

High Frequency Guided Wave Propagation and Scattering in Silicon Wafers

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Thin monocrystalline silicon wafers are employed for the manufacture of solar cells with high conversion efficiency. Micro-cracks can be induced by the wafer cutting process, leading to breakage of the fragile wafers. High frequency guided waves allow for the monitoring of wafers and detection and characterization of surface defects. The material anisotropy of the monocrystalline silicon leads to variations of the guided wave characteristics, depending on the guided wave mode and propagation direction relative to the crystal orientation. Selective excitation of the first anti-symmetric A_0 wave mode at 5 MHz center frequency was achieved experimentally using a custom-made wedge transducer. Strong wave pulses with limited beam skewing and widening were measured using non-contact laser interferometer measurements. This allowed the accurate characterization of the Lamb wave propagation and scattering at small artificial surface defects with a size of less than 100 μm . The surface extent of the defects of varying size was characterized using an optical microscope. The scattered guided wave field was evaluated, and characteristic parameters extracted and correlated to the defect size, allowing in principle detection of small defects. Further investigations are required to explain the systematic asymmetry of the guided wave field in the vicinity of the indents.

Keywords: Monocrystalline Silicon, Lamb Waves, Anisotropy, Ultrasonics, Surface Cracks

25 **1. INTRODUCTION**

26 Monocrystalline silicon wafers are used for the manufacture of solar photovoltaic panels with
27 high conversion efficiency [1]. Thin wafers are beneficial to increase conversion efficiency and
28 reduce manufacturing costs. However, in practice the minimum thickness of the brittle wafers
29 is limited by wafer breakage rates due to surface micro-cracks induced by the wafer cutting
30 process [2]. A range of nondestructive testing (NDT) techniques have been developed and
31 proposed for the detection of micro-cracks in silicon wafers [3]. These include
32 electroluminescence, thermography, optical transmission and interferometry imaging, impact
33 testing, and ultrasonic wave propagation [4]. Resonance ultrasonic vibration techniques
34 monitoring frequency shift and bandwidth of longitudinal vibration modes were proposed as
35 indicators of millimeter sized cracks in silicon wafers [5].

36 Long propagation distances relative to the thickness of thin plate-like structures can be
37 achieved employing guided ultrasonic waves [6]. For in-process monitoring this could provide
38 the required full area coverage of silicon wafers during manufacture. The sensitivity of high
39 frequency guided waves for crack detection in metallic plate structures has been demonstrated
40 experimentally [7, 8], and mode conversion and nonlinear effects due to rough crack surfaces
41 considered [9]. Guided wave scattering at horizontal cracks (disbonds) has been studied using
42 modal expansion and Finite Element Analysis (FEA) [10]. The effects of anisotropy on guided
43 wave propagation and defect detection have been widely investigated for composite structures,
44 including the energy concentration along the fiber direction [11], modal focusing [12], and
45 scattering at defects [13]. For an anisotropic multilayered structure, it has been demonstrated
46 that a guided wave beam can deviate with respect to the sagittal plane of excitation towards the
47 stiffer direction of the anisotropic structure [14]. For anisotropic silicon wafers the ultrasonic
48 and guided wave propagation are direction dependent, but only limited experimental
49 measurements have been reported. The variation of the arrival time and amplitude of the S_0

50 and SH modes with propagation direction was measured for silicon wafers and composite
51 plates [15]. The focusing effect of silicon material anisotropy on the ultrasonic wave energy
52 was predicted theoretically and measured experimentally [16, 17]. From an inversion of
53 experimental data, the silicon material properties could be obtained [18]. The variation of zero
54 group velocity, cut-off frequency, and amplitude with crystallographic orientation were
55 measured using a line laser source [19]. The slowness surface of leaky Lamb waves for
56 immersed silicon wafers was measured [20]. For monocrystalline silicon wafers, effects of the
57 anisotropy on the propagation of the fundamental guided modes (A_0 and S_0) were measured
58 and compared to Finite Element (FE) simulations and theoretical predictions [21]. Due to the
59 anisotropy, significant guided wave skewing in the non-principal crystallographic directions
60 was observed [22]. Especially for the S_0 mode, beam widening in directions away from the
61 stiffest crystal orientation ($\langle 110 \rangle$) and significant variation of the phase slowness occurs [21].
62 Laser interferometry measurements of the fundamental guided wave modes allowed crack
63 detection in silicon wafers [23]. Using a B-scan configuration, the guided wave amplitude drop
64 due to cracks in monocrystalline and polycrystalline silicon wafers was measured using air-
65 coupled transducers [24]. For a circular measurement arrangement with air-coupled
66 transducers, the amplitude profile of the A_0 mode at 200 kHz was recorded for different
67 propagation directions [25]. This allowed the detection of a 20 mm long through-thickness
68 notch in a wafer.

