

1 The role of chemical design on the performance of organic semiconductors Hugo Bronstein¹,
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9 **Introduction**

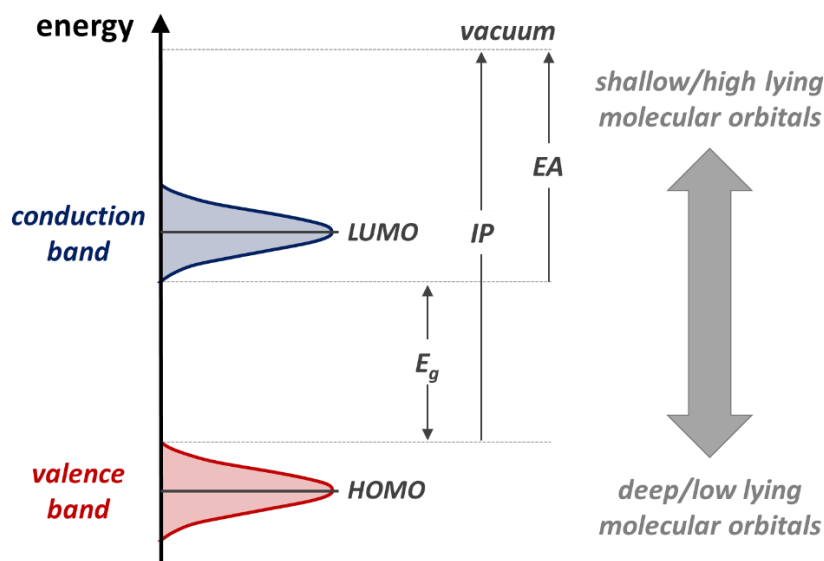
10 Organic semiconductors are increasingly utilised as the active layer in a wide range of new
11 technologies, arising from their solution processability, light-weight and flexibility. The versatility of
12 synthetic organic chemistry drives the development of applications benefiting from such properties,
13 including biological sensing, wearable electronics, semi-transparent photovoltaics and flexible
14 displays. Advances in these applications are underpinned not only by developments in chemistry but
15 also by a deeper understanding of the factors that govern the properties of the material through the
16 continuous improvement of analytical and computational techniques. Judicious molecular design
17 provides control of the semiconductor frontier molecular orbital energy distribution and guides the
18 hierarchical assembly of organic semiconductors into functional films where we can control the
19 properties and motion of charges and excited states. This review aims to illustrate how chemistry and
20 molecular design play integral roles in advancing electronic device properties in the main and
21 emerging applications of organic semiconductors.

22

23 **Molecular Orbital Design Considerations**

24 The frontier molecular orbital energy levels and their distribution along the conjugated pi electron
25 skeleton of the molecule (both small molecule and polymer) play a critical role in intra and
26 intermolecular charge transport, light absorption, light emission, charge injection/extraction, charge
27 trapping, and ambient electrochemistry. The highest occupied molecular orbital (HOMO) energy level
28 is largely dependent on the electron density and delocalisation of the conjugated π -electron system.
29 Functional groups which can directly donate electron density either mesomerically (i.e. lone pair
30 donation from nitrogen, oxygen, sulfur) or inductively (i.e. alkyl chains) can contribute to raising the
31 HOMO energy level, i.e. decreasing the solid-state ionisation potential. Conversely electron
32 withdrawing groups, such as fluorine, carbonyl groups or cyano groups, can act to both lower the
33 HOMO energy level, whilst also stabilising the lowest unoccupied molecular orbital (LUMO) energy
34 level, corresponding to an increase in the solid-state electron affinity. When considering
35 intermolecular electronic coupling, the spatial distribution of the molecular orbital density is essential,
36 in order to maximise π -orbital overlap along the conjugated system. A practical illustration of the
37 manipulation of molecular orbital energy levels is shown in the Tutorial Box, with the example of the
38 systematic functionalisation of a fluorene benzothiadiazole rhodanine conjugated molecule.

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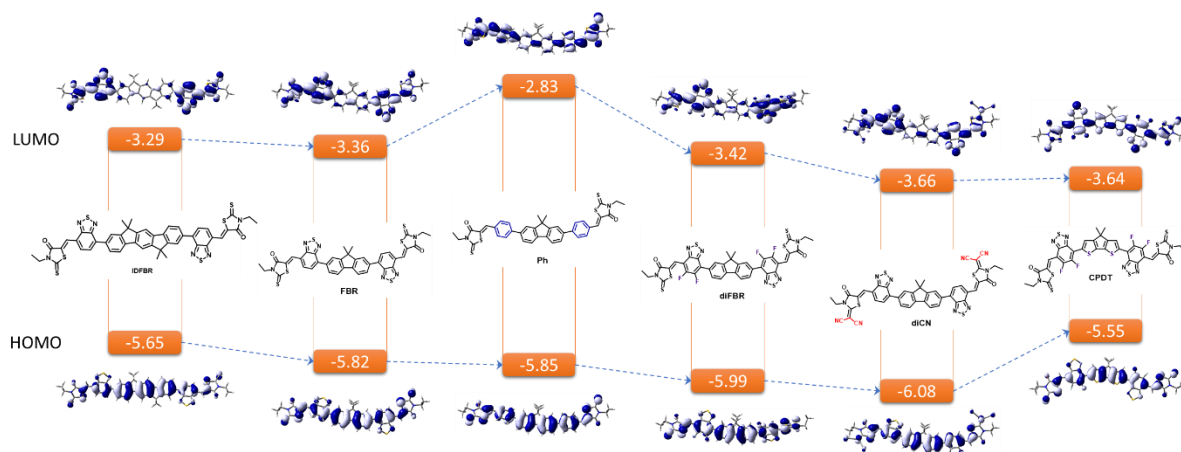
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41 **TUTORIAL BOX 1.** Definition of terminology used to describe molecular orbital energy levels.

