i. We also note that Mg migration across the shared edge
 of $NbO_{6/7}$ units (Figure 7h) had a high barrier of ~ 1.72 eV
 and these pathways were inaccessible. For Mg-ions to ac-
 cess a higher fraction of the lattice, passage through desta-
 bilized F sites in empty 5-membered rings, with higher en-
 ergy barriers, is required.

1 **Table 3. The lattice parameters for $V_4Nb_{18}O_{55}$ from the DFT models and the LeBail refinement of the experimental**
 2 **material, compared with the previously reported values from literature.**

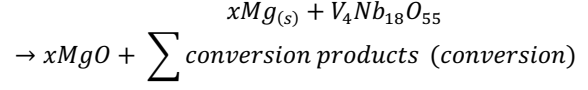
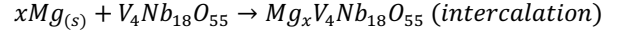
	Stoichiometry	a (Å)	Δa	b (Å)	Δb	c (Å)	Δc
Börrnert et al.	$V_4Nb_{18}O_{55}$	17.355	-	17.698	-	3.971	-
C	$V_4Nb_{22}O_{65}$	17.3918	0.2%	18.0077	1.7%	3.9748	0.1%
P	$V_3Nb_{11}O_{34}$	17.5097	0.9%	17.9076	1.2%	4.0013	0.8%
This work	$V_4Nb_{18}O_{55}$	17.3658	0.06%	17.6500	-0.3%	3.9789	0.2%
This work	$Mg_{3.5}V_4Nb_{18}O_{55}$	17.4045	0.3%	17.7385	0.2%	3.9924	0.5%

3 In the present case, for Mg-ions to penetrate the center of 52
 4 the $V_4Nb_{18}O_{55}$ particles, they must travel along a diffusion 53
 5 length (L) which can be estimated to be ~ 25 nm, corre- 54
 6 sponding to half the distance along the ab plane for the 54
 7 $V_4Nb_{18}O_{55}$ rods in this study (~ 50 nm width). The diffusion 55
 8 length L scales with \sqrt{Dt} , where D is the diffusivity and t is 56
 9 the discharge time (in this case 3 h, corresponding to a rate
 10 of C/3 or 15 mA g⁻¹ as used in the electrochemical experi- 57
 11 ments). Therefore, a diffusivity of $\sim 5.8 \times 10^{-20}$ m² s⁻¹ is re- 58
 12 quired for Mg insertion to the center of a $V_4Nb_{18}O_{55}$ particle 59
 13 during discharge. The diffusivity D can be approximated us-
 14 ing:

$$D \approx \nu \cdot a^2 \cdot \exp\left(\frac{-E_m}{kT}\right)$$

15 Where ν is the atomic jump frequency, and a is the jump 62
 16 distance. If we assume reasonable values of $\nu = 10^{12}$ s⁻¹ and 62
 17 $a = 3$ Å,⁴⁰ the maximum tolerable activation barrier is ~ 720 63
 18 meV at 25 °C (298.15 K), and 930 meV at the elevated temper- 63
 19 ature used in this study (110 °C). These computational 64
 20 results provide a framework for rationalizing the electro-
 21 chemical observations in Figure 5. The low overpotentials 65
 22 observed for shallow charge/discharge cycles may corre- 66
 23 spond to Mg migration through the regions of the structure 67
 24 accessible along low energy pathways with barriers of up to 68
 25 ~ 650 meV. Much of the structure will, however, be inacces- 69
 26 sible for Mg in this way. This is due to the low fraction 70
 27 (~ 0.35) of 5-membered rings occupied by VO_5 units, some 71
 28 of which will block Mg mobility rather than facilitate it. Un- 72
 29 der these shallow discharge conditions, and limited regions 73
 30 of the structure accessible, a low capacity was expected, in 74
 31 agreement with the 15 - 30 mA h g⁻¹ achieved. Under deeper 75
 32 discharge conditions at elevated temperatures, Mg-ions 76
 33 may be able to pass through kinetically limited high-energy 77
 34 regions of unoccupied 5-membered rings, with higher bar- 78
 35 riers in the range 0.92 - 1.05 eV for mobility. Under these 79
 36 conditions, we would expect Mg to access a greater propor- 80
 37 tion of the structure and thus provide a higher capacity, but 81
 38 at the cost of significantly increased overpotentials, as is ob- 82
 39 served experimentally. We note that Mg-ions may be re- 83
 40 tained in stable locations in the structure under these deep 84
 41 cycling conditions, which would result in capacity fade and 85
 42 irreversible crystallographic changes within the $V_4Nb_{18}O_{55}$ 86
 43 structure.

