

1 **Revealing the microstructural stability of a three-phase soft solid**  
2 **(ice cream) by 4D synchrotron X-ray tomography**

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## Abstract

35  
36 Understanding the microstructural stability of soft solids is key to optimizing  
37 formulations and processing parameters to improve the materials' properties. In this  
38 study, *in situ* synchrotron X-ray tomography is used to determine the temperature  
39 dependence of ice-cream's microstructural evolution, together with the underlying  
40 physical mechanisms that control microstructural stability. A new tomographic data  
41 processing method was developed, enabling the features to be segmented and quantified.  
42 The time-resolved results revealed that the melting-recrystallization mechanism is  
43 responsible for the evolution of ice crystal size and morphology during thermal cycling  
44 between -15 and -5 °C, while coalescence of air cells is the dominant coarsening  
45 mechanism controlling air bubble size and interconnectivity. This work also revealed  
46 other interesting phenomena, including the role of the unfrozen matrix in maintaining  
47 the ice cream's microstructural stability and the complex interactions between ice  
48 crystals and air structures, e.g. the melting and recrystallization of ice crystals  
49 significantly affect the air cell's morphology and the behavior of the unfrozen matrix.  
50 The results provide crucial information enhancing our understanding of microstructural  
51 evolution in multi-phase multi-state complex foodstuffs and other soft solids.

52

53 **Keywords:** Ice cream; Microstructure; Tomography; Ice crystals; Coarsening; Soft  
54 solid.

55

56 **1. Introduction**

57 Soft solids are important composites that are characterized by complex multi-phase  
58 structures and possess inherently complex non-Newtonian rheological properties under  
59 external stress [1-3]. Soft solids exist either in nature, e.g. muds, or in many artificially  
60 manufactured products such as emulsions, biopolymers, fresh concrete and domestic  
61 baking materials. Many soft solids, such as soft foams (e.g. ice cream) and aerated  
62 desserts, contain porous phases within a viscous matrix [4-6].

63 Structural stability is desired for many soft foams and microstructural instability  
64 greatly influences the materials' properties and their applications. Take ice cream for  
65 example. The microstructure of ice cream, including the size distribution and  
66 connectivity of each phase, plays a critical role in determining product quality (e.g.  
67 mouthfeel, taste, appearance, etc.) and the product's shelf-life [6-10]. For example, the  
68 microstructural change at different storage temperatures has been shown to alter ice  
69 cream's viscoelastic properties and hence our oral sensory perception of it [10].

70 However, irreversible microstructural changes often occur in ice creams (over a range  
71 of different timescales), as well as in other similar foam structures that contain air,  
72 above a certain temperature ( $\sim -30$  °C for ice cream [4]) which may occur during  
73 shipping, storage at the grocery store and in domestic freezers (ca.  $-18$  °C), and on the  
74 consumer's table.

75 For ice cream, the structural instability is affected by many factors, including  
76 compositions [11, 12] and thermal variations [5, 13]. One of the well-recognized  
77 phenomena due to thermal instability is coarsening of microstructure [4, 5, 13-15]. This  
78 was initially examined in light microscopy [14, 16-19] and cryo-scanning electron

79 microscopy [20, 21], and transmission electron microscopy [22], all of which provide  
80 only 2D information of the surface or of cuts through the ice cream sample. These  
81 phenomena were recently studied in 3D with X-ray tomography on *ex situ* coarsened  
82 samples [4, 5]. Using a synchrotron X-ray tomography technique, we revealed in our  
83 previous work [5] that after thermal cycling (or thermal ‘abuse’) between -15 and -5 °C  
84 for a number of cycles/days, both ice crystals and air cells grew in size creating ice  
85 dominated structures within a deteriorated ice cream microstructure. More specifically,  
86 the size of ice crystals was observed to continuously increase up to 14 cycles; however,  
87 the growth rate significantly decreased after 7 thermal cycles. The air cells also  
88 increased in size, and they continued to grow into interconnected irregular shapes with  
89 long continuous channels after 14 thermal cycles. However, for up to 7 thermal cycles  
90 the air cells seemed to remain more or less spherical.

91 From that *ex situ* study, we also observed that the ice crystals within the unfrozen  
92 matrix tended to align along the boundary between air cells and the matrix, minimizing  
93 surface energy. Our findings in those 3D experiments provided valuable insight into the  
94 structural changes of ice cream upon thermal cycling. However, the study was  
95 performed on *ex situ* thermally cycled samples and thus, the interactions between the  
96 microstructural features could not be elucidated. The detailed mechanisms that control  
97 the microstructural evolution, which are only available via *in situ* studies, still remain  
98 to be explored. Questions regarding the growth mechanisms and relative movement of  
99 the phases and their exact interactions during thermal cycling, need to be answered to  
100 be able to improve the stability of the ice cream’s structure [5, 14, 23]. For example,

101 what are the dynamics of the changes for each phase and how do they impose on each  
102 other due to thermal variations, in order to maintain the integrity of ice cream's structure?  
103 What are the dominant coarsening mechanisms that control the ice cream's  
104 microstructural evolution?

105 This work non-destructively studies the thermal stability of the ice cream  
106 microstructure via 4D (3D plus time) synchrotron X-ray tomography to reveal the  
107 dynamics of the microstructural changes. This technique has become increasingly used  
108 in the study of opaque materials systems to study both coarsening [24] and rheology  
109 during deformation [25, 26]. X-ray tomographs were continuously acquired on a  
110 sample during a heating and cooling cycle at a well-controlled slow ramp rate of 0.05  
111 °C/min. To analyze the acquired data, an iterative tomographic data reconstruction and  
112 image processing method was developed. Through quantitative analysis, the physical  
113 mechanisms which dominate the degradation of ice cream's microstructure due to  
114 temperature variation, are examined and discussed in detail.

115

## 116 **2. Materials and methods**

### 117 **2.1 Sample and experimental methods**

118 Fresh ice cream containing 5% fat was manufactured by Unilever R&D (U. K). A  
119 500 ml block of fresh ice cream was initially thermally cycled between -15 and -5 °C  
120 for seven times (1 week) before it was used for the *in situ* synchrotron experiment. The  
121 seven cycles created a larger scale of microstructure, enabling easy identification of  
122 phases for quantification, and also represent a transition point between the observations  
123 made in the first seven cycles where change was relatively rapid and the next seven

124 cycles where the size of ice crystals became more stable as reported in our *ex situ* studies  
125 [5]. Small ice cream samples, each contained in a 3 mm inner diameter kapton tube (67  
126  $\mu\text{m}$  thick, American Durafilm Co. Inc, Holliston, U.S), were cut from the 500 ml block.  
127 Details of the sample preparation method are described in [5].

128 The *in situ* synchrotron experiment was conducted on the Diamond Manchester  
129 Beamline (I13-2) of the Diamond Light Source (DLS, U.K) using a pink beam. The set-  
130 up for running the beamline experiment, together with the cold stage used to provide  
131 the sample temperature, is described in [5, 27]. During the *in situ* thermal cycling  
132 experiment, the sample was loaded in the cold stage at  $-15\text{ }^{\circ}\text{C}$  and stabilized for 10 min.  
133 Then, the sample was heated to  $-5\text{ }^{\circ}\text{C}$  at a ramp rate of  $0.05\text{ K/min}$  and held there for  
134 10 min. After that, the sample was cooled back to  $-15\text{ }^{\circ}\text{C}$  at the same ramp rate as the  
135 heating stage. A schematic of the thermal cycle history is shown in the inset of Fig. 1.  
136 The tomographic scans were acquired using a  $2560 \times 2160$  pixel PCO Edge 5.5 CMOS  
137 camera that was optically coupled to a single crystal  $\text{CdWO}_4$  scintillator during the  
138 thermal cycle. For each tomographic scan, 900 projections were recorded with an  
139 exposure time of 100 ms (90 s for each scan) and a pixel size of  $0.8\text{ }\mu\text{m}$ . However, at  
140 the end of each tomographic scan, the sample stage was rotated back to the initial  
141 position for system re-initiation to start the next tomographic scan, adding an additional  
142 delay of  $\sim 51\text{ s}$ , for a cycle time of  $\sim 141\text{ s}$ . In total, 178 tomographic scans were acquired  
143 during a thermal cycle.