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43 The figure of Tutorial Box 1. Illustrates the energies relevant to the molecular orbital design of organic
 44 semiconductors. The discrete HOMO and LUMO molecular energy levels observed in isolated organic
 45 semiconductor molecules, are broadened into electronic bands, the valence and conduction band.
 46 The optical band gap (E_g) is defined as the energy gap between the highest energy of the valence band
 47 and the lowest energy of the conduction band. The ionisation potential (IP) is the energy needed to
 48 remove one electron from the top of the valence band, whereas the electron affinity (EA) is defined
 49 as the energy from the bottom of the conduction band to the vacuum level.

50



51

52 **TUTORIAL BOX 2.** Case study in frontier energy level manipulation through molecular design.

53 [FOOTNOTE: Frontier energy levels (all in eV) and molecular orbital distributions have been estimated
 54 using density functional theory. The ground state geometries were obtained by B3LYP density
 55 functional method with basis set def2-SVP. The dispersion correction was conducted by Grimme's D3
 56 version with the BJ damping function. The excited states of target molecules on ground state
 57 structures were calculated by TD-DFT with the same level. The systematic name of FBR is 5,5'-((9,9-

58 dimethyl-9*H*-fluorene-2,7-diyl)bis(benzo[*c*][1,2,5]thiadiazole-7,4-diyl))bis(methanylylidene))bis(3-
59 ethyl-2-thioxothiazolidin-4-one)]

60 TEXT FOR TUTORIAL BOX: As illustrated on the left of tutorial box 2, the HOMO energy level of the
61 non-fullerene acceptor IDFBR ((5*Z*,5'*Z*)-5,5'-(((6,6,12,12-tetramethyl-6,12-dihydroindeno[1,2-
62 *b*]fluorene-2,8-diyl)bis(benzo[*c*][1,2,5]thiadiazole-7,4-diyl))bis(methaneylylidene))bis(3-ethyl-2-
63 thioxothiazolidin-4-one)) is delocalised along the entire π -conjugated backbone, whereas the LUMO
64 is predominantly localised on the electron-deficient 2,1,3-benzothiadiazole (BT) and rhodanine
65 moieties.¹ Consequently, chemical modifications to modify the LUMO of the molecule need to be on
66 the periphery, while the HOMO is potentially tuned with chemical modifications on any conjugated
67 part of the molecule. Reducing the conjugation length therefore of the central indenofluorene core
68 to a fluorene unit (FBR) predominantly raises the HOMO energy level (by almost 0.2 eV) while there is
69 little shallowing of the LUMO (less than 0.1 eV). Substituting the BT moieties on FBR with phenylenes
70 (Ph) raises the LUMO significantly (~0.5 eV) due to the less electron-deficient nature of the phenylene
71 groups. The HOMO, on the other hand, is only marginally affected. Difluorination of the BT units (diFBR)
72 lowers both the HOMO and LUMO by 0.1 – 0.2 eV compared with FBR due to the electron withdrawing
73 inductive effects from the electronegative fluorine atoms. Introduction of the much stronger electron-
74 withdrawing dicyanovinyl moiety than the thiocarbonyl on the rhodanine (diCN) has a stronger effect,
75 lowering both HOMO and LUMO by approximately 0.3 eV relative to FBR. As well as electron density
76 considerations, the delocalisation of the pi orbital skeleton plays a significant role in the orbital energy
77 levels. FBR has a non-coplanar backbone conformation due to repulsive torsional interactions
78 between neighbouring hydrogens on the BT and fluorene units; these steric effects can be alleviated
79 with the replacement of the fluorene core with cyclopentadithiophene (CPDT), allowing more orbital
80 overlap and thus further raising the HOMO (in addition to the effect of the electron rich thiophene
81 rings) as well as lowering the LUMO. Further planarising S-F through electrostatic interactions are
82 possible between the lone pairs of electrons on the fluorine atoms of the diFBT, and the
83 electropositive sulfur atoms on the peripheral thiophene rings of the CPDT unit, with a consequent
84 additional lowering of the LUMO. An additional consequence of pi orbital delocalisation, is that the
85 LUMO molecular orbital now is now distributed more over the core of the molecule, and hence any
86 functionalisation of the core will affect the LUMO as well as the HOMO energy level. Although not
87 illustrated in the tutorial box, heteroatom substitution of the carbon bridging atom of the fluorene
88 unit (with silicon or germanium) surprisingly did not appear to influence either of the frontier
89 molecular orbital energies. Typically, incorporation of bridging heteroatoms promotes an increase in
90 the respective bond lengths of the central ring to accommodate the larger atoms. The resultant larger
91 C-C bond length linking the two aromatic rings, reduces the antibonding interactions that arise from
92 the HOMO orbital node between the rings, thus deepening the HOMO energy level. Similarly, there
93 was also not a noticeable effect of chalcogen substitution of the sulphur atoms on the CPDT unit. In
94 general, increasing the chalcogen atom size reduces ring aromaticity, through decreasing participation
95 of the chalcogen lone pair in the ring, as the chalcogen – carbon bond length increases. This enhances
96 the diene character of the ring, reducing aromatic stabilisation and thus increasing the HOMO energy
97 level, and a reduced bandgap².

98 Charge transport is sensitive to traps within the bulk,³ therefore p-type materials benefit from a small
99 ionisation potential, where filling of deep traps is thermodynamically unfavourable. Similarly, n-type
100 materials benefit from a large electron affinity, where again there are less accessible charge traps.
101 Injection of holes from an electrode into the HOMO of a semiconductor is more energetically facile
102 when the electrode work function is close to, or preferably larger than, the ionisation potential of the
103 semiconductor. This allows for ohmic contact, with low contact resistance at the electrode-
104 semiconductor interface. Correspondingly, for electron injection, the electron affinity should be as

