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The Role of Anisotropy in Distinguishing Domination of Néel or Brownian Relaxation Contribution to Magnetic Inductive Heating: Orientations for Biomedical Applications

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Abstract: Magnetic inductive heating (MIH) has been a topic of great interest because of its potential applications, especially in biomedicine. In this paper, the parameters characteristic for magnetic inductive heating power including maximum specific loss power (SLP_{max}), optimal nanoparticle diameter (D_c) and its width (ΔD_c) are considered as being dependent on magnetic nanoparticle anisotropy (K). The calculated results suggest 3 different Néel-domination (N), overlapped Néel/Brownian (NB), and Brownian-domination (B) regions. The transition from NB- to B-region changes abruptly around critical anisotropy K_c . For magnetic nanoparticles with low K ($K < K_c$), the feature of SLP peaks is determined by a high value of D_c and small ΔD_c while those of the high K ($K > K_c$) are opposite. The decreases of the SLP_{max} when increasing polydispersity and viscosity are characterized by different rates of $d(SLP_{max})/d\sigma$ and $d(SLP_{max})/d\eta$ depending on each domination region. The critical anisotropy K_c varies with the frequency of an alternating magnetic field. A possibility to improve heating power via increasing anisotropy is analyzed and deduced for Fe_3O_4 magnetic nanoparticles. For MIH application, the monodispersity requirement for magnetic nanoparticles in the B-region is less stringent, while materials in the N- and/or NB-regions are much more favorable in high viscous media. Experimental results on viscosity dependence of SLP for $CoFe_2O_4$ and $MnFe_2O_4$ ferrofluids are in good agreement with the calculations. These results indicated that magnetic nanoparticles in the N- and/or NB-regions are in general better for application in elevated viscosity media.

Keywords: magnetic heating; Néel & Brownian relaxation; particle anisotropy; polydispersity; ferrofluid viscosity

1. Introduction

MIH of magnetic nanoparticles (MNPs) while exposed to an alternating magnetic field (AFM) has attracted a lot of attention due to its potential applications in various domains, in particular cancer hyperthermia treatment and drug delivery. In cancer hyperthermia, it uses MIH on the principle that the cancer cells can be killed due to their higher thermal sensitivity than that of healthy ones [1–11]. For drug delivery, the MIH can be used to trigger remotely or on-site release of drug molecules being tagged to the MNPs [1–3,8,11–13].

Recently, new interesting biomedical applications were proposed including tuning the cellular gate for regulation of plasma glucose [14], biomaterials devitrification [15–17] and determination of MNPs accumulation in various organs [18]. In order to optimize the amount of MNPs used in the MIH, many studies focused on the conditions for maximizing heating efficiency, which is described theoretically by SLP or experimentally by specific absorption rate (SAR). The MIH power of MNPs has been known to consist of 3 main contributions: magnetization hysteresis, Néel relaxation generated by changing direction of the magnetic moments, and Brownian relaxation due to physical movement of MNPs. According to the linear response theory (LRT) [19,20], the Néel and Brownian relaxation processes are the main mechanisms to generate heat by the fluid containing superparamagnetic nanoparticles (NPs). Similar to other physical properties, the MIH power depends strongly on the diameter (D) of MNPs. Finding optimal particle diameter at which the SLP exhibits maximum value is frequently the purpose of both theoretical and experimental MNPs-MIH research.

Theoretical calculations indicated that, as a result of the competition between Néel and Brownian dissipation processes, the SLP versus D curve exhibits a peak shape with maximum value (SLP_{max}) at a characteristic particle diameter D_c for different material parameters and field conditions [21–26]. From the studies of various Fe_3O_4 (FO) MNPs with different size at frequency $f = 376$ kHz, the maximum SLP_{max} was found for a 16 nm (size standard deviation, $\sigma = 0.175$) sample [27–29]; for iron oxide nanocubes at $f = 520$ kHz, it is 19 ± 3 nm [30,31]. An analysis of various data after taking a normalization of field factor gave an elevation of SLP in the diameter range from 14–18 nm [25]. For $\gamma-Fe_2O_3$, elevation of SLP for 5 sample sizes from 5.3 nm to 16.5 nm ($\sigma = 0.19$ – 0.43) is in agreement with the theoretical prediction for the peak appearance at $D_c = 14.5$ nm [22]. Although these results have confirmed the existence of optimal particle size D_c of MNPs in MIH, finding the value of D_c is not a simple task in experimental works because of the impact of other parameters such as the magnetic anisotropy, viscosity of fluid, and the size distribution.

Firstly, while the magnetic anisotropy of MNPs can be considered as a constant in theoretical works, it depends strongly on the particle size. Additionally, it was showed that the SLP_{max} and D_c decreases monotonically with increasing anisotropy [23,24]. It is noted that, due to the difference in morphology, the synthesized MNPs of the same material can have quite different polydispersity, anisotropy K values (e.g., FO MNPs having K ranging from about 20 to 550 kJ/m³ [32,33]) which can influence the SLP_{max} and D_c parameters. Secondly, the value of SLP_{max} was found to reduce with expansion of size distribution [22]. In practice, there will be some size distribution of the MNPs regardless of the synthesis method and producing MNPs with a sufficiently narrow size distribution is a difficult task. Thus far there have been few reports on experimental verification of calculated optimal D which would require the ability to synthesize MNPs with high monodispersity and to tune D precisely [22,25,28,34]. Thirdly, there might be some relationships between the value of D_c and SAR and ferrofluid viscosity. The size dependences of heating power with some variations of D and ferrofluid viscosity was studied experimentally and theoretically which showed for example for $\gamma-Fe_2O_3$ and $CoFe_2O_4$ MNPs SAR decreased about 20% and 80% with an increase of the viscosity, respectively [22]. To be able to obtain MNPs with optimal heating power, it is desirable to develop a complete correlation of SLP versus D and to consider impacts of different parameters, e.g., size distribution and shape of the NPs [20–22].

