Magnetic particle hyperthermia modelling - different core sizes and magnetic interactions

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ABSTRACT

We present experimental intrinsic loss power (ILP) values, measured at an excitation frequency of 1 MHz and at relatively low field amplitudes of 3.4 to 9.9 kA/m, as a function of the mean core diameter, for selected magnetic nanoparticle (MNP) samples synthesized in the recent EU-funded NanoMag project. The mean core sizes ranged from ca. 8 nm to 31 nm. Transmission electron microscopy indicated that those with smaller core sizes (less than ca. 22 nm) were single-core MNPs, while those with larger core sizes (ca. 29 nm to 31 nm) were multi-core MNPs. The ILP data showed a peak at ca. 20 nm. We show here that this behaviour correlates well with the predicted ILP values obtained using either a non-interacting Debye model, or via dynamic Monte-Carlo simulations, the latter including core-core magnetic interactions for the multi-core particles. We show that this alignment of the models is a consequence of the low field amplitudes used.

Keywords: magnetic nanoparticles, magnetic interactions, magnetic relaxation, Monte-Carlo simulations, multi-core particles, single-core particles.

1. Introduction

Iron oxide based magnetic nanoparticles (MNPs) can be found in several biomedical applications such as in the areas of diagnosis, actuation, imaging and therapy [1, 2]. One interesting and promising *in-vivo* MNP application is magnetic hyperthermia for cancer therapy [3], for which the figures of merit are the specific loss power (SLP) or intrinsic loss power (ILP) parameters, both of which characterise the heating performance of a given MNP sample.

In a recent EU-funded project (NanoMag), we have studied ca. 50 different iron oxide MNP systems with the goal of improving our metrological understanding of this class of materials. These have been both single-core and multi-core MNP systems[†], with different core sizes, core packing fractions and number of cores per particle.

The data obtained on both single-core and multi-core MNPs is interesting in the context of magnetic hyperthermia, give that it has been reported previously that ILP values may be dependent on the strength of magnetic core-core interactions, with stronger interactions leading to lower ILPs [4]. This implies that the ILP should be lower in multi-core MNPs than in corresponding single-core MNPs, due to the core-core interactions.

In this study we compare the measured ILP values in selected particles with both a non-interacting Debye model and with dynamic Monte-Carlo simulations including, for multi-core particles, core-core magnetic interactions.

We have earlier made substantial static Monte-Carlo simulations where we introduced anisotropy energies and dipolar-dipolar interactions in multi-core particles and where we also included core size distribution as well as varying the easy-axis distribution [5-7]. In these studies, the energies can be independently introduced in order to study their different contribution to the magnetic response.

In this study we have further developed this static equilibrium simulation to a dynamic Monte-Carlo model by studying the probabilities of core magnetization switching compared to the frequency of an applied AC field.

Several other studies on dynamic Monte-Carlo simulations on interacting magnetic nanoparticles have been performed [9-12], but in this study we have compared the simulation results with experimental data for an especially large core size range (ca. 8 nm to ca. 31 nm), and we have paid particular attention to the particle type (single-core or multi-core).

2. Material and methods

[†] For further details on the 'single-core' and 'multi-core' terminology, see Wells et al., *Standardisation of magnetic nanoparticles in liquid suspension*, J. Phys. D **50**, 383003 (2017).

2.1 Magnetic nanoparticles

Iron oxide based magnetic nanoparticles where synthesized with different core sizes and different types of magnetic nanoparticles (single- and multi-core particles). For the single-core particles (CSIC-01, CSIC-02 and CSIC-03), thermal decomposition of iron oleate in organic media was used as the synthesis route. The CSIC-01 sample was coated with silica (silica layer of about 9 nm). CSIC-02 and CSIC-03 MNPs were coated with meso-2,3-dimercaptosuccinic acid (DMSA). The multi-core particles CSIC-04, CSIC-05 and CSIC-06 were synthesized using oxidative precipitation including an acid treatment. These particles were coated with dextran under high pressure homogenization (HPH) conditions.

The mean core diameter (D_m), log-normal distribution standard deviation (σ), and type of MNP system, as determined by transmission electron microscopy (TEM), are given in Table 1.

Table 1. Structural details of the studied MNP systems.