144

## 145 **2.2 Image reconstruction and three-phase segmentation approach**

146 Initially, the acquired projection data were reconstructed using the conventional

147 Filtered Back Projection (FBP) algorithm [28], producing extremely noisy  
148 reconstructions with low contrast and ring artifacts (see Fig. 2(a)). The poor quality of  
149 FBP images was due to angular under-sampling (only 900 projections for a  $2k \times 2k \times$   
150  $2k$  volume), short exposure time and low attenuation contrast between ice and water. In  
151 order to improve the image quality suitable for segmentation, we applied a Model-  
152 Based Iterative Reconstruction (MBIR) approach.

153 Our MBIR algorithm is based on the Group-Huber data fidelity function to  
154 minimize ring artifacts and 3D total variation (TV) regularization penalty [29-31]. The  
155 TV-related regularization sub-problem has been solved using the Split-Bregman  
156 method in order to enhance the weak contrast between ice crystals and unfrozen matrix  
157 (see Fig. 2(b)).

158 Although MBIR reconstruction substantially improves the contrast and also  
159 removes noise, the reconstructed images suffer from the visible intensity  
160 inhomogeneities within various ice-crystals (indicated by arrows in Fig. 2(b)). These  
161 artifacts can be a result of the combined effects of strong noise and beam hardening.  
162 The latter is possible due to abrupt changes in attenuation coefficients between unfrozen  
163 matrix (highly attenuated) and ice crystals (poorly attenuated), therefore introducing  
164 non-linearity in a beam. Intensity inhomogeneity within a single crystal restricts  
165 successful segmentation by histogram thresholding. One can use more sophisticated  
166 segmentation methods (e.g. 3D snake contours can successfully segment features with  
167 intensity inhomogeneity using supervised seeding). However, due to the large data size,  
168 a computationally efficient approach is required.

169 Here we applied an additional post-processing step to equalize intensity within  
170 phases by means of gradient-constrained nonlinear isotropic diffusion [5, 32]. The mask  
171 to terminate the diffusion process across selected boundaries was acquired from the  
172 image in Fig. 2(b) by thresholding the magnitude of the gradient. Since the gradient  
173 magnitude between phases is large and the variations of intensity within regions are  
174 gradual, one can run the constrained diffusion until the regions become fully  
175 homogeneous. Therefore, we run 1000 diffusion iterations on GPU to equalize  
176 intensities within enclosed regions of the image in Fig. 2(b). The result (Fig. 2(c)) is  
177 sufficient to implement a simple histogram thresholding operation and generate images  
178 with homogeneous contrast within three phases (Fig. 2(d)).

179 Finally, the processed volume was cropped into a smaller volume for 3D analysis.  
180 3D rendering of the features, as well as quantification of size/volume, was performed  
181 using Avizo® (FEI, Thermo Fisher Scientific, U. S). The thickness of unfrozen matrix  
182 was measured using BoneJ in ImageJ [33].

183

### 184 **3. Results and Discussion**

#### 185 **3.1 General microstructural evolution during thermal cycling**

186 Fig. 3 and supplementary video 1 show the 2D tomographic slices of ice cream's  
187 microstructural evolution during thermal cycling between -15 and -5 °C. A few salient  
188 observations can be made based on these images. First, ice crystals were continuously  
189 melting, decreasing in size, during the heating stage (Fig. 3(a-c)) and their  
190 morphologies became more spherical at “high” temperature (e.g. -5 °C, Fig. 3(c)). It



191 appears that the most significant morphologic change took place in the temperature  
192 range -7.5 to -5 °C. During cooling, the relatively round ice crystals grow into a more  
193 irregular morphology again (Fig. 3(d-f)), as expected for ice which has a high  
194 anisotropy in interfacial energy. Second, some of the air cells tended to coalesce. One  
195 example is indicated by the arrows in Fig. 3(a) and (b), where two neighboring large  
196 air cells gradually merged into one.

197 Third, it is observed that during the heating stage the microstructural features  
198 moved to the upper right corner of the region in Fig. 3, and observed more clearly in  
199 supplementary video 1, because of the volume shrinkage when ice crystals were melted  
200 into the unfrozen matrix. An air cell in Fig. 3(b) (indicated by a red arrow) shows the  
201 example of this movement. Note that the microstructure seemed to be compressed  
202 slightly due to the decreasing volume. It is mentioned that the expansion of volume  
203 seemed to be less significant during cooling as compared to the reduction in volume  
204 upon heating. The degree of expansion and relative movement during cooling also  
205 suggests the unfrozen matrix is quite flexible in response to external thermal change at  
206 high temperatures.

207 Microstructures in the longitudinal section were extracted for further examination  
208 and the results are shown in supplementary Fig. S1. The microstructural features (e.g.  
209 ice crystals and air cells) move less in the sample axis direction as compared to that in  
210 the cross section (Fig. 3). In other words, flotation of air cells along the sample axis  
211 does not occur, or we did not observe it occurring. This observation suggests that  
212 macro-fluid flow was largely inhibited due to the increased viscosity of the unfrozen

213 matrix due to the stabilizers present in the ice cream samples, which would inhibit the  
214 mobility of air cells [14]. In addition, it is also difficult to identify if drainage, which  
215 can play an important role in the instability of the air structure at high temperature [14],  
216 occurred in this work. However, it is also mentioned that microscopic fluid flow in the  
217 channels within the unfrozen matrix probably still occurs which cannot be resolved by  
218 the technique used in this work.

219

### 220 **3.2 3D Ice crystal evolution**

221 The volume fraction of ice crystals was examined first (Fig. 1), together with the  
222 calculated data based on the thermal properties (i.e. extrapolated melting points at  
223 various concentrations) of the ice cream formulation. Generally, the results measured  
224 by tomography compare well with the theoretical predictions. Minor errors might be  
225 caused by shrinkage and expansion, which would change some of the features measured.

226 Fig. 4 shows the 3D evolution of ice crystals during a thermal cycle. Ice crystals  
227 are individually color-rendered according to the equivalent diameter of each ice crystal.  
228 In the figure, the blue color corresponds to small size. As expected, ice crystals  
229 gradually decreased in size when the sample was heated, reaching the minimum size at  
230  $-5\text{ }^{\circ}\text{C}$ . This is consistent with the observation that more small blue ice crystals were  
231 present at  $-5\text{ }^{\circ}\text{C}$ , as shown in Fig. 4(c). The ice crystals then continuously grew in size  
232 when the sample was cooled again to  $-15\text{ }^{\circ}\text{C}$  (Fig. 4(d-e)).

233 Detailed examination of the position change for the 3D ice crystals after heating  
234 shows that the ice crystals moved upwards slightly, by less than  $20\text{ }\mu\text{m}$ . This movement

235 of the ice crystals might be caused by compression due to volume shrinkage, which  
236 drives the ice crystals to move upwards towards the center of the sample when the  
237 unfrozen matrix is less viscous at warmer temperatures [34].

238 The ice crystals in Fig. 4 were quantified in terms of equivalent diameter and size  
239 distribution (Fig. 5). The average equivalent diameter of ice crystals decreases upon  
240 heating from 101  $\mu\text{m}$  at  $-15\text{ }^\circ\text{C}$  to 87  $\mu\text{m}$  at  $-5\text{ }^\circ\text{C}$ , and then increases again to  $\sim 103\text{ }\mu\text{m}$   
241 at  $-15\text{ }^\circ\text{C}$  after cooling (Fig. 5(a)). Interestingly, the average equivalent diameter of the  
242 ice crystals is increased by  $\sim 2\text{ }\mu\text{m}$  at  $-15\text{ }^\circ\text{C}$  after a thermal cycle (Fig. 5(a)), suggesting  
243 a mild coarsening of ice crystals after thermal cycling. This coarsening of ice crystals  
244 at a storage temperature (e.g.  $-5\text{ }^\circ\text{C}$   $\sim -18\text{ }^\circ\text{C}$ ) with imposed oscillation temperatures (e.g.  
245  $\pm 2.5\text{ }^\circ\text{C}$ ) have been previously observed [10, 23, 35, 36]. It is pointed out that the  
246 performed thermal cycling in this study leads to much higher coarsening rates than  
247 isothermal storage at a certain temperature due to the temperature oscillations which,  
248 through observed melting-recrystallization process, increases the rate of the coarsening  
249 process [23].

250 The size distribution of the ice crystals at different temperatures was analyzed and  
251 the results are plotted in Fig. 5(b-d). In general, the size distribution curves reflect the  
252 temperature change, where the curves shift to the left upon heating (Fig. 5(c)), while  
253 they shift to the right during the cooling stage (Fig. 5(d)). It is also shown in Fig. 5(b)  
254 that the distribution of the two curves at  $-15\text{ }^\circ\text{C}$ , before and after the thermal cycle, is  
255 similar indicating only a minor change in the overall size after the thermal cycle was  
256 completed. This observation confirms that the ice crystals follow mainly a “melting-

257 recrystallization” mechanism and that other proposed mechanisms [37] play a small or  
258 negligible role.