In this work, assessments of the impact of K on overall behavior of the heating power quantity have been systematically analyzed. The characteristics of SLP versus D for different MNPs have been studied by considering the impacts of two experimental parameters of the MNPs ferrofluid: σ and media viscosity (η). The different shapes of $SLP(D)$ or the different rates of $d(SLP_{max})/d\sigma$ and $d(SLP_{max})/d\eta$ for FO MNPs with various anisotropy have been studied and explained as a result of the competition between Néel and Brownian dissipation processes. Here, the experimental results of MNPs of manganese ferrite (MFO), and cobalt ferrite (CFO) are considered as representative of low K (Néel

domination) N type, and high K (Brownian domination) B type, which have confirmed the theoretical predictions.

2. Theoretical Basis

The SLP of the MNPs can be calculated as follows [4,20–26]:

$$\text{SLP} = \frac{P}{\Phi\rho}, \quad (1)$$

where Φ is the volume fraction (%), ρ the mean mass density (g/m^3) and P the loss power density (W/m^3). The calculations were carried out for particles with size standard deviation (or size polydispersity) parameter $\sigma = 0\text{--}0.5$, and viscosity $\eta = 1\text{--}10$ mPa·s, assuming the volume fraction $\Phi = 0.1\%$ and the surface ligand layer thickness of 1 nm.

For the superparamagnetic NPs, it is generally known that Néel and Brownian relaxation losses mainly contribute to the MIH of materials. The Néel relaxation loss is originated from relaxation effects of magnetization in a magnetic field, the Brownian relaxation loss is due to the rotation of the NPs as a whole in a ferrofluid. In the LRT, χ is assumed to remain constant with increasing H . This approach was shown to be valid for small magnetic fields. Based on theoretical results, it was found that the condition of validity for the LRT is $\chi < 1$. The dimensionless parameter, ξ , can be given by ref. [19]:

$$\xi = \frac{\mu_0 M_s V H}{k_B T}, \quad (2)$$

Therefore, LRT based on Néel and Brownian relaxation losses is suitable for superparamagnetic NPs or $H \ll H_C$ [19]. The correlation between SLP (based on LRT), the field amplitude (H), frequency (f) and the MNPs imaginary susceptibility (χ'') is ref. [20,24]:

$$\text{SLP} = \pi\mu_0\chi''H^2f, \quad (3)$$

with μ_0 being the permeability of free space; where

$$\chi''(f) = \chi_0 \frac{f\tau}{1 + (f\tau)^2}, \quad (4)$$

The effective relaxation time (τ) can be determined from the compositional relaxation of Néel (τ_N) and Brownian (τ_B) processes [20,24]:

$$\tau = \tau_N \cdot \tau_B / (\tau_N + \tau_B), \quad (5)$$

with

$$\tau_N = \frac{\sqrt{\pi}}{2} \tau_0 \frac{\exp(KV/k_B T)}{\sqrt{(KV/k_B T)}}, \quad (6)$$

$$\tau_B = \frac{3\eta V_H}{k_B T}, \quad (7)$$

where V and V_H are the volume of the core magnetic, and the whole capped particle, respectively.

For monodisperse superparamagnetic NPs ($\sigma = 0$), the value of SLP is calculated based on above equations. Considering the impact of particle size deviation ($\sigma > 0$), the mean loss power density can be described as ref. [20]:

$$\bar{P} = \int_0^{\infty} P g(D) dD, \quad (8)$$

with

$$g(D) = \frac{1}{\sqrt{2\pi}\sigma D} \exp\left[-\frac{(\ln(D/D_0))^2}{2\sigma^2}\right], \quad (9)$$

and $g(D)$ the distribution function

$$\int_0^{\infty} g(D)dD = 1, \quad (10)$$

3. Results and Discussion

3.1. Characteristics of Optimal Parameters in Domination Regions of Néel or Brownian Relaxations

Magnetite is one of the extensively studied and important in a wide range of medical applications because of its biocompatibility. The anisotropy value of FO MNPs was reported between a few tens to higher than 500 kJ/m^3 [32,33]. The calculations are, therefore, made for different K up to 50 kJ/m^3 and the results of SLP versus particle diameter at K equal to 9 kJ/m^3 , 20 kJ/m^3 , 41 kJ/m^3 are presented in Figure 1.

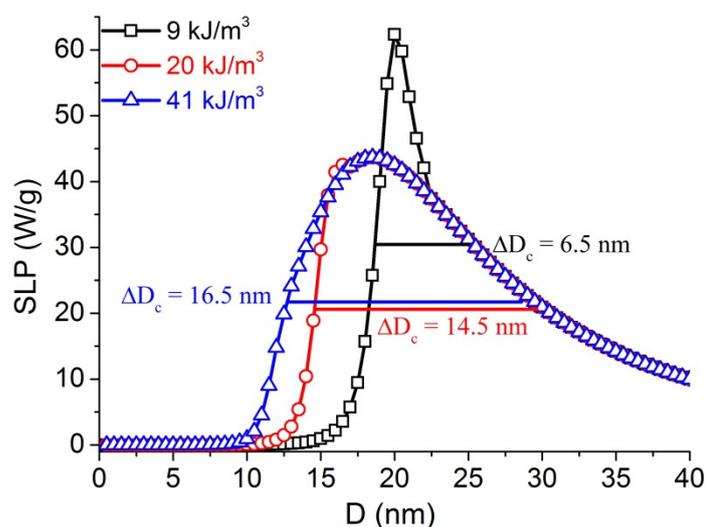


Figure 1. SLP versus D , calculated at $f = 100 \text{ kHz}$, $H = 5.18 \text{ kA/m}$, $\sigma = 0$, $\eta = 1 \text{ mPa}\cdot\text{s}$ for monodisperse FO MNPs with 3 different K of 9, 20, and 41 kJ/m^3 . The bars denote the (full-width-half-maximum (FWHM)) ΔD_c .