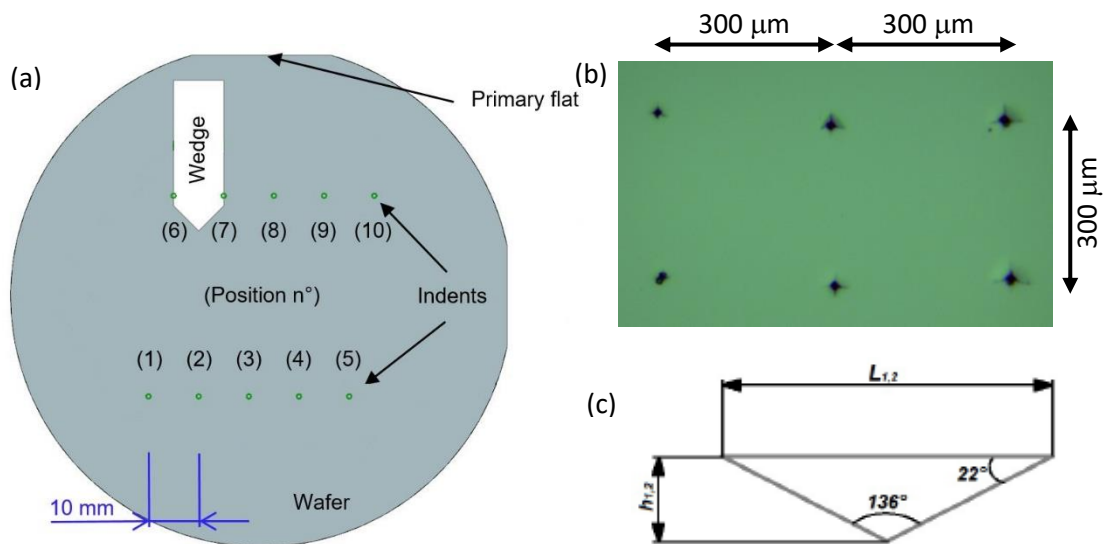
69 In this contribution, the scattering of the A_0 guided wave mode at 5 MHz center frequency was
70 measured using a non-contact laser interferometer in the near field around small surface defects
71 in thin monocrystalline silicon wafers. Using an indenter with different force levels, micro-
72 defects of increasing severity were created and characterized using an optical microscope [26].
73 Section 2 provides details of the experimental measurements, including silicon wafer
74 specimens and selective guided wave mode excitation and measurement. Section 3 provides a

75 brief overview of guided wave propagation in the anisotropic, monocrystalline silicon wafers.
 76 Section 4 shows and compares the experimental results for the scattered wave field variation
 77 at defects of increasing severity, while section 5 concludes the paper.

78 **2. EXPERIMENTS**

79 **2.1 Monocrystalline silicon specimens with artificial surface defects**

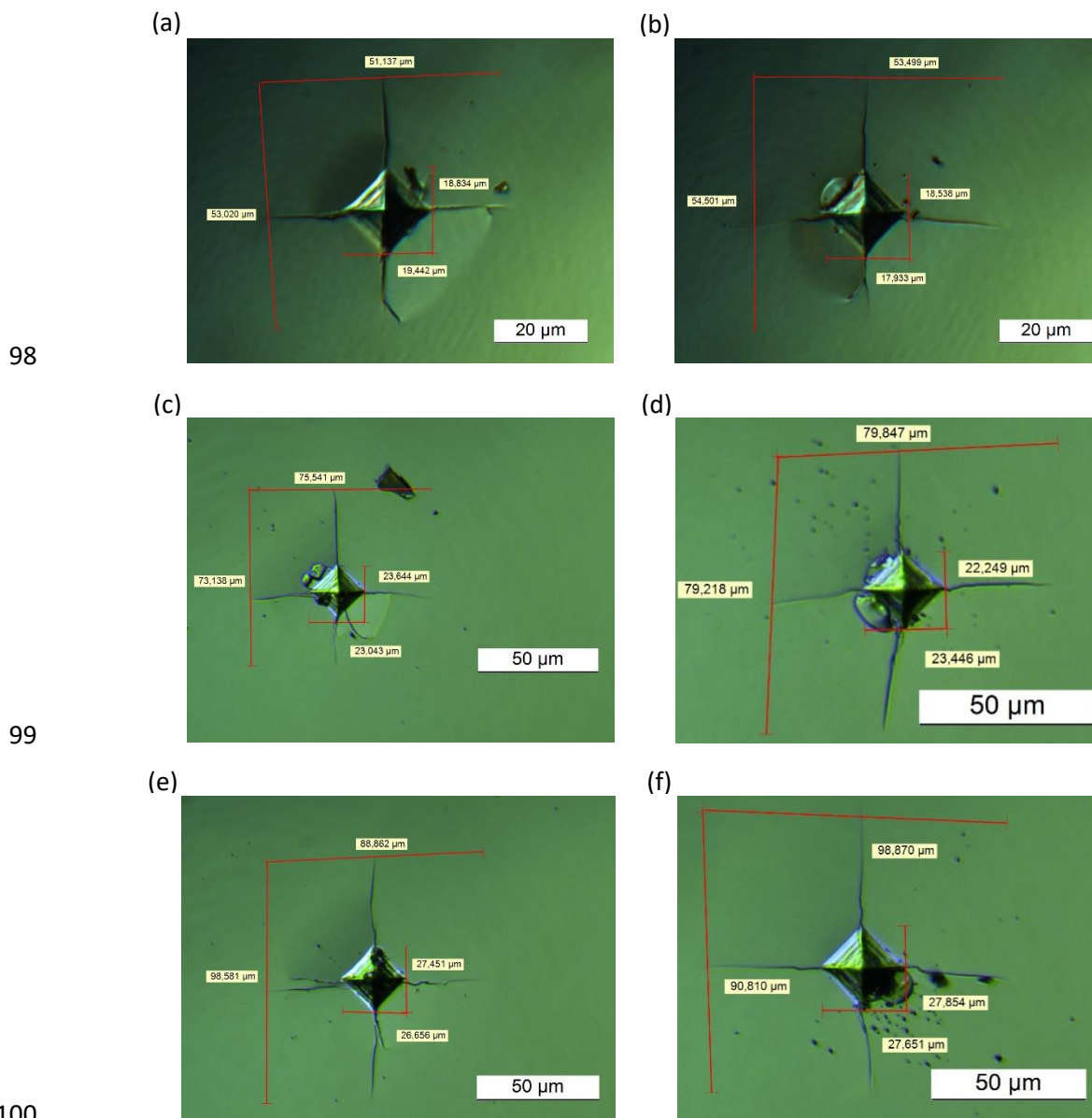
80 Polished, boron doped P-type monocrystalline silicon wafers (001-cut) with a nominal
 81 thickness of $380\ \mu\text{m}$ and $100\ \text{mm}$ (4'') diameter were used in this study. An undamaged wafer
 82 was used to measure the guided wave excitation and propagation characteristics. In a second
 83 wafer, a cluster of 6 indents was made using a Vickers indenter at controlled speed and
 84 specified force. The indents were located on a predefined grid with a spacing of approximately
 85 $300\ \mu\text{m}$ in both directions, with 2 indents per force level of 1 N, 1.5 N, and 2 N, respectively
 86 (Fig. 1b). Three additional silicon wafers were used for the investigation of the wave scattering
 87 at individual indents.