259 Fig. 6 shows the morphological evolution of five individually separated ice crystals  
260 extracted for detailed examination. The morphology of the ice crystals changed during  
261 thermal cycling being more irregular at low temperature, e.g. -15 °C (Fig. 6(a) and (f)),  
262 while they became more spherical at “high” temperature, e.g. -5 °C (Fig. 6(c) and (d)).

263 The five ice crystals shown in Fig. 6 are quantified in terms of volume, volume  
264 change, specific surface area and sphericity, and the results are presented in Fig. 7. It is  
265 seen that the volume of each ice crystal keeps decreasing during heating and then  
266 increasing during cooling. The volume change (as compared to -15 °C before thermal  
267 cycling) shows that the volume of four of the ice crystals increased by 5.5-11%,  
268 indicating an increase in size of ice crystals, which is consistent with the observation of  
269 the overall increased equivalent diameter of ice crystals due to coarsening (Fig. 5(a)).  
270 However, the volume of ice crystal 3 decreased slightly, by ~4%. This ice crystal is one  
271 of the smallest of the 5, and most likely the other four ice crystals grew at the expense  
272 of ice crystal 3 via an Ostwald Ripening mechanism [37].

273 In this study, complete melting of ice crystals was not observed when the sample  
274 was heated to -5 °C since the smallest crystals will have dissolved in the seven thermal  
275 cycles prior to the *in situ* experiment. After seven cycles the ice crystals have grown to  
276 a large enough size such that they don’t completely melt during heating to -5 °C. This  
277 critical size can be extracted from Fig. 5(b). Note that the smallest crystals measured  
278 are about 60 μm diameter at -15 °C, and reduce in size by ~25 μm when heated to -5

279 °C. This observation suggests that ice crystals with an equivalent diameter less than 25  
280  $\mu\text{m}$  have a high probability of being completely melted during the applied thermal  
281 cycling. This is further supported by the observation that there is no significant number  
282 of crystals less than 30  $\mu\text{m}$  diameter at -5 °C. This finding confirms the measurements  
283 made in our previous study (i.e. Fig. 9(d) in ref. [5]) showing a significant reduction in  
284 the number of ice crystals during the first 7 cycles.

285 As mentioned, the morphology of ice crystals also changes during thermal cycling  
286 (Figs. 3 and 6). This change is quantified using measures of specific surface area and  
287 sphericity in Fig. 7 (c) and (d). For example, both the values of specific surface area  
288 and sphericity increased continuously as the ice crystals melt, decreasing as the ice  
289 crystals recrystallized. These changes support the observation that ice crystals become  
290 more spherical during melting (to minimize interfacial energy) and then became more  
291 faceted (irregular) during the recrystallization stage as they strive to reach the Wulff  
292 shape driven by anisotropy in interfacial energy [38, 39]. However, the ice crystals  
293 show only minor morphological changes after a thermal cycle, as indicated by the  
294 average sphericity of more than 300 ice crystals where the value increased from ~0.80  
295 at -15 °C to 0.84 at -5 °C during the heating stage, and then decreased to ~0.80 when  
296 cooled back to -15 °C. A similar trend was observed for the specific surface area. The  
297 minor changes to the size and morphology of ice crystals after one thermal cycle  
298 support our previous findings that only small differences in size and morphology of ice  
299 crystals were observed between the sample thermally cycled for 7 days and the sample  
300 cycled for 14 days [5].

301

### 302 **3.3 Unfrozen matrix evolution**

303 Fig. 8 and supplementary Fig. S2 (larger volume) show the 3D morphological  
304 evolution of the unfrozen matrix during a thermal cycle. The unfrozen matrix forms a  
305 very complex 3D network-like shape with ice crystals and air bubbles dispersed within  
306 the matrix. The 3D images in Fig. 8(a-f) show that the unfrozen matrix appeared thicker  
307 between air cells upon heating, while they became thinner as the sample cooled down  
308 presumably due to the melting of the ice crystals during heating and recrystallization  
309 during cooling. This is reflected by the quantified thickness of the unfrozen matrix  
310 analyzed in a  $721 \times 721 \times 504 \mu\text{m}^3$  volume (supplementary Fig. S2), as plotted in Fig. 8(g)  
311 where the thickness monotonically increased from  $\sim 12.6 \mu\text{m}$  at  $-15 \text{ }^\circ\text{C}$  to  $27.0 \mu\text{m}$  at  $-5$   
312  $^\circ\text{C}$ , and then decreased to  $\sim 19.6 \mu\text{m}$  at  $-15 \text{ }^\circ\text{C}$  at the end of the thermal cycle.

313 Figure 8 shows that the thickness of the unfrozen matrix is greater during the cooling  
314 stage than during the heating stage, such that the thickness increases by  $\sim 6 \mu\text{m}$  after  
315 the thermal cycle was completed. After detailed examination, we propose that the  
316 formation of a local region of the matrix (e.g. upper right corner in Fig. 3) that is  
317 concentrated with more water molecules, is responsible for this change. The shrinkage  
318 of the sample, and the associated macro-flow induced by the compression effect, might  
319 have accelerated the formation of a larger region of low viscosity matrix. During the  
320 cooling stage, no additional new ice crystals were nucleated under the cooling rate  
321 studied. Thus, the measured thickness of the unfrozen matrix in this region is higher  
322 compared to that before thermal cycling.

323 In addition, the ‘strength’ (or viscosity) of the unfrozen matrix would decrease at the  
324 higher temperatures during heating due to a reduction in viscosity allowing more  
325 movement of the matrix to accommodate the overall shrinkage of the sample (see  
326 above). Thus, a shift of the structure was observed during the heating stage, suggesting  
327 a significant negative impact of the ice melting process on structural stability.

328

### 329 **3.4 3D air cell evolution**

330 A few 3D air cells were extracted to examine the coarsening mechanism. Fig. 9  
331 shows one example where two separate air cells gradually merge into one. It is observed  
332 that the air cells at -15 °C are not necessarily round, instead, they have many concave  
333 regions (indicated by an arrow in Fig. 9(a)), or even an elongated shape for some cases  
334 as seen in Fig. 3. Upon heating, the two air cells merged by creating a bridge between  
335 them (Fig. 9(b)), and then the bridge (or neck) continued to thicken with increasing  
336 temperature during the heating stage. Meanwhile, some of the concave regions on the  
337 air cell surface gradually disappear forming a smooth or spherical surface. The gradual  
338 rounding was driven by the reduction in surface energy of the air/unfrozen matrix  
339 interface [40].

340 Decreased viscosity of the unfrozen matrix at “warmer” temperatures upon heating  
341 increases the diffusion rate of gas between air cells and promotes coalescence of air  
342 cells. It is mentioned that adding stabilizers and emulsifier to ice cream helps reduce  
343 air cell coarsening, due to the increased extent of fat destabilization and the increased  
344 viscosity of the matrix phase, respectively [14]. It is also noticed that the surrounding

345 ice crystals significantly affect the shape of the area around the neck between merging  
346 cells. One example is shown in Fig. 9(d-2). It is likely that those ice crystals in the  
347 vicinity of the neck limit further coalescence of the two air cells due to the constraint  
348 imposed by the ice crystal imbedded unfrozen matrix (also see supplementary Fig. S3).

349       Upon cooling, some of the phenomena observed during the heating stage act in  
350 reverse. That is, the surface of air cells became rough or even distorted again at “cooler”  
351 temperatures, e.g. -12 °C in Fig. 9(g). This is most obvious at the lowest temperature of  
352 -15 °C (Fig. 9(h) and (h-2)). The changes are likely to be caused by two main factors.  
353 One is that the growing ice crystals continue to push towards the air cells through the  
354 unfrozen matrix. This is realized more easily when the unfrozen matrix becomes thinner  
355 and thinner as more water molecules are attached to the recrystallizing ice crystals. In  
356 total, twenty ice crystals were observed to grow around the air cells shown in Fig. 9  
357 (see supplementary Fig. S3). The second factor is that the pressure within the air cell  
358 decreases with the decreasing temperature according to the ideal gas law ( $PV=nRT$ ),  
359 releasing some of the force on the surface that resists morphological change. It is  
360 mentioned that the final morphological change is a result of competition between the  
361 force imposed on the air cell surface by the growing ice crystals and the surface tension  
362 of the air cell/matrix interface. It seems that for the case in Fig. 9 the force imposed by  
363 the growing ice crystals through the matrix was greater than the surface tension at  
364 temperatures lower than  $\sim -12$  °C, under which the surface started to deform  
365 significantly. In addition, the coalescence process seemed to be inhibited by the  
366 increased viscosity at low temperatures during the cooling stage, indicating a



367 significantly reduced rate of morphological change of the air cells than during the  
368 heating stage.

