



Search for a Higgs boson in the diphoton final state using the full CDF data set from $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV

CDF Collaboration

T. Aaltonen^w, B. Álvarez González^{j,27}, S. Amerio^{ap}, D. Amidei^{ah}, A. Anastassov^{q,25}, A. Annovi^s, J. Antos^{m,n}, G. Apollinari^q, J.A. Appel^q, T. Arisawa^{bj}, A. Artikov^o, J. Asaadi^{bd}, W. Ashmanskas^q, B. Auerbach^{bm}, A. Aurisano^{bd}, F. Azfar^{ao}, W. Badgett^q, T. Bae^{aa}, A. Barbaro-Galtieri^{ab}, V.E. Barnes^{ax}, B.A. Barnett^y, P. Barria^{as,au}, P. Bartos^{m,n}, M. Baucé^{ap,aq}, F. Bedeschi^{as}, S. Behari^y, G. Bellettini^{as,at}, J. Bellinger^{bl}, D. Benjamin^p, A. Beretvas^q, A. Bhatti^{az}, D. Bisello^{ap,aq}, I. Bizjak^{ad}, K.R. Bland^e, B. Blumenfeld^y, A. Bocci^p, A. Bodek^{ay}, D. Bortoletto^{ax}, J. Boudreau^{aw}, A. Boveia^l, L. Brigliadori^{f,g}, C. Bromberg^{ai}, E. Brucken^w, J. Budagov^o, H.S. Budd^{ay}, K. Burkett^q, G. Busetto^{ap,aq}, P. Bussey^u, A. Buzatu^{ag}, A. Calamba^k, C. Calancha^{ae}, S. Camarda^d, M. Campanelli^{ad}, M. Campbell^{ah}, F. Canelli^{l,q}, B. Carls^x, D. Carlsmith^{bl}, R. Carosi^{as}, S. Carrillo^{r,14}, S. Carron^q, B. Casal^{j,12}, M. Casarsa^{be}, A. Castro^{f,g}, P. Catastini^v, D. Cauz^{be}, V. Cavaliere^x, M. Cavalli-Sforza^d, A. Cerri^{ab,7}, L. Cerrito^{ad,20}, Y.C. Chen^a, M. Chertok^h, G. Chiarelli^{as}, G. Chlachidze^q, F. Chlebana^q, K. Cho^{aa}, D. Chokheli^o, W.H. Chung^{bl}, Y.S. Chung^{ay}, M.A. Ciocci^{as,au}, A. Clark^t, C. Clarke^{bk}, G. Compostella^{ap,aq}, M.E. Convery^q, J. Conway^h, M. Corbo^q, M. Cordelli^s, C.A. Cox^h, D.J. Cox^h, F. Crescioli^{as,at}, J. Cuevas^{j,27}, R. Culbertson^q, D. Dagenhart^q, N. d'Ascenzo^{q,24}, M. Datta^q, P. de Barbaro^{ay}, M. Dell'Orso^{as,at}, L. Demortier^{az}, M. Deninno^f, F. Devoto^w, M. d'Errico^{ap,aq}, A. Di Canto^{as,at}, B. Di Ruzza^q, J.R. Dittmann^e, M. D'Onofrio^{ac}, S. Donati^{as,at}, P. Dong^q, M. Dorigo^{be}, T. Dorigo^{ap}, K. Ebina^{bj}, A. Elagin^{bd}, A. Eppig^{ah}, R. Erbacher^h, S. Errede^x, N. Ershaidat^{q,31}, R. Eusebi^{bd}, S. Farrington^{ao}, M. Feindt^z, J.P. Fernandez^{ae}, R. Field^r, G. Flanagan^{q,22}, R. Forrest^h, M.J. Frank^e, M. Franklin^v, J.C. Freeman^q, Y. Funakoshi^{bj}, I. Furic^r, M. Gallinaro^{az}, J.E. Garcia^t, A.F. Garfinkel^{ax}, P. Garosi^{as,au}, H. Gerberich^x, E. Gerchtein^q, S. Giagu^{ba}, V. Giakoumopoulou^c, P. Giannetti^{as}, K. Gibson^{aw}, C.M. Ginsburg^q, N. Giokaris^c, P. Giromini^s, G. Giurgiu^y, V. Glagolev^o, D. Glenzinski^q, M. Gold^{ak}, D. Goldin^{bd}, N. Goldschmidt^r, A. Golossanov^q, G. Gomez^j, G. Gomez-Ceballos^{af}, M. Goncharov^{af}, O. González^{ae}, I. Gorelov^{ak}, A.T. Goshaw^p, K. Goulianos^{az}, S. Grinstein^d, C. Grosso-Pilcher^l, R.C. Group^{q,bi}, J. Guimaraes da Costa^v, S.R. Hahn^q, E. Halkiadakis^{bc}, A. Hamaguchi^{an}, J.Y. Han^{ay}, F. Happacher^s, K. Hara^{bg}, D. Hare^{bc}, M. Hare^{bh}, R.F. Harr^{bk}, K. Hatakeyama^e, C. Hays^{ao}, M. Heck^z, J. Heinrich^{ar}, M. Herndon^{bl}, S. Hewamanage^e, A. Hocker^q, W. Hopkins^{q,8}, D. Horn^z, S. Hou^a, R.E. Hughes^{al}, M. Hurwitz^l, U. Husemann^{bm}, N. Hussain^{ag}, M. Hussein^{ai}, J. Huston^{ai}, G. Introzzi^{as}, M. Iori^{ba,bb}, A. Ivanov^{h,17}, E. James^q, D. Jang^k, B. Jayatilaka^p, E.J. Jeon^{aa}, S. Jindariani^q, M. Jones^{ax}, K.K. Joo^{aa}, S.Y. Jun^k, T.R. Junk^q, T. Kamon^{bd,aa}, P.E. Karchin^{bk}, A. Kasmi^e, Y. Kato^{an,16}, W. Ketchum^l, J. Keung^{ar}, V. Khotilovich^{bd}, B. Kilminster^q, D.H. Kim^{aa}, H.S. Kim^{aa}, J.E. Kim^{aa}, M.J. Kim^s, S.B. Kim^{aa}, S.H. Kim^{bg}, Y.K. Kim^l, Y.J. Kim^{aa}, N. Kimura^{bj}, M. Kirby^q, S. Klimentenko^r, K. Knoepfel^q, K. Kondo^{bj,1}, D.J. Kong^{aa}, J. Konigsberg^r, A.V. Kotwal^p, M. Kreps^z, J. Kroll^{ar}, D. Krop^l, M. Kruse^p, V. Krutelyov^{bd,4}, T. Kuhr^z, M. Kurata^{al,26}, S. Kwang^l, A.T. Laasanen^{ax}, S. Lami^{as}, S. Lammel^q, M. Lancaster^{ad}, R.L. Lander^h, K. Lannon^{al,26}, A. Lath^{bc}, G. Latino^{as,au}, T. LeCompte^b, E. Lee^{bd}, H.S. Lee^{l,18}, J.S. Lee^{aa}, S.W. Lee^{bd,29}, S. Leo^{as,at}, S. Leone^{as}, J.D. Lewis^q, A. Limosani^{p,21}, C.-J. Lin^{ab}, M. Lindgren^q, E. Lipeles^{ar}, A. Lister^t, D.O. Litvintsev^q, C. Liu^{aw}, H. Liu^{bi}, Q. Liu^{ax}, T. Liu^q, S. Lockwitz^{bm}, A. Loginov^{bm}, D. Lucchesi^{ap,aq}, J. Lueck^z, P. Lujan^{ab}, P. Lukens^q, G. Lungu^{az}, J. Lys^{ab}, R. Lysak^{m,n,6}, R. Madrak^q, K. Maeshima^q, P. Maestro^{as,au}, S. Malik^{az},