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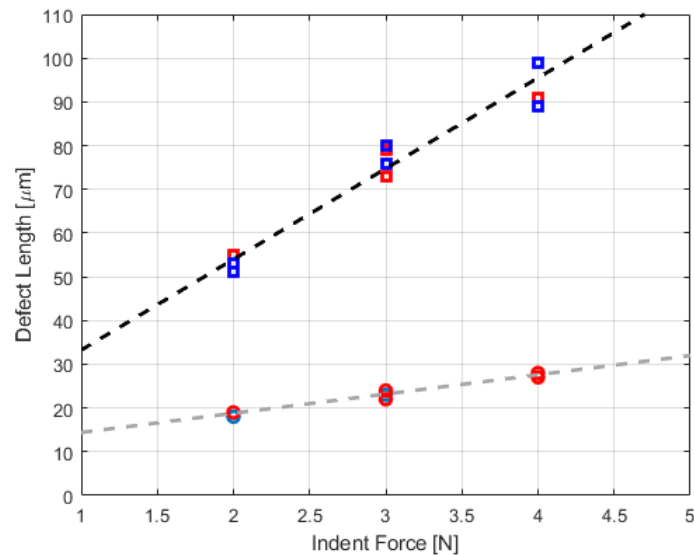
89 **Figure 1:** (a) Schematic of individual indent location and wedge position on wafer, 10 indents
 90 per wafer; (b) microscopy image of clustered pattern of 6 indents; (c) schematic of indenter
 91 geometry.

92 Each wafer contains 10 indents generated with the same force level (2 N, 3 N, and 4 N). The
93 geometrical arrangement of the indents on a wafer is shown in Fig. 1a, with the indents placed
94 in two lines with 10 mm separation between adjacent indents. All defects were made with a
95 square (four-sided) pyramid indenter tip with an angle of 136° (Fig. 1c). The tip was oriented
96 such that the four corners were along the highest stiffness directions $\langle 110 \rangle$, which repeats
97 every 90° for monocrystalline silicon.



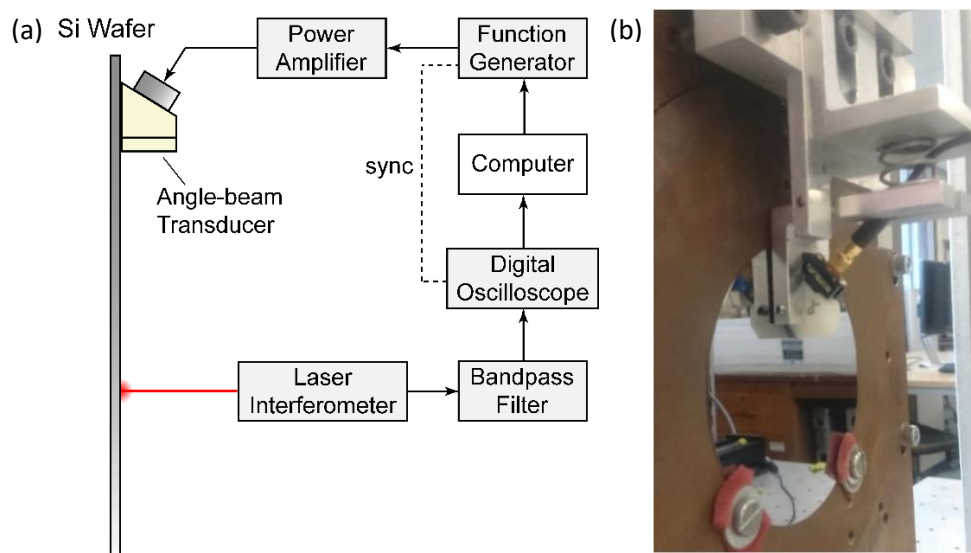
101 **Figure 2:** Microscopy photographs of individual defects with optical measurement of indent
102 and overall defect size; (a), (b): 2N indent force; (c), (d): 3N indent force; (e), (f): 4N indent
103 force.

