# **Diffusion Behavior of Methane in 3D Kerogen Models**

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**ABSTRACT:** As global energy demand increases, natural gas recovery from source rocks is attracting considerable attention since recent development in shale extraction techniques has made the recovery process economically viable. Kerogen is thought to play an important role in gas recovery, however, the interactions between trapped shale gas and kerogen remain poorly understood due to the complex, heterogeneous microporous structure of kerogens. This study examines the diffusive behavior of methane molecules in kerogen matrices of different types (Type I, II, and II) and maturity levels (A to D for Type II kerogens) on a molecular scale. Models of each kerogen type were developed using simulated annealing. We employed grand canonical Monte Carlo (GCMC) simulations to predict the methane loadings of the kerogen models and then used equilibrium molecular dynamics (EMD) simulations to compute the mean square displacement (MSD) of methane molecules within the kerogen matrices at reservoir-relevant conditions, i.e., 365 K and 275 bar. Our results show that methane self-diffusivity exhibits some degree of anisotropy in all kerogen types examined here except for Type I-A kerogen, where diffusion is the fastest and isotropic diffusion is observed. Self-diffusivity appears to correlate positively with pore volume for Type II kerogens, where an increase in diffusivity is observed with increasing maturity. Swelling of the kerogen matrix up to a 3% volume change is also observed upon methane adsorption. The findings contribute to a better understanding of hydrocarbon transport mechanisms in shale and may lead to further development of extraction techniques, fracturing fluids, and recovery predictions.

# 1. Introduction

Shale gas is attracting considerable attention as an unconventional energy source due to the abundance and geographic distribution of shale as well as the everdemand.<sup>1,2</sup> global energy Recent increasing advancements in shale gas extraction techniques such as horizontal drilling and hydraulic fracturing, which enhance the extraction volume and shale permeability, respectively, have allowed successful commercial exploitation of shale hydrocarbon resources. However, a recent assessment suggests only up to a fifth of the hydrocarbons in a shale matrix can be recovered using the aforementioned methods.<sup>3</sup> During production, shale gas migrates from the shale matrix to a production wellbore through a network of fractures created by the injection of pressurised fluid;<sup>4-8</sup> typically water, though interest in supercritical carbon dioxide based approaches is growing.<sup>9,10</sup> It has been established that the rate-limiting step in shale gas production lies in the gas migration from the shale matrix of low permeability into the fracture network. Thus, it is of crucial importance to understand the transport of small organic molecules in shale matrices to explain the relatively low yield and rationalize production decline over time.<sup>11–13</sup> Such studies are also needed to identify ways to optimize the shale gas extraction techniques.

Within shale rocks, the majority of the organic material is present as kerogen - a waxy, complex, heterogeneous mixture of hydrocarbons found in organic matter-rich sedimentary rocks deposited in lacustrine, marine, and terrestrial environments around the world. In such an environment, kerogens formed at the end of diagenesis are slowly buried more deeply in the Earth and undergo thermal maturation (i.e., catagenesis and metagenesis),



Figure 1: The van Krevelen diagram.<sup>14</sup> The arrows indicate the transformations of kerogens over geological timescale due to thermal maturation, and the crosses mark the kerogen types considered in this work.

during which subsurface thermal stress causes kerogen fragments to break down chemically, eventually (producing oil and gas.<sup>15</sup> Kerogens are insoluble in common organic solvents, such as carbon tetrachloride.<sup>16</sup> Because of their chemical heterogeneity, kerogens are typically classified into three main types based on overall C/H/N ratios determined via elemental analysis. Each type Type I, II, and III) occupies a specific region on the van Krevelen diagram (Figure 1),<sup>14</sup> a plot of H/C ratio against oxygen-to-carbon O/C ratio. The type of kerogen present in a shale depends on numerous factors, including but not limited to its biological origin, depositional environment, burial depth and age/maturity level. Previous research<sup>17-19</sup> reports that the main factor dictating the adsorption capacity of shale samples for CH<sub>4</sub> and CO<sub>2</sub> is their total organic content (TOC), much of which is kerogen. Further studies comparing the adsorption capacity of shale and isolated kerogen samples confirmed the uptake of CH<sub>4</sub> and CO<sub>2</sub> is significantly higher in pure kerogen than that of shale.<sup>20-23</sup> Observing and quantifying the adsorption and transport processes in such systems presents numerous experimental challenges, such as the difficulty in obtaining shale samples and isolating pure, unaltered kerogen as well as the need to examine behaviour over many length and time scales. To mitigate these limitations, classical methods of atomistic computer simulations, such as molecular dynamics (MD) and Monte Carlo (MC) or hybrid MD/MC simulations, provide a convenient way of studying kerogen at a molecular level, providing unique insight into the behaviour of these fluids in nanometerscale confined spaces. Molecular simulations can provide detailed insight into the kerogen structure and how that structure responds upon fluid uptake/desorption, the rate at which the uptake/desorption occurs, and the effect of fluid-kerogen interactions on adsorption and transport.