The typical peak behaviors of the calculated SLP against D obtained for FO MNPs can be seen in Figure 1. Here, the difference in the SLP (D) graph shape of FO MNPs depends on the value of K [20,23]. We have introduced full-width-half-maximum (FWHM) ΔD_c for describing the peaks. The D_c , ΔD_c and SLP_{\max} obtained at $f = 100 \text{ kHz}$, $H = 5.18 \text{ kA/m}$ (65 Oe), $\sigma = 0$, $\eta = 1 \text{ mPa}\cdot\text{s}$ for monodisperse FO ferrofluids with different anisotropy K are listed in Table 1.

The results showed that the high K (50 kJ/m^3) FO sample (18.5 nm) has much smaller value of D_c compared with that of the low K (3 kJ/m^3) one (28.5 nm). In contrast, there is a large increase in the value of FWHM ΔD_c when the value of K goes from 3 kJ/m^3 (3.5 nm) to 50 kJ/m^3 (16.5 nm). The fact that the SLP peak is narrow for low K but broad for high K can be attributed to Néel or Brownian domination, respectively [9,20,22,24].

Based on Equation (5), it is clear that the effective relaxation time is dominated by the shorter of the two components. For example, A. E. Deatsch et al. found that the Brownian relaxation becomes more significant at 13 nm for high anisotropy particles (40 kJ/m^3) [25]. We calculated the value of relaxation times at D_c for all samples. As can be seen in Table 1, the value of effective relaxation time is approximately equal to the Brownian relaxation

time at D_c when the value of K is higher than 19 kJ/m^3 . The calculations indicated that the Brownian relaxation dominated at D_c when the value of K is higher than 19 kJ/m^3 . Therefore, the difference in peak behaviors of the SLP against D depends on the domination of Néel or Brownian relaxations. As can be seen in Table 1, the value of D_c changes with an increase of K from 3 kJ/m^3 to 20 kJ/m^3 , because the Néel relaxation loss still affects this parameter. When the value of K for FO is higher than 20 kJ/m^3 , D_c and SLP_{\max} are unchanged due to the domination of Brownian relaxation loss. Therefore, it changes abruptly at some critical anisotropy $K_c = 20 \text{ kJ/m}^3$.

Table 1. D_c , ΔD_c , SLP_{\max} , τ_N , τ_B , and τ obtained for monodisperse FO ferrofluids with different anisotropy constant K . ($f = 100 \text{ kHz}$, $H = 5.18 \text{ kA/m}$, $\sigma = 0$, $\eta = 1 \text{ mPa}\cdot\text{s}$).

$K \text{ (kJ/m}^3\text{)}$	$D_c \text{ (nm)}$	$\Delta D_c \text{ (nm)}$	$SLP_{\max} \text{ (W/g)}$	$\tau_N \text{ (s)}$	$\tau_B \text{ (s)}$	$\tau \text{ (s)}$
3	28.5	3.5	76.9	1.95×10^{-6}	10.77×10^{-6}	1.65×10^{-6}
5	24	3	71.4	1.87×10^{-6}	6.67×10^{-6}	1.47×10^{-6}
7	21.5	4	66.2	1.98×10^{-6}	4.92×10^{-6}	1.42×10^{-6}
9	20	6.5	62.3	2.65×10^{-6}	4×10^{-6}	1.6×10^{-6}
12	18.5	9.5	58.3	4.26×10^{-6}	3.27×10^{-6}	1.85×10^{-6}
15	17.5	12.5	50.9	3.24×10^{-6}	2.81×10^{-6}	2.02×10^{-6}
19	17	14	46.9	3.46×10^{-6}	2.6×10^{-6}	1.52×10^{-6}
20	17	14.5	45.9	2×10^{-3}	2.6×10^{-6}	2.6×10^{-6}
21	18.5	14.5	43.6	4×10^{-3}	3.27×10^{-6}	3.27×10^{-6}
30	18.5	15	43.6	4.903	3.27×10^{-6}	3.27×10^{-6}
41	18.5	16	43.6	2.8×10^4	3.27×10^{-6}	3.27×10^{-6}
50	18.5	16.5	43.6	3.4×10^7	3.27×10^{-6}	3.27×10^{-6}

Changes of characteristic parameters, SLP_{\max} , D_c and ΔD_c as a function of K for $f = 100 \text{ kHz}$, $\eta = 1 \text{ mPa}\cdot\text{s}$ are presented in Figure 2, where it suggests 3 different regions. In the lowest K (region I: $K \leq 5 \text{ kJ/m}^3$): SLP_{\max} , D_c and ΔD_c decrease with increasing K . The middle region (region II: $5 \text{ kJ/m}^3 \leq K \leq 20 \text{ kJ/m}^3$) is characterized by a tendency of the decrease of D_c to slow down while ΔD_c starts to increase with increasing K . In the high K region (region III: $K > 20 \text{ kJ/m}^3$), all the 3 parameters become almost constant. Relating to the dissipation mechanisms, the regions I, II and III can be correspondingly assigned as Néel-domination (N-region), Néel/Brownian overlap (NB-region) and Brownian-domination (B-region). While the transition from N- to NB-region is quite smooth, the transition to the B region is distinctly of the first order and in case of FO MNPs it corresponds to a critical anisotropy value K_c of about 20 kJ/m^3 .