G. Manca^{ac,2}, A. Manousakis-Katsikakis^c, F. Margaroli^{ba}, C. Marino^z, M. Martínez^d, P. Mastrandrea^{ba}, K. Matera^x, M.E. Mattson^{bk}, A. Mazzacane^q, P. Mazzanti^f, K.S. McFarland^{ay}, P. McIntyre^{bd}, R. McNulty^{ac,11}, A. Mehta^{ac}, P. Mehtala^w, C. Mesropian^{az}, T. Miao^q, D. Mietlicki^{ah}, A. Mitra^a, H. Miyake^{bg}, S. Moed^q, N. Moggi^f, M.N. Mondragon^{q,14}, C.S. Moon^{aa}, R. Moore^q, M.J. Morello^{as,av}, J. Morlock^z, P. Movilla Fernandez^q, A. Mukherjee^q, Th. Muller^z, P. Murat^q, M. Mussini^{f,g}, J. Nachtman^{q,15}, Y. Nagai^{bg}, J. Naganoma^{bj}, I. Nakano^{am}, A. Napier^{bh}, J. Nett^{bd}, C. Neu^{bi}, M.S. Neubauer^x, J. Nielsen^{ab,5}, L. Nodulman^b, S.Y. Noh^{aa}, O. Norniella^x, L. Oakes^{ao}, S.H. Oh^p, Y.D. Oh^{aa}, I. Oksuzian^{bi}, T. Okusawa^{an}, R. Orava^w, L. Ortolan^d, S. Pagan Griso^{ap,aq}, C. Pagliarone^{be}, E. Palencia^{j,7}, V. Papadimitriou^q, A.A. Paramonov^b, J. Patrick^q, G. Pauletta^{be,bf}, M. Paulini^k, C. Paus^{af}, D.E. Pellett^h, A. Penzo^{be}, T.J. Phillips^p, G. Piacentino^{as}, E. Pianori^{ar}, J. Pilot^{al}, K. Pitts^x, C. Plagerⁱ, L. Pondrom^{bl}, S. Poprocki^{q,8}, K. Potamianos^{ax}, F. Prokoshin^{o,30}, A. Pranko^{ab}, F. Ptohos^{s,9}, G. Punzi^{as,at}, A. Rahaman^{aw}, V. Ramakrishnan^{bl}, N. Ranjan^{ax}, I. Redondo^{ae}, P. Renton^{ao}, M. Rescigno^{ba}, T. Riddick^{ad}, F. Rimondi^{f,g}, L. Ristori^{q,as}, A. Robson^u, T. Rodrigo^j, T. Rodriguez^{ar}, E. Rogers^x, S. Rolli^{bh,10}, R. Roser^q, F. Ruffini^{as,au}, A. Ruiz^j, J. Russ^k, V. Rusu^q, A. Safonov^{bd}, W.K. Sakumoto^{ay}, Y. Sakurai^{bj}, L. Santi^{be,bf}, K. Sato^{bg}, V. Saveliev^{q,24}, A. Savoy-Navarro^{q,28}, P. Schlabach^q, A. Schmidt^z, E.E. Schmidt^q, T. Schwarz^q, L. Scodellaro^j, A. Scribano^{as,au}, F. Scuri^{as}, S. Seidel^{ak}, Y. Seiya^{an}, A. Semenov^o, F. Sforza^{as,au}, S.Z. Shalhout^h, T. Shears^{ac}, P.F. Shepard^{aw}, M. Shimojima^{bg,23}, M. Shochet^l, I. Shreyber-Tecker^{aj}, A. Simonenko^o, P. Sinervo^{ag}, K. Sliwa^{bh}, J.R. Smith^h, F.D. Snider^q, A. Soha^q, V. Sorin^d, H. Song^{aw}, P. Squillacioti^{as,au}, M. Stancari^q, R. St. Denis^u, B. Stelzer^{ag}, O. Stelzer-Chilton^{ag}, D. Stentz^{q,25}, J. Strologas^{ak}, G.L. Strycker^{ah}, Y. Sudo^{bg}, A. Sukhanov^q, I. Suslov^o, K. Takemasa^{bg}, Y. Takeuchi^{bg}, J. Tang^l, M. Tecchio^{ah}, P.K. Teng^a, J. Thom^{q,8}, J. Thome^k, G.A. Thompson^x, E. Thomson^{ar}, D. Toback^{bd}, S. Tokar^{m,n}, K. Tollefson^{ai}, T. Tomura^{bg}, D. Tonelli^q, S. Torre^s, D. Torretta^q, P. Totaro^{ap}, M. Trovato^{as,av}, F. Ukegawa^{bg}, S. Uozumi^{aa}, A. Varganov^{ah}, F. Vázquez^{r,14}, G. Velev^q, C. Vellidis^{q,*}, M. Vidal^{ax}, I. Vila^j, R. Vilar^j, J. Vizán^j, M. Vogel^{ak}, G. Volpi^s, P. Wagner^{ar}, R.L. Wagner^q, T. Wakisaka^{an}, R. Wallnyⁱ, S.M. Wang^a, A. Warburton^{ag}, D. Waters^{ad}, W.C. Wester III^q, D. Whiteson^{ar,3}, A.B. Wicklund^b, E. Wicklund^q, S. Wilbur^l, F. Wick^z, H.H. Williams^{ar}, J.S. Wilson^{al}, P. Wilson^q, B.L. Winer^{al}, P. Wittich^{q,8}, S. Wolbers^q, H. Wolfe^{al}, T. Wright^{ah}, X. Wu^t, Z. Wu^e, K. Yamamoto^{an}, D. Yamato^{an}, T. Yang^q, U.K. Yang^{l,19}, Y.C. Yang^{aa}, W.-M. Yao^{ab}, G.P. Yeh^q, K. Yi^{q,15}, J. Yoh^q, K. Yorita^{bj}, T. Yoshida^{an,13}, G.B. Yu^p, I. Yu^{aa}, S.S. Yu^q, J.C. Yun^q, A. Zanetti^{be}, Y. Zeng^p, C. Zhou^p, S. Zucchelli^{f,g}

^a Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, ROC

^b Argonne National Laboratory, Argonne, IL 60439, USA

^c University of Athens, 157 71 Athens, Greece

^d Institut de Física d'Altes Energies, ICREA, Universitat Autònoma de Barcelona, E-08193, Bellaterra (Barcelona), Spain

^e Baylor University, Waco, TX 76798, USA

^f Istituto Nazionale di Fisica Nucleare Bologna, Italy

^g University of Bologna, I-40127 Bologna, Italy

^h University of California Davis, Davis, CA 95616, USA

ⁱ University of California Los Angeles, Los Angeles, CA 90024, USA

^j Instituto de Física de Cantabria, CSIC–University of Cantabria, 39005 Santander, Spain

^k Carnegie Mellon University, Pittsburgh, PA 15213, USA

^l Enrico Fermi Institute, University of Chicago, Chicago, IL 60637, USA

^m Comenius University, 842 48 Bratislava, Slovakia

ⁿ Institute of Experimental Physics, 040 01 Kosice, Slovakia

^o Joint Institute for Nuclear Research, RU-141980 Dubna, Russia

^p Duke University, Durham, NC 27708, USA

^q Fermi National Accelerator Laboratory, Batavia, IL 60510, USA

^r University of Florida, Gainesville, FL 32611, USA

^s Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy

^t University of Geneva, CH-1211 Geneva 4, Switzerland

^u Glasgow University, Glasgow G12 8QQ, United Kingdom

^v Harvard University, Cambridge, MA 02138, USA

^w Division of High Energy Physics, Department of Physics, University of Helsinki and Helsinki Institute of Physics, FIN-00014, Helsinki, Finland

^x University of Illinois, Urbana, IL 61801, USA

^y The Johns Hopkins University, Baltimore, MD 21218, USA

^z Institut für Experimentelle Kernphysik, Karlsruhe Institute of Technology, D-76131 Karlsruhe, Germany

^{aa} Center for High Energy Physics: Kyungpook National University, Daegu 702-701, Republic of Korea; Seoul National University, Seoul 151-742, Republic of Korea; Sungkyunkwan University, Suwon 440-746, Republic of Korea; Korea Institute of Science and Technology Information, Daejeon 305-806, Republic of Korea; Chonnam National University, Gwangju 500-757, Republic of Korea; Chonbuk National University, Jeonju 561-756, Republic of Korea

^{ab} Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA

^{ac} University of Liverpool, Liverpool L69 7ZE, United Kingdom

^{ad} University College London, London WC1E 6BT, United Kingdom

^{ae} Centro de Investigaciones Energéticas Medioambientales y Tecnológicas, E-28040 Madrid, Spain