However, the complete molecular structure of isolated kerogens is not known, meaning that construction of realistic molecular models remains a challenge. Several experimental techniques, such as solid-state <sup>13</sup>C NMR spectroscopy, sulphur X-ray absorption near edge structure (S-XANES), and X-ray photoelectron spectroscopy (XPS), have been used to characterize kerogen samples.<sup>24</sup> These studies identify the elemental composition, the fundamental functional groups and their distribution in the structure, and chemical/structural features present in a given kerogen sample. From these data, realistic two-dimensional (2D) kerogen models can be created with appropriate overall chemical composition and functionality. Thanks to several important studies, detailed 2D molecular diagrams of kerogen have been created for several types and maturity levels. Behar and Vandenbrouke<sup>25</sup> constructed detailed 2D molecular diagrams representative of the three main kerogen types at various stages of their evolution. The model structures they proposed have a high starting molecular weight for the least mature kerogens with decreasing molecular weight as kerogen maturity increases. Their pioneering work enabled Zhang and LeBoeuf<sup>26</sup> to modify the models in their study of the volumetric properties of immature Green River Shale kerogen. In particular, Zhang and LeBoeuf's modification brought the kerogen model in closer alignment with the chemical composition of an extracted immature Green River Shale kerogen. Following a similar approach, Siskin et al.<sup>27</sup> also proposed a 2D Green River oil shale kerogen model consisting of seven unique molecules based on selective chemical derivatizations and NMR spectroscopy. More recently, 3D configurations of kerogen based on Siskin et al.'s model have been built through a combination of ab initio and molecular mechanics calculations.<sup>28,29</sup> One common feature of all these models is that the kerogen molecules contain a large number of atoms to duplicate the elemental analysis data. Ungerer et al.<sup>30</sup> overcame this problem by creating six realistic and relatively smaller kerogen models of different types and maturity levels, covering a broad region of the van Krevelen diagram.<sup>14</sup> In addition to constructing models with computational efficiency in mind, incorporating prior knowledge about molecular fragments allowed them to improve kerogen models such that geometrical constraints imposed by (tetrahedral) alkane and (planar) aromatic fragments can be satisfied in 3D. The six kerogen model units are representative of kerogen samples obtained in various depositional environments, where Type I, Type II, and Type III correspond to kerogens of lacustrine, marine, and terrestrial origins, respectively. These individual

molecular units can be used as building blocks to generate larger models of nanoporous kerogen matrices. Since Ungerer et al.,<sup>30</sup> a considerable amount of literature has been published based on their models. More recently, in the same vein, Lee et al.<sup>31</sup> developed three new molecular models for Type II kerogen based on a sample extracted from the Bakken formation.

In a 2016 study, Ho et al.<sup>32</sup> demonstrated that methane release in kerogen matrix is a two-step process, with an initial pressure-driven fast release of free gas followed by a slow release of adsorbed gas through desorption and diffusion across the low permeability kerogen matrix.<sup>33</sup> The second desorption/diffusion step is the rate-limiting step, with diffusion behavior of methane in pores differing from that of bulk methane. Furthermore, Ho and colleagues show that the same packing procedure can give rise to kerogen matrices with different pore connectivity and that some methane might be trapped in isolated pores and non-recoverable. In an investigation into the swelling properties of Type II-D kerogen,<sup>34</sup> they observed swelling of the kerogen matrix following gas adsorption leading to increased surface area, porosity and pore size of the kerogen model. Kazemi et al.<sup>35</sup> calculated both self- and transport diffusion coefficients of methane in a Type II-C kerogen matrix, showing that both converge to a similar value (low diffusivity) as the pressure increases. Michalec and Lísal<sup>36</sup> controlled the microporosity of a Type II-D kerogen model by introducing spherical dummy particles of varying sizes (up to 15 Å) as the kerogen matrices were packed. Subsequently, they compared the adsorption of shale gas on rigid kerogen structures of different microporosity through GCMC as well as MD simulations. The simulated kerogen exhibited lower gas uptake due to lower pore space accessibility, which can be improved by modelling kerogen as a flexible matrix. They also showed that the kerogen matrix preferably adsorbs CO<sub>2</sub> rather than CH<sub>4</sub> due to stronger van der Waals and electrostatic interactions. This result corroborates earlier findings by Sui and Yao<sup>37</sup> and Wang et al.<sup>38</sup> who explored the adsorption of CH<sub>4</sub> and CO<sub>2</sub> in Type II-A kerogen matrix computationally. Through MD simulations, Pathak et al.<sup>39</sup> studied the swelling of flexible Type II-C kerogen



Figure 2: A schematic diagram of the annealing/relaxation procedure. (A) A single Type I-A kerogen unit. (B) The initial configuration of the system before annealing. (C) A kerogen matrix of Type I-A after annealing.

by simulated annealing of kerogen models with a series of 17 fixed mass liquid hydrocarbons. They demonstrated that kerogen exhibits variations in swelling with various organic liquids. Vasileiadis et al.<sup>40</sup> manipulated the porosity of Type II-D kerogen matrices of different system sizes through introduction of dummy particles of different sizes (up to 40 Å) and devised a new algorithm to characterize porosity in kerogen. Their results indicate that system size effects appear to affect the system density and pore size distribution, both of which are also affected by the choice of the force field for the simulations. Pore characteristics, e.g., methane accessible area and volume, increase with increasing number of kerogen units as well as the number of dummy particles