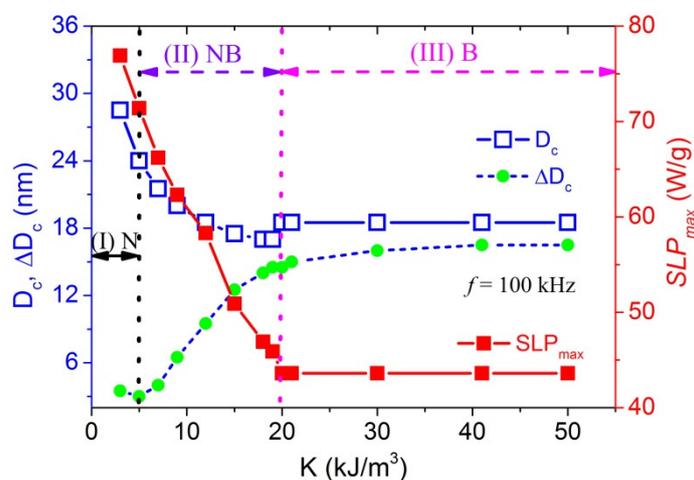


Figure 2. Changes of characteristic parameters, SLP_{\max} , D_c and ΔD_c from Néel-domination to Brownian-domination region calculated for frequency $f = 100 \text{ kHz}$, viscosity $\eta = 1 \text{ mPa}\cdot\text{s}$.

3.2. Dependence of Critical Anisotropy K_c on Frequency (f)

Based on LRT, the value of SLP reaches the maximum SLP_{max} at the optimal particle size D_c when the condition $\omega\tau = 1$ is satisfied. Therefore, the anisotropy boundary of the transition from NB- to B-region might depend on the frequency of AMF.

Based on the data of FO MNPs at $\eta = 1$ mPa·s for the expanded K range and at various frequencies, the SLP was calculated and results of D_c versus K are presented in Figure S1a. For frequencies up to 1000 kHz can be seen an abrupt change of D_c at K_c associated with the transition to the Brownian domination region (for example with $f = 500$ kHz, the value of K_c is about 60 kJ/m³, and the three N, NB and B regions are presented in Figure S1b. The dependence of the critical anisotropy K_c on the field frequency is presented in Figure 3, which follows the function:

$$K_c(f) = A(1 - e^{-B \times (f+f_0)}), \tag{11}$$

where $A = 214.63$ kJ/m³, $B = 0.0000814$ ms, and $f_0 = 81.27$ kHz are fitting constants.

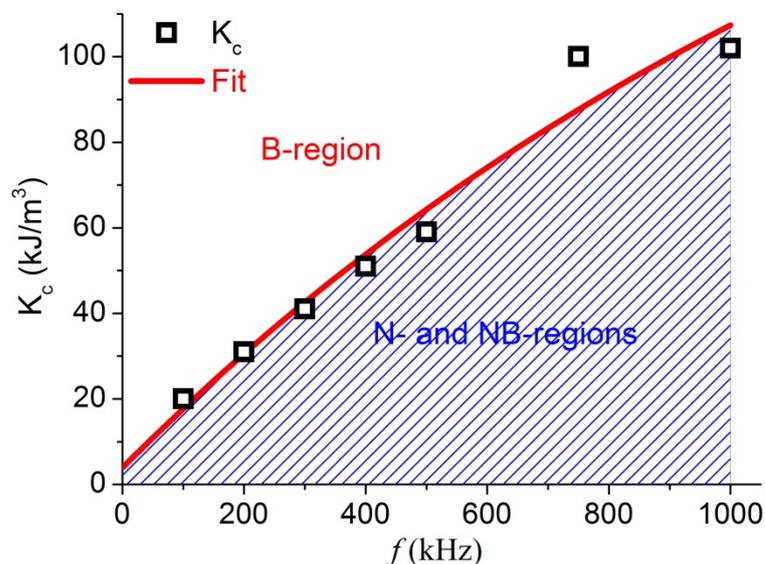


Figure 3. K_c versus f at $\eta = 1$ mPa·s, the fitting gave $R^2 = 0.94121$.

The values of critical anisotropy K_c at different frequencies for fluids with various viscosities are listed in Table 2. For MFO fluid with $K = 3$ kJ/m³, the Brownian relaxation dominates when $\eta = 1$ mPa·s and $f = 100$ kHz (low frequency) but it becomes the Néel relaxation domination when $\eta \geq 1$ mPa·s or $f \geq 100$ kHz. For CFO fluid, the Néel relaxation dominates when $\eta \geq 4$ mPa·s or $f \geq 500$ kHz, but it is the Brownian relaxation domination when $\eta \leq 2$ mPa·s or $f \leq 1000$ kHz.