MNP system	Type	D_m (nm)	σ (nm)
CSIC-01	Single-core	11.6	1.2
CSIC-03	Single-core	7.9	1.6
CSIC-11	Single-core	14.6	1.9
CSIC-12	Single-core	21.7	3.8
CSIC-04	Multi-core	28.7	8.7
CSIC-05	Multi-core	30.7	9.2
CSIC-06	Multi-core	28.7	7.9

The mean number of cores in each multi-core particle was ca. 2.5, 4.6 and 2.1 for CSIC-04, CSIC-05 and CSIC-06 respectively, as determined by TEM.

2.2 Experimental

ILP values were measured using a Materials Characterisation MACH system (Resonant Circuits Limited, UK) at an excitation frequency of 1 MHz and at field amplitudes ranging from 3.4 to 9.9 kA/m. The temperature change was measured using fluoroptic probes, and the data were analyzed using the 'corrected slope method' for non-adiabatic systems [13].

TEM was performed using a FEI Titan 80-300 equipped with a field emission gun and operating at 80 or 300 kV.

2.3 Theory

A non-interacting ILP model based on the Debye model and inspired by Rosensweig [8] has been developed. The model takes into account the core size distribution, the field dependence of the relaxation time as well as an approximation of the non-linear field effects in the out-of-phase magnetization. The ILP is given by:

$$ILP = \frac{\pi \mu_0 M_s f}{H_0 \rho < V_c >} \int V_c L\left(\frac{M_S V_c \mu_0 H_0}{kT}\right) \frac{2\pi f \tau_{eff}}{1 + (2\pi f \tau_{eff})^2} f_{LN}(D_c) dD_c \quad (1)$$

where

$$\frac{1}{\tau_{eff}} = \frac{1}{\tau_N} + \frac{1}{\tau_B}$$

$$\tau_N = \tau_0 exp\left(\frac{KV_c}{kT}\left(1 - \frac{H_0}{H_{SW}}\right)^2\right)$$

$$\tau_B = \frac{3V_H \eta}{kT}$$

and $L(x)$ is the Langevin function, H_{sw} the switching field equal to $2K/M_S$, K the anisotropy constant, H_0 the field amplitude, f the excitation frequency, M_s the saturation magnetization, f_{LN} the core size distribution (log-normal), ρ density of the core, V_c the core volume, V_H the hydrodynamic volume of the particle, η is the viscosity of the fluid where the MNPs are dispersed, k the Boltzmann constant and T is the temperature.

A dynamic Monte-Carlo model has also been developed to study the dynamic magnetic response of single- and multi-core particles. We use a two-level approximation of the energy states cores [10,11] and a unit time of simulation is equal to one Monte-Carlo step (1 mcs) where all cores are updated. 1 mcs corresponds to τ_0 (typically 10^{-9}–10^{-11} s). Multicore particles are built up as clusters of N spherical domains defined by number of cores, core size distribution, packing density, random or aligned easy axis, and dead layer thickness [6]. Anisotropy, dipole-dipole, and exchange interactions are all included in the model. The total energy of a spherical domain i can be written as:

$$E_i = -KV_i(\boldsymbol{\mu}_i \cdot \boldsymbol{n}_i)^2 - M_s V_i \boldsymbol{\mu}_i \cdot \boldsymbol{B}_i \quad (2)$$

where K is the anisotropy constant, V_i the core volume, $\boldsymbol{\mu}_i$ the magnetic moment of the core, $\boldsymbol{n}_i$ the unit vector of the easy axis and M_s the saturation magnetization.

The total magnetic field $\boldsymbol{B}_i$ acting on a domain i is a sum of the external applied field, the dipolar field from all surrounding domains and the exchange interaction field from the very close neighbours. The total field H_i is given by:

$$H_i = H_{ext} + \sum_j J_{ij} \boldsymbol{\mu}_j - \frac{1}{4\pi} \sum_j \left(\frac{\boldsymbol{\mu}_j}{|\boldsymbol{r}_{ij}|^3} - \frac{3(\boldsymbol{\mu}_j \cdot \boldsymbol{r}_{ij})\boldsymbol{r}_{ij}}{|\boldsymbol{r}_{ij}|^5}\right) \quad (3)$$