^{af} Massachusetts Institute of Technology, Cambridge, MA 02139, USA

- ^{ag} Institute of Particle Physics: McGill University, Montréal, Québec, H3A 2T8, Canada; Simon Fraser University, Burnaby, British Columbia, V5A 1S6, Canada; University of Toronto, Toronto, Ontario, M5S 1A7, Canada; TRIUMF, Vancouver, British Columbia, V6T 2A3, Canada
- ^{ah} University of Michigan, Ann Arbor, MI 48109, USA
- ^{ai} Michigan State University, East Lansing, MI 48824, USA
- ^{aj} Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia
- ^{ak} University of New Mexico, Albuquerque, NM 87131, USA
- ^{al} The Ohio State University, Columbus, OH 43210, USA
- ^{am} Okayama University, Okayama 700-8530, Japan
- ^{an} Osaka City University, Osaka 588, Japan
- ^{ao} University of Oxford, Oxford OX1 3RH, United Kingdom
- ^{ap} Istituto Nazionale di Fisica Nucleare, Sezione di Padova–Trento, Italy
- ^{aq} University of Padova, I-35131 Padova, Italy
- ^{ar} University of Pennsylvania, Philadelphia, PA 19104, USA
- ^{as} Istituto Nazionale di Fisica Nucleare Pisa, Italy
- ^{at} University of Pisa, Italy
- ^{au} University of Siena, Italy
- ^{av} Scuola Normale Superiore, I-56127 Pisa, Italy
- ^{aw} University of Pittsburgh, Pittsburgh, PA 15260, USA
- ^{ax} Purdue University, West Lafayette, IN 47907, USA
- ^{ay} University of Rochester, Rochester, NY 14627, USA
- ^{az} The Rockefeller University, New York, NY 10065, USA
- ^{ba} Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1, Italy
- ^{bb} Sapienza Università di Roma, I-00185 Roma, Italy
- ^{bc} Rutgers University, Piscataway, NJ 08855, USA
- ^{bd} Texas A&M University, College Station, TX 77843, USA
- ^{be} Istituto Nazionale di Fisica Nucleare Trieste/Udine, I-34100 Trieste, Italy
- ^{bf} University of Udine, I-33100 Udine, Italy
- ^{bg} University of Tsukuba, Tsukuba, Ibaraki 305, Japan
- ^{bh} Tufts University, Medford, MA 02155, USA
- ^{bi} University of Virginia, Charlottesville, VA 22906, USA
- ^{bj} Waseda University, Tokyo 169, Japan
- ^{bk} Wayne State University, Detroit, MI 48201, USA
- ^{bl} University of Wisconsin, Madison, WI 53706, USA
- ^{bm} Yale University, New Haven, CT 06520, USA

ARTICLE INFO

Article history:

Received 27 July 2012

Received in revised form 25 August 2012

Accepted 27 August 2012

Available online 30 August 2012

Editor: M. Doser

ABSTRACT

A search for a narrow Higgs boson resonance in the diphoton mass spectrum is presented based on data corresponding to 10 fb^{-1} of integrated luminosity collected by the CDF experiment from proton–antiproton collisions at $\sqrt{s} = 1.96 \text{ TeV}$. To increase the sensitivity of the search, we employ a multivariate discriminant technique for the first time in this channel at CDF. No evidence of signal is observed, and upper limits are set on the cross section times branching ratio of the resonant state as a function of the Higgs boson mass. The limits are interpreted in the context of the standard model with an expected (observed) limit on the cross section times branching ratio of 9.9 (17.0) times the standard model prediction at the 95% credibility level for a Higgs boson mass of $125 \text{ GeV}/c^2$. Moreover, a Higgs boson with suppressed couplings to fermions is excluded for masses below $114 \text{ GeV}/c^2$ at the 95% credibility level.

© 2012 Elsevier B.V. All rights reserved.

* Corresponding author.

E-mail address: vellidis@fnal.gov (C. Vellidis).¹ Deceased.² Visitor from Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy.³ Visitor from University of CA Irvine, Irvine, CA 92697, USA.⁴ Visitor from University of CA Santa Barbara, Santa Barbara, CA 93106, USA.⁵ Visitor from University of CA Santa Cruz, Santa Cruz, CA 95064, USA.⁶ Visitor from Institute of Physics, Academy of Sciences of the Czech Republic, Czech Republic.⁷ Visitor from CERN, CH-1211 Geneva, Switzerland.⁸ Visitor from Cornell University, Ithaca, NY 14853, USA.⁹ Visitor from University of Cyprus, Nicosia CY-1678, Cyprus.¹⁰ Visitor from Office of Science, U.S. Department of Energy, Washington, DC 20585, USA.¹¹ Visitor from University College Dublin, Dublin 4, Ireland.¹² Visitor from ETH, 8092 Zurich, Switzerland.¹³ Visitor from University of Fukui, Fukui City, Fukui Prefecture, 910-0017, Japan.¹⁴ Visitor from Universidad Iberoamericana, Mexico D.F., Mexico.¹⁵ Visitor from University of Iowa, Iowa City, IA 52242, USA.¹⁶ Visitor from Kinki University, Higashi-Osaka City, 577-8502, Japan.¹⁷ Visitor from Kansas State University, Manhattan, KS 66506, USA.¹⁸ Visitor from Ewha Womans University, Seoul, 120-750, Republic of Korea.¹⁹ Visitor from University of Manchester, Manchester M13 9PL, United Kingdom.²⁰ Visitor from Queen Mary, University of London, London, E1 4NS, United Kingdom.²¹ Visitor from University of Melbourne, Victoria 3010, Australia.

1. Introduction

The standard model (SM) of particle physics has proven to be a robust theory that accurately describes the properties of elementary particles and the forces of interaction between them. However, the origin of mass has remained an unsolved mystery for decades. The SM suggests that particles acquire mass due to interactions with the Higgs field via spontaneous symmetry breaking [1]. In direct searches at the Large Electron–Positron Collider (LEP) [2] and recent search results from the Tevatron [3] and the Large Hadron Collider (LHC) [4], all potential SM Higgs boson masses outside the ranges 116.6–119.4 GeV/ c^2 and 122.1–127.0 GeV/ c^2 are excluded by at least one experiment.

In the SM, the branching ratio for a Higgs boson decaying into a photon pair $\mathcal{B}(H \rightarrow \gamma\gamma)$ is maximal for Higgs boson masses between about 110 and 140 GeV/ c^2 . This is a mass range that is most suitable for Higgs boson searches at the Fermilab Tevatron [3] and is favored by indirect constraints from electroweak observables [5]. The SM $H \rightarrow \gamma\gamma$ branching ratio peaks at a value of about 0.23% for a Higgs boson mass $m_H = 125$ GeV/ c^2 [6]. This is a very small branching ratio; however, the distinctive signal that photons produce in the detector makes $H \rightarrow \gamma\gamma$ an appealing search mode. Compared to the dominant decay modes involving b quarks, a larger fraction of $H \rightarrow \gamma\gamma$ events can be identified and the diphoton invariant mass of these events would cluster in a narrower range, thus providing a better discriminator against the smoothly distributed background. There are also theories beyond the standard model that predict a suppressed coupling of a Higgs boson to fermions. In these “fermiophobic” Higgs boson models, the diphoton decay can be greatly enhanced [7].

The Collider Detector at Fermilab (CDF) and D0 experiments at the Tevatron have searched for both a SM Higgs boson, H , and a fermiophobic Higgs boson, h_f , decaying to two photons [8]. The CDF and D0 experiments recently set 95% credibility level (C.L.) upper limits on the cross section times branching ratio $\sigma \times \mathcal{B}(H \rightarrow \gamma\gamma)$ relative to the SM prediction and on $\mathcal{B}(h_f \rightarrow \gamma\gamma)$ using data corresponding to integrated luminosities \mathcal{L} of 7.0 fb $^{-1}$ [9] and 8.2 fb $^{-1}$ [10], respectively. The h_f result sets a lower limit on m_{h_f} of 114 GeV/ c^2 and 112.9 GeV/ c^2 , respectively. These results surpassed for the first time the 109.7 GeV/ c^2 mass limit obtained from combined searches at the LEP collider at CERN [11].