Table 2. Values of K_c at different frequencies for fluids with various viscosities (η (mPa·s)).

f (kHz)	K_c (kJ/m ³)				
	$\eta = 1$	$\eta = 2$	$\eta = 4$	$\eta = 6$	$\eta = 8$
100	20	33	60	83	102
200	31	56	103	147	205
300	41	74	141	207	261
400	51	95	188	252	322
500	59	112	227	308	399
750	100	153	292	>400	>400
1000	102	205	364	>400	>400

3.3. Orientations to Choose Proper Regions in Biomedical Applications

For biomedical applications, the size distribution and the viscosity of ferrofluid must be considered. In the above calculation of the relaxation times and MIH power of MNPs, the particles were assumed to have the same size, in other words the MNPs are monodisperse with $\sigma = 0$. However, in reality, this assumption cannot be satisfied because usually there would exist some size distribution of the NPs regardless of the synthetic method used. In practice, producing monodisperse MNPs ($\sigma = 0$) is almost impossible task. For example, J. P. Fortin et al. synthesized 16.5 nm γ - Fe_2O_3 with $\sigma = 0.43$, which could be reduced to 0.19 ($d = 5.3$ nm) after size sorting by successive phase separations [22]. M.P. Pileni et al. synthesized 2–5 nm CoFe_2O_4 with 30% polydispersity in size distribution by using functionalized surfactants [35]. When A. K. Gupta et al. compared the different characteristic features of the iron oxide NPs fabricated through different methods, they were always polydisperse NPs with narrow or broad size distributions [3]. Therefore, the assumption ($\sigma = 0$) appears not to be realized. In fact, in the synthesis of MNPs using seed mediated-growth or size sorting route, the common value of σ was reported in between 0.3–0.5 but sometime it can reduce to 0.15–0.2 [2,22,36,37]. The polydispersity characterized by the σ of MNPs can give rise to the reduction of SLP [20]. For γ - Fe_2O_3 MNPs, σ in the range of 0.08–0.15 could result in a decrease of SLP_{\max} to half of the monodisperse value [22,24]. Calculations were carried out for various anisotropy (K) and standard deviation (σ) of FO MNPs at $f = 100$ kHz, $H = 5.18$ kA/m, $\eta = 1$ mPa·s and the results of relative $\text{SLP}_{\max}(\sigma)/\text{SLP}_{\max}(\sigma = 0)$ are represented in Figure 4.

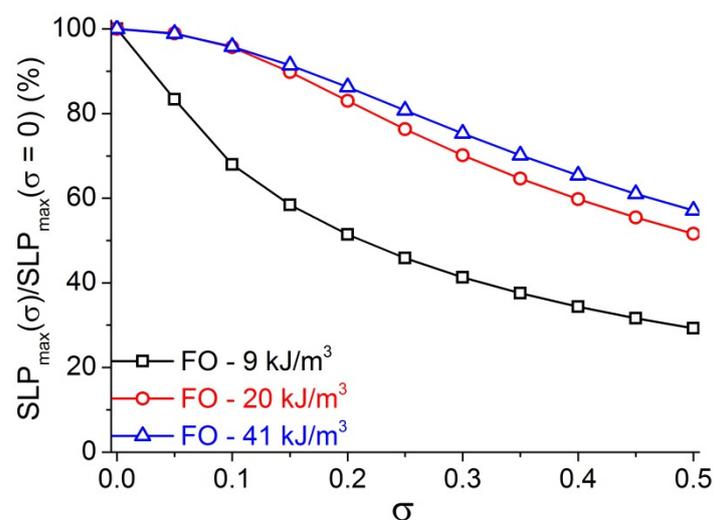


Figure 4. $\text{SLP}_{\max}(\sigma)/\text{SLP}_{\max}(\sigma = 0)$ calculated at $f = 100$ kHz, $H = 5.18$ kA/m, $\eta = 1$ mPa·s for FO samples with $K = 9, 20$, and 41 kJ/m³.

The change of heating power behavior with varying K has been reflected clearly also in the dependence of relative SLP on size standard variation. It can be seen that the domination region strongly influences the behavior of SLP besides material composition and particle size. Our results in Figure 4 indicate that, in all the curves, the SLP_{\max} decreases with increasing σ and there is a clear trend of increasing relative SLP_{\max} when the value of K changes from the NB-region (9 kJ/m³) to K_c (20 kJ/m³) and then the B-region (41 kJ/m³). In other words, for the similar polydispersity degree σ , the MNPs in the B-region would have higher heating power than that of those in the NB-region. These results indicated that the monodispersity requirement for getting the same heating power is much less strict for the MNPs in B-region as compared with the case of MNPs in N- or NB-region. Quantitatively, the polydispersity-caused SLP reduction for iso-dispersity of $\sigma = 0.2$ FO MNPs and σ_{50} (defined as at the position of $\frac{1}{2} \text{SLP}_{\max}$) with various anisotropies at $f = 100$ kHz, $H = 5.18$ kA/m, $\eta = 1$ mPa·s are listed in Table 3. The parameter σ_{50} is seen to increase from 0.21 to 0.58 with K from 9 to 32 kJ/m³, then becomes constant for $K \geq$

36 kJ/m³ indicating that the monodispersity requirement for obtaining the same heating power is much less strict for the high K as compared with the case of low K MNPs. The results in Table 3 also show that the polydispersity-caused SLP reduction decreased with an increase of magnetic anisotropy. From the slope of the SLP graph in Figure S2, it is estimated that an increase of K by about 4 times can result in saving as much as 35% of the heating power. These results seem to be consistent with the observation on cubic-shaped FO MNPs of about 20 nm (K = 180 kJ/m³) having higher heating power than spherical ones (K = 20 kJ/m³) [32]. The improvement of heating power by increasing K is an interesting topic and it has been recently reported for different materials, while the elevation of K was possible by various approaches including surface morphology, exchange-bias or stacking of particles in chains [32,33,36–39].