369 Pelan et al. [41] and Rohenkohl and Kohlus [42] both suggested that the  
370 coalescence of air cells to create large coarsened air pockets was the major destabilizing  
371 mechanism in the ice cream they studied. A previous study revealed that the storage of  
372 ice cream without emulsifier or stabilizer at -15 °C for 16 days lead to interconnected  
373 channels [14]. Our recent observations also showed that thermal cycling of ice cream  
374 between -15 and -5 °C for 14 days resulted in a very complex interconnected air  
375 structure [5]. Although Ostwald ripening was observed in the aerated emulsions [43],  
376 the *in situ* observations in this study strongly suggest that for ice cream that was cycled  
377 for seven times coalescence is the dominant mechanism responsible for the creation of  
378 complex interconnected air structures. It should be noted that gas formation can occur  
379 due to radiation damage, resulting in molecular bond cleavage (H-H and O-O) or water  
380 photolysis, as reported in water under high pressures [44]. If this is occurring, it could  
381 explain the increase in bubble volume fraction and the coarsening of bubbles. However,  
382 the increase in bubble volume will not have a significant impact on the coarsening of  
383 the ice crystals. Gas formation due to irradiation is an open question, as is how this  
384 might affect bubble coarsening.

385 Apart from the coarsening of air cells, another interesting phenomenon was  
386 observed, i.e. the reduced volume of some of the air cells after thermal cycling. Fig. 10  
387 shows the evolution of three individually extracted air cells during thermal cycling, and  
388 their corresponding quantified volume changes are plotted in Fig. 11. The overall

389 volume fraction of air cells was also analyzed and the result was observed to decrease  
390 monotonically during the heating stage, and continued to decrease until 0.256 at ~ -7  
391 °C during cooling before it started to rise upon further freezing (Fig. 11(a)). The volume  
392 fraction after the thermal cycle (~ 0.285) was lower compared to that before thermal  
393 cycling began (~0.330). This corresponds to a reduction of volume fraction by ~ 13.8%.  
394 The trend of the volume change of the three individual air cells is consistent with that  
395 for the overall volume change. It is also noticed that the change of volume is even more  
396 than 50% for Air 2 and Air 3, and that those two air cells did not grow in size during  
397 the cooling stage. Detailed mechanisms here are still unknown. It is unlikely that the  
398 hydrostatic pressure causes such a large change, as the sample height is quite small.  
399 The shift of the sample during the thermal cycle might contribute to some measurement  
400 errors; however, we believe the volume (thus size) change of the air cells is the main  
401 contribution, which is supported by the volume change of all three air cell cases (Fig.  
402 11(c-d)). We suspect that the diffusion of gas into the matrix and the surrounding air  
403 cells, as well as out of the whole sample, at the warm temperatures might have  
404 contributed to this change. The detailed mechanisms will be investigated in our future  
405 study.

406

### 407 **3.5 Summary of microstructural evolution mechanisms**

408 Here, we summarize the mechanisms that control the microstructural evolution of  
409 ice cream as observed in this *in situ* study (Table 1). Generally, the microstructural  
410 evolution of ice cream during thermal cycling is controlled by the interaction of three  
411 phases.

412        Regarding ice crystals, nucleation of new ice crystals does not occur under the  
413 cooling condition studied in this work. The melting-recrystallization mechanism  
414 hypothesized as an important mechanism in our previous study [5] was quantified by  
415 analyzing the 4-D tomographs during the thermal cycle. The melting and  
416 recrystallization of ice crystals also affects the air cell's morphology, as well as its  
417 coarsening process, through the unfrozen matrix layer between the ice crystals and the  
418 air structures. For the air phase, coalescence of air cells is clearly observed to be  
419 responsible for the coarsening mechanism. For the sample that was initially thermally  
420 cycled for seven times, Ostwald ripening takes a less important role in the coarsening  
421 of both ice crystals and air cells during thermal cycling. The continuous reduction of  
422 air cell volume needs further investigation. The third phase, the unfrozen matrix, is a  
423 crucial component controlling the microstructural stability of ice cream. It acts as the  
424 reservoir for the water from dissolving ice during heating and releases water for  
425 recrystallization of the ice crystals during the cooling cycle. The network of unfrozen  
426 matrix, reinforced by the distributed ice crystals (and dissolved hydrocolloids), holds  
427 the whole structure together and greatly influences the structural stability of ice cream  
428 when subjected to external temperature variations.

429

#### 430 **4. Conclusions**

431        Using 4D synchrotron X-ray tomography, we investigated the effect of thermal  
432 variation on the microstructural stability of ice cream during a heating and cooling cycle  
433 between -15 °C and -5 °C, at a ramp rate of 0.05 °C/min. A new data reconstruction and

434 image processing method was developed, enabling the large 4D data sets to be  
435 segmented and quantified. The experimental set-up, as well as the image processing  
436 routine developed, can be applied to a wide range of soft materials.

437 The dynamic evolution of individual microstructural features, i.e. an ice crystal, air  
438 cell, and unfrozen matrix, was quantitatively analyzed. The findings integrate the *ex*  
439 *situ* observations made in our previous work enhancing our understanding of the  
440 mechanisms controlling ice cream's microstructural evolution. The experimental  
441 results in this study reveal important physical mechanisms that influence  
442 microstructural instability: that is, the coarsening of air cells takes place mainly through  
443 the coalescence of neighboring air cells, while ice crystal growth results from the  
444 melting-recrystallization mechanism during thermal cycling, both of which lead to  
445 degradation of ice cream's microstructure. The unfrozen matrix plays an important role  
446 in maintaining the integrity of the structure of ice cream while being flexible enough at  
447 the higher temperatures to reduce the stresses imposed during heating and then cooling  
448 by the melting and recrystallization of the ice crystals.

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455 and group members for technical support.

456

457 **Data statement**

458 Representative samples of the research data are shown in the figures. Other datasets  
459 generated during and/or analysed during this study are not publicly available due to  
460 their large size but are available from the corresponding author on reasonable request.

461

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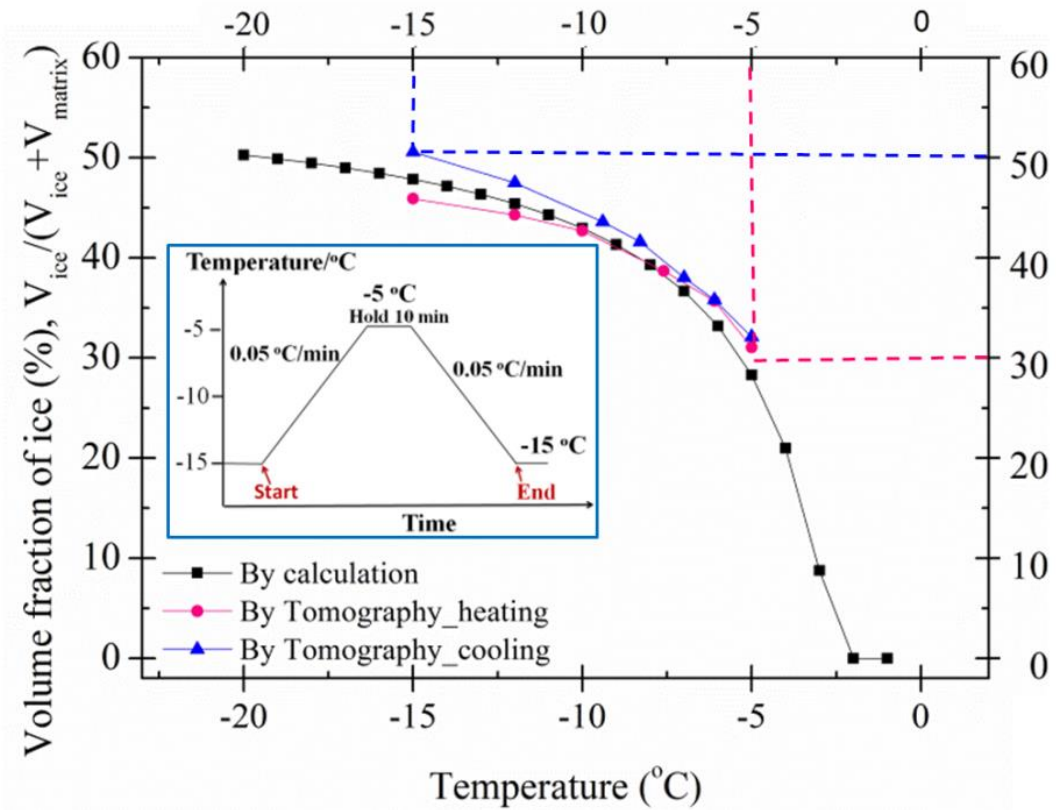
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562 induced dissociation of H<sub>2</sub>O and formation of an O-2-H-2 alloy at high pressure, *Science* 314(5799)  
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568 Fig. 1 Change in ice volume fraction as a function of temperature. The inset shows the