104 From each of the corners a surface crack was generated along the highest stiffness directions
 105 $\langle 110 \rangle$ as shown in Fig. 2. Especially for the higher forces, chipping of the wafers adjacent to
 106 the indent was observed, and these defects were not investigated further. For each indenter
 107 force, at least two indents without visual indication of chipping at the surface were available.
 108 The surface extent of the indents and cracks was measured using an optical microscope, but no
 109 accurate depth information about the defects was available. The surface defect size showed a
 110 linear correlation with the indent force, with a maximum variation of the overall defect size of
 111 up to 10 μm for indents generated using the same force level (Fig. 3). From the measurement
 112 of the indent dimensions and the known angle of the indenter tip, the indent depth was
 113 estimated as 4 μm , 5 μm , and 6 μm for indent forces of 2 N, 3 N, and 4 N, respectively.



114

115 **Figure 3:** Optically measured defect size; square pyramid indent base (circles) and overall
 116 length including surface cracks (squares) against indent force (horizontal - blue; vertical - red);
 117 linear fit (dashed lines).



118

119 **Figure 4:** (a) Schematic diagram of the experimental setup; (b) photograph of monocrystalline
 120 silicon wafer on holder with spring-mounted wedge transducer.

121 2.2. Guided wave measurement

122 For the guided wave measurements, the silicon wafers were held in a custom-made wafer
 123 holder with a defined orientation against the flat (Fig. 4). This helped to reduce the risk of wafer
 124 breakage and allowed accurate angular orientation of the wafer with an accuracy of 1° . The A_0
 125 guided wave mode (first anti-symmetric Lamb wave mode) was excited selectively at a center
 126 frequency of 5 MHz. A commercial piezoelectric transducer (width: 8 mm) was placed on a
 127 custom-made nylon wedge with an angle of 41° to match the wavelength for propagation along
 128 the $\langle 110 \rangle$ direction according to Snell's law and nominal phase velocity values. Using springs,
 129 the wedge was pressed against the silicon wafer with a controlled force of approximately 10 N
 130 [27]. Standard ultrasonic couplant was used on the contact surface and excess couplant
 131 removed before measurements. The excitation signal was defined in Labview as a narrowband
 132 Hanning-windowed sinusoidal pulse (12 cycles) with a center frequency of 5 MHz. The signal
 133 was generated using an arbitrary function generator and amplified to $400 V_{pp}$ using a power
 134 amplifier.

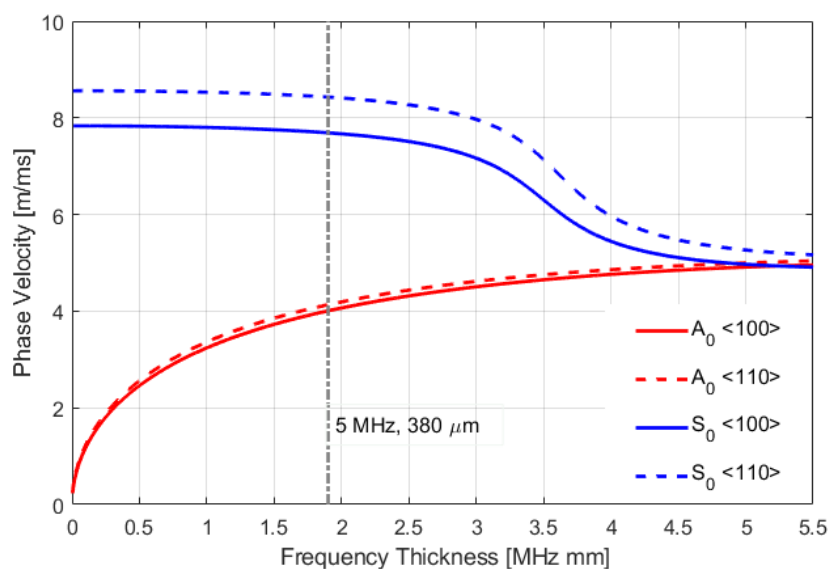
135 The out-of-plane surface displacement was measured using a non-contact, commercial laser
136 interferometer. Very good laser beam reflection and signal to noise ratio (SNR) was obtained
137 due to the mirror-like, polished silicon wafer surface, but care had to be taken to position the
138 wafer exactly perpendicular to the laser beam. The laser head was moved parallel to the silicon
139 wafer and holder using a positioning system with an accuracy and repeatability better than 1
140 μm . The measured out-of-plane displacement signal (output of laser demodulation) was
141 bandpass filtered (2-8 MHz, 4th order Butterworth filter) and recorded with a sampling
142 frequency of 100 MHz using a digital storage oscilloscope. 40 averages at each measurement
143 point were taken and the signal was transferred to a PC for evaluation in Matlab. The recorded
144 time trace at each point was time-gated and the amplitude and phase values were extracted at
145 the center frequency of 5 MHz using Fast Fourier Transform (FFT).

146 Initial measurements were conducted on the undamaged silicon wafer to characterize the
147 guided wave field excited by the wedge transducer [18]. The chosen measurement grid had 9
148 steps of 5 mm in the wave propagation direction and 21 steps of 1 mm perpendicular to the
149 wave propagation direction. Measurements were conducted for the wedge transducer aligned
150 in both the $\langle 110 \rangle$ and $\langle 100 \rangle$ directions. For comparison, the field for the S_0 guided wave mode
151 excited using the same setup, but with a custom-made wedge with an angle of 19° according
152 to Snell's law are shown [21]. For the silicon wafer with the clustered defect, the guided wave
153 field over an area of $1000 \mu\text{m}$ by $600 \mu\text{m}$, containing the 6 indents, was measured with a step
154 size of $20 \mu\text{m}$ in horizontal and vertical directions. For the 3 silicon wafers containing
155 individual defects, the scattered guided wave field was measured around two of the defects
156 with the same indenter force and with no evidence of chipping. In order to accurately capture
157 local amplitude variation close to each indent, an area of $400 \mu\text{m}$ by $400 \mu\text{m}$ was scanned with
158 a step size of $5 \mu\text{m}$ in the horizontal and vertical direction. For the guided wave measurements,
159 the exact defect location relative to the laser beam positioning could not be verified