and the dummy particle size. In the absence of dummy particles, the pore limiting diameter (PLD) of kerogen matrices is smaller than TraPPE-UA model of CH<sub>4</sub>.<sup>41</sup> They suggested using a large system size when simulating kerogen models of this kind. In a separate study, they found that a linear relationship between porosity and adsorption capacity, and that diffusion across kerogen matrices is anisotropic.<sup>42</sup> Zhao et al.<sup>43</sup> and Huang et al.<sup>44</sup> considered the effects of maturity and moisture content and on methane adsorption in Type II kerogen matrices. In both studies, a positive correlation has been shown between methane adsorption and kerogen maturity. This is further supported by a later study by Alafnan et al.<sup>45</sup> Tesson and Firoozabadi<sup>46</sup> reported

methane adsorption in both rigid and flexible Type II-A kerogen and concluded that kerogen flexibility has limited effect on the self-diffusion of methane. It is worth pointing out that while ensuring flexibility of the kerogen matrix, their MD simulations were conducted in the NVT ensemble, where deformation was not accounted for. Wu and Firoozabadi47 conducted boundary-driven nonequilibrium molecular dynamics (BD-NEMD) simulations to study the transport of methane across a matrix comprises of 60 Type II-A kerogen molecules. They noted that a reduction in the CH<sub>4</sub> flux is primarily driven by the change in pore size and shape as the kerogen matrix was flexible. He et al.48 developed a new model to quantify the diffusive tortuosity of kerogens based on Type II-C kerogen. In order to realistically model mature Type II kerogens, Rezlerová et al.49 modelled embedding of various molecules in the kerogen matrices of Type II-C and Type II-D through annealing dynamics, thereby inducing microporosity. They also subsequently introduced mesoporosity by creating a slitshaped mesopore of varying sizes (i.e., 20 or 30 Å) between the replicated matrices to model a multi-scale pore network. They computed the adsorption isotherms of pure CH<sub>4</sub> and a binary equimolar mixture of CH<sub>4</sub> and CO<sub>2</sub> for their models and also evaluated the selfdiffusivity of CH<sub>4</sub> and CO<sub>2</sub>. They noted that an overmature Type II-D kerogen has a higher accessible surface area than a mature Type II-C kerogen due to the preferential parallel stacking of the overmature kerogen macromolecules. In addition, this stacking arrangement in the overmature kerogen creates a less tortuous micropore network, such that the self-diffusivity of CH<sub>4</sub> is higher than that in the mature kerogen. Their study also complements the conclusions from other studies<sup>36–38</sup> that kerogens are relatively selective towards CO2 vs. CH4. Comparably, Sun et al.<sup>50</sup> represented fractures in kerogen using slit-shaped nanopores of different sizes (up to 15 Å) and studied adsorption and diffusion of pure gas and a binary mixture of CH<sub>4</sub> and CO<sub>2</sub> in Type II-D kerogen. They found that the kerogen matrix preferably adsorbs CO<sub>2</sub> more than the slit mesopores. Both studies observed that the CH4 adsorption isotherm in kerogen follows type-I Langmuir adsorption behavior. Li et al.<sup>51</sup> conducted MD simulations to study CO<sub>2</sub> storage in water-filled slitshaped nanopores of different sizes through a kerogen matrix of the Type II series. Chong et al.<sup>52</sup> cluster size analysis showed the existence of pore discontinuity upon adsorption of CH4 and CO2, whereas with water, continuity can be observed in the micropores of Type II-A kerogen matrix. In addition, results from multiple investigations<sup>37,43,52–54</sup> into immature Type II-A kerogens with different starting number of kerogen macromolecules have demonstrated that despite achieving similar final simulated density, the differences in the porosity of the resulting packed molecular structures can be as high as an order of magnitude. More recently, Li et al.<sup>55</sup> attempted to shed some light on Schroeder's Paradox,<sup>56</sup> which suggests a discrepancy in solute sorption exists when a polymer is exposed to solvent in saturated vapor and pure liquid states, concerning kerogen systems. By employing hybrid MD-GCMC simulations, they reported that liquid-induced kerogen swelling is greater than that of their gas counterparts.

Together, these studies highlight the complexity of kerogens and that they assume a broad spectrum of pore (connectivity, characteristics constrictivity and tortuosity).<sup>57,58</sup> Adding to the complexity, the use of different kerogen model molecules and packing/annealing procedures may lead to different results. Whilst all the models described here are chemically representative, it is unclear how structurally descriptive they are compared to pure isolated kerogen or kerogen as it would be found in a field sample of shale. Generating a reliable and controlled distribution of 3D kerogen structures remains challenging, particularly given the limited experimental data to use as blueprints for guidance. Given these challenges, in this study we set out to qualitatively assess the diffusion behavior of methane in kerogen matrices constructed from model kerogen molecules representing different types and maturity at conditions relevant to a geological shale-gas reservoir (365 K and 275 bar). A sufficiently large number of macromolecules (i.e., 50) were used to avoid system size effects noted in the literature and no dummy particles were involved to introduce porosity. Unlike earlier studies, in order to mimic reservoir conditions and allow realistic deformation (swelling) of the kerogen matrices, here we computed the MSD of CH4 in hundred nanosecond long MD simulations via both NVT and NPT statistical ensembles. To the best of our knowledge, this is the first study to include all major kerogen types, enabling this work to provide a broader overview of CH<sub>4</sub> diffusion in kerogens.

# 2. Methods

Whilst the exact structures of bulk kerogen are unknown, many researchers have utilized simulated annealing procedures to construct model kerogen matrices out of several model kerogen molecules (Figure 2), where a number of molecular units of a kerogen type are first placed in a large simulation box (low initial kerogen density). The box is subsequently equilibrated through a series of cooling and heating cycles, often in a stepwise manner and at high pressure, resulting in a condensed structure, i.e., a matrix. The density of the final configuration is then compared against experimental data of the corresponding kerogen type to validate the packed structure.