The ATLAS and CMS experiments at the LHC at CERN have searched for a SM Higgs boson decaying to two photons using $\mathcal{L} = 4.9$ fb $^{-1}$ [12] and 4.8 fb $^{-1}$ [13], respectively. In the low mass range, rates corresponding to more than twice the SM cross section are excluded at the 95% C.L. An excess of 1.8σ is present in the CMS result and of 1.5σ in the ATLAS result, accounting for the look-elsewhere effect, which could be consistent with a SM Higgs boson with a mass near 125 GeV/ c^2 . Recent updates confirm these excesses at the 5σ level [14]. Searches for a fermiophobic Higgs boson in the ATLAS diphoton data exclude the ranges 110.0–118.0 GeV/ c^2 and 119.5–121.0 GeV/ c^2 at the 95% C.L. [15], and in the CMS diphoton data the ranges 110.0–124.5 GeV/ c^2 and 127.0–137.5 at the 95% C.L. [16].

In this Letter, we present a search for a Higgs boson decaying to two photons using the final CDF diphoton data set, corresponding to an integrated luminosity of 10 fb $^{-1}$. This analysis searches the diphoton mass distribution for a narrow resonance that could reveal the presence of a SM or fermiophobic Higgs boson, updating the previous CDF result [9] with more than 40% additional integrated luminosity. We furthermore implement a new multivariate technique for events that contain two central photons, using both diphoton and jet kinematic variables to improve the sensitivity for identifying a Higgs boson signal from the diphoton backgrounds.

2. Higgs boson signal model

For the SM search, we consider the three most likely production mechanisms at the Tevatron: gluon fusion (GF); associated production (VH), where a Higgs boson is produced in association with a W or Z boson; and vector boson fusion (VBF), where a Higgs boson is produced alongside two quark jets. As an example, the SM cross sections [3] for $m_H = 125$ GeV/ c^2 are 949.3 fb [17], 208.0 fb [18], and 65.3 fb [19] respectively. In the fermiophobic search, we consider a benchmark model in which a Higgs boson does not couple to fermions, yet retains its SM couplings to bosons [7]. In this model, the GF loop process is suppressed and fermiophobic Higgs boson production is dominated by VH and VBF. With $\mathcal{L} = 10$ fb $^{-1}$, about 28 (43) $H \rightarrow \gamma\gamma$ ($h_f \rightarrow \gamma\gamma$) events are predicted to be produced for $m_H = 125$ GeV/ c^2 .

The acceptance of these events in well-instrumented regions of the CDF detector times the efficiency of passing the full diphoton selection discussed in Section 3 [9] is only about 25%. This fraction, along with the predicted distributions of kinematic variables, is obtained from a simulation of Higgs boson decays into diphotons. For each Higgs boson mass hypothesis tested in the range 100–150 GeV/ c^2 , in 5 GeV/ c^2 steps, signal samples are developed from the PYTHIA 6.2 [20] Monte Carlo (MC) event generator and a parametrized response of the CDF II detector [21]. All PYTHIA samples were made with CTEQ5L [22] parton distribution functions, where the PYTHIA underlying event model is tuned to CDF jet data [23]. Each signal sample is corrected for multiple interactions and differences between the identification of photons in the simulation and the data [9,24]. The GF signal is furthermore corrected based on a higher-order theoretical prediction of the transverse momentum distribution [25].

3. Detector and event selection

We use the CDF II detector [26] to identify photon candidate events produced in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV. The silicon vertex tracker [27] and the central outer tracker [28], contained within a 1.4 T axial magnetic field, measure the trajectories of charged particles and determine their momenta. Particles that pass through the outer tracker reach the electromagnetic (EM) and hadronic calorimeters [29–31], which are divided into two regions: central ($|\eta| < 1.1$) and forward or “plug” ($1.1 < |\eta| < 3.6$). The EM calorimeters contain fine-grained shower maximum detectors [32], which measure the shower shape and centroid position in the plane transverse to the direction of the shower development.

The event selection is the same as in the previous $H \rightarrow \gamma\gamma$ search [9]. Events with two photon candidates are selected and the data are divided into four independent categories according to the position and type of the photons. In central–central (CC) events with non-converted photons, both photon candidates are detected within the fiducial region of the central EM calorimeter ($|\eta| < 1.05$); in central–plug (CP) events with non-converted pho-

²² Visitor from Muons, Inc., Batavia, IL 60510, USA.

²³ Visitor from Nagasaki Institute of Applied Science, Nagasaki, Japan.

²⁴ Visitor from National Research Nuclear University, Moscow, Russia.

²⁵ Visitor from Northwestern University, Evanston, IL 60208, USA.

²⁶ Visitor from University of Notre Dame, Notre Dame, IN 46556, USA.

²⁷ Visitor from Universidad de Oviedo, E-33007 Oviedo, Spain.

²⁸ Visitor from CNRS-IN2P3, Paris, F-75205 France.

²⁹ Visitor from Texas Tech University, Lubbock, TX 79609, USA.

³⁰ Visitor from Universidad Tecnica Federico Santa Maria, 110v Valparaiso, Chile.

³¹ Visitor from Yarmouk University, Irbid 211-63, Jordan.

tons, one photon candidate is detected in this region and the other is in the fiducial region of the plug calorimeter ($1.2 < |\eta| < 2.8$); in central–central events with a conversion (C’C), both photon candidates are in the central region, but one photon converts and is reconstructed from its e^+e^- decay products; in central–plug events with a conversion (C’P), there is one central conversion candidate together with a plug photon candidate.

For the diphoton resonance technique described in Section 4, the event selection in the fermiophobic Higgs boson search is extended by taking advantage of the final-state features present in the VH and VBF processes. Because the Higgs boson from these processes will be produced in association with a W or Z boson, or with two jets, the transverse momentum of the diphoton system $p_T^{\gamma\gamma}$ is generally higher relative to the diphoton backgrounds. A requirement of $p_T^{\gamma\gamma} > 75$ GeV/ c isolates a region of high h_f sensitivity, retaining roughly 30% of the signal while removing 99.5% of the background [8]. Two lower- $p_T^{\gamma\gamma}$ regions, $p_T^{\gamma\gamma} < 35$ GeV/ c and 35 GeV/ $c < p_T^{\gamma\gamma} < 75$ GeV/ c , are additionally included and provide about 15% more sensitivity to the h_f signal.

4. Diphoton resonance search

The decay of a Higgs boson into a diphoton pair would appear as a very narrow peak in the distribution of the invariant mass $m_{\gamma\gamma}$ of the two photons. The diphoton mass resolution as determined from simulation is better than 3% for the Higgs boson mass region studied here and is limited by the energy resolution of the electromagnetic calorimeters [33] and the ability to identify the primary interaction vertex [9,24]. The diphoton invariant mass distribution for the most sensitive search category in the SM and fermiophobic scenarios is provided in Fig. 1, with an inset showing the signal shape expected from simulation. In each diphoton category, we perform a search of the $m_{\gamma\gamma}$ spectrum for signs of a resonance.

For this search, the total diphoton background is modeled from a fit to the binned diphoton mass spectrum of the data using a log-likelihood method, as described in [9,24]. The fit is performed independently for each diphoton category and includes only the sideband region for each m_H hypothesis, which is the control region excluding a mass window centered on the Higgs boson mass being tested. The full width of the mass window is chosen to be approximately ± 2 standard deviations of the expected Higgs boson mass resolution, which amounts to 12 GeV/ c^2 , 16 GeV/ c^2 , and 20 GeV/ c^2 for mass hypotheses of 100–115 GeV/ c^2 , 120–135 GeV/ c^2 , and 140–150 GeV/ c^2 , respectively. Example fits for the CC category for $m_H = 125$ GeV/ c^2 are shown in Fig. 1.