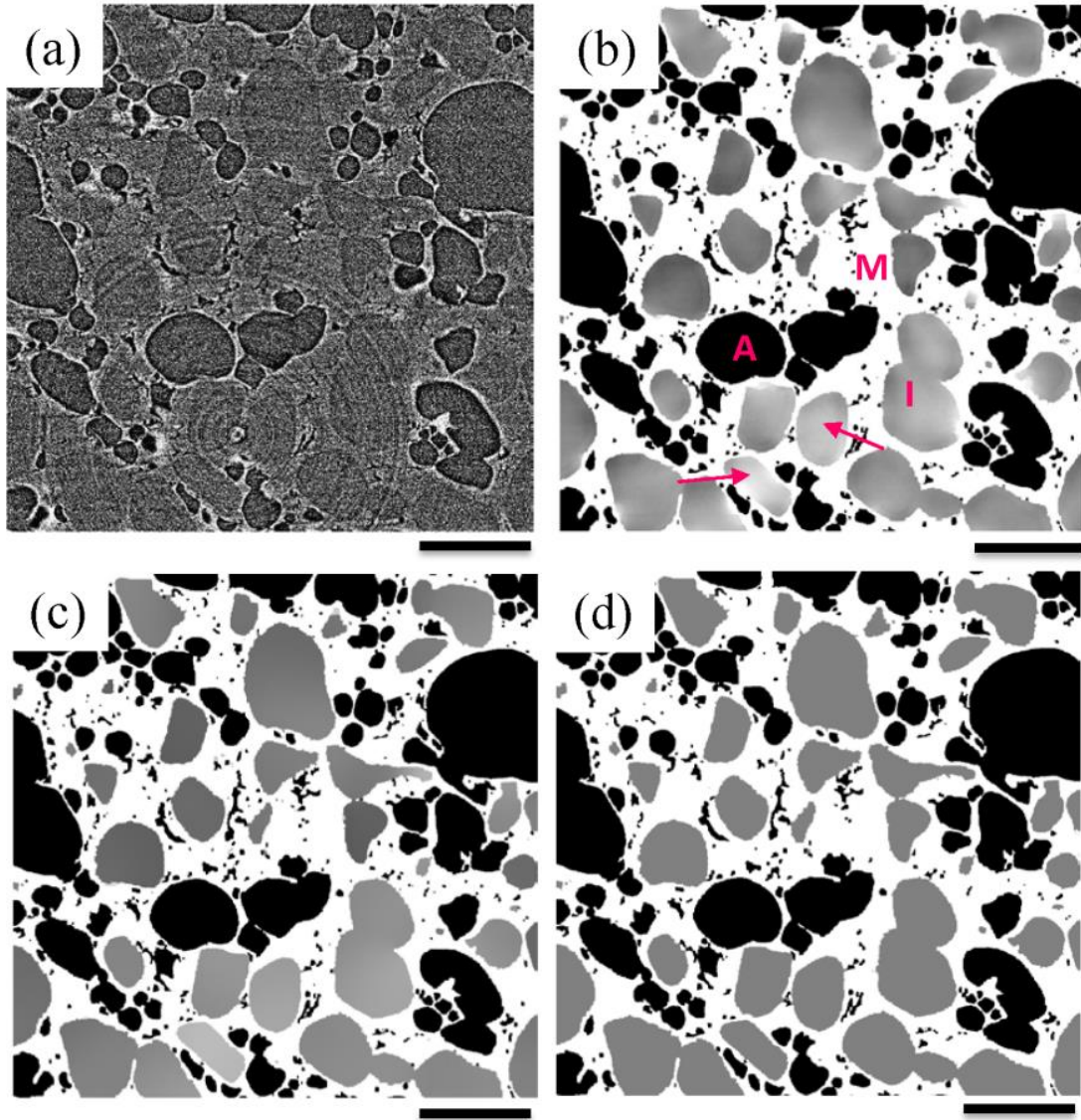
569 thermal cycling history of the ice cream sample for the *in situ* synchrotron experiment.

570 The measured phase diagram (i.e. extrapolated melting points at various

571 concentrations) of the ice cream formulation is presented in [5].

572





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574 Fig. 2 Reconstructed images using (a) conventional FBP reconstruction method and (b)

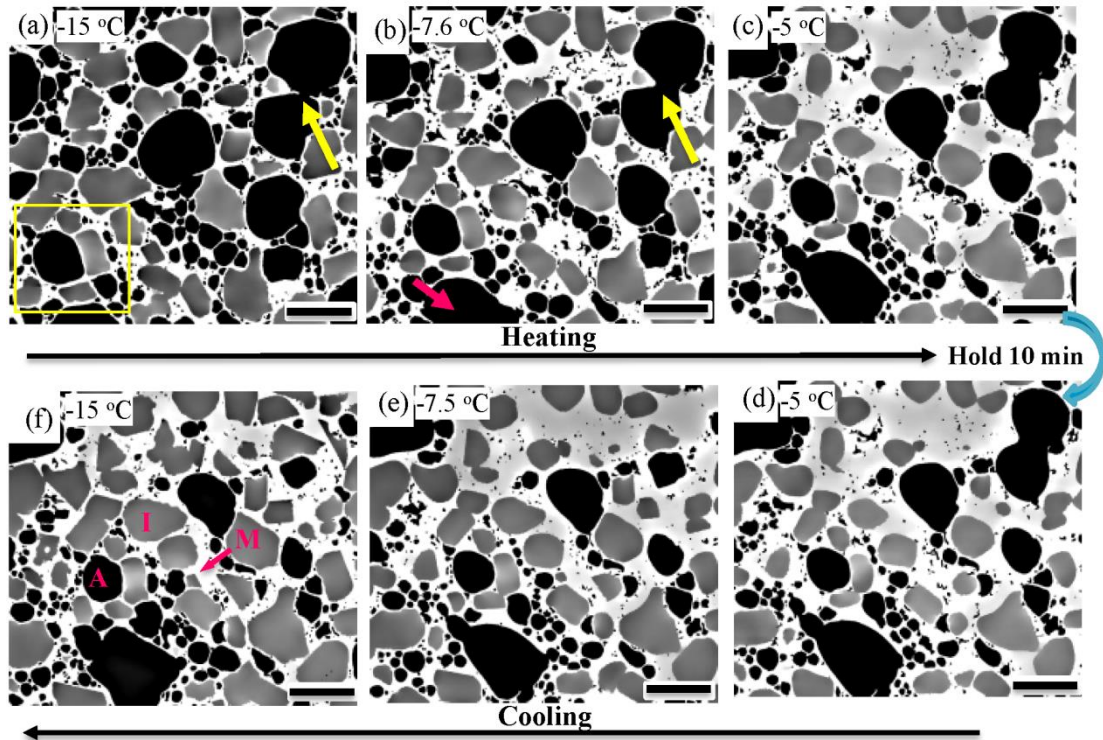
575 MBIR reconstruction method; (c) post-processed image of reconstructed MBIR image

576 of (b), note more homogeneous (equalized) intensities compared to (b) within ice-

577 crystals; (d) 3-phases segmentation using (c). A, I and M in (b) stand for air cell, ice

578 crystal and unfrozen matrix, respectively. Scale bar =150  $\mu\text{m}$ .