105 large as possible. Achieving *in operando* stability of OS devices relies on the suppression of reactions
106 of both the neutral and charged semiconductor species.⁴ In order to prevent the most
107 thermodynamically favourable reactions (involving a combination of oxygen and water), it has been
108 predicted that a neutral p-type semiconductor should have an ionisation potential of greater than 4.9
109 eV.⁴ One such example of ambient electrochemistry is the reduction of oxygen in the presence of
110 water to form hydroxide anions from the transfer of electrons from shallow HOMO semiconductors.
111 However, in operation, i.e. the hole polaron formed from deep HOMO semiconductors has the
112 opposite trend in reactivity and can oxidise ambient water for example. The activation barriers of
113 these forward and reverse reactions fortunately lead to overpotentials which allow many organic
114 semiconductors to exhibit reasonable shelf and operational stabilities. Electron transport is
115 particularly affected by ambient reactions, and in the same way as p-type semiconductors, preventing
116 electrochemical oxidation reactions of excited electrons delocalized within the LUMO is essential. To
117 accomplish this, the LUMO energy level must be deep enough, i.e. the electron affinity must be large
118 enough, to prevent the formation of, for example, superoxide radical anion from reaction of excited
119 electrons with hydrated oxygen complexes⁵ (one of the most favourable electrochemical reactions),
120 or the reduction of water. These unwanted electrochemical processes can reduce charge transport,
121 but moreover also lead to further irreversible chemical reactions within the semiconducting structure.
122 Defining a precise electron affinity value that needs to be exceeded to prevent these redox reactions
123 requires consideration of the electrochemical overpotential of the reaction, and is also influenced by
124 morphology. There is literature evidence however, to suggest that an electron affinity of greater than
125 4 eV will suffice to suppress oxidation.⁶

126

127 **Organic Field Effect Transistors**

128 Charge transport in organic semiconducting polymers relies on a combination of intrachain polaron
129 conjugation, facilitated by the polymer backbone π -electron delocalisation, and intermolecular charge
130 hopping between adjacent chains, facilitated by thin film microstructure. A close packed hierarchical
131 assembly of polymer backbones, promoted by intermolecular interactions, is often the dominant
132 factor in the optimisation of charge carrier mobility. Many design motifs have been proposed to
133 facilitate intermolecular contacts. An early example was the micron scale, three dimensional ordering
134 achieved in thin films of poly[2,5-bis(3-alkylthiophen-2-yl)thieno[3,2-*b*]thiophene] (pBTTT), illustrated
135 in Figure 1a⁷. This morphology was described as ordered lamellar sheets comprised of π -stacked
136 conjugated pBTTT backbones, with registration in the out of plane direction directed by the side chain
137 interdigitation of vertically adjacent polymer chains. Interdigitation was made possible by the
138 conformationally tolerant, regular spacing of the side chains along the backbone, with optimal special
139 separation to ensure an ordered and close packed side chain density on interdigitation. This approach
140 was also subsequently exploited in a series of isoindigo polymers⁸, where the analogy with biological
141 interactions invoked the description “molecular docking” to describe the possible locking of alkyl
142 chains into sterically empty regions along the polymer chain. Indeed, the cumulative Van der Waals
143 interactions between “inert” aliphatic side chains are likely the dominant interaction that templates
144 morphology and short contacts. Even though the degree of interdigitation exhibited by pBTTT is
145 extremely unusual, a close packed and regular alkyl distribution between backbones can act to not
146 only create ordered morphology, but reduce the local free volume, where unwanted species such as
147 water can reside, leading to charge trapping and subsequently a suppression of carrier mobility.

148 More recently, it has been identified that efficient intramolecular charge transport can be optimally
149 facilitated when the polymer backbone exhibits as low intrinsic energetic disorder as possible, and is
150 resilient to torsional fluctuations between adjacent monomer units^{9 10}. This is typically achieved by

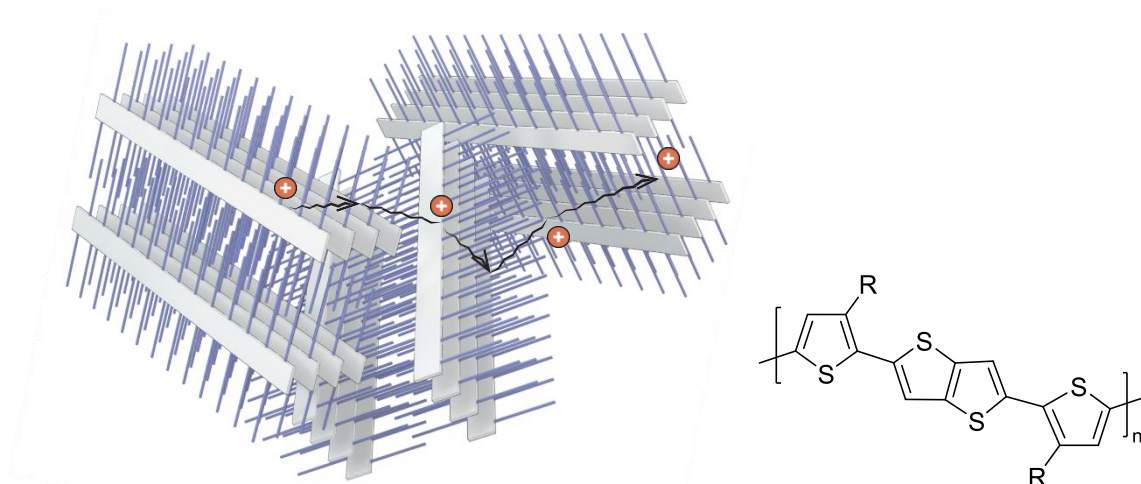
151 maximising the energetic barrier to rotation between aromatic repeat units along the backbone. There
152 are several molecular design strategies demonstrated to be successful in imparting co-planarity to a
153 conjugated backbone. The use of non-covalent through space interactions between adjacent ring
154 structures that can act as a barrier to rotation, thus promoting co-planarity. Common examples of this
155 approach include the electrostatic attractions between heteroatoms in close proximity to the linking
156 bond between aromatic units along the backbone. The sulphur atom on a linking thieno group, despite
157 the presence of two lone pairs of electrons, presents a partially positive charge due to the donation
158 of one lone pair to complete an aromatic ring. As a result, heteroatoms with available lone pairs of
159 electrons, such as oxygen and nitrogen can form strong attractive interactions with the sulphur, thus
160 rigidifying the linking bond.¹¹⁻¹³ This has specifically been exploited in the coupling of a thiophene ring
161 with a 3-alkoxythiothene unit. Bis-lactam based building blocks, like diketopyrrolopyrrole (DPP) and
162 isoindigo, on the other hand rely on more traditional hydrogen bonding-like interactions as evidenced
163 from their crystal structures.^{14,15} The oxygen in the electron deficient lactam core and the hydrogen
164 atoms in close proximity to introduce conformational locks along the polymer backbone thereby
165 reducing rotational disorder and maximising orbital overlap. (Figure 2) Non-covalent stabilising
166 approaches were also exploited to design highly coplanar copolymers with low energetic disorder such
167 as the copolymer indacenodithiophene-co-benzothiadiazole (IDT-BT)⁹, illustrated in Figure 1b. The
168 dominant planarising interaction in this case is a non-traditional hydrogen bond between the nitrogen
169 of the benzothiadiazole and the alpha hydrogen of the adjacent indacenodithiophene group. These
170 strong interactions overcome the repulsive steric exchange energies, giving rise to an extremely low
171 energetic disorder in this polymer. Despite a lack of long-range order, the assembly of side chains
172 promotes intermittent short contacts along the backbone. Transport is primarily one dimensional
173 along the backbone, with occasional hops between chains. As a result, these “short contact” polymers
174 exhibit very high charge carrier mobilities.