5. Multivariate discriminator

The diphoton mass distribution is the most powerful variable for separating a Higgs boson signal from the diphoton backgrounds. However, other information is available that can be used to further distinguish this signal. We improve the most sensitive search category (CC) by replacing the diphoton mass shape as a final discriminator with a “Multi-Layer Perceptron” neural network (NN) [34] output distribution. The NN combines the information of several well-modeled kinematic variables into a single discriminator, optimized to separate signal and background events. Four diphoton kinematic variables are included: $m_{\gamma\gamma}$, $p_T^{\gamma\gamma}$, the difference between the azimuthal angles of the two photons, and the cosine of the photon emission angle relative to the colliding hadrons in the diphoton rest frame (the Collins–Soper angle) [35]. For events with jets, we also include four variables related to the jet activity, which are particularly useful for identifying VBF

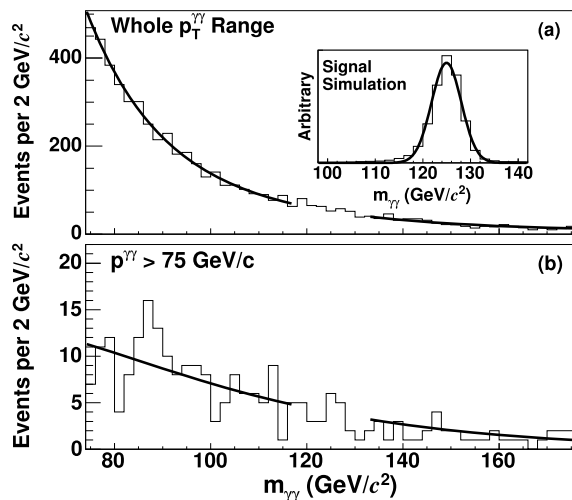


Fig. 1. The invariant mass distribution of CC photon pairs in the data is shown for (a) the entire $p_T^{\gamma\gamma}$ region used in the SM Higgs boson diphoton resonance search and (b) the highest- $p_T^{\gamma\gamma}$ region (the most sensitive region) used in the h_f diphoton resonance search. Each distribution shows a fit to the data for the hypothesis of $m_H = 125$ GeV/ c^2 , for which the signal region centered at 125 GeV/ c^2 is excluded from the fit. The expected shape of the signal from simulation is shown in the inset of (a). A Gaussian fit to the 125 GeV/ c^2 signal simulation yields a 1σ resolution of less than 4 GeV/ c^2 .

and VH signal events. These variables are the number of jets in the event, the sum of the jet transverse energies, and the event sphericity and aplanarity [36]. Jets are reconstructed from tower clusters in the hadronic calorimeter within a cone of radius 0.4 in the η – ϕ plane [37]. Each jet is required to have $|\eta| < 2$ and a transverse energy $E_T > 20$ GeV, where the energy is corrected for calorimeter response, multiple interactions, and absolute energy scale.

In order to optimize the performance of the method, we divide the CC category into two independent subsamples of events: the CC0 category for events with no jets and the CCJ category for events with at least one jet. The CC0 category uses a network trained with only the four diphoton variables; the CCJ category uses a network trained with the four diphoton and four jet variables.

The sideband fit used in the diphoton resonance search provides an estimate of the total background prediction in each signal mass window; however, the multivariate analysis requires a more detailed background model. Specifically, we divide the background into its distinct components in order to best model all input variables used by the discriminator, which is also sensitive to correlations. There are two main background components in the CC data sample: a prompt diphoton ($\gamma\gamma$) background produced from the hard parton scattering or from hard photon bremsstrahlung from energetic quarks, and a background comprised of γ –jet and jet–jet events ($\gamma j + jj$) in which the jets are misidentified as photons [38]. To model the shape of kinematic variables in the $\gamma\gamma$ background, we use a PYTHIA MC sample developed and studied in a measurement of the diphoton cross section [35]. To model the variable shapes in the $\gamma j + jj$ background, we obtain a data sample enriched in misidentified photons by selecting events for which one or both photon candidates fail the NN photon ID requirement [9].

In the diphoton cross section analysis [35] it was found that a $p_T^{\gamma\gamma}$ -dependent correction was needed for the PYTHIA modeling of prompt diphoton events. We adopt the correction for this analysis, reweighting the $p_T^{\gamma\gamma}$ distribution from PYTHIA to match the $p_T^{\gamma\gamma}$ distribution from control regions in prompt diphoton data. For each category, CC0 and CCJ, and for each Higgs boson mass hy-

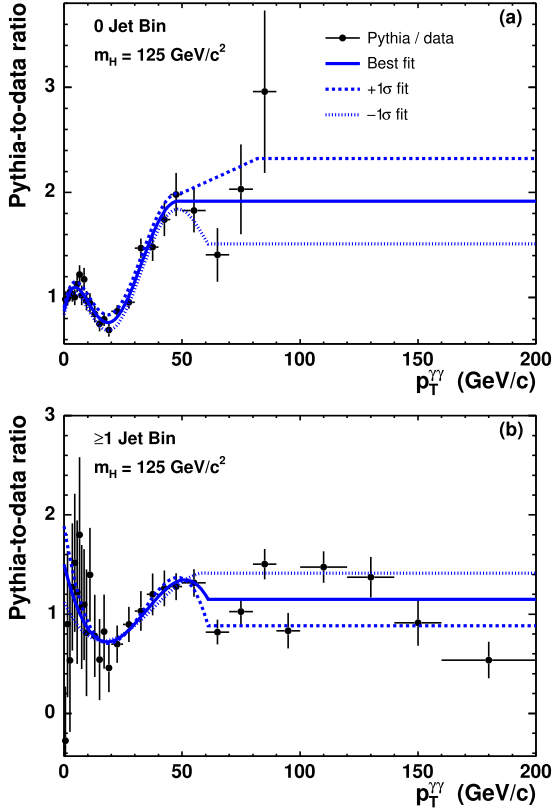


Fig. 2. For a Higgs boson mass of $125 \text{ GeV}/c^2$, the reweighting function obtained from the ratio of the $p_T^{\gamma\gamma}$ distribution in PYTHIA to the $p_T^{\gamma\gamma}$ distribution in prompt diphoton data, for events with (a) zero jets and (b) at least one jet. In both plots, the best fit to the PYTHIA-to-data ratio points is given by a solid curve. The other two curves show the systematic uncertainty of the fit.

pothesis, event weights are derived for the $\gamma\gamma$ background based on the sideband regions, excluding the signal mass window. The weights are derived by fitting a smooth function to the ratio of the $p_T^{\gamma\gamma}$ distribution from the data to that from the PYTHIA prediction. The best fit in the CC0 category is obtained from a polynomial (constant) function for $p_T^{\gamma\gamma} < 50 \text{ GeV}/c$ ($p_T^{\gamma\gamma} > 50 \text{ GeV}/c$). A different polynomial (constant) function provides the best fit in the CCJ category for $p_T^{\gamma\gamma} < 60 \text{ GeV}/c$ ($p_T^{\gamma\gamma} > 60 \text{ GeV}/c$). Fig. 2 shows the reweighting function for a Higgs boson mass hypothesis of $125 \text{ GeV}/c^2$. The solid curve shows the best fit to the data and the other two curves show the variations induced by propagating the 68% C.L. fit uncertainties to the fitting function. The rise of the reweighting function from $p_T^{\gamma\gamma} \sim 20 \text{ GeV}/c$ to $p_T^{\gamma\gamma} \sim 50 \text{ GeV}/c$ in both the CC0 and CCJ categories is interpreted in Ref. [35] as an effect of parton fragmentation not modeled in PYTHIA, which contributes to the prompt diphoton production cross section in that range.