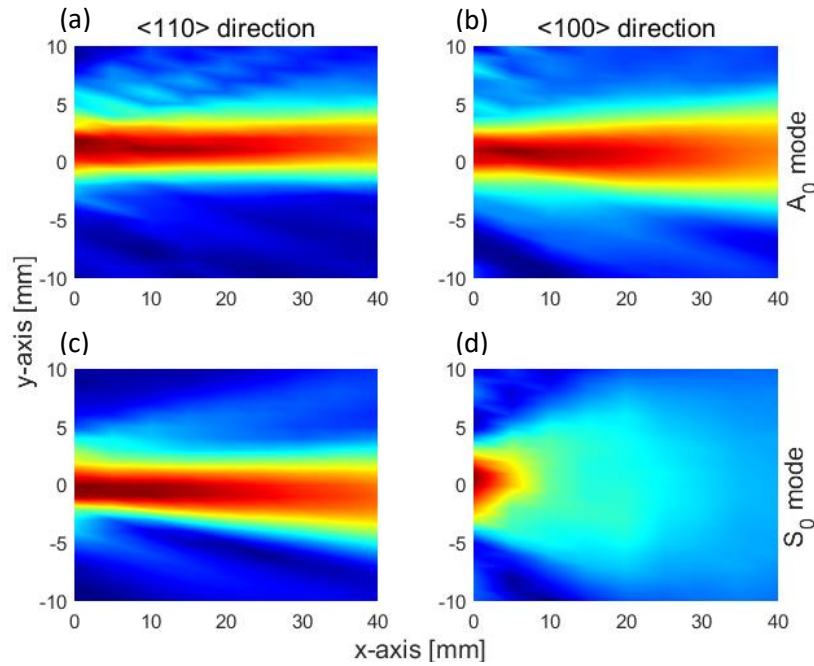
160 independently, and several steps were required to center the small measurement area
 161 approximately around the defect location [27, 28].

162 3. GUIDED WAVE PROPAGATION IN SILICON WAFERS

163 For the nominal silicon material properties (cubic symmetric stiffness constants: $C_{11} = 165.70$
 164 GPa, $C_{12} = 63.90$ GPa, $C_{44} = 79.56$ GPa, density $\rho = 2390$ kg/m³) the guided wave propagation
 165 characteristics were predicted using the Disperse software [29]. The phase velocity of the
 166 fundamental S_0 mode shows a variation of approximately 10% between the propagation in the
 167 direction of highest stiffness $\langle 110 \rangle$ and lowest stiffness $\langle 100 \rangle$. For the flexural A_0 mode a
 168 lower variation of approximately 3% was predicted. Good agreement with experimental
 169 measurements of the phase slowness was found in a previous publication [21]. The chosen
 170 excitation frequency of 5 MHz for a wafer thickness of 380 μm is marked in the dispersion
 171 diagram (Fig. 5). Good separation between the phase velocities of the fundamental A_0 and S_0
 172 modes can be observed and both modes show reasonably low dispersion for the chosen
 173 frequency thickness product of approximately 2 MHz mm.



174
 175 **Figure 5:** Dispersion diagram for silicon wafer, product of nominal thickness and excitation
 176 frequency marked.



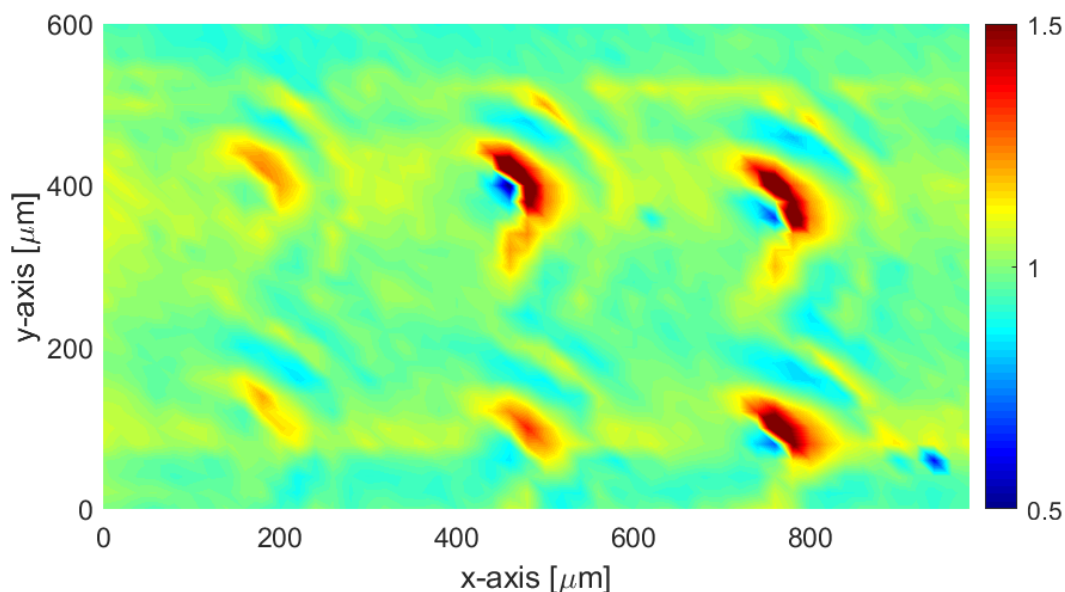
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178 **Figure 6:** Measured guided wave field amplitude in front of the wedge transducer; (a) A_0 mode
 179 $\langle 110 \rangle$ direction; (b) A_0 mode $\langle 100 \rangle$ direction; (c) S_0 mode $\langle 110 \rangle$ direction; (d) S_0 mode
 180 $\langle 100 \rangle$ direction.