Table 3. Polydispersity-caused SLP reduction for iso-dispersity $\sigma = 0.2$ FO MNPs and σ_{50} , defined as at the position of $\frac{1}{2} SLP_{max}$, with various anisotropy K (f = 100 kHz, H = 5.18 kA/m, $\eta = 1$ mPa·s).

K (kJ/m ³)	SLP Reduction (%)	σ_{50}
9	48.6	0.21
12	38.8	0.3
16	25.3	0.43
20	17	0.52
24	15.6	0.54
28	14.8	0.57
32	14.3	0.58
36	13.9	0.59
40	13.9	0.59

The medium in which MNPs are dispersed also can affect the heating efficacy of MNPs. Based on Equations (3)–(7), Rosensweig’s calculations demonstrated a dependence of the heating rate $\Delta T/\Delta t$ on the medium viscosity at 3 different frequencies [20]. We performed systematic calculations of SLP_{max} in the viscosity range from 1 mPa·s to 10 mPa·s for MNPs with various K. The results of the relative $SLP_{max}(\eta)/SLP_{max}(\eta = 1 \text{ mPa}\cdot\text{s})$ for the fluid of FO MNPs are presented in Figure 5. As can be seen in Figure 5, for FO MNPs with the value of K in NB-region (9 kJ/m³), SLP_{max} decreases very little with increasing η . In contrast, there is a large decrease in the SLP_{max} when the value of η is going from 1 kJ/m³ to 4 mPa·s for 20 kJ/m³, and 41 kJ/m³ FO MNPs. It is noted that the value of critical anisotropy at 4 mPa·s for the field frequency of 100 kHz is 60 kJ/m³ (Table 2). Therefore, the decrease of SLP_{max} with increasing η might occur when the value of anisotropy is smaller than K_c . In other words, this behavior happens when the K of FO MNPs is in the B-region. These calculations indicated that the FO MNPs with the K in the N- or NB-regions are much better for MIH applications in highly viscous conditions. The observation of weak and strong decrease of the heating power with increasing viscosity, will be discussed as compared with experimental results in the following section.

3.4. A Comparison with Experimental Results

To compare the results from the experiments with the theoretical ones in a simple way, we have investigated SLP of the same MNPs but varying viscosity. Previously we reported the SAR measurements carried out at the field of 178 kHz and 5.18 kA/m (65 Oe) for bare CFO and MFO MNPs suspended in water fluid with added agar forming media with viscosity up to about 2 mPa·s, and observed strong and weak decrease of SAR with an increase of η , respectively for CFO and MFO, which is in agreement with the LRT prediction [40]. For comparing the results from the experiments with the theoretical ones, we studied two ferrofluids of 20 nm low-K MnFe₂O₄ and 15 nm high-K CoFe₂O₄ MNPs

coated with chitosan. They were prepared by co-precipitation. To create the magnetic fluids with different viscosity, the MNPs suspension was mixed with various agar concentrations. The value of viscosity extended up to about 8 mPa·s by adding more agar amounts. The fabrication of samples with various viscosities as well as the method of viscosity measurement were described earlier [40] as were the experimental details of the synthesis and the coating procedure by chitosan [41]. As discussed in our previous reports, chitosan has been widely used in biomedicine due to its biocompatibility and biodegradability [41]. Based on the measured magnetization curves shown in Figure S3, the thickness of chitosan coating layer was estimated to be around 4 nm. The value of M_S , H_C and K_{eff} for MFO and CFO NPs (shown details in Supplementary Section S3) are presented in Table 4.

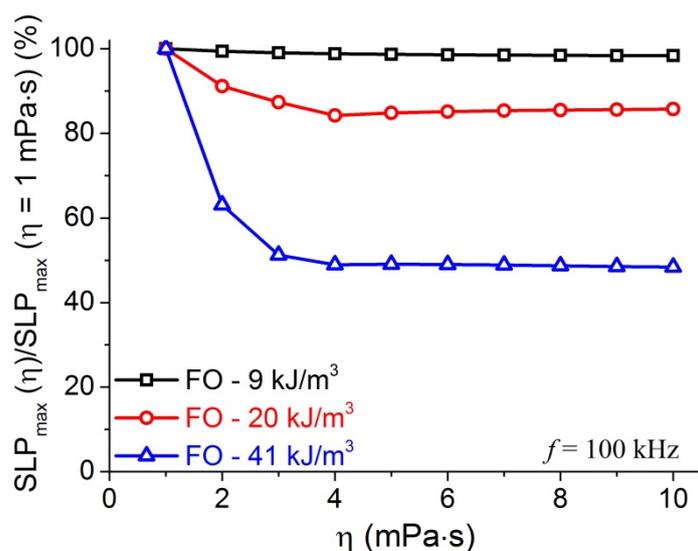


Figure 5. $SLP_{\text{max}}(\eta)/SLP_{\text{max}}(\eta = 1 \text{ mPa}\cdot\text{s})$ calculated at $f = 100 \text{ kHz}$, $H = 5.18 \text{ kA/m}$ for FO samples with $K = 9, 20$, and 41 kJ/m^3 .

Table 4. Values of K_c at different frequencies for fluids with various viscosities (η (mPa·s)).

Sample	M_S (emu/g)	H_C (Oe)	K_{eff} (kJ/m ³)
MFO (uncoated)	60.9	72	-
MFO (coated)	51.9	72	11
CFO (uncoated)	67.2	875	-
CFO (coated)	57.4	875	62

In addition to the field of 178 kHz, 5.18 kA/m (provided by a commercial generator RDO HFI 5 kW [40]), the calorimetric measurements were additionally conducted at two fields of the same amplitude of 15.9 kA/m (200 Oe) with frequencies of 340 kHz and 450 kHz (provided by a commercial generator UHF-20A [42]), to seek evidence of SAR behavior changing with increasing field frequency. The temperature elevation with time, $T(t)$, was monitored by using the same commercial optical sensor (GaAs-Opsens) for both MIH sets up.