### **Creation of Bulk Kerogen Matrices**

The relaxation procedure set out by Michalec and Lísal<sup>36</sup> was closely followed to prepare the initial configuration of bulk kerogen matrices. 50 kerogen macromolecules were randomly inserted into a large cubic simulation box with an initial system density of ~0.05 g/cm<sup>3</sup> to avoid interactions with copies of their images through the imposed periodic boundary conditions. This was done for each kerogen type reported by Ungerer et al.<sup>30</sup> Then the systems were relaxed through energy minimization via steepest descent, followed by annealing dynamics. A total of 12 cycles were executed as described in Table 1. The choice of final temperature and pressure represents a typical shale reservoir at ~3-4 km depth. Figure 3 shows the changes in system density during annealing for Type I-A system, as an example. To



Figure 3: Heating and cooling profile during annealing dynamics for Type I-A kerogen.

reduce the statistical uncertainty, this relaxation procedure was carried out with ten independent runs, and to check for convergence, the final system density of 10 matrices was then averaged and compared to experimental data. We then randomly selected three of the 10 relaxed matrix structures for the subsequent grand canonical Monte Carlo simulations (GCMC).

### GCMC

GCMC simulations were employed to compute CH<sub>4</sub> loading in selected matrices for each kerogen type. In the grand canonical ensemble, the chemical potential ( $\mu$ ), volume (V) and temperature (T) of the systems are fixed whereas the number of molecules fluctuate. The GCMC simulations allowed us to compare variations in CH<sub>4</sub> loading due to packing as well as to determine the number of CH<sub>4</sub> molecules to insert in the kerogen matrices for MD simulations.

#### **Molecular Dynamics**

In order to generate the initial configurations for MD simulations, one of the three kerogen matrices used in the GCMC simulations was selected for each type of kerogen. We randomly inserted the number of CH<sub>4</sub> molecules corresponding to the predicted CH<sub>4</sub> loadings from GCMC simulations into the kerogen matrices. The steepest descent algorithm was used to relax the CH<sub>4</sub> loaded systems to avoid bad contacts. Then, the systems were subjected to an NVT MD run and successively an NPT run, both for 100 ns where N, V, T and P represent the number of atoms, volume, temperature, and pressure, respectively. In order to minimize statistical error, we performed five independent MD simulations for each kerogen type, with randomly determined different initial CH<sub>4</sub> positions in the matrix. Finally, we computed the mean square displacement (MSD) of methane molecules for each kerogen matrix in all three directions of Cartesian space. The entire workflow for simulations is illustrated in Figure S1.

Table 1: The anneal	ling procedure	for the	creation	of
kerogen matrices.				

Ensemble	<i>T</i> (K)	P (bar)	<i>t</i> (ns)
NVT	900	-	0.2
NPT	900	10	0.2
NPT	$900 \rightarrow 700$	10	0.1
NPT	700	10	0.2
NPT	$700 \rightarrow 500$	10	0.1
NPT	500	10	0.2
NPT	$500 \rightarrow 365$	10	0.1
NPT	365	10	0.2
NPT	365	$10 \rightarrow 275$	0.2
NPT	365	275	0.5
NPT	$1000 \rightarrow 365$	275	1.0
NPT	365	275	1.0

### **Kerogen and Methane Interaction Potentials**

A number of force fields have been used to model kerogen, including but not limited to, COMPASS,<sup>59</sup> CVFF,<sup>60</sup> DREIDING,<sup>61</sup> GAFF,<sup>62</sup> and PCFF+.<sup>30</sup> PCFF+ was originally used in the development of kerogen model units by Ungerer et al.<sup>30</sup> However, the short-range non bonded interactions in PCFF+ are described by a repulsive–attractive 9-6 Lennard-Jones (LJ) potential, which is incompatible with many other force fields that include a 12-6 LJ potential, such that it is difficult to simulate kerogen with other compounds. CVFF is one of the most commonly used force fields in kerogen simulations,<sup>32–34,36,51,52</sup> and the 12-6 LJ potential in its

functional form is compatible with other force fields. CVFF correctly reproduces the experimental density of the Type II kerogen series<sup>51</sup> and provides a reasonable description of kerogen interactions with its constituents, e.g., carbon dioxide (EPM2),<sup>63</sup> methane (TraPPE-UA),<sup>41</sup> and water (SPC).<sup>64</sup> The partial atomic charges of CVFF atoms were assigned using a bond increment scheme. In this study, we further extended the application of CVFF to model two additional immature kerogen types, namely, Type I-A and Type III-A. Methane molecules were represented by the TraPPE-UA model,<sup>41</sup> which has been shown to correctly predict thermodynamic properties and reproduce experimental phase equilibrium data. The relevant simulation input files and force field parameters used in this work can be found in the supporting information section.