where H_{ext} is the applied external field, $\boldsymbol{J}_{ij}$ the nearest-neighbor exchange coupling constant between domain i and j, $\boldsymbol{\mu}_i$ the magnetic moment and $\boldsymbol{r}_{ij}$ the vector between domain i and j. It has been shown that the minimum energy is found when all three vectors, $\boldsymbol{\mu}_i$, $\boldsymbol{n}_i$, and $\boldsymbol{B}_i$ are in the same plane and hence the energy function can be written as:

$$E_i = -KV_i\left(\cos^2(\alpha_i) + \left(M_s/K\right)|\boldsymbol{B}_i|\cos(\varphi_i - \alpha_i)\right) \quad (4)$$

where α_i is the angle between $\boldsymbol{n}_i$ and $\boldsymbol{\mu}_i$, and φ_i the angle between $\boldsymbol{B}_i$ and $\boldsymbol{n}_i$.

For a given applied field we calculate the energy minima and the saddle point for each domain i in the cluster. We assume that the magnetic moment must be aligned with one of the minima and the probability for the moment to flip between the two extreme points is given by the energy barrier related to the thermal energy, according to:

$$p_i = e^{-\left(E_i^{sad} - E_i^{min}\right)/k_B T} \quad (5)$$

where E_i^{min} is the minimum energy where the moment is located, E_i^{sad} is the saddle point energy, k_B the Boltzmann constant and T is the temperature. The magnetic moment flips if the probability p_i is greater than a drawn random number between 0 and 1.

To simulate the dynamic response from a particle we apply an AC-field, $H_0 \sin(\omega t)$. One period of the signal is divided into n discrete points and the mean particle magnetisation is calculated for each field point. The number of Monte-Carlo steps m for each field point is given by the frequency of the applied field f, τ_0 and n. The magnetization is averaged over a large number of clusters (typ. order of 1000) and a number of periods (typ. 10 periods). The results are typically presented as hysteresis loops and the loop area is used to calculate the ILP value.

3. Results and discussion

Using the non-interacting model and K=10 kJ/m^3, T=300 K and a log-normal size distribution width (i.e. standard deviation) of σ=1.2 the ILP are calculated and plotted versus the field amplitude and the median core diameter, as visualized in the 3D plot in Figure 1 below.

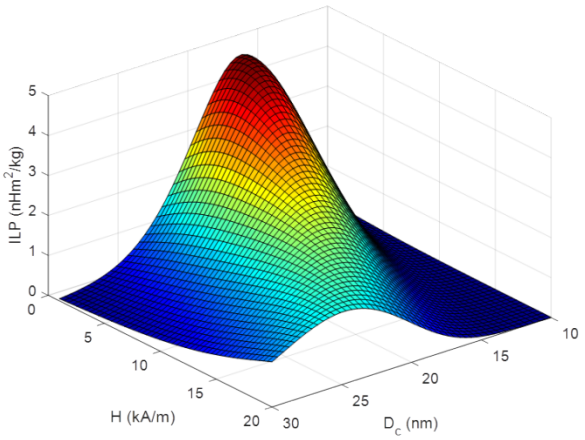

Fig. 1 ILP versus field amplitude and mean core diameter for the following parameters K=10 kJ/m3, T=300 K, size distribution width (σ=1.2) and τ_0 = 10^{-10} s, only taking the Néel contribution into account.

As can be seen from Figure 1 there is maximum in ILP at a specific core diameter when varying the core size. At this specific core diameter, the mean core relaxation rate is in the same range as the excitation frequency. In the non-interacting model, we assume that the non-linear field behavior of the out-of-phase magnetization can be approximated by the Langevin function and we approximate the field dependence of the Néel relaxation by Equation 1.

To test these assumptions, we used the same parameters in the non-interacting model and in the Monte-Carlo simulations for a monodispersed non-interacting MNP case‡. The result is plotted in Figure 2 below.