The heating curves, measured at the lower field, for the two series of MFO and CFO samples in solutions of various agar concentrations are presented in Figure 6. The values of SAR_{exp} were calculated from the fitting of the temperature rising curves $T(t)$ for the whole time range following the method described in reference [40].

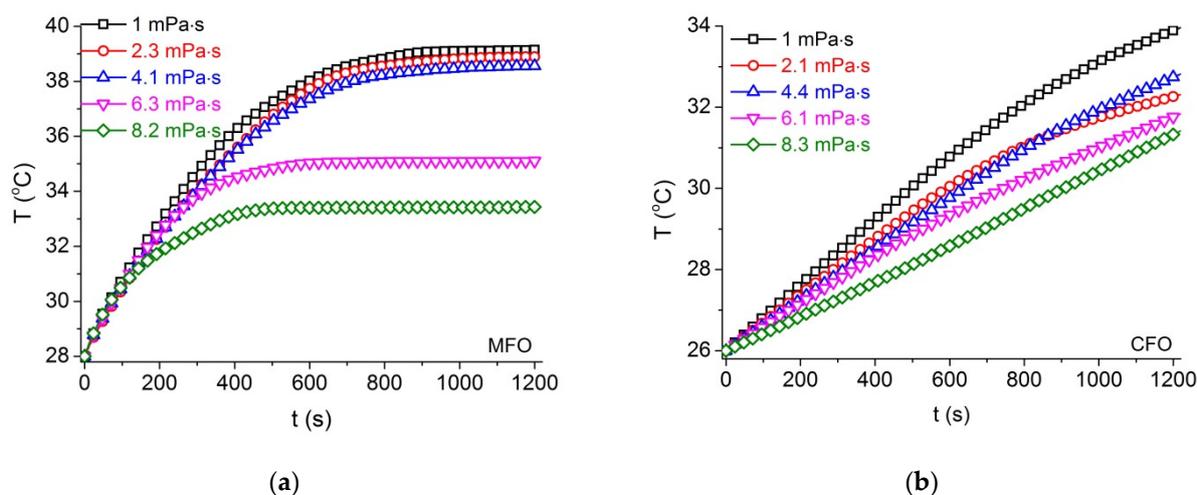


Figure 6. Hyperthermia curves measured at field frequency $f = 178$ kHz, $H = 5.18$ kA/m for (a) MFO and (b) CFO ferrofluids of various viscosities.

In order to compare the experimental and theoretical SAR, we calculated the Néel-Brownian relaxation experimental SAR_{exp}^{LRT} by using the Equation:

$$SAR_{exp}^{LRT} = SAR_{exp} - SAR_{exp}^{hys} \quad (12)$$

where SAR_{exp} is the SAR contribution by hysteresis mechanism, whose estimation is described in Supplementary Section S3, similarly to the approach reported in [42]. The results of calculated SAR_{exp}^{LRT} obtained from the measurements at 340 kHz and 450 kHz. They are gathered in Supplementary Section S4. As can be seen, for the MFO samples with $K = 11$ kJ/m³ and H_c were close to or smaller than the AFM amplitudes, the contribution of SAR_{exp}^{hys} is less than 2.5% of the whole heating power. However, for the CFO samples with $K = 62$ kJ/m³ and H_c much larger than the field amplitude used, the maximum SAR_{exp}^{hys} contributions were 17%, 16% and 9% for 178, 340 and 145 kHz, respectively.

In order to discuss about the impact of AFM frequency on the SAR characteristics relevant to the relaxation mechanisms, we calculated the ratio of $SAR(\eta)/SAR(\eta = 1 \text{ mPa}\cdot\text{s})$ for all the MFO and CFO samples of various viscosities measured at the 3 frequencies (see also Tables S1–S3), the dependences of which are shown in Figure 7.

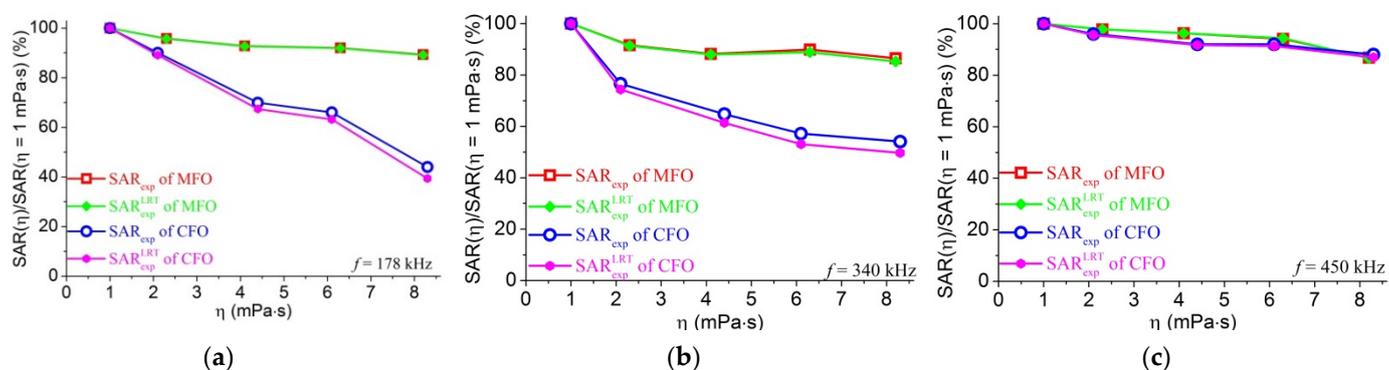


Figure 7. Relative SAR_{exp} and SAR_{exp}^{LRT} measured at field of (a) 5.18 kA/m, 178 kHz; (b) 15.9 kA/m, 340 kHz; and (c) 15.9 kA/m, 450 kHz.