## **Simulation Settings**

RASPA molecular simulation package<sup>65</sup> (2.0.39) was used for GCMC simulations. Each GCMC run included a  $5 \times 10^5$  initialization cycle followed by a  $5 \times 10^5$ production cycle, where each cycle is N steps. N is equal to the number of particles present in the system. Owing to the enormous number of atoms involved, for computational efficiency, kerogen structures at this stage were treated as a rigid framework, and the methanekerogen interactions were pretabulated. During the GCMC simulations, insertion/deletion, translation and reinsertion of methane molecules were sampled with equal probability (approx. 33 %). The acceptance rules for insertion and deletion in the grand canonical ensemble included fugacity, a measure of chemical potential, and was calculated using the Peng-Robinson equation of state for methane.66

All MD simulations were performed using GROMACS 2020.4 molecular dynamics simulation software.67 A time step of 0.5 fs was used in the integration of Newton's equations of motion via the leapfrog algorithm. During the annealing procedure, a velocity rescaling thermostat<sup>68</sup> was used to control the system temperature, whereas during the MD production runs, a Nosé–Hoover thermostat<sup>69,70</sup> with a coupling time of 0.1 ps was used. In the NPT simulations, the pressure was maintained using an anisotropic Parrinello-Rahman barostat<sup>71</sup> with a coupling time of 0.5 ps, such that all three dimensions of the systems were allowed to fluctuate independently to allow realistic deformation of kerogen matrices. Long range Coulombic interactions were calculated using a smooth particle mesh Ewald (PME) method<sup>72</sup> of a fourth order polynomial with a mesh width of 0.12 nm.

In both MD and MC simulations, periodic boundary conditions were applied in all three directions. LJ potential was used to describe short-range non-bonded interactions. The LJ interactions of unlike atom pairs in different molecules or further than 3 bonds of the same molecule were calculated using the Lorentz-Berthelot mixing rules. A cut-off radius of 14 Å was used for the LJ interactions and the real part of the Ewald summation. Long-range dispersion corrections were not applied to energy or pressure. All simulations were carried out at 365 K and 275 bar unless stated otherwise. Sample input files, including force field parameters, from MD and GCMC simulations are provided in the Supporting Information.

Utilising PoreBlazer v4.0,<sup>73</sup> a grid-based algorithm was used to calculate and characterize the porosity of the kerogen matrices, before and after the MD simulations. Settings in PoreBlazer were modified such that CVFF interaction parameters, 14 Å cut-off radius, and the cubelet size of 0.2 Å were used in the calculations. Pore limiting diameter, which is defined as the maximum penetrant diameter where a pore network remains percolated, was calculated. In addition, accessible pore volume was estimated with the Widom's ghost atom insertion method<sup>74</sup> by using a helium probe, with parameters taken from Hirschfelder et al.,<sup>75</sup> where  $\sigma_{He} =$ 2.64 Å and  $\varepsilon/k_{B_{He}} = 10.9$  K. Accessible geometric surface area of the kerogen matrices was calculated using a N<sub>2</sub> probe of size 3.314 Å.<sup>76</sup>

# Results and Discussion

#### Kerogen Model Validation and Characterization

Before investigating the effects of anisotropy and heterogeneity of the different kerogen matrices, we validated the CVFF force field to ensure the calculated densities of the kerogen matrices are comparable to experimental data. For each kerogen type, we created 10 condensed kerogen matrices, each starting from a random initial configuration of kerogen molecules, using the simulated annealing procedure outlined in the methods. Table 2 shows the average simulated density of six kerogen types in comparison to relevant experimental data. The density of Type I-A compares quite well with the Green River Shale sample. However, the densities of the Type II kerogen series appear to be slightly lower or close to the lower bound of experimentally measured Kimmeridge Clay kerogen densities, except for Type II-B for which we found no experimental comparison set in the literature. A possible explanation for this might be that Ungerer et al.'s Type II kerogen model molecules<sup>30</sup> are based on Duvernay kerogen samples. It is well understood that kerogens of the same maturity may exhibit different chemical compositions,<sup>24</sup> even when extracted from the same geological site. In addition, their composition is also dependent on the evolutionary history of the sedimentary formation that they formed in and the heterogeneity in composition of kerogen samples. Experimental data<sup>77</sup> suggest that Duvernay kerogens have a density range of  $1.28 \pm 0.3$  g/cm<sup>3</sup>, which our results fit

well. Similarly, the same explanation applies to Type III-A kerogen,<sup>24</sup> since the calculated density is reasonably close to the lower boundary of the reported range. A comparison of our results with the density of Type II

Kerogen	Source	Density (g/cm <sup>3</sup> )		
Туре		Experimental	Simulated	
I-A	Green River Shale	0.95 <sup>78</sup>	$0.964 \pm 0.004$	
II-A	Kimmeridge Clay Formation	1.18–1.29 <sup>79</sup>	$1.117\pm0.008$	
II-B	-	-	$1.080\pm0.017$	
II-C	Kimmeridge Clay Formation	1.18–1.25 <sup>79</sup>	$1.138\pm0.022$	
II-D	Kimmeridge Clay Formation	1.30–1.40 <sup>79</sup>	$1.305 \pm 0.015$	
III-A	Blanzy–Montçeau Basin	$1.16 - 1.20^{80}$	$1.119 \pm 0.060$	

Table 2: Comparison of experimental and averaged simulated kerogen densities at 365 K and 275 bar.

kerogens from other modelling studies shows that our results are slightly lower than the reported values. It is important to highlight that during the initial step of our packing procedure, the kerogen units were inserted randomly to account for a wide range of structural variations, whereas most researchers<sup>36,46,52</sup> considered limited cases, where the kerogen macromolecules were oriented in the same direction, with the assumption that they will ultimately evolve into a uniformly layered kerogen matrix. Furthermore, the discrepancy between our results and those of other Type II modelling studies could be attributed to the packing procedure and conditions or the number of kerogen molecules involved. Vasileiadis et al.<sup>40</sup> reported system size effect when using kerogen models of the same kind, where system density decreases as the number of kerogen molecules increases. It seems possible that these results are related to the favorable stacking of the polyaromatic units of kerogen macromolecules,<sup>81</sup> where pores in one direction may be elongated and cause reduction in system density. Overall, the simulated density of kerogen matrices is in the range of 0.96 - 1.20 g/cm<sup>3</sup> which is in line with the general trend of the available experimental and modelling data.