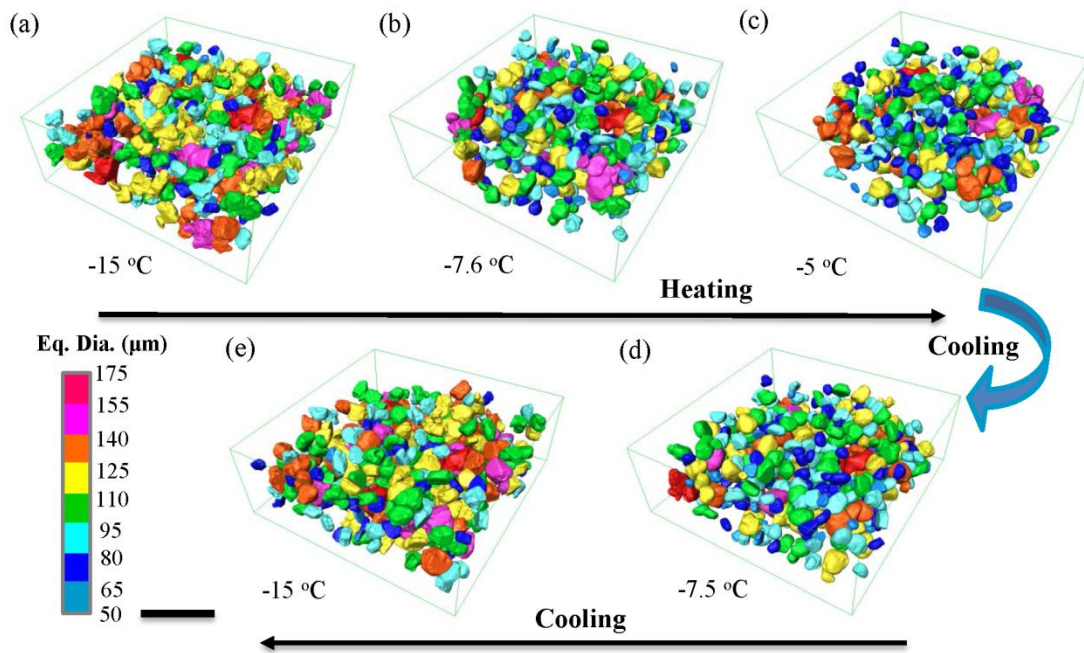
579



580

581 Fig. 3 Reconstructed tomographic slices (prior to the equalization step) showing the  
 582 overall microstructural evolution of ice cream during a thermal cycle: (a) -15 °C, (b) -  
 583 7.6 °C, and (c) -5 °C during heating, (d) after holding at -5 °C for 10 min, (e) -7.5 °C,  
 584 and (f) -15 °C during refreezing. A, I and M in (f) stand for air cell, ice crystal and  
 585 unfrozen matrix, respectively. Scale bar equals 150  $\mu\text{m}$ .

586



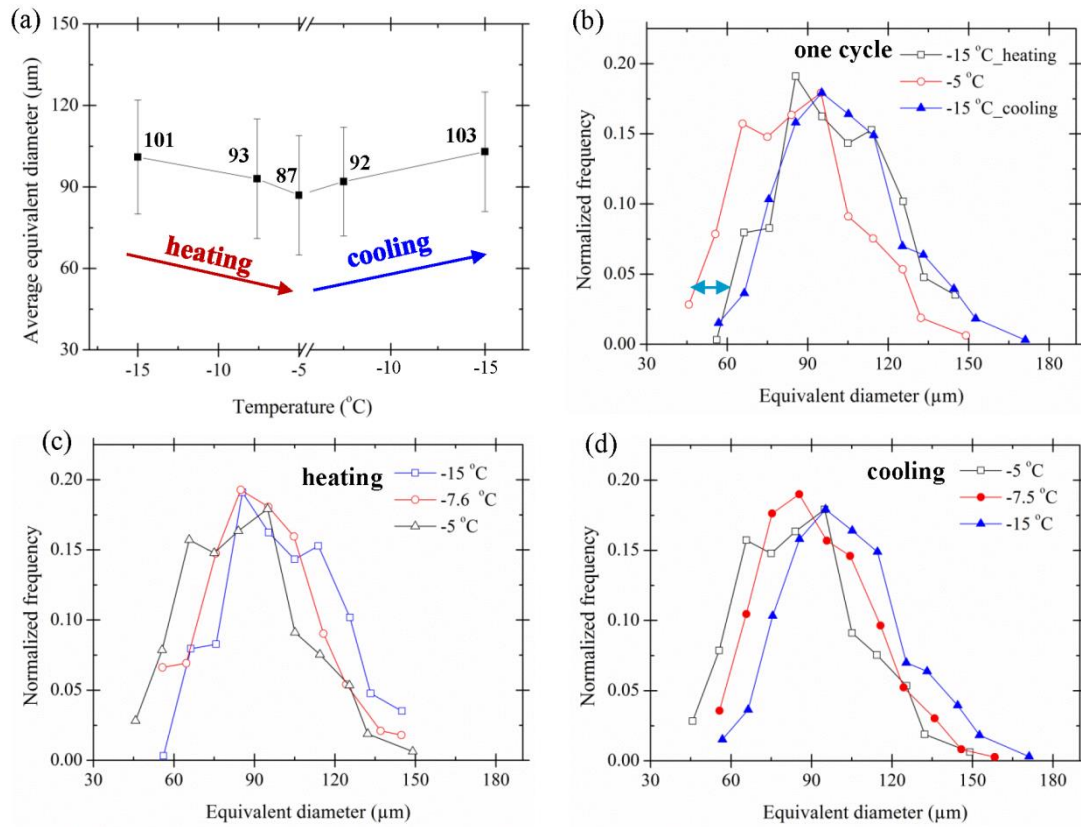
587

588 Fig. 4 3D ice crystal evolution in a  $1416 \times 1416 \times 504 \mu\text{m}^3$  volume during a thermal cycle:

589 (a) -15 °C, (b) -7.6 °C, (c) -5 °C, (d) -7.5 °C, and (e) -15 °C. Ice crystals are size-colored

590 using the equivalent diameter. Scale bar equals  $500 \mu\text{m}$ .

591



592

593 Fig. 5 Quantified ice crystal size during a thermal cycle: (a) change of average

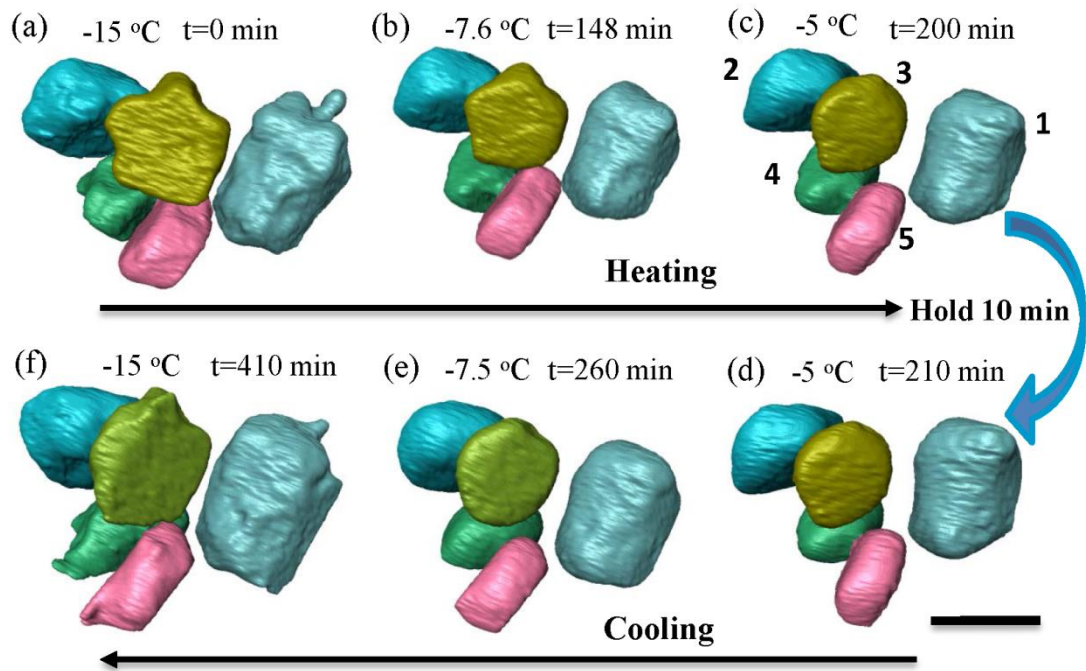
594 equivalent diameter of ice crystals during a thermal cycle; (b-d) size distribution of ice

595 crystals during (b) a complete thermal cycle, (c) heating stage and (d) cooling stage.

596 The arrow in (b) indicates the size shift of the curves. Note, more than 300 ice crystals

597 were analyzed.