Firstly, it can be noted that the decrease of SAR with increasing viscosity is experimentally observed in all the cases, so that the impact of the hysteresis contribution is insignificant.

Secondly, Figure 7a,b suggest that the strong and weak decrease of SAR with increasing viscosity, depicted respectively for CFO and MFO, confirm our previous experiment reported for $\eta \leq 2.12$ mPa·s [40] and the agreement with the LRT theory prediction shown in Figure 5. While such a behavior is natural for the case of the field (178 kHz, 65 Oe) with the parameter ξ of about 0.7, it is worth noting that such behavior ‘survived’ to the field of 340 kHz and 15.9 kA/m with ξ of about 2. The trend of decreasing SAR with increasing viscosity η in our experiments is also in agreement with the observation for γ -Fe₂O₃ (D = 7.1 nm, σ = 0.37) and CoFe₂O₄ (D = 9.7 nm, σ = 0.35) MNPs where SAR decreased, respectively, about 20% and 80% when the viscosity increased from 0.75 to 335 mPa·s [22]. In fact, it was found that SLP value from MNPs dispersed in water is higher than in glycerol for the same type of MNPs, MNP dose, and AMF condition [43]. A similar phenomenon was observed in the cellular environment as well, in which the measured SLP value decreased by half associated with attenuation of the Brownian relaxation in cellular conditions [44]. Thirdly, while the MFO ferrofluids remain almost not impacted by the field frequency used, it is interesting to observe for CFO samples a clear shift from high K toward low-K characteristic behavior when the used field frequency increases from 178 kHz to 450 kHz. In particular, Figure 7c for the case of $f = 450$ kHz indicates that the relative SAR versus η for the CFO MNPs approaches to the type of very weak decrease of the heating power with increasing the media viscosities, similarly to that of MFO MNPs. This observation is in good agreement with the scheme shown in Figure S7, according to which the MFO MNPs of $K = 11$ kJ/m³ are predicted to exhibit the SAR contributed mainly from the Néel or overlapped Néel-Brownian relaxation losses depending not on the used field frequency, and the dissipation mechanism contributing to the SAR for the CFO MNPs with $K = 62$ kJ/m³ belongs to the pure Brownian region for AMF with f below about 460 kHz and changes to the Néel-Brown or Néel type above this frequency.

4. Conclusions

In summary, in addition to the peak position D_c and its amplitude SLP_{max} , we have introduced a peak width ΔD_c and performed systematic analyses of the three parameters against the magnetic anisotropy of MNPs, fluid viscosities and the frequency of the alternating magnetic field. High K MNPs have been shown to exhibit domination of Brownian loss while it becomes that of Néel loss in low K ones. To the best of our knowledge, our results have demonstrated for the first time the transition from Néel to Brownian loss region not to occur in a continuous way, but at critical anisotropy K_c which increases with the frequency of the alternating magnetic field and the viscosity of the ferrofluid. We have pointed out that, for a given material, fabrication of MNPs with higher anisotropy up to K_c can improve the heating power as much as 35% thanks to the offsetting of the polydispersity-caused reduction. From the calculation and experimental results, it appears that low K MNPs are in general better for MIH applications in highly viscous conditions than those of high K ones.

Supplementary Materials: The following are available online at <https://www.mdpi.com/article/10.3390/ma14081875/s1>, Figure S1: (a) D_c versus K for various AMF frequencies. (b) The width ΔD_c , and D_c versus K at $f = 500$ kHz showing 3 characteristic N, NB and B regions; Figure S2: Polydispersity-caused SLP reduction calculated at $f = 100$ kHz for iso-dispersity $\sigma = 0.2$ FO MNPs as a function of anisotropy K, Figure S3: Magnetization curves measured for as-synthesized and chitosan coated (a) MFO, and (b) CFO MNPs; Figure S4: The initial magnetization curves of MFO and CFO MNPs. The solid lines represent the fitting curve assuming “the law of approach to saturation”; Figure S5: Hyperthermia curves measured at fields of frequency $f = 340$ kHz, $H = 15.9$ kA/m (200 Oe) for (a) MFO and (b) CFO ferrofluids of various viscosities; Figure S6: Hyperthermia curves measured at fields of frequency $f = 450$ kHz, $H = 15.9$ kA/m for (a) MFO and (b) CFO ferrofluids of various viscosities; Figure S7: Illustration scheme for the MIH experiments for CFO and MFO MNPs, Table S1: Values of SAR_{exp} , SAR_{exp}^{hys} , SAR_{exp}^{LRT} , and $\frac{SAR_{exp}^{LRT}}{SAR_{exp}^{LRT}(\eta=1 \text{ mPa}\cdot\text{s})}$ at 5.18 kA/m, 178 kHz; Table S2: Values

of SAR_{exp} , SAR_{exp}^{hys} , SAR_{exp}^{LRT} , and SLP^{LRT} at 15.9 kA/m, 340 kHz; Table S3: Values of SAR_{exp} , SAR_{exp}^{hys} , SAR_{exp}^{LRT} , and SLP^{LRT} at 15.9 kA/m, 450 kHz.

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