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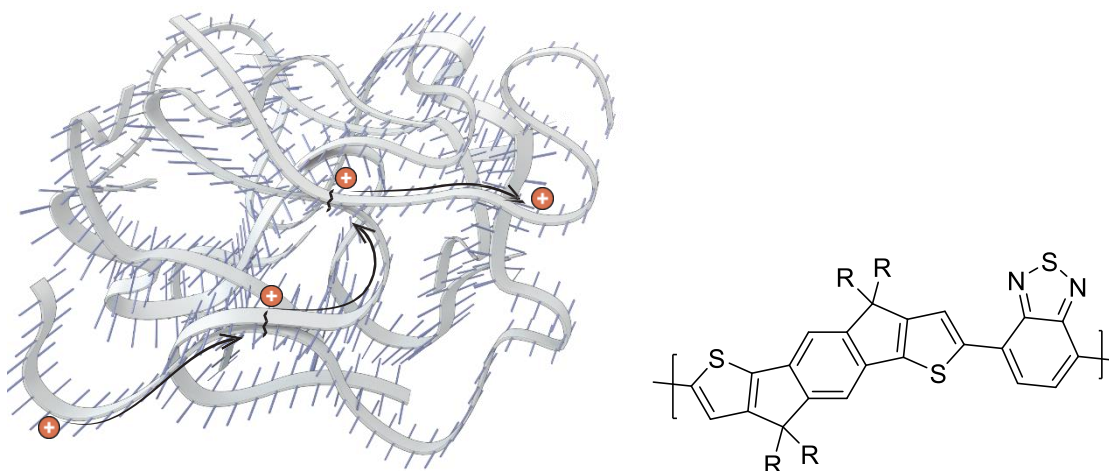
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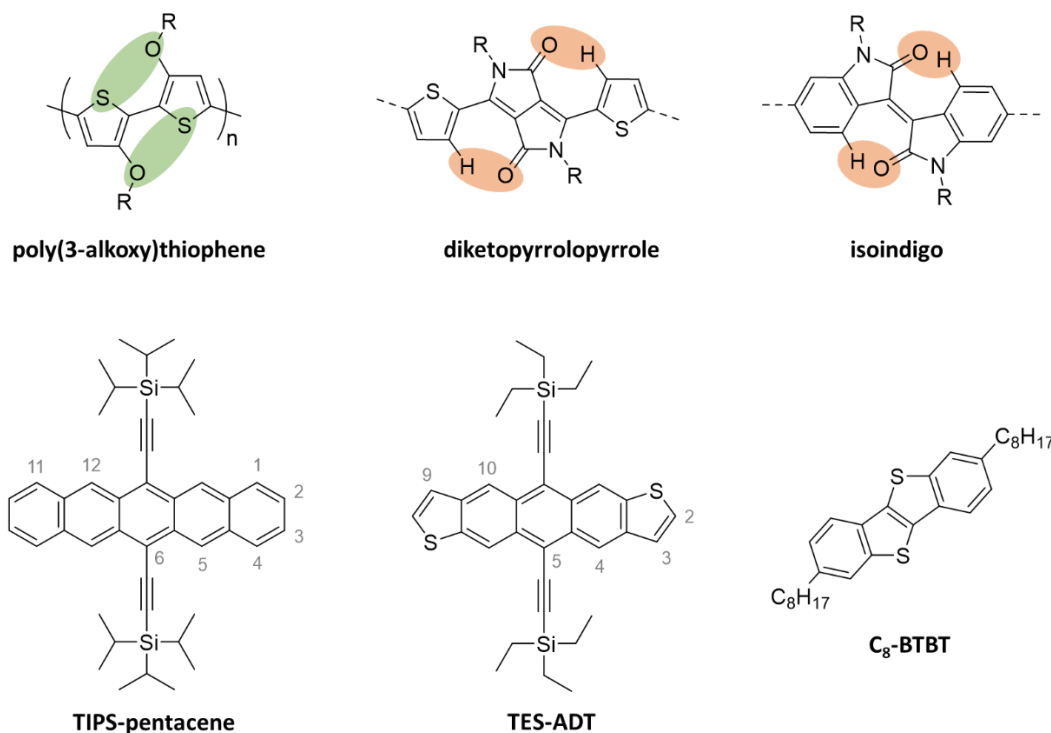
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182 Figure 1. Illustration of the three-dimensional packing arrangement (left) of pBTTT , where R is
183 typically a 14, 16 or 18 carbon linear alkyl chain and the optimal transport direction across and
184 between domains, and (right) the intermittent short contact packing arrangement of IDT-BT where R
185 is typically a 16 carbon linear alkyl chain and the optimal transport direction along the chains.