598

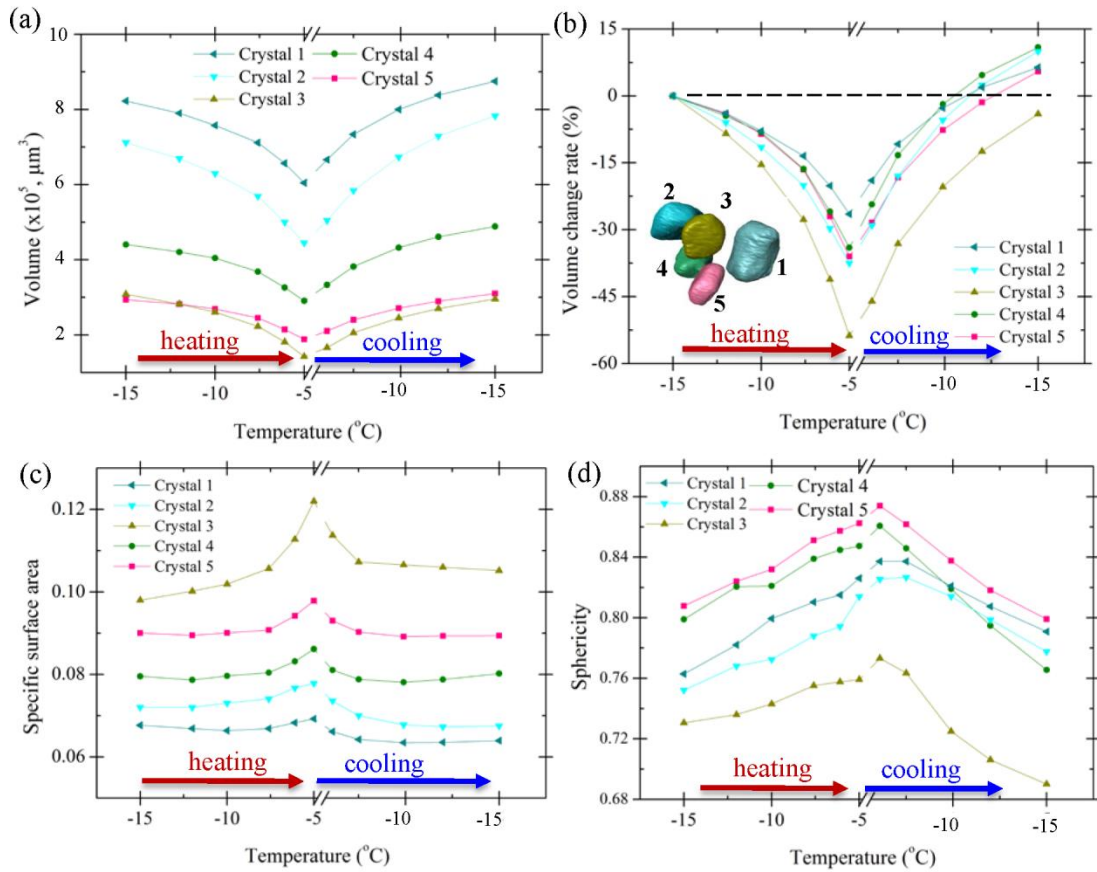


599

600 Fig. 6 3D morphological evolution of five ice crystals during a thermal cycle: (a) -15  
 601 °C, (b) -7.6 °C, (c) -5 °C, (d) after holding at -5 °C for 10 min, (e) -7.5 °C, and (f) -15  
 602 °C. The time is indicated in each figure during the thermal cycle. Scale bar 150 μm for

603 all images. Numbers in (c) match the ice crystals analyzed in Fig. 7.

604



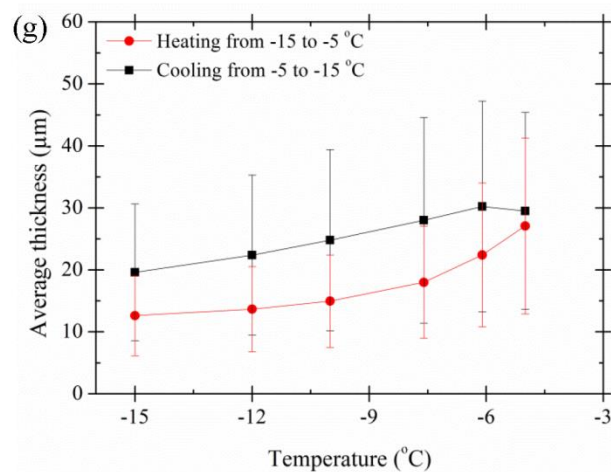
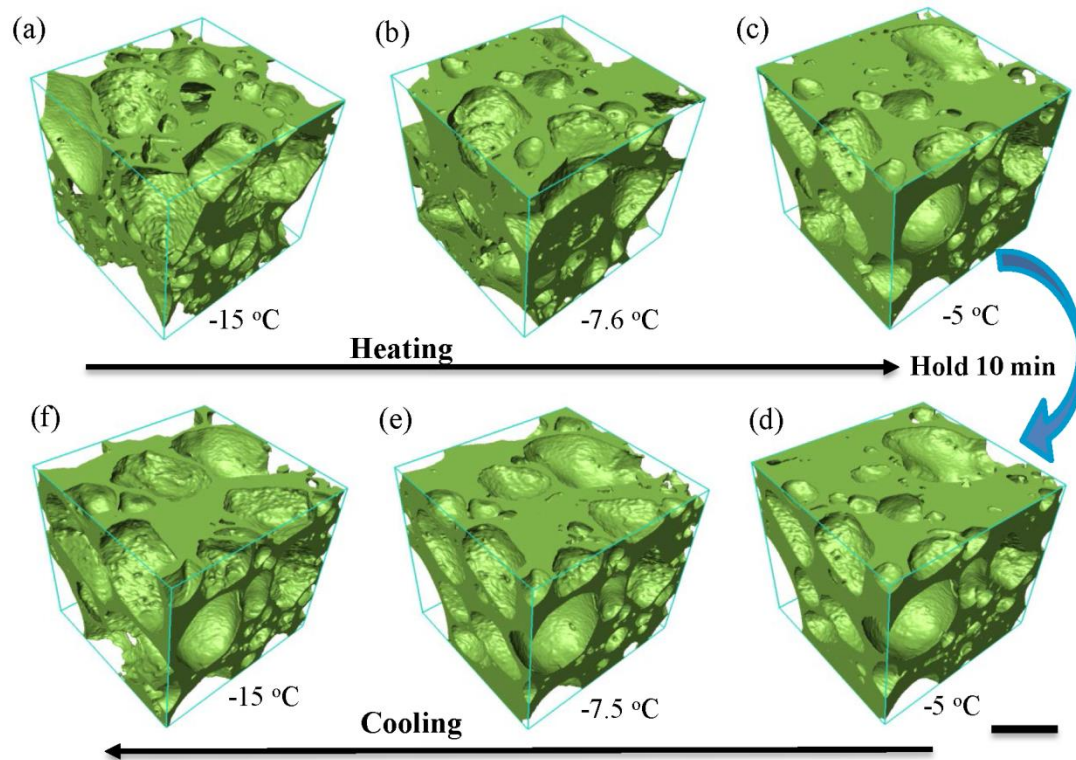
605

606 Fig. 7 Quantified results of five ice crystals during a thermal cycle: (a) volume, (b)

607 volume change, (c) specific surface area, and (d) sphericity. Note the colours of the

608 plots in each figure are identical to the colour-rendered ice crystals in (b).