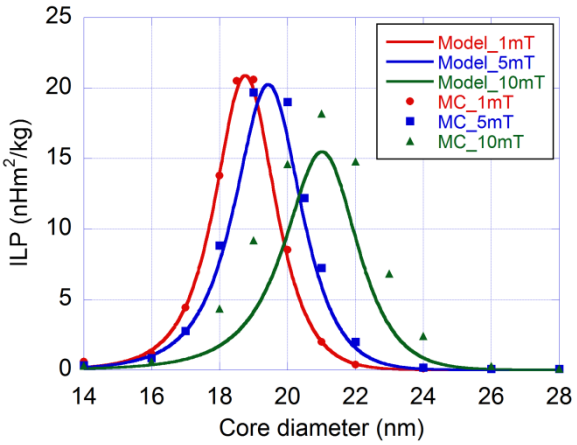

Fig. 2 ILP versus core diameter using the non-interacting model together with the results from the dynamic MC-simulations, for mono-dispersed cores at different field amplitudes (only taking the Néel contribution into account),

‡ Note that using the ILP parameter to describe magnetic heating in a monodisperse MNP system is rather an oxymoron, as the ILP is strictly only valid in polydisperse ($\sigma > 0.1$) systems [14]. It is used here for illustration only.

red: B_0 = 1 mT, blue: B_0 = 5 mT and green: B_0 = 10 mT. f = 1 MHz, K = 10 kJ/m^3 and T = 300 K. "Model" in the legend indicates the result from the non-interaction model and "MC" the result from the dynamic Monte-Carlo simulations.

As can be seen in Figure 2 there is good resemblance between the non-interacting model (Equation 1) and the dynamic Monte-Carlo simulations. At higher field amplitudes we have some deviations, probably due to the assumption of using the Langevin function that describes the non-linear field effects for this dynamic magnetic case.

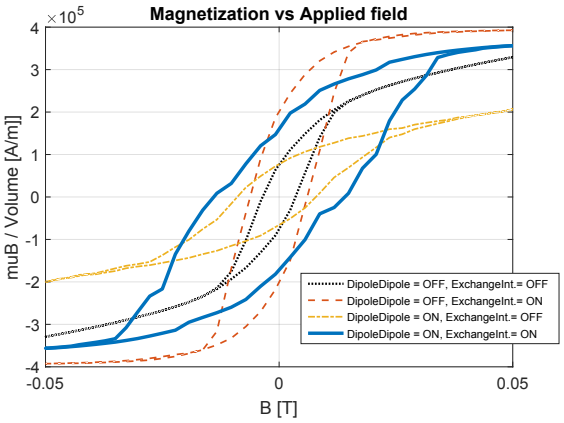

Fig. 3 Hysteresis loops (magnetization vs field) at different interactions by switching ON and OFF dipole-dipole and exchange contributions; Number of cores N_c = 20, random closed packed (RCP), K = 10 kJ/m^3, M_S = 400 kA/m, T = 300 K, D_c = 15 nm. Applied field amplitude and frequency is B_0 = 50 mT and f = 1 MHz.

As can be seen from Figure 3 the different types of interactions give different behaviour in the hysteresis loop. Dipole-dipole interactions stretch out the hysteresis in field and the exchange interactions stretch out the hysteresis loop in magnetization (for instance increase of the remanence).

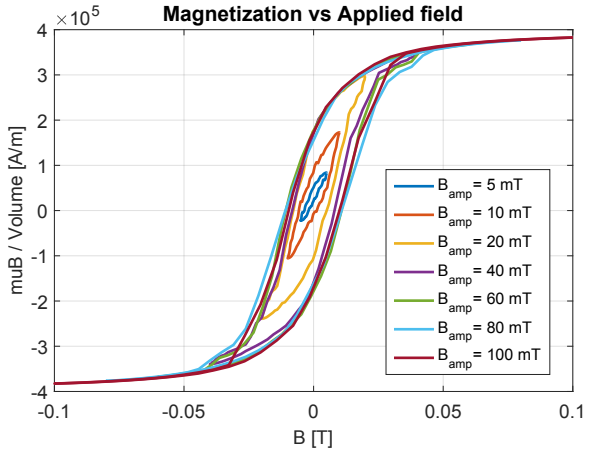

Fig. 4 Hysteresis loops (magnetization vs field) at different field amplitudes as indicated in the figure; using N_c = 20 (RCP), K = 10 kJ/m^3, M_S = 400 kA/m, T = 300 K, D_c = 15 nm. Applied frequency is f = 1 MHz.

From the result showing in Figure 4 the enclosed area (absorbed energy) increase with field amplitude, as expected.