181 The guided wave propagation in front of the wedge transducer was characterized from the non-
 182 contact laser measurements. For the A_0 mode a strong wave beam with high amplitude was
 183 observed for the two principal directions (Fig. 6a/b), with some slight beam widening for the
 184 lower stiffness $\langle 100 \rangle$ direction (Fig. 6b). For the S_0 mode a strong wave beam was measured
 185 in the stiffer $\langle 110 \rangle$ principal direction (Fig. 6c), but significant beam widening in the lower
 186 stiffness $\langle 100 \rangle$ direction was observed, which significantly reduces the range for defect
 187 detection. Together with the significant wave skew angle of the S_0 mode in the non-principal
 188 directions due to the variation of the phase slowness [21] and a two times larger wavelength
 189 for the selected center frequency, the S_0 mode has some disadvantages for the monitoring of
 190 defects compared to the A_0 mode.

191 **4. DEFECT SCATTERING**

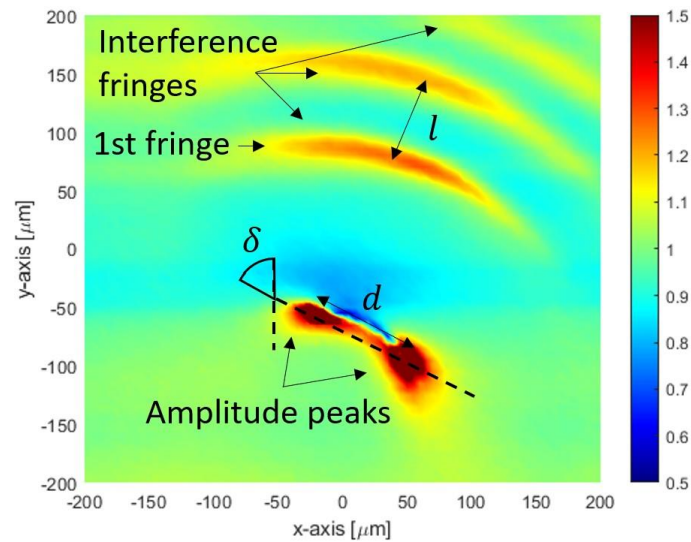
192 A first measurement of the guided wave interaction with a cluster of defects was conducted for
193 the silicon wafer containing 6 indents with a spacing of approximately 300 μm (Fig. 7). The
194 measurement was conducted with a step size of 20 μm in both directions to cover the area
195 containing the defects. The amplitude of the guided wave field was normalized and corrected
196 for the gradient of the incident wave by fitting and subtracting a plane with a linear gradient in
197 the y-direction (propagation direction). The wave field shows some variability of up to 10%
198 of the amplitude of the incident wave, which might likely be caused by changes in the coupling
199 of the wedge transducer over the several hours of measurement duration. At the approximate
200 location of each of the clustered defects, a clear amplitude peak can be observed. The maximum
201 peak amplitude (normalized with the amplitude of the incident wave) increases in general with
202 the severity of the defect, from approximately 1.2 for the 1 N defects to approximately 1.8 for
203 the 2 N defects. Interestingly, for the 1.5 N defects, a variability in amplitude was observed,
204 with normalized amplitudes of 1.3 and 1.9. Around each defect a roughly quarter-circular area
205 of higher and lower amplitude similar to interference fringes was observed, but not of
206 significantly higher amplitude than the variability of the measured guided wave field
207 amplitude.



208

209 **Figure 7:** Measured scattered field around clustered defects (20 μm step size), incident wave
 210 top to bottom, A_0 mode at 5 MHz center frequency, amplitude from FFT, normalized, corrected
 211 for gradient; $x = 200 \mu\text{m}$: indent force 1 N; $x = 500 \mu\text{m}$: indent force 1.5 N; $x = 800 \mu\text{m}$: indent
 212 force 2 N.

213 In order to better understand the interaction of the high frequency guided wave with defects,
 214 the three wafers containing well separated defects with higher indent force (2 N, 3N and 4 N)
 215 were investigated. The amplitude of the scattered field was measured in a smaller area of 400
 216 μm by 400 μm , but with higher resolution step size of 5 μm . This new measurement shows
 217 similar patterns in the vicinity of the indent, but with more details. To evaluate the correlation
 218 of this pattern with the indent size and force, several parameters were chosen for analysis (Fig.
 219 8). The maximum normalized amplitude of the peaks and their orientation relative to the y-axis
 220 (angle δ) were computed. Where two peaks were visible in the scattered field, the peak spacing
 221 d was calculated. For smaller indents, where the two peaks merged to a single peak, the peak
 222 spacing was replaced by the FWHM (full width at half maximum). The interference length l
 223 and the normalized amplitude of the 1st fringe were extracted from the interference-like fringes
 224 in the upper right area of the scans.



225

226 **Figure 8:** Typical pattern (FFT amplitude measurement, incident wave top to bottom) observed
 227 around individual indent at 5MHz with evaluation parameters marked; amplitude of peaks and
 228 angle δ with y-axis; peak separation distance d ; interference length l between fringes and
 229 amplitude of first fringe.