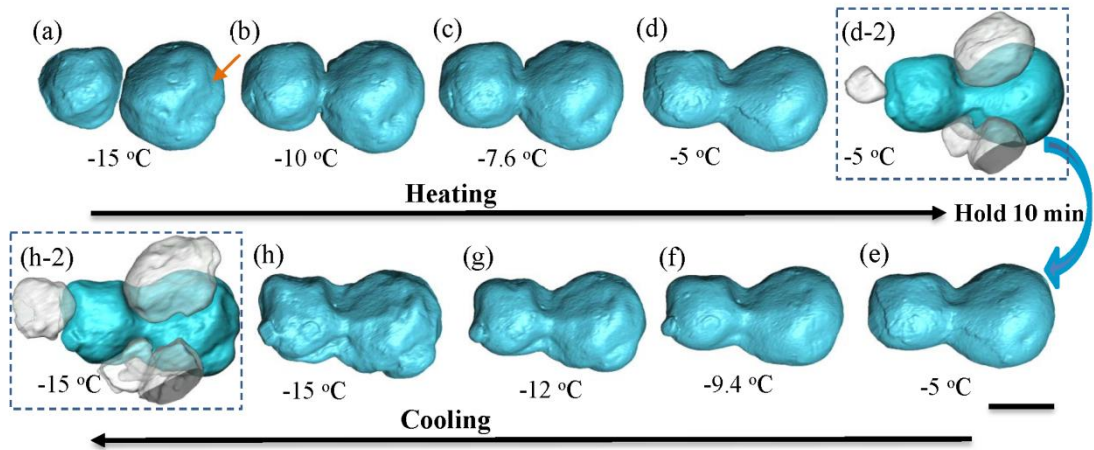
609



610

611 Fig. 8 3D morphological evolution of unfrozen matrix within a  $259 \times 243 \times 243 \mu\text{m}^3$   
 612 volume during a thermal cycle: (a)  $-15 \text{ }^\circ\text{C}$ , (b)  $-7.6 \text{ }^\circ\text{C}$ , (c)  $-5 \text{ }^\circ\text{C}$ , (d) after holding at  $-5$   
 613  $^\circ\text{C}$  for 10 min, (e)  $-7.5 \text{ }^\circ\text{C}$ , and (f)  $-15 \text{ }^\circ\text{C}$ ; (g) Average thickness of the unfrozen matrix  
 614 as a function of temperature. Note, the thickness is measured within a  $721 \times 721 \times 504$   
 615  $\mu\text{m}^3$  volume, identical to the domain as shown in supplementary Fig. S2. Figures (a-f)  
 616 share the same scale bar. Scale bar equals  $100 \mu\text{m}$ .

617



618

619 Fig. 9 Coalescence of two air cells during the heating stage (a-d) and cooling stage (e-

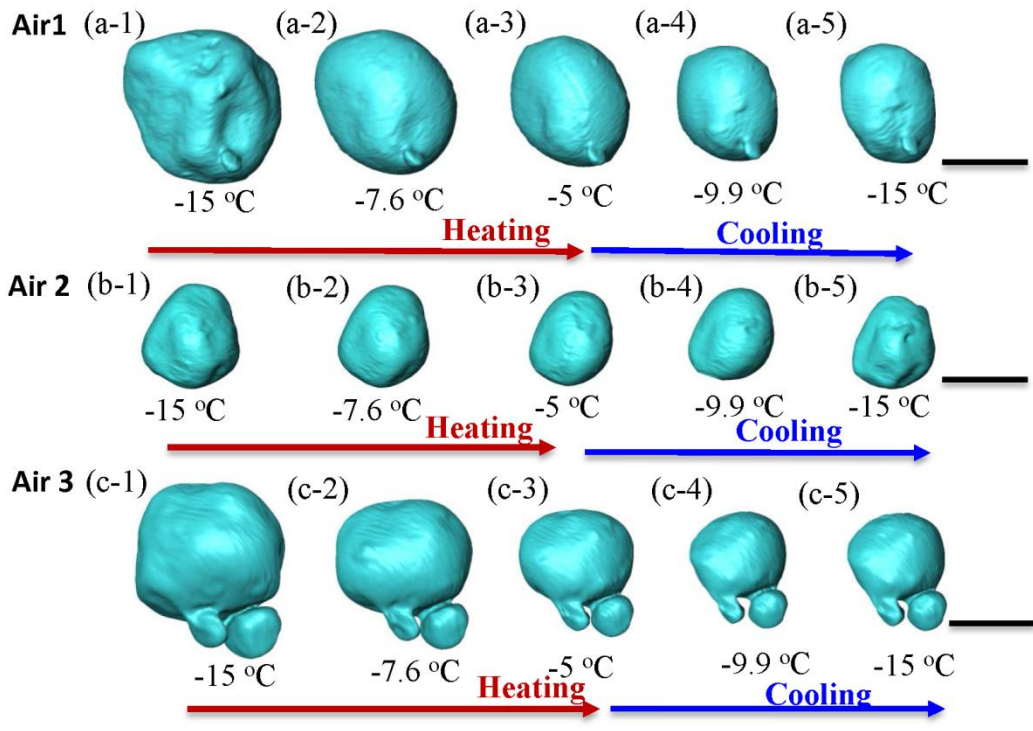
620 h) of a thermal cycle. (d-2) and (h-2) show the morphological relationship between the

621 surrounding ice crystals and the air cell at -5 °C and -15 °C, respectively. Scale bar 100

622 μm for all images.

623



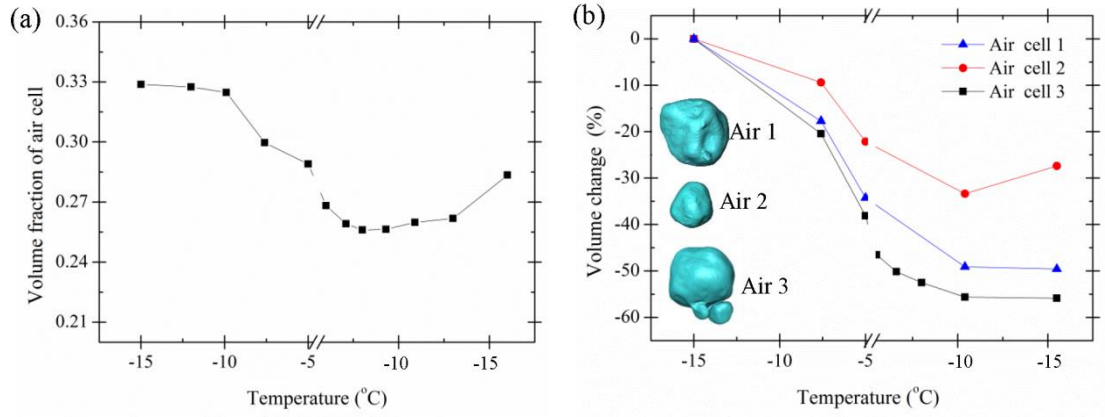


624

625 Fig. 10 Morphological evolution of three individual air cell cases during a thermal cycle.

626 Scale bar 100  $\mu\text{m}$  for all images.

627



628

629 Fig. 11 Volume change of air cells as a function of temperature during a thermal cycle:

630 (a) overall volume change of air cells in a  $1416 \times 1416 \times 504 \mu\text{m}^3$  volume, (b) volume



631 change of three individual air cells.

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**Table**

Table 1 A summary of the microstructural changes that occur during a thermal cycling from -5 to -15 °C highlighting the differences between what occurs during the first seven cycles and the following seven cycles. The arrows indicate that part of the cycle over which most change occurs.

	Heat to -5 °C	Hold at -5 °C	Cool to -15 °C	Hold at -15 °C
<b>Air cells</b>				
Coalescence of neighbouring air cells. 				
1 to 7 cycles	Size increases although some small air cells remain leading to a bimodal distribution and remain equiaxed.		The air cells shrink.	Cells remain relatively spherical.
8 to 14 cycles	Coalescence of neighbouring air cells.		The air cells become irregular due to constraint by matrix and ice crystals.	Air cells become interconnected and form channels within the matrix network.
<i>Air cells continue to grow into a large interconnected network of irregular shapes as the number of thermal cycles increase. The morphology is constrained by the network of unfrozen matrix and ice crystals.</i>				
<b>Ice crystals</b>				
	Melt by ~40%		Grow by ~66%	
Dissolution  Recrystallization				
1 to 7 cycles	Size of ice crystals decrease and those < 25 µm melt completely. The morphology becomes rounded.	The size and morphology of crystals change little.	Size of ice crystals increase and no nucleation of new crystals occurs. The morphology becomes irregular during recrystallization.	Over the cycle, the size increases significantly and the number decrease significantly.
8 to 14 cycles	Size decreases by ~ 25 µm dissolving and the number changes little.	The size of crystals change little.	Size of ice crystals increase by about 25 µm. The number remains unchanged.	Over the cycle, the size increases by a small amount and the number do not increase.
<i>After 7 thermal cycles, the number changes little and the size of ice crystals increase slowly. The ice crystals form networks within the unfrozen matrix network.</i>				
<b>Unfrozen matrix</b>				
Water content	Increases	High	Decreases	Low
Viscosity	Viscosity decreases	Low	Viscosity increases	High
Mechanical response	Matrix becomes and remains flexible, reducing residual stresses		Matrix becomes less flexible	Matrix is effectively rigid
Total volume	Ice cream expands	May shrink somewhat	Ice cream shrinks	Relatively constant
<i>Alignment of ice crystals with the unfrozen matrix network occurs to minimise surface energy and reduce local stress with each additional thermal cycle. At the warmer temperatures the matrix</i>				

*becomes flexible also reducing stresses developed by the constraint of ice crystal and air cells.*

639