44 While we saw no evidence of conversion reactions taking 87
 45 place during cycling, we performed DFT calculations follow- 88
 46 ing the approach detailed in Hanna *et al.* to calculate volt- 88
 47 ages for intercalation (V_{int}) and conversion (V_{conv}) and 89
 48 thereby assess the relative favorability of both reaction 90
 49 pathways.⁶³ The competing reactions in this case can be de- 91
 50 fined as:



and the calculated voltages can be determined by as:

$$V_{int} = -\frac{G_{Mg_xV_4Nb_{18}O_{55}} - G_{xMg_{(s)}} - G_{V_4Nb_{18}O_{55}}}{zF}$$

and

$$V_{conv} = -\frac{\sum_i G_i^{conv} - G_{xMg_{(s)}} - G_{V_4Nb_{22}O_{65}}}{zF}$$

where

$$\sum_i G_i^{conv}$$

is the sum of the conversion products, containing x moles of
 MgO. G is the Gibbs free energy, which can be reasonably
 approximated by the 0 K enthalpies obtained from the calcu-
 lations. Since there is no experimental suggestion of the
 potential conversion products, we considered the reduced
 oxides V_2O_3 and NbO_2 , given the lowest oxidation states for
 V and Nb were 3+ and 4+, respectively, in discharged
 $V_4Nb_{18}O_{55}$ from XAS analysis.

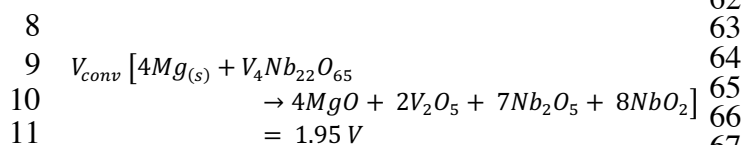
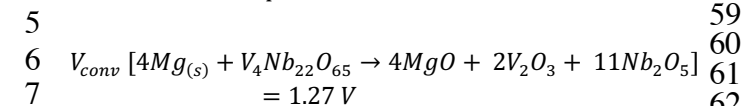
Our DFT calculations necessitated the use of an off-stoichi-
 ometric model of the conventional unit cell to account for
 the disorder of V ions in the 5-membered rings in the real
 material. We used a cell with composition $V_4Nb_{22}O_{65}$. The
 approximate composition of the discharged electrode is
 $Mg_4V_4Nb_{18}O_{55}$. We therefore chose to investigate the differ-
 ence in conversion vs intercalation voltages at stoichi-
 ometries of $MgV_4Nb_{22}O_{65}$ and $Mg_4V_4Nb_{22}O_{65}$. For
 $MgV_4Nb_{22}O_{65}$ the potential to form conversion products was
 calculated to be:

$$\begin{aligned} V_{conv} [Mg_{(s)} + V_4Nb_{22}O_{65} \\ \rightarrow MgO + 2V_2O_5 + 10Nb_2O_5 + 2NbO_2] \\ = 1.98 V \end{aligned}$$

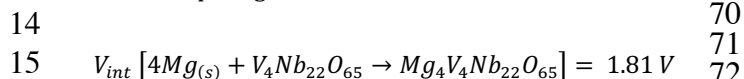
and the intercalation voltage was:

$$V_{int} [Mg_{(s)} + V_4Nb_{22}O_{65} \rightarrow MgV_4Nb_{22}O_{65}] = 2.16 V$$

1 Therefore, intercalation is thermodynamically favoured in 55
2 this case. For the electrode discharged to a stoichiometry of 56
3 $Mg_4V_4Nb_{22}O_{65}$, we can calculate pathways to two difference 57
4 sets of conversion products: 58



13 and the competing intercalation reaction is: 69



17 In this case, the conversion reaction is slightly thermody- 73
18 namically favoured over intercalation, but only by 0.16 V. 74
19 Considering the thermodynamic driving force for conver- 75
20 sion is small, the $V_4Nb_{18}O_{55}$ structure is expected to show 76
21 high kinetic stability, and there is no evidence from experi- 77
22 ment for conversion products such as MgO , V_2O_3 , Nb_2O_5 and 78
23 NbO_2 , we can exclude significant conversion reactions from 79
24 taking place. 80

25 Overall, these DFT results suggested that bronze-type 81
26 frameworks can be used to achieve good Mg^{2+} migration in 82
27 oxide materials, but diffusion is highly dependent on local 83
28 coordination geometry of the migrating Mg^{2+} . Considering 84
29 the above discussion, we can suggest some preliminary de- 85
30 sign rules for Mg -ion mobility in bronze-type structures. 86
31 Frameworks containing 5-membered rings should be 87
32 avoided, since these provide regions limiting Mg mobility. 88
33 Since the instability of the Mg -ions can be assigned to the 89
34 unfavorable 4-fold coordination they achieve in the 5-mem- 90
35 bered ring windows (Figure 6f,h), it follows that larger rings 91
36 (i.e. 6 and 7-membered) will be similarly unfavorable re- 92
37 gions for Mg to occupy and traverse. Structures comprised 93
38 of only 3 and 4-membered rings are therefore expected to 94
39 provide more favorable energy pathways. An alternative 95
40 strategy to achieve enhanced Mg mobility can be deduced 96
41 from the facilitation of migration through the 5-membered 97
42 rings by the VO_5 units. The geometry of these interlayer V 98
43 ions offers an additional O-ion coordination site for the mi- 99
44 grating Mg , and provide stabilization. The use of extra- 100
45 framework, interstitial or mobile O-ions may therefore pro- 101
46 vide a strategy for enhancing Mg -ion mobility. It should be 102
47 noted that the observed experimental kinetics for Mg inter- 103
48 calation in $V_4Nb_{18}O_{55}$ are still relatively inhibited in compar- 104
49 ison to Mo_6S_8 , despite possessing similar computed activa- 105
50 tion energy values (0.65 eV vs 0.44-0.51 eV respectively), 106
51 suggesting that other limiting factors, such as transfer from 107
52 the electrolyte to the electrode, present kinetic barriers in 108
53 oxide electrode magnesiation. 109