## **GCMC Simulations**

To account for structural variation that arises during the annealing procedure caused by the random initial configurations of the kerogen units used, we first randomly chose three of the 10 packed structures for each kerogen type for a consistency check. First, we perform GCMC simulations to determine their  $CH_4$  adsorption capacity. Table 3 shows an overview of methane adsorption for the kerogen matrices estimated via GCMC simulations. It is apparent that immature kerogen of Type I has the lowest methane uptake, followed by relatively moderate and high methane adsorption capacity for Type II-A and Type III-A immature kerogens, respectively. In addition, the CH<sub>4</sub> adsorption capacity increases with thermal maturity in Type II kerogen series. These results match those observed in earlier modelling studies<sup>43–45,49,53</sup> and experimental measurements.<sup>82–86</sup> The CH<sub>4</sub> adsorption capacity was found to be highest in Type II-D kerogen compared to the other kerogen types.

The amount of CH<sub>4</sub> loading has a positive correlation with increasing aromatic content of the kerogen units, where Type II-D kerogen has the highest percentage of aromatic content of all.<sup>30</sup> This can be explained by the preferential parallel stacking of the polyaromatic parts of the kerogen macromolecules,<sup>81</sup> and hence larger accessible surface area, resulting in increased CH<sub>4</sub>kerogen interactions.<sup>49</sup> The results of the aforementioned relationships are depicted in a 3D scatter plot in Figure 5.

Table 3: The amount of CH<sub>4</sub> loading in various rigid microporous kerogen matrices obtained via GCMC simulations at 365 K and 275 bar.

Kerogen Type	Loading (mmol/g)
I-A	$0.430 \pm 0.002$
II-A	$0.683 \pm 0.001$
II-B	$1.353 \pm 0.001$
II-C	$1.363 \pm 0.002$
II-D	$2.034 \pm 0.004$
III-A	$1.969 \pm 0.002$



Figure 4: A schematic diagram of the preparation of kerogen matrix loaded with methane. (A) An unfilled kerogen matrix. (B) A random initial configuration of methane molecules, where the capacity was pre-determined through GCMC. (C) A filled kerogen matrix.



Figure 5: The amount of CH<sub>4</sub> loading versus the surface area of 3D kerogen models of various types and the fraction of its aromatic carbon content.

From the chart, there is a clear trend of increasing CH<sub>4</sub> loading with the increase of the aromatic carbon content and surface area of the kerogen matrices. Closer inspection of the figure shows that the surface areas of the considered kerogen matrices are not uniform, and this is likely due to the packing procedure that we employed. The computed methane uptake in our overmature Type II-D kerogen models, which did not have any artificially induced microporosity, is 2.5 times higher than thatreported by Michalec and Lísal.<sup>36</sup> A possible explanation for this might be that Michalec and Lísal<sup>36</sup> used a smaller number of kerogen molecules and are therefore more susceptible to finite size effects. It has been reported that structural properties such as methane accessible area changes with the number of macromolecules as well as the packing procedure used.<sup>37,40,43,52-54</sup> As a result, the pore network characteristics of the packed structures from our studies are likely to differ from other studies whether artificial porosity is introduced in those studies or not. As a result of these size effects and a limited sample size, the relationship between pi stacking and increased surface area might not generalize for other kerogen types or packing methodologies.

### Methane Self-Diffusivity

Diffusion is the primary mode of gas transport in shale rocks, making diffusion rates crucial for understanding hydrocarbon recovery and the factors that influence it. Anisotropic behavior of shale rocks as a result of shale bedding has been found to influence permeability. A number of recent studies<sup>38,42</sup> have reported anisotropy in Type-II kerogen structures. To improve our understanding of shale gas transport mechanisms (i.e., mainly CH<sub>4</sub>), we computed the MSD of pure CH<sub>4</sub> along all three Cartesian directions in kerogen matrices of different types and maturity levels. Isolated and connected pores exist in kerogen and will almost certainly affect methane uptake and mobility; thus, we randomly inserted  $CH_4$  into the selected kerogen matrices to evaluate the effect or pore network accessibility on  $CH_4$  self-diffusivity. For each kerogen type, five  $CH_4$ -



Figure 6: The averaged MSD plots of CH<sub>4</sub> in the x, y and z directions from five independent NPT ensemble MD simulations in six kerogen types at 365 K and 275 bar. Refer to Figure S3 for unscaled version of each plot and Figure S6 for analysis of anisotropy in MSD of methane molecules.