The relative contributions of the two background components are obtained from a fit to the diphoton data. Three histograms for each NN input variable are constructed: one from the $\gamma\gamma$ background sample after reweighting, one from the $\gamma j + jj$ background sample, and one from the diphoton data. Events used for the fit are required to have diphoton mass values greater than $70 \text{ GeV}/c^2$ and to be outside of the signal mass window. The histograms are then used to build a χ^2 function defined by

$$\chi^2 = \sum_{i=1}^{N_{\text{bins}}} \sum_{j=1}^{N_{\text{variables}}} \left[\frac{(\alpha g_{ij} + \beta f_{ij} - d_{ij})^2}{d_{ij}} \right] \quad (1)$$

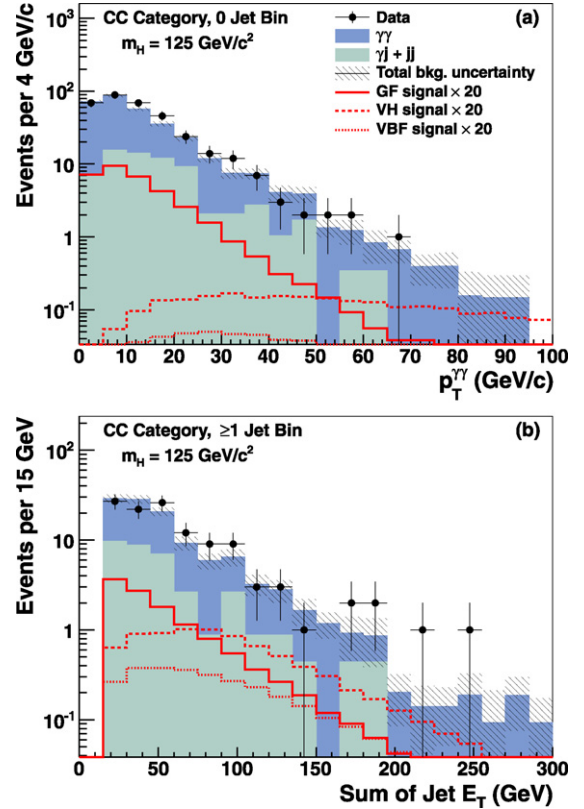


Fig. 3. For a Higgs boson mass of $125 \text{ GeV}/c^2$, a comparison of the data to the background prediction in (a) the $p_T^{\gamma\gamma}$ distribution for the CC0 category and (b) the distribution of the sum of the reconstructed jet E_T for the CCJ category. The expected SM Higgs boson signal for the three production processes is multiplied by a factor of 20.

where g_{ij} , f_{ij} , and d_{ij} refer to the number of events in the i th bin of the j th input variable for the prompt $\gamma\gamma$ background, $\gamma j + jj$ background, and diphoton data samples, respectively. The sums are over all bins of each input variable for which there are at least 5 events in the data, and the global α and β coefficients are determined by minimizing the χ^2 function. This function is defined and minimized separately for each Higgs boson mass hypothesis and for each category (CC0 and CCJ).

A neural network discriminant is trained separately for each mass hypothesis using signal and background events. The signal events used in the training are optimized for the SM scenario and are composed of GF, VH, and VBF MC samples so that the corresponding total numbers are proportional to their SM cross section predictions. The background sample is made by taking a portion of the $\gamma j + jj$ sample available for each mass hypothesis and adding $\gamma\gamma$ events from PYTHIA weighted by the ratio α/β from the χ^2 fit for the given mass hypothesis.

After training, the NN is applied to diphoton data events with mass inside the signal window. Fig. 3 shows input variables such as the $p_T^{\gamma\gamma}$ distribution for events with no reconstructed jets and the sum of the jet E_T for events with ≥ 1 reconstructed jet. The signal shapes are scaled to 20 times the expected number of reconstructed events in the SM scenario. The background prediction is also provided. While the χ^2 fit described by Eq. (1) is used to fix the relative composition of the $\gamma\gamma$ and $\gamma j + jj$ background components, the total expected number of background events is more accurately determined from the sideband mass fits described in Section 4. The resulting NN shapes for $m_H = 125 \text{ GeV}/c^2$ are provided in Fig. 4.

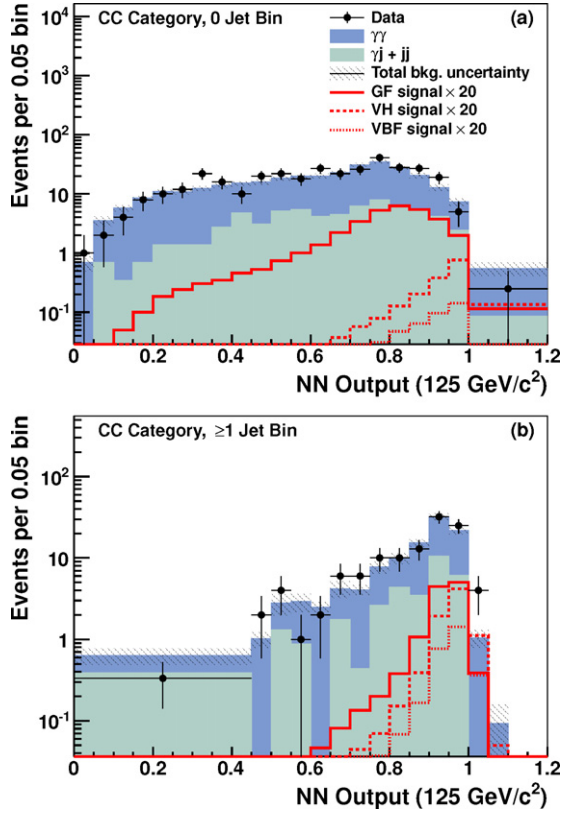


Fig. 4. For a Higgs boson mass of $125 \text{ GeV}/c^2$, a comparison of the NN output distributions for the data and the background prediction for (a) the CC0 category and (b) the CCJ category. The expected SM Higgs boson signal for the three production processes is multiplied by a factor of 20.

6. Systematic uncertainties

The sources of systematic uncertainties on the expected number of signal events are the same as in the previous $H \rightarrow \gamma\gamma$ search [9,24]. They arise from the conversion ID efficiency (7%), the integrated luminosity measurement (6%), varying the parton distribution functions used in PYTHIA (up to 5%) [39,40], varying the parameters that control the amount of initial- and final-state radiation from the parton shower model of PYTHIA (about 4%), and the PYTHIA modeling of the shape of the $p_T^{\gamma\gamma}$ distribution for the h_f signal (up to 4%) [41]. Finally, we include uncertainties from the photon ID efficiency (up to 4%), the trigger efficiency (less than 3%), and the EM energy scale (less than 1%). The signal rate uncertainties that arise from a common source are treated as correlated when combining results from each category.

The statistical uncertainties on the total background rate in the signal region are determined by the $m_{\gamma\gamma}$ fits. They are 4% or less for the channels associated with the SM diphoton resonance search and are less than 7% for the CC0 and CCJ categories used in the multivariate technique. For the channels associated with the fermiophobic Higgs boson diphoton resonance search, the background rate uncertainty is 12% or less, except for the high- $p_T^{\gamma\gamma}$ bins with conversion photons, where it is 20%. These background rate uncertainties are treated as uncorrelated between categories because the $m_{\gamma\gamma}$ fits are determined from exclusive data samples.

For the search using the multivariate technique, in addition to the rate uncertainties summarized above, we consider shape uncertainties and bin-by-bin statistical uncertainties of the NN discriminant. The signal shape uncertainties are associated with initial- and final-state radiation and the jet energy scale [37], and the background shape uncertainties are associated with the PYTHIA

$p_T^{\gamma\gamma}$ -correction and the jet energy scale. The PYTHIA shape uncertainties due to the $p_T^{\gamma\gamma}$ fits are taken as uncorrelated between the CC0 and the CCJ categories because the fits determining the corrections for each category are done independently. The jet energy scale shape uncertainties are correlated between the two categories in order to take into account event migration between categories. The dominant uncertainty in the multivariate analysis is the bin-by-bin statistical uncertainty of the $\gamma j + j j$ background histograms.

7. Results

No evidence of a narrow peak or any other structure is visible in the diphoton mass spectrum or the NN output distribution. We calculate a Bayesian C.L. limit for each Higgs boson mass hypothesis based on a combination of likelihoods from the discriminant distributions for all channels in the corresponding mass signal region. The combined limits for the SM search use the NN discriminants of the CC0 and CCJ categories and the mass discriminants from the CP, C'C, and C'P categories. The fermiophobic limits use the NN discriminants of the CC0 and CCJ categories and the mass discriminants from the CP, C'C, and C'P categories divided into $p_T^{\gamma\gamma}$ regions. For the limit calculation, we assume a flat prior (truncated at zero) for the signal rate and a truncated Gaussian prior for each of the systematic uncertainties. A 95% C.L. limit is determined such that 95% of the posterior density for $\sigma \times \mathcal{B}(H \rightarrow \gamma\gamma)$ falls below the limit [42]. The expected 95% C.L. limits are calculated assuming no signal, based on expected backgrounds only, as the median of 2000 simulated experiments. The observed 95% C.L. limits on $\sigma \times \mathcal{B}(H \rightarrow \gamma\gamma)$ are calculated from the data.