186

187 Small molecule semiconductor thin film transport properties have been notoriously difficult to
188 optimise through rational design at the molecular level. Crystal packing motifs, which are extremely
189 sensitive to even the smallest molecular perturbations, play a dominant role, with both a herringbone
190 and 2D slip stack arrangement being shown to have the required combination of electronic coupling
191 and sufficient isotropy for charge propagation. Two key building blocks have emerged as exemplary
192 aromatic cores, namely pentacene (and its analogue anthradithiophene (ADT)) and
193 benzothienobenzothiophene (BTBT). The packing motifs of these semiconducting molecules have
194 been manipulated by the addition of specific groups either on the central core (in the case of
195 pentacene and ADT) or periphery (in the case of BTBT). Both pentacene and ADT can be readily
196 centrally functionalised, on the 6 and 13 positions of the ring for pentacene¹⁶ (most successfully with
197 the triisopropylsilylethynyl (TIPS) group), and on the 5 and 11 positions for ADT (with
198 triethylsilylethynyl (TES)). (Figure 1) This functionalisation serves three principal functions. It renders
199 the molecules soluble for solution processing. It suppresses a potential oxidation site for degradation.
200 However, most importantly, it acts as a contributing molecular template to promote a 2D

201 interdigitated π -stack packing motif which dramatically enhances thin film charge carrier mobilities in
 202 comparison to most other packing motifs. In contrast, replacing the isopropyl groups on the silicon
 203 atom with slightly less bulky ethyl groups (TES pentacene) dramatically changes the crystal structure,
 204 templating a lamellar 1D slip stack with a correspondingly much larger anisotropy of charge transport,
 205 which in the length-scale of a transistor device, leads to much lower charge carrier mobilities. In the
 206 case of TES-ADT, the inclusion of a peripheral fluorine atom can also promote non-covalent short
 207 contacts, further enhancing electronic coupling and subsequently carrier mobility. BTBT was identified
 208 as a suitable aromatic core for charge transport and has been most optimally substituted in the 2 and
 209 7 positions with linear octyl chains ¹⁷. This both increases solubility, but critically, with a highly
 210 calamitic shape, induces a herringbone packing motif. The short contacts between the conjugated
 211 cores of adjacent molecules in the crystal cell give rise to excellent, and relatively isotropic, electronic
 212 coupling, thus enhancing charge carrier mobility. Historically, the charge carrier mobility of small
 213 molecules has been larger than polymers in thin films, typically attributed to a lower energetic
 214 disorder, narrower density of states, and less defects. However, semiconducting polymers with short
 215 intermolecular contacts are catching up, and generally exhibit less anisotropy of charge transport, a
 216 feature essential for integrated circuitry.



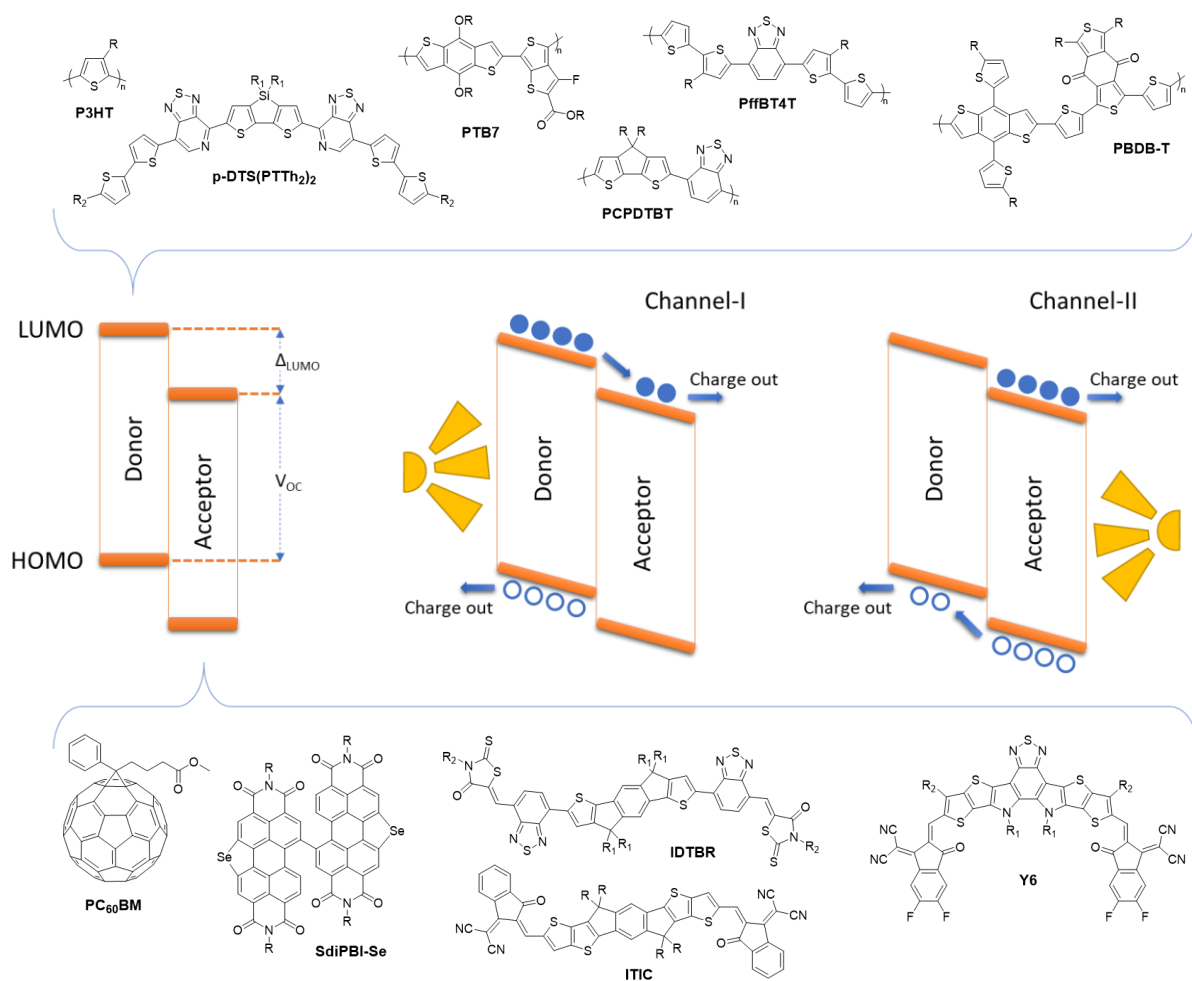
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 218 Figure 2. (top) Chemical structures of poly(3-alkoxy)thiophene , diketopyrrolopyrrole and isoindigo,
 219 with attractive planarizing interactions highlighted in green and orange. (bottom) chemical
 220 structures of highperforming molecular semiconductors TIPS-pentacene, TES-ADT and C₈-BTBT.