Figure 7: Diffusion coefficients of CH4 in kerogen matrices of various types versus its density and helium accessible pore volume.

loaded matrix configurations were generated such that each configuration had different initial positions of CH<sub>4</sub>. A total of 30 EMD simulations were performed, i.e., five for each kerogen type. Figure 6 displays the averaged MSD plots for CH<sub>4</sub> in various kerogen matrices over the last 50 ns in the NPT ensemble at 365 K and 275 bar. Similar to the earlier observation, Type I-A kerogens exhibit isotropic behavior with no appreciable difference in the MSD in any direction (Figure S6). Looking at the MSD plots for Type II-A and Type II-B kerogens, there appears to be 2D planar regions where the diffusion rate is similar and differs from that perpendicular to the plane. MSD plots of Type II-C show a large standard deviation for the x-component but a low standard deviation for their y-component, suggesting a more uniform interconnected pore structure such as a channel in the y-direction. Both MSD plots for the Type II-D and Type III-A kerogen models show a similar behavior to that of Type II-C kerogen. Furthermore, radial distribution function plots methane with atoms of different functional groups present in kerogen models show that predominant methane binding sites may vary from one kerogen model to another. Due to the specific composition of the model and the specific conformations limiting access to certain types of sites for some models. But overall, it appears that sulfur, nitrogen and oxygen atoms of different functional groups provide the majority of preferred binding sites.

Figure 7 presents the 3D scatter plot of the relationship between the diffusion coefficients of  $CH_4$  and the density and pore volume of the swollen kerogen matrices. The most interesting aspect of this graph is that, apart from Type I-A kerogens, a positive correlation is found

between the diffusion coefficients of CH<sub>4</sub> and the pore volume of kerogens. The most unexpected result is the anomalous CH<sub>4</sub> diffusion rate in Type I-A kerogens, which suggests that there are other factors at play that affect the CH<sub>4</sub> diffusion in kerogen matrices. The observed increase in CH4 diffusion could be attributed to the lack of aromatic carbon in Type I-A kerogens, where less stacking could potentially create more diffusion paths, such that CH<sub>4</sub> could permeate easily in all three directions, i.e., isotropic diffusion. The percentage of aromatic groups in kerogen molecules increases from Type I to Type III kerogen, with the highest in Type II-D kerogen. The increase in polyaromatic regions of the kerogen units and subsequent stacking of these regions may result in a matrix that is not amorphous and hence create a system with anisotropy. The MSD of CH<sub>4</sub> in Type I-A kerogen is an order of magnitude higher than the other kerogen types, which suggests that in Type I-A, CH4 has the fastest diffusion rate compared to other kerogen types. This makes sense in light of the lower density of this material and the larger size of the molecular units, leading to a packed structure with greater pore volume and connectivity compared to the other packed models. The lowest diffusion rate is observed for the Type II-A kerogen. The results from the Type II kerogen series further support previous findings<sup>38,42</sup> that diffusion behavior of CH<sub>4</sub> in Type II kerogens is anisotropic, although many of the observed differences occur within the error bars on each plot. The magnitude of the standard deviations can be attributed to different pore space accessibility as a result of several factors outlined above. In general, the magnitude of the standard

deviations is similar between different plots (typically fractions of a  $nm^2$  at early stages of the simulations and roughly 2-4  $nm^2$  - after 50 ns) and is on the same order of magnitude for each direction.

#### Volume Changes up on CH<sub>4</sub> Adsorption

Table 4 compares the structural properties of the kerogen matrices before and after MD simulations, showing that kerogens experience volumetric strain upon CH<sub>4</sub> adsorption. This is consistent with other literature data.<sup>34,37</sup> In addition, the pore limiting diameter of all kerogen matrices appears to be smaller than the methane molecular diameter in the TraPPE-UA force field, i.e., 3.73 Å. This suggests methane should not be able to diffuse through the percolated pore network. Yet some diffusion is observed, most notably in the Type I-A model (Figure 5). This finding is consistent with the trend reported in an earlier study<sup>40</sup> for Type II-D kerogen. Diffusion of CH<sub>4</sub> through pores that are too small to permit methane motion suggests that the CH<sub>4</sub> molecules themselves may actively change pore accessibility, opening up micropores that were previously inaccessible. Indeed, pore size distribution analyses of kerogen models before and after the MD simulations suggest that methane imbibition has caused a shift of pore size distribution from the initial peak at approx. 3.2 Å to 3.8 Å, which is about the size of a TraPPE united atom methane molecule (Figure S4). Slow diffusion observed in kerogen types other than Type I-A could be due to CH<sub>4</sub> molecules being confined and localized in small nanospaces, i.e., isolated pores with no diffusion pathways. Finally, total helium volume of the kerogen matrices appears to have increased after CH<sub>4</sub> adsorption, which suggest that CH<sub>4</sub> adsorption increases the porosity of the matrices, consistent with the hypothesis that CH<sub>4</sub> actively opens diffusion pathways. Additional studies that can provide energetic insight into this process are warranted. Vasileiadis et al.42 reported that the self-diffusion coefficient has a positive correlation with the pore limiting diameter. However, the findings of this study do not support their results. However, the findings of this study do not support their results. In any case, construction of 3D kerogen matrices from relatively low molecular weight kerogen macromolecules should be handled with care as the lack of cross-linking between kerogen units may allow increased flexibility which is not present in real kerogen.36