For the SM Higgs boson search, the results are given relative to the theory prediction, where theoretical cross section uncertainties of 14% on the inclusive GF process, 7% on the VH process, and 5% on the VBF process are included in the limit calculation [3, 43]. Since the NN technique divides the CC category into separate channels based on the number of reconstructed jets, different GF cross section uncertainties are assigned to the CC0 and CCJ channels [3,44]. For the h_f model, SM cross sections and uncertainties are assumed (GF excluded) and used to convert limits on $\sigma \times \mathcal{B}(h_f \rightarrow \gamma\gamma)$ into limits on $\mathcal{B}(h_f \rightarrow \gamma\gamma)$. The SM and fermiophobic limit results for the CC category alone are provided in Table 1, showing the gain obtained by incorporating a multivariate technique for this category. The combined limit results for both searches are displayed in Table 2 and graphically in Fig. 5. Limits are also provided on $\sigma \times \mathcal{B}(H \rightarrow \gamma\gamma)$ for the SM search without including theoretical cross section uncertainties. For the SM limit at $m_H = 120 \text{ GeV}/c^2$, we observe a deviation of greater than 2.5σ from the expectation. After accounting for the look-elsewhere effect associated with performing the search at 11 mass points, the significance of this discrepancy decreases to less than 2σ . When the analysis is optimized for the fermiophobic benchmark model, no excess is observed. For the h_f model, we obtain a limit of $m_{h_f} < 114 \text{ GeV}/c^2$ by linear interpolation between the sampled values of m_{h_f} based on the intersection of the observed limit and the model prediction.

8. Summary and conclusions

This Letter presents the results of a search for a narrow resonance in the diphoton mass spectrum using data taken by the CDF II detector at the Tevatron. We have improved upon the previous CDF analysis by implementing a neural network discriminant to increase sensitivity in the most sensitive diphoton category by as much as 13% (17%) for the SM (fermiophobic) scenario. In addition, we have included the full CDF diphoton data set, which adds more

Table 1

Expected and observed 95% C.L. upper limits on the production cross section multiplied by the $H \rightarrow \gamma\gamma$ branching ratio relative to the SM prediction for the most sensitive category (CC) using the NN discriminant. For comparison, values for the CC category are also provided based on the diphoton resonance technique, which uses the $m_{\gamma\gamma}$ shape as a discriminant for setting limits. The expected and observed 95% C.L. upper limits on the h_f branching ratio (in %) are provided in parentheses, based on both the NN discriminant and diphoton resonance technique for the CC category.

m_H (GeV/ c^2)	NN discriminant		$m_{\gamma\gamma}$ discriminant	
	Expected	Observed	Expected	Observed
100	13.9 (4.6)	10.6 (4.7)	15.1 (5.1)	11.3 (3.5)
105	12.6 (4.6)	13.0 (6.1)	14.1 (5.5)	10.6 (5.1)
110	11.9 (5.2)	11.8 (5.5)	13.5 (5.8)	11.4 (6.3)
115	11.4 (5.2)	14.1 (6.7)	12.9 (6.2)	15.4 (6.0)
120	11.3 (5.5)	23.2 (9.2)	12.8 (6.6)	22.2 (7.3)
125	11.7 (6.4)	20.5 (10.2)	12.9 (6.9)	21.2 (8.0)
130	12.5 (7.0)	13.1 (6.5)	13.9 (7.3)	16.0 (6.0)
135	13.7 (7.7)	15.0 (6.0)	15.3 (7.9)	17.2 (4.9)
140	16.5 (8.2)	20.4 (8.1)	17.5 (8.3)	25.4 (5.9)
145	18.5 (8.4)	27.4 (11.8)	21.2 (8.6)	24.3 (8.8)
150	25.7 (8.7)	17.1 (7.0)	28.2 (9.0)	15.1 (8.4)

Table 2

Expected and observed 95% C.L. upper limits on the production cross section times branching ratio relative to the SM prediction, the production cross section times branching ratio with theoretical cross section uncertainties removed, and the h_f branching ratio. The fermiophobic benchmark model prediction for $\mathcal{B}(h_f \rightarrow \gamma\gamma)$ is also shown for comparison.

	m_H (GeV/ c^2)	100	105	110	115	120	125	130	135	140	145	150
$\sigma \times \mathcal{B}(H \rightarrow \gamma\gamma)/SM$	Expected	12.2	10.9	10.6	9.7	9.7	9.9	10.5	11.6	14.0	16.0	21.3
	Observed	10.4	11.0	7.7	10.9	21.3	17.0	12.9	12.9	18.3	21.2	14.9
$\sigma \times \mathcal{B}(H \rightarrow \gamma\gamma)$ (fb)	Expected	45.1	39.0	37.2	31.8	29.7	27.2	25.5	24.0	23.0	20.4	20.2
	Observed	37.9	40.6	26.8	35.9	66.6	47.7	31.5	26.5	30.7	27.2	13.9
$\mathcal{B}(h_f \rightarrow \gamma\gamma)$ (%)	Expected	3.7	3.8	4.3	4.3	4.6	5.3	5.7	6.1	6.6	6.7	7.1
	Observed	4.9	5.1	3.5	4.8	5.9	4.9	5.3	7.9	8.4	8.3	5.0
	Fermiophobic prediction	18.5	10.4	6.0	3.7	2.3	1.6	1.1	0.8	0.5	0.4	0.3

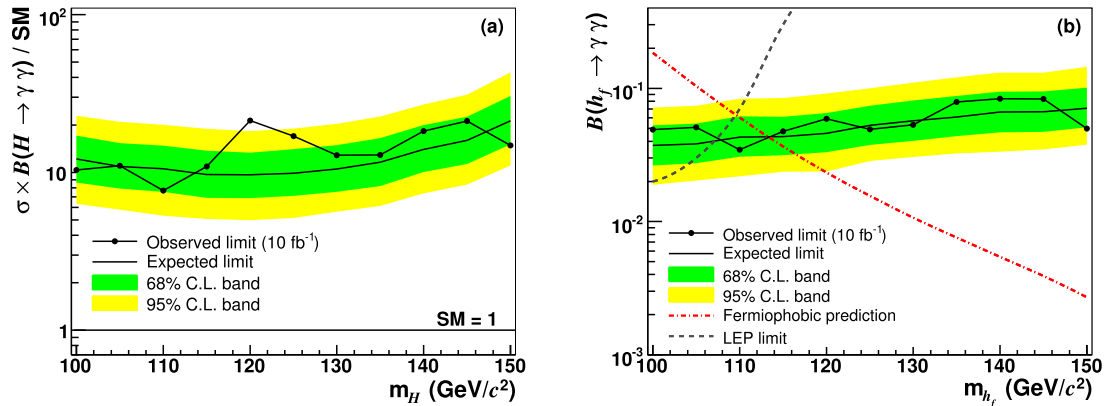


Fig. 5. (a) As a function of m_H , the 95% C.L. upper limit on the cross section times branching ratio for the SM Higgs boson decay to two photons, relative to the SM prediction. (b) The 95% C.L. upper limit on the branching ratio for the fermiophobic Higgs boson decay to two photons as a function of m_{h_f} . For reference, the 95% C.L. limits from LEP are also included. The shaded regions represent the 68% C.L. and 95% C.L. bands for the observed limit with respect to the expected limit based on the distribution of simulated experimental outcomes.

than 40% additional integrated luminosity relative to the previous diphoton Higgs boson search. There is no significant evidence of a resonance in the data. Limits are placed on the production cross section times branching ratio for Higgs boson decay into a photon pair and compared to the predictions of the standard model and a benchmark fermiophobic model. The latter results in a limit on the fermiophobic Higgs boson mass of $m_{h_f} < 114$ GeV/ c^2 at the 95% C.L.

Acknowledgements

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was

supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium für Bildung und Forschung, Germany; the Korean World Class University Program, the National Research Foundation of Korea; the Science and Technology Facilities Council and the Royal Society, UK; the Russian Foundation for Basic Research; the Ministerio de Ciencia e Innovación, and Programa Consolider-Ingenio 2010, Spain; the Slovak R&D Agency; the Academy of Finland; and the Australian Research Council (ARC).