230 Figure 9 presents the 6 FFT amplitude high-resolution scans for defects with a surface extent
 231 from approximately 50 μm to 100 μm (including the cracks). The exact localization of the
 232 defect center for the laser measurement proved difficult, and several steps as described in [27]
 233 were required to define the origin. Although a repeatable pattern was observed for the scattered
 234 wave field at the 6 defects, the guided wave amplitude after normalization shows some
 235 variability. The 6 measurements shown in Fig. 9 were evaluated to understand which of the
 236 parameters illustrated in Fig. 8 correlate with the defect size. Values are presented in Table 1.

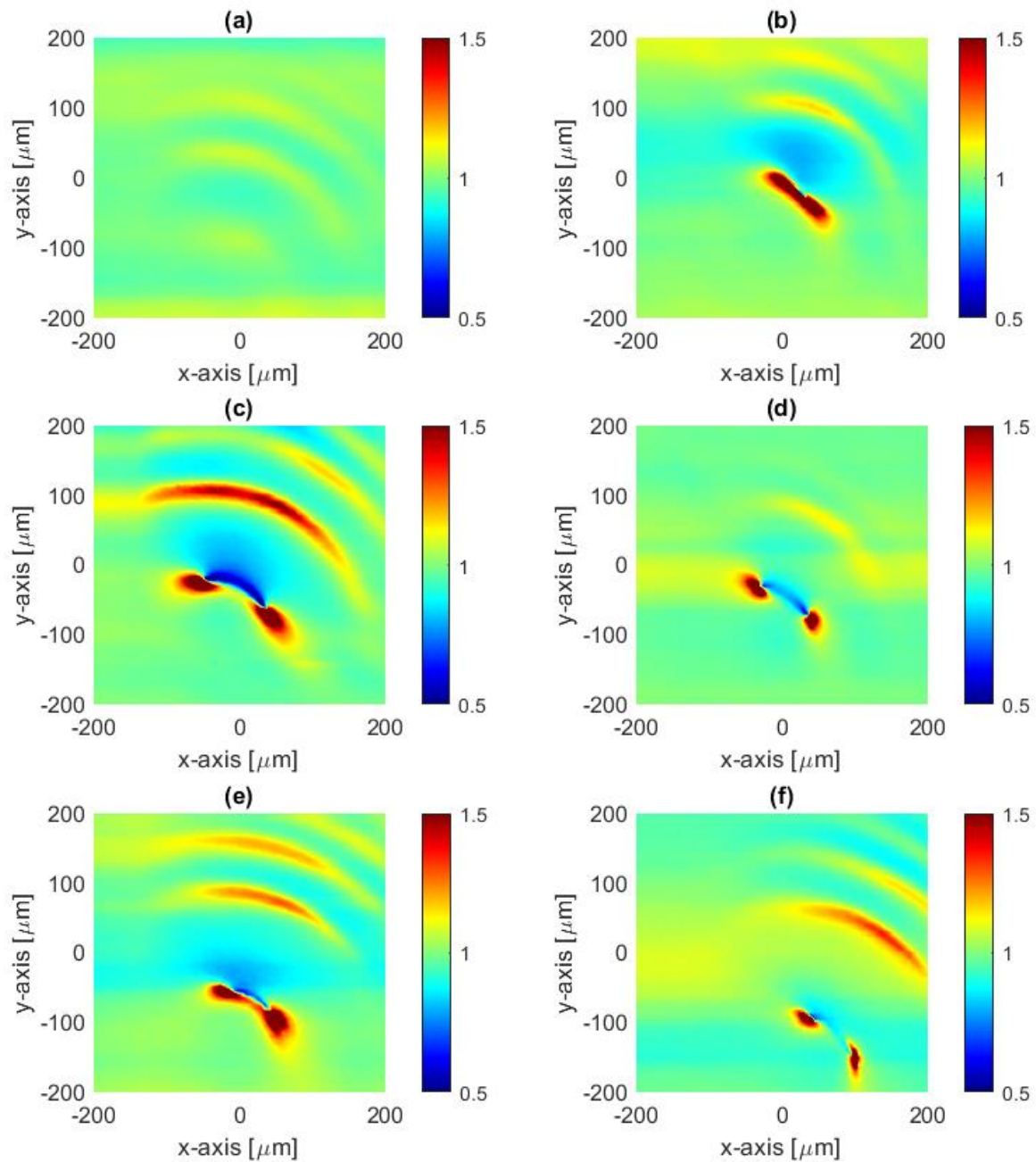
237 High amplitude was observed at the center of each defect, as shown in Fig. 9. For the smaller
 238 defects (indent force 2 N: Fig. 9a, b), only a single peak was visible with a large amplitude
 239 variation between 1.1 and 4.8 of relative amplitude. For the larger defects (indent force 3 N:
 240 Fig. 9c, d, 4 N: Fig. 9e, f), two peaks with (normalized) amplitudes varying between 4.7 and
 241 8.3 were observed. The peak separation d is similar to the defect surface extent, potentially

242 showing that these quantities are correlated. However, it should be noted that the largest peak
243 separation was actually observed for one of the 3N defects (Fig. 9c).