Property	Kerogen Type	Initial	Final*	Change (%)
Volume (nm <sup>3</sup> )	I-A	328.58	$330.06\pm0.97$	0.45
	II-A	289.11	$290.30\pm0.91$	0.41
	II-B	262.33	265.13 ± 1.25	1.07
	II-C	250.45	$253.56 \pm 0.54$	1.24
	II-D	161.04	$165.25 \pm 0.77$	2.62
	III-A	249.08	$252.08\pm0.74$	1.20
Pore limiting diameter (Å)	I-A	1.580	$1.578\pm0.183$	-0.13
	II-A	1.540	$1.542\pm0.098$	0.13
	II-B	1.890	$1.892 \pm 0.190$	0.11
	II-C	2.460	$2.460 \pm 0.284$	0.00
	II-D	2.930	$2.932 \pm 0.544$	0.07
	III-A	2.960	$2.956\pm0.445$	-0.14
Helium accessible pore volume (cm <sup>3</sup> /g)	I-A	0.061	$0.074 \pm 0.002$	21.64
	II-A	0.071	$0.085 \pm 0.002$	19.72
	II-B	0.117	$0.137 \pm 0.004$	17.44
	II-C	0.118	$0.142\pm0.002$	20.00
	II-D	0.176	$0.216\pm0.006$	22.84
	III-A	0.167	$0.193 \pm 0.004$	15.57

\* Average volume of the final frames of 5 independent MD runs.

# 4. Conclusions

The primary motivation of this study was to examine CH<sub>4</sub> diffusion in 3D kerogen matrices built from model kerogen molecules that span the recognized types and maturity levels of natural kerogens. The MD and GCMC simulation results reported here show that swelling of the 3D kerogen matrix occurs upon CH<sub>4</sub> adsorption. Methane adsorption also induces other changes in the pore network such as variations in the pore limiting diameter and accessibility between pores in the network. Importantly, all kerogen structures exhibit some degree of anisotropy in the CH<sub>4</sub> diffusion behavior with the exception of Type I-A kerogen (Figure S6), the least mature and lowest density model. As has been observed for fluid diffusion in other microporous to nanoporous materials, the diffusion of CH<sub>4</sub> in kerogen matrices is slow compared to bulk CH<sub>4</sub>.

This study has many broader implications for understanding CH<sub>4</sub> diffusion and recovery in kerogens and shales. The scientific community already knows that kerogens contain a significant fraction of the methane in shale-gas reservoirs<sup>87-89</sup> and that the recovery of methane from shale-gas reservoirs is low compared to other types of gas reservoirs (a few 10s of %).90,91 What is less clear is why the recovery is so low and what can be done about it. The solutions to both of these questions rely on a detailed molecular-scale understanding of methane behavior in shales and all shale components, including kerogen. The ability to extract kerogen and know the extracted kerogen has the same physical and chemical characteristics as kerogen in the native rocks is currently limited. Thus, computational approaches are likely key to developing a molecular-scale understanding of what is happening in shale-gas reservoirs, and will also be needed to help verify experimental data once the community at large is able to collect detailed molecular-scale experimental data for methane adsorption and dynamics in kerogens. Thus, there is a pressing need to have good computational models of kerogens, to understand how to build realistic computational models of kerogens, and to do both of these for a broad range of kerogen types. Due to the scope of the kerogen models examined in this study and the clear procedures for generating kerogen matrices, one of the more important broader implications of the present study is that it forms an important baseline data set for future research into gas recovery from kerogen. These results also highlight several operational considerations that must go into creating computational models of kerogen and interpreting and comparing the associated molecular modelling data sets. Operationally, the annealing procedure employed when creating a 3D kerogen matrix from small representative model pore characteristics of the dense kerogen structures. More research is needed to understand the relationship between the annealing method and the pore accessibility of the resulting packed structures. To improve consistency between modelling studies and applicability to the field, our results and those in the literature also suggest that more experimental data are required regarding the pore network and network connectivity and that future modelling studies must map the model kerogen matrix pore network to match these data. Based on the volume expansion and potential active role of CH<sub>4</sub> in the pore network, future work should be conducted under NPT conditions to enable swelling of the kerogen matrix. Likewise, additional modelling studies that estimate the energetics of methane adsorption and kerogen matrix swelling are warranted to improve our understanding of CH4 transport thermodynamics in this important phase of shale gas reservoirs. The results presented here make it possible to conduct new and much needed modeling calculations of thermodynamic parameters regarding methane adsorption in a broad array of kerogen types, transport calculations, and computational models of methane/water and methane/carbon dioxide competition for binding sites in kerogens. The results presented here show that mature kerogens have more anisotropic diffusion pathways than immature kerogens and that it should be easier to recover methane from kerogen types with higher methane diffusivities and when the kerogens have the physical space required to swell since the flexibility of the kerogen network seems key to enabling methane diffusion. The latter point is particularly important given the known volume expansion of shales due to water uptake by clay minerals, which may alter the pore network and kerogen density/porosity.<sup>4</sup> Ultimately, a detailed molecular-scale understanding of methanekerogen interactions from this and future studies using similar methods of constructing kerogen models will contribute to practical in-field solutions for increasing the methane recovery from kerogens in shales.

molecules plays an important role in determining the final

# 5. Supporting Information

A workflow diagram of the simulation set up procedure. The average MSD of  $CH_4$  in the x, y and z directions in six kerogen types in NVT and NPT ensembles. Normalised MSD of  $CH_4$  in kerogens. Pore size distributions of kerogens before and after methane loading. RDFs between  $CH_4$  and between atoms of different functional groups in kerogen macromolecules of various types. Files containing simulation input parameters.

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