54 **Conclusions** 110

The tetragonal tungsten bronze $V_4Nb_{18}O_{55}$ was evaluated as a member of a new class of Mg -ion intercalation hosts for Mg batteries. Partially reversible Mg -ion insertion was achieved at 110 °C, with a maximum capacity of 75 mA h g⁻¹, corresponding to 4 Mg^{2+} ions per formula unit. The electrochemical response could be sustained at moderate efficiency over tens of cycles, indicating some reversibility. The insertion and partial removal of Mg was corroborated with elemental, crystallographic, and spectroscopic methodologies. Powder X-ray diffraction revealed crystal lattice expansion on Mg insertion, whilst X-ray absorption spectroscopy showed reduction of V^{5+} to V^{3+} and partial reduction of Nb^{5+} to Nb^{4+} upon discharge. These changes were shown to be partially reversed upon charge, suggesting some irreversible Mg insertion, consistent with the electrochemical measurements. DFT calculations and analysis suggested the structure of $V_4Nb_{18}O_{55}$ facilitates relatively rapid Mg^{2+} diffusion along the *ab* crystallographic plane. Therefore, it is suggested that other tetragonal tungsten bronze frameworks may hold promise as Mg electrode materials in the future.

ASSOCIATED CONTENT

Supporting information. V^{4+} solution synthesis protocol, Continuous Hydrothermal Flow Synthesis apparatus, LeBail refinement plots, electrochemical figures, EDX analysis, XAS analysis, DFT computational methodology. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Author Contributions

‡ The contribution of IDJ and GN to this work was equal. IDJ and GN co-wrote the manuscript. IDJ, GN and HYD performed the electrochemical testing and analysis. IDJ and DT performed the synthesis of $V_4Nb_{18}O_{55}$. GN and JWF performed the XAS analysis. KM and FC performed the DFT analysis and interpretation. IDJ and JKC performed the XRD and LeBail analysis. YAW, LH and RFK performed the HAADF imaging and EDX measurements. JC, JAD and IPP edited the manuscript. JC is the academic lead of the Mg battery research team at the University of Illinois at Chicago. JAD is the academic lead of the synthesis team at University College London, and co-developed the CHFS process used within this work. All authors have given approval to the final version of the manuscript.

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15 **ABBREVIATIONS** 82

16 ACC, activated carbon cloth; BPR, back-pressure regulator; 84
 17 CJM, confined jet mixer; CHFS, continuous hydrothermal flow 85
 18 synthesis; DFT, density functional theory; EDS, energy-disper- 86
 19 sive X-ray spectroscopy; FWHM, full width at half maximum; 87
 20 HAADF-STEM, High-Angle Annular Dark-Field Scanning Trans- 88
 21 mission Electron Microscopy; ICSD, International Crystal Struc- 89
 22 ture Database; MgTFSI₂-PY₁₄TFSI, Mg[N(SO₂)₂(CF₃)₂]₂- 90
 23 (C₉H₂₀N)(N(SO₂)₂(CF₃)₂); NMP, N-methyl pyrrolidone; NMR, 91
 24 nuclear magnetic resonance; PDF, pair distribution function; 92
 25 PVDF, polyvinylidene difluoride; PXRD, powder X-ray diffrac- 93
 26 tion; SDD, silicon drift detector; TEM, transmission electron mi- 94
 27 croscopy; TEY, total electron yield; TFY, total fluorescence 95
 28 yield; XAS, X-ray Absorption Spectroscopy; XRD, X-ray diffrac- 96
 29 tion. 97

30 **REFERENCES** 100

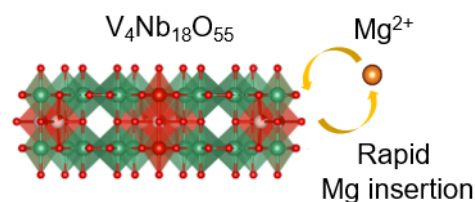
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78 The tetragonal tungsten bronze, $V_4Nb_{18}O_{55}$, dis-
79 plays relatively rapid Mg^{2+} intercalation, due to
80 rapid Mg diffusion pathways present in the struc-
81 ture. This family of compounds are therefore ex-
82 citing candidate electrodes for next-generation re-
83 chargeable Mg batteries.