221
 222 **Organic Photovoltaics**

223 In an OPV device, the photoactive layer consists of an electron donor and an electron acceptor
 224 material which can both be molecularly tailored to absorb a large proportion of the incoming solar
 225 spectrum. An appropriate energetic offset between the donor and acceptor subsequently drives the
 226 splitting of excitons to free charge carriers which can be collected at the electrodes. To absorb a large
 227 proportion of the incident photons, an active layer significantly thicker than the exciton diffusion

228 length is required and this necessitates an intermixed donor:acceptor bulk heterojunction morphology
 229 (BHJ).¹⁸⁻²⁰ Optimisation of the frontier molecular energy levels of both donor and acceptor such that
 230 the majority of incident light is absorbed, whilst retaining sufficient energetic driving force for charge
 231 separation and a high open circuit voltage (V_{oc}) is required. Additionally, the donor:acceptor
 232 heterojunction must be controlled judiciously throughout the active layer to ensure efficient charge
 233 separation and collection with a high degree of morphological stability.

234 In the extensively studied P3HT:PCBM heterojunction the LUMO energy levels are \sim -3.2 and -4.2 eV
 235 for P3HT and PCBM respectively, providing a large energetic offset to overcome the P3HT exciton
 236 binding energy (typically 0.3 eV) and drive charge separation following a Channel-I mechanism
 237 illustrated in Figure 3. Combined with the fact that P3HT absorbs strongly in the visible spectrum and
 238 that the donor:acceptor BHJ comprises a favourable intermixed morphology, OPV devices with power
 239 conversion efficiencies over 4% have been achieved.²¹



240
 241 Figure 3. Energy diagrams (middle) outlining donor-acceptor energy alignment and photocurrent
 242 generation, with evolution of donor (top) and acceptor (bottom) materials design.

243 In order to harvest a greater proportion of the solar spectrum, OPV research rapidly shifted from P3HT
 244 towards molecular orbital hybridised donor materials by conjugating electron rich and electron poor
 245 repeat units. This powerful and highly modular tool to narrow the HOMO-LUMO gap has enabled the
 246 development of vast numbers of “push-pull” type polymers with hybridised molecular orbitals. While
 247 a narrower band gap donor material will absorb a greater proportion of the incoming photons and
 248 thus increase the extracted current (J_{sc}), it can compromise either the V_{oc} (shallower donor HOMO

249 leading to smaller LUMO(acceptor)-HOMO(donor) gap) or the charge separation (deeper donor LUMO
250 leading to smaller LUMO(donor)-LUMO(acceptor) offset) according to the energy diagram in Figure 3.
251 Computational modelling suggests that this competition between a high voltage and efficient light
252 absorption is best negotiated in a PC₆₀BM-based device using a donor material with a band gap around
253 1.5 eV.^{22,23} One such material is PCPDTBT; despite having a near ideal band gap, it performed poorly
254 in bulk heterojunction devices. However, in 2007, the efficiency of a PCPDTBT:PC₇₀BM device was
255 nearly doubled from 2.8 to 5.5% using a small amount of molecular additive (1,8-octanedithiol) when
256 casting the active layer.²⁴ Being able to optimise the bulk heterojunction morphology in this manner,
257 by preferentially improving the solubility of one of the photoactive components during film casting,
258 has subsequently allowed the field to more reliably assess new photoactive materials without
259 performance being overshadowed by poor morphology. Focusing still on push-pull type narrow band
260 gap donor polymers in conjunction with fullerene-based acceptors, OPV device efficiencies increased
261 steadily to around 11%, often exploiting the non-covalent interactions discussed earlier, to promote
262 polymer aggregation and facilitate phase separation at the exciton diffusion lengthscale.²⁵ From a
263 frontier energy level perspective, the HOMO level of these polymers, such as PTB7 (-5.15 eV), PTB7-
264 Th (-5.24 eV) and PffBT4T (-5.34 eV), have been lowered to improve the V_{oc} , while the band gap has
265 been increased slightly to ~1.6 eV. It should be noted that most high-performing PCBM devices are
266 prepared with PC₇₀BM; its asymmetry allows it to absorb visible light more strongly than PC₆₀BM, thus
267 allowing also for more Channel-II photocurrent generation (Figure 3).²⁶

268 Polymer donors have dominated the OPV landscape, but there are nevertheless some notable small
269 molecule donors. The pyridalthiadiazole-based small molecule p-DTS(PTTh₂)₂ (Figure 3) with frontier
270 energy levels of -3.6 eV (LUMO) and -5.2 eV (HOMO) absorbs broadly from 500 to 800 nm and has a
271 large LUMO-LUMO offset when paired with PC₇₀BM.²⁷ Controlling the blend morphology in a two
272 molecule blend is particularly challenging, which is highlighted by a remarkable sensitivity to
273 processing conditions for the p-DTS(PTTh₂)₂:PC₇₀BM devices; for a 70:30 weight ratio blend with 0.25%
274 solvent additive, a power conversion efficiency of 6.7% is achieved while small deviations from these
275 conditions sees the performance drop precipitously.