244 The values in Table 1 exhibit the variability visible in Fig. 9 for the peak measurement
245 amplitude between indents of similar size (indent size increasing from left to right). The
246 amplitudes of the main peak increase in general with the indent size, with the exception of
247 measurement (c), which shows higher values for all parameters. The peaks of the amplitude
248 maximum are extremely sharp with typically very high amplitude value (except for
249 measurement (a)), even relative to the small scan step of 5 μm . Therefore, one can expect the
250 peak amplitude to show high variability, since its value is based on a single measurement point.
251 The peak orientation δ is similar for all indents ($\delta = 56 \pm 6^\circ$, mean and standard deviation)
252 and correlates to the left-right asymmetry in the scattered amplitude pattern, which was also
253 present in the measurement shown in Fig. 7. The origin of this asymmetry is not understood so
254 far. Even though silicon has material anisotropy, for the geometry and orientation of the
255 specimens and defects, symmetric scattered fields were expected. This might be correlated to
256 effects of the laser measurement, as the laser spot size on the polished silicon wafer is quite
257 large.

258 In front of the defects, the quarter-circular pattern of high and low amplitudes could be
259 observed with more detail. The amplitudes of the 1st fringe increase in general with the indent
260 size, again with the exception of measurement (c). It is important to stress that the fringes have
261 different shapes compared to the main amplitude peaks, but a similar asymmetry relative to the
262 incident wave direction can be seen. The interference fringes have smooth features, with
263 limited amplitude variation over multiple measurement points, making the measurement of the
264 amplitude evaluation quite robust. Similar to the amplitude of the main peak, significant
265 variation of the fringe amplitude for the same indent force was observed. This is especially
266 evident for the defects generated using the 3N indent force, where Fig. 9d shows a weak fringe

267 pattern comparable to the 2N indents with only one fringe clearly visible, while the fringe
 268 pattern in Fig. 9c has the highest amplitude of all measured defects (see table 1).



269

270 **Figure 9:** Measured scattered field around individual defects (5 μm step size), incident wave
 271 top to bottom, A_0 mode at 5 MHz center frequency, amplitude from FFT, normalized, corrected
 272 for gradient: (a), (b): 2N indent force; (c), (d) 3N indent force; (e), (f): 4N indent force.

273 **Table 1:** Characterization of the 6 measurements at individual defects shown in Fig. 9. Data
 274 with * shows FWHM for single amplitude peak; defect size average combined size of indent
 275 with cracks (horizontal and vertical).

	2N (a)	2N (b)	3N (c)	3N (d)	4N (e)	4N (f)
Defect size (μm)	52	54	76	80	94	95
Normalized peak amplitude	1.1*	4.8*	5.4	4.7	8.3	5.3
Peak orientation δ ($^\circ$)	-	49*	64	59	59	49
Peak separation d (μm)	81*	71*	105	87	82	93
1 st fringe amplitude	1.1	1.2	1.6	1.2	1.3	1.4
Interference length l (μm)	75	64	76	55	70	70

276

277 The distance between areas of high amplitudes (fringes) does not match the expected
 278 correlation to half of the wavelength of the A_0 wave mode for constructive and destructive
 279 interference of incident and scattered waves. The interference length l appears not to be
 280 correlated to the defect size and has similar values for all indents. The average interference
 281 length is $l = 68 \pm 7 \mu\text{m}$ (mean and standard deviation) which is approximately 10 times
 282 smaller than the A_0 mode wavelength at 5 MHz ($\lambda = 800 \mu\text{m}$). As for the asymmetric wave
 283 field, parts of the observed pattern could be due to a combination or interference of the

284 ultrasonic signal and optical artefacts from the laser system. Further investigations are required
285 to clarify this hypothesis.

286 **5. CONCLUSIONS**

287 The applicability of high frequency guided ultrasonic waves for the detection of small surface
288 defects in thin monocrystalline silicon wafers was investigated. In order to improve the
289 efficiency and reduce the cost of renewable electricity generation using solar panels, reduction
290 of wafer thickness and reliable detection of small surface cracking during the manufacture
291 process are required. The propagation of guided waves in the anisotropic wafers was
292 investigated, and good results were achieved experimentally for the A_0 Lamb wave mode at
293 approximately 2 MHz mm frequency-thickness product with a strong beam and propagation
294 range, allowing in principle inspection of complete wafers from a single excitation position.

295 A Vickers indenter with controlled force was used to create defects on the wafer surface. While
296 penetrating into the brittle material, the indenter tip generated surface cracks propagating in the
297 highest stiffness directions from each corner of the pyramid-like tip with a maximum surface
298 extent of 100 μm . The scattering at these surface defects was measured using a non-contact
299 laser interferometer. This showed a repeatable pattern consisting of sharp amplitude peaks at
300 the center and fringes similar to the interference of incident and scattered waves, but with a
301 smaller interference length than expected for the A_0 Lamb wave mode. A systematic
302 asymmetry was clearly visible in the scattering patterns, which could not be explained from
303 the silicon crystalline properties or the surface geometry of the defects. Parts of the observed
304 pattern could be due to a combination or interference of the ultrasonic signal and optical
305 artefacts from the laser system. Further investigations are required to clarify this hypothesis.
306 The amplitude of the peaks and fringes in general increases with indent size and the main peak
307 amplitude is significantly larger than the incident wave amplitude, making defect detection

308 possible. It should be noted that the depth of the defects, except for the indent (up to 6 μm), is
309 not known, and this might account for some of the observed variability.

310 Better characterisation of the depth profile of the generated micro-cracks and improved
311 understanding of the scattered wave amplitude patterns will be required and could be obtained
312 from numerical simulations or further measurements to better quantify the interaction and
313 sensitivity of the proposed measurement methodology. Further investigation must also be done
314 to characterize possible optical artefacts of the non-contact laser measurement for very small
315 defects.

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