References

- [1] P.W. Higgs, Phys. Rev. Lett. 13 (1964) 508;
G.S. Guralnik, C.R. Hagen, T.W.B. Kibble, Phys. Rev. Lett. 13 (1964) 585;
F. Englert, R. Brout, Phys. Rev. Lett. 13 (1964) 321.
- [2] LEP Working Group for Higgs Boson Searches, LEP Collaborations, Phys. Lett. B 565 (2003) 61.
- [3] Tevatron New Phenomena and Higgs Working Group, CDF, D0 Collaborations, arXiv:1207.0449.
- [4] G. Aad, et al., ATLAS Collaboration, Phys. Rev. D 86 (2012) 032003;
S. Chatrchyan, et al., CMS Collaboration, Phys. Lett. B 710 (2012) 26.
- [5] LEP, Tevatron, SLD Electroweak Working Groups, arXiv:1012.2367.
- [6] A. Djouadi, J. Kalinowski, M. Spira, Comput. Phys. Commun. 108 (1998) 56.
- [7] A.G. Akeroyd, Phys. Lett. B 368 (1996) 89;
A. Barroso, L. Brucher, R. Santos, Phys. Rev. D 60 (1999) 035005;
S. Mrenna, J. Wells, Phys. Rev. D 63 (2000) 015006;
H.E. Haber, G.L. Kane, T. Sterling, Nucl. Phys. B 161 (1979) 493;
J.F. Gunion, R. Vega, J. Wudka, Phys. Rev. D 42 (1990) 1673;
V. Barger, N.G. Deshpande, J.L. Hewett, T.G. Rizzo, arXiv:hep-ph/9211234;
J.-L. Basdevant, E.L. Berger, D. Dicus, C. Kao, S. Willenbrock, Phys. Lett. B 313 (1993) 402;
A. Stange, W. Marciano, S. Willenbrock, Phys. Rev. D 49 (1994) 1354;
M.A. Díaz, T.J. Weiler, arXiv:hep-ph/9401259;
L. Brucher, R. Santos, Eur. Phys. J. C 12 (2000) 87;
G. Landsberg, K.T. Matchev, Phys. Rev. D 62 (2000) 035004.
- [8] T. Affolder, et al., CDF Collaboration, Phys. Rev. D 64 (2001) 092002;
B. Abbott, et al., D0 Collaboration, Phys. Rev. Lett. 82 (1999) 2244;
V.M. Abazov, et al., D0 Collaboration, Phys. Rev. Lett. 101 (2008) 051801;
T. Aaltonen, et al., CDF Collaboration, Phys. Rev. Lett. 103 (2009) 061803.
- [9] T. Aaltonen, et al., CDF Collaboration, Phys. Rev. Lett. 108 (2012) 011801.
- [10] V.M. Abazov, et al., D0 Collaboration, Phys. Rev. Lett. 107 (2011) 151801.
- [11] A. Rosca, LEP Collaborations, Nucl. Phys. B Proc. Suppl. 117 (2003) 743.
- [12] G. Aad, et al., ATLAS Collaboration, Phys. Rev. Lett. 108 (2012) 111803.
- [13] D. Chatrchyan, et al., CMS Collaboration, Phys. Lett. B 710 (2012) 403.
- [14] G. Aad, et al., ATLAS Collaboration, Phys. Lett. B 716 (2012) 1;
S. Chatrchyan, et al., CMS Collaboration, Phys. Lett. B 716 (2012) 30.
- [15] G. Aad, et al., ATLAS Collaboration, arXiv:1205.0701, Eur. Phys. J. C, submitted for publication.
- [16] S. Chatrchyan, et al., CMS Collaboration, arXiv:1207.1130, J. High Energy Phys., submitted for publication.
- [17] C. Anastasiou, R. Boughezal, F. Petriello, J. High Energy Phys. 0904 (2009) 003;
D. de Florian, M. Grazzini, Phys. Lett. B 674 (2009) 291.
- [18] J. Baglio, A. Djouadi, J. High Energy Phys. 1010 (2010) 064;
O. Brein, A. Djouadi, R. Harlander, Phys. Lett. B 579 (2004) 149;
M.L. Ciccolini, S. Dittmaier, M. Krämer, Phys. Rev. D 68 (2003) 073003.
- [19] P. Bolzoni, F. Maltoni, S.-O. Moch, M. Zaro, Phys. Rev. Lett. 105 (2010) 011801;
M. Ciccolini, A. Denner, S. Dittmaier, Phys. Rev. Lett. 99 (2007) 161803;
M. Ciccolini, A. Denner, S. Dittmaier, Phys. Rev. D 77 (2008) 013002.
- [20] T. Sjöstrand, P. Edén, C. Friberg, L. Lönnblad, G. Miu, S. Mrenna, E. Norrbin, Comput. Phys. Commun. 135 (2001) 238.
- [21] R. Brun, et al., CERN Report No. CERN-DD-EE-84-01, 1987;
G. Grindhammer, M. Rudowicz, S. Peters, Nucl. Instrum. Methods Phys. Res. Sect. A 290 (1990) 469.
- [22] H.L. Lai, et al., CTEQ Collaboration, Eur. Phys. J. C 12 (2000) 375.
- [23] T. Aaltonen, et al., CDF Collaboration, Phys. Rev. D 82 (2010) 034001.
- [24] K.R. Bland, Ph.D. thesis, Baylor University, Fermilab Report No. FERMILAB-THESIS-2012-11, 2012.
- [25] G. Bozzi, S. Catani, D. de Florian, M. Grazzini, Phys. Lett. B 564 (2003) 65;
G. Bozzi, S. Catani, D. de Florian, M. Grazzini, Nucl. Phys. B 737 (2006) 73;
D. de Florian, G. Ferrera, M. Grazzini, D. Tommasini, J. High Energy Phys. 2011 (2011) 064.
- [26] D. Acosta, et al., CDF Collaboration, Phys. Rev. D 71 (2005) 032001.
- [27] A. Sill, CDF Collaboration, Nucl. Instrum. Methods Phys. Res. Sect. A 447 (2000) 1.
- [28] T. Affolder, et al., Nucl. Instrum. Methods Phys. Res. Sect. A 526 (2004) 249.
- [29] L. Balka, et al., Nucl. Instrum. Methods Phys. Res. Sect. A 267 (1988) 272.
- [30] S. Bertolucci, et al., Nucl. Instrum. Methods Phys. Res. Sect. A 267 (1988) 301.
- [31] M. Albrow, et al., Nucl. Instrum. Methods Phys. Res. Sect. A 480 (2002) 524.
- [32] G. Apollinari, K. Goulianos, P. Melese, M. Lindgren, Nucl. Instrum. Methods Phys. Res. Sect. A 412 (1998) 515.
- [33] The energy resolution for photons is $13.5\%/\sqrt{E_T}$ in the central region and $(16\%/\sqrt{E_T}) \oplus 1\%$ in the plug region.
- [34] A. Hoecker, et al., arXiv:physics/0703039.
- [35] T. Aaltonen, et al., CDF Collaboration, Phys. Rev. D 84 (2011) 052006.
- [36] D. Acosta, et al., CDF Collaboration, Phys. Rev. D 72 (2005) 052003.
- [37] A. Bhatti, et al., Nucl. Instrum. Methods Phys. Res. Sect. A 566 (2006) 375.
- [38] Typically, this occurs when a jet fragments into a π^0 or η particle that subsequently decays to multiple photons, which are then reconstructed as a single photon.
- [39] D. Stump, et al., J. High Energy Phys. 0310 (2003) 046.
- [40] D. Bourilkov, R.C. Group, M.R. Whalley, arXiv:hep-ph/0605240.
- [41] S. Mrenna, C.-P. Yuan, Phys. Lett. B 416 (1998) 200.
- [42] J. Beringer, et al., Particle Data Group, Phys. Rev. D 86 (2012) 010001.
- [43] M. Botje, et al., arXiv:1101.0538;
S. Dittmaier, et al., LHC Higgs Cross Section Working Group, arXiv:1101.0593.
- [44] I.W. Stewart, F.J. Tackmann, Phys. Rev. D 85 (2012) 034011;
S. Dittmaier, et al., LHC Higgs Cross Section Working Group, arXiv:1201.3084;
C. Anastasiou, G. Dissertori, M. Grazzini, F. Stöckli, B.R. Webber, J. High Energy Phys. 0908 (2009) 099.