276 The V_{oc} loss is significant for Channel-I dominated fullerene-based OPV devices; defined as the
277 difference between the optical energy gap and the obtained open-circuit voltage, V_{oc} losses in high-
278 performing PCBM cells are typically on the order of 0.8 eV and thus significantly larger than the exciton
279 binding energy.²⁸ The development of non-fullerene acceptors (NFAs) have proven to address this
280 crucial barrier for further improving OPV device efficiencies. While it is relatively straightforward from
281 a molecular energy level perspective to design and synthesise NFAs with appropriate HOMO and
282 LUMO energy levels, early work highlighted the difficulties achieving both an intermixed blend
283 morphology and sufficiently high charge carrier mobilities. The high crystallinity of rylene diimides has
284 most successfully been mitigated by twisted dimer-type acceptor structures, for example illustrated
285 by the selenium-annulated PDI motif used in SdiPBI-Se to afford a PCE of 8.4%, and four-ring propeller-
286 type structures such as FTTB-PDI4 giving rise to a high PCE of 10.6% with a low V_{oc} loss of 0.53 V.^{29,30}
287 These examples also illustrate the advantage of NFAs over fullerenes in terms of controlling the
288 frontier energy levels with relative ease; SdiPBI-Se and FTTB-PDI4 have LUMO energy levels of -3.9
289 and -3.6 eV and optical band gaps of 2.2 and 1.9 eV, respectively. Not only can the acceptor strength
290 be adjusted to match the chosen donor material, the optical band gap can likewise be tuned to ensure
291 complementary absorption with the donor. The control over frontier energy levels is further
292 emphasised by the fact that NFAs often have their HOMOs and LUMOs spatially separated meaning
293 that they can be adjusted independently of each other through molecular design (see tutorial box for
294 more detail). Additionally, NFAs typically show markedly enhanced light absorption compared to

295 fullerene acceptors; coupled with efficient hole transfer from acceptor to donor, this provides high
296 Channel-II photocurrent generation and high J_{sc} values.

297 A further leap for NFAs has been manifested through the recent development of calamitic π -
298 conjugated molecules characterised by a central fused electron-rich (hetero)aromatic core flanked by
299 terminal electron-deficient units such as indanedione or rhodanine; further modularity in the
300 molecular design can be incorporated by inserting electron-rich or deficient π -conjugated spacers
301 between the core and the flanking groups. IDTBR, with an indacenodithiophene core flanked on each
302 side by 2,1,3-benzothiadiazole and rhodanine, is one such calamitic acceptor.³¹ Its co-planar backbone
303 results in a relatively narrow band gap (1.6 eV); a P3HT:IDTBR active layer consequently absorbs
304 strongly over the entire visible spectrum and affords device efficiencies above 6%. In addition to
305 pushing the OPV performance for a simple and commercially scalable polymer such as P3HT, this work
306 also illustrates that the NFAs can be used as the primary low band gap light absorber, a role which was
307 hitherto reserved for the donor material. Further improvements towards 8% power conversion
308 efficiency with IDTBR has been reported for ternary blends using two complementary NFAs in
309 conjunction with P3HT.³² IDTBR has a LUMO energy level of -3.9 eV, making it compatible in terms of
310 LUMO-LUMO offset with polymers such as PffBT4T; this is clearly demonstrated with a power
311 conversion efficiency of 10% and a very small V_{oc} loss around 0.5 V for a PffBT4T:IDTBR based device.³³
312 ITIC is another calamitic acceptor with comparable frontier energy levels to those of IDTBR.³⁴ It has
313 likewise been used in several high efficiency OPV devices, but similarly to IDTBR it has an optical band
314 gap around 1.6 eV which is comparable to most high-performing donor polymers developed
315 specifically for fullerene-based devices. As such, to cover more of the solar spectrum and increase J_{sc} ,
316 it has proven beneficial to widen the band gap of the donor polymer to prevent a large spectral overlap
317 with the NFA, and possibly provide a pathway for Forster Energy transfer from the donor to acceptor.
318 PBDB-T, with a HOMO level of -5.3 eV and an optical band gap around 1.8 eV, comprises the weaker
319 acceptor unit benzodithiophenedione as the electron-deficient unit in conjunction with the well-
320 known benzodithiophene donor motif. Devices with a PBDB-T:ITIC active layer have gradually been
321 improved to over 14% device efficiency.^{35,36} PBDB-T has been modified by substitution of alkyl chains
322 with thioalkyls, by fluorination and chlorination, while ITIC similarly has been fluorinated and
323 chlorinated. Halogenation of the ITIC electron acceptor increases its electron withdrawing potential,
324 and stabilises both the HOMO and LUMO, but more so the LUMO, as there is a higher LUMO orbital
325 distribution on the ITIC unit, which results in a slightly narrower band gap. The red-shifted absorption
326 increases the J_{sc} while enhanced intra- and intermolecular interactions concurrently improve the
327 extinction coefficient and charge transport properties. The HOMO energy level of dichlorinated PBDB-
328 T-2Cl is lowered 0.2 eV relative to its non-chlorinated analogue PBDB-T; this maintains a high V_{oc} in
329 OPV devices with halogenated ITIC derivatives, while the stabilised frontier energy levels likewise
330 improve the device oxidative stability.