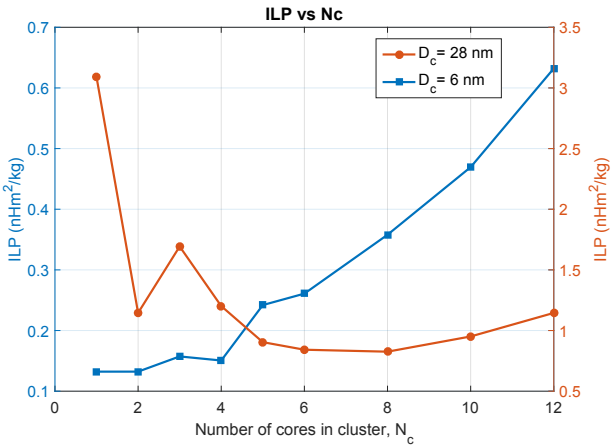

Fig. 5 ILP versus number of cores in multi-core MNP for two core sizes, below and above the core size that gives the maxima in ILP. Red curve D_c = 28 nm, blue curve D_c = 6 nm, K = 10 kJ/m^3, M_S = 400 kA/m, T = 300 K, Applied field amplitude and frequency is B_0 = 10 mT and f = 1 MHz. Only dipole-dipole interactions included.

In Figure 5 we can see an interesting behaviour. For core sizes above the size that gives the maxima in ILP, the ILP value decrease with increasing number of cores, N_c. This is expected and has also been determined from experiment [4]. For core sizes below the size that gives the ILP maxima the ILP increases instead. Both observations may be explained by the change in dipole-dipole interactions between the cores when the number of cores in cluster change and thereby changing the Néel relaxation time with respect to the excitation frequency.

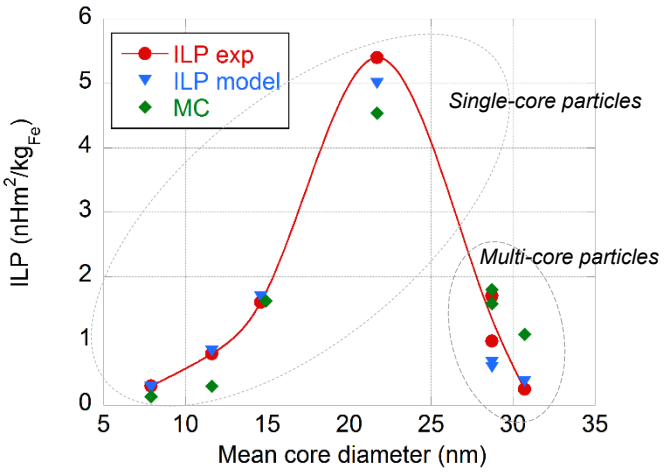

Fig. 6 Experimental ILP data determined from SAR/ILP analysis (field amplitudes between 3.4-9.9 kA/m and at frequency 1 MHz) (red), calculated ILP values using the non-interacting ILP model (blue) and simulated ILP using the dynamic Monte-Carlo model at 1 MHz with dipole-dipole interactions for the multi-core MNPs (green). MNP parameters used in the analysis are determined from TEM, and magnetization vs field (e.g. core size, core size

distribution and saturation magnetization) and magnetic anisotropies from [15].

Bringing this all together, Figure 6 shows the experimental ILP data superimposed with the calculated ILP values obtained using the non-interacting model and the MC simulations at 1 MHz, as a function of core size. It is clear that there is very good agreement across the full range of core sizes, using both models.

4. Conclusions

We have investigated experimental ILP values for different iron-oxide based cores sizes both for single- and multi-core particles.

We have found that a non-interacting ILP model can be used when the field is not too high (as compared to $2K/M_s$). In order to include interactions (dipole-dipole and exchange interactions) and non-linear field effects, a dynamic Monte-Carlo should be used.

Both models can explain how the ILP varies with core sizes and the core-core interactions in the multi-core particles, for instance the peak in ILP at a specific core size (about 20 nm) for 1 MHz.

From the analysis we have found that the core-core interactions give a different behaviour in ILP dependent on if the core size is above or below the core size that gives the maxima in ILP. This is something we will investigate further in an upcoming paper [16].

Acknowledgements

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