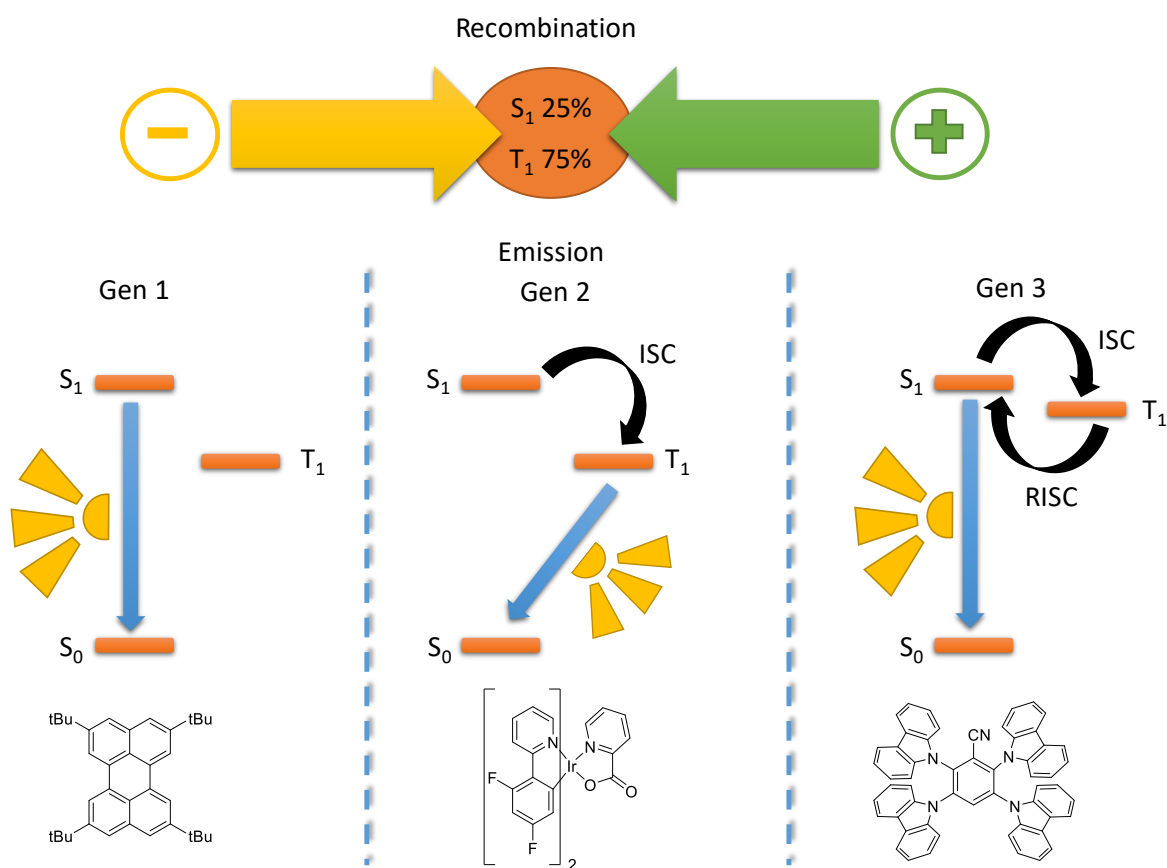
331 Most recently, attention has broadened from linear calamitic motifs to a curved structure in the form
332 of Y6.³⁷ Introducing as the central π -conjugated chromophore a 2,1,3-benzothiadiazole flanked by
333 thieno[3,2-*b*]thiophene units that are fused onto the benzothiadiazole with alkyl-bearing nitrogen
334 atoms, a narrow band gap NFA absorbing strongly beyond 900 nm is obtained. The two *N*-alkyl side
335 chains facing each other in Y6 twist the molecule slightly (~17 degrees), ensuring good solubility and
336 favourable control of aggregation. Paired with a fluorinated derivative of PBDB-T devices obtained
337 high power conversion efficiencies above 15% with V_{oc} values above 0.8 V. The small V_{oc} loss for this
338 system is further improved using a chlorinated analogue of Y6, giving a device efficiency above 16%
339 and a V_{oc} loss of only 0.53 V with a small non-radiative energy loss of 0.2 V.³⁸ It is hypothesised that
340 the much larger dipole moments in these bent NFAs compared to ITIC-type NFAs, and high electron
341 mobility, contribute to the efficient charge separation and high fill factors around 75%.

342 While it is clear that molecule structure and frontier energy level design has driven the field of OPV
343 forward continuously over the last decade, related design approaches are likewise being explored to
344 sensitise other PV applications. This is most prominently exemplified with the development of singlet
345 fission materials that convert one high-energy singlet exciton to two lower energy triplet excitons,
346 which can in turn generate additional electron-hole pairs in, for example, low bandgap silicon-based
347 PV devices. With an S_1 singlet exciton energy of 2.4 eV, tetracene can absorb high-energy photons that
348 would otherwise generate high thermalisation losses in silicon cells and subsequently through the
349 process of singlet fission convert this singlet exciton to two triplet excitons. The T_1 triplet exciton
350 energy of tetracene around 1.25 eV matches nicely the 1.1 eV bandgap of silicon. Although it is proving
351 challenging to effectively transfer the triplet excitons to the silicon cell, this technology holds promise
352 to exceed the Shockley-Queisser limit of 29% efficiency for a single junction PV device.

353

354 **Organic Light Emitting Diodes**

355 Despite organic light emitting diodes (OLEDs) being at an advanced stage of commercialisation, it
356 remains an extremely active area of research with many fascinating developments over the last 5-10
357 years. The fundamental operating principle of an OLED is extremely simple; charges injected from
358 electrodes recombine within an organic layer to form an excited state which decays radiatively.
359 However, spin statistics dictates that non-geminate recombination of charges results in a 1:3 ratio of
360 singlet and triplet excited states. The vast majority of research since the birth of the OLED has been to
361 overcome the issue associated with the non-emissive nature of triplet excited states which limits a
362 maximum of 25% internal quantum efficiency (IQE). The first successful approach to overcome this
363 was to move away from fluorescence (Gen 1) to phosphorescence (Gen 2) by exploiting the efficient
364 heavy-metal induced intersystem crossing and phosphorescence of organo-platinum and iridium
365 complexes (Figure 4). This allowed all of the electrically generated excited states to be used radiatively
366 and 100% internal quantum efficiency devices were demonstrated (with a maximum of 20-25%
367 external quantum efficiency due to outcoupling losses). However, issues associated with both colour
368 purity, and stability particularly in the blue region of the electromagnetic spectrum motivated
369 researchers to continue developing new materials for OLEDs. The 3rd generation of OLED materials
370 relies on the principle of reverse intersystem crossing (rISC) to utilise the dark triplet states. By
371 reducing the energy gap between the lowest excited singlet and triplet states, rISC can take place
372 whereby triplets can be converted into emissive singlets using thermal energy, i.e. thermally activated
373 delayed fluorescence (TADF).



374

375 **Figure 4: Evolution of OLED Emitters**

376 The most commonly used phosphorescent emitters are cyclometallated iridium (III) complexes which
 377 are dispersed within an organic host. It has long been known that the emitting dipole orientation of a
 378 dye influences the outcoupling efficiency of OLEDs.³⁹ Despite this, the dipole orientation of
 379 phosphorescent dopants has not been studied extensively, presumably due to the lack of any obvious
 380 driving force for alignment (in typically amorphous host materials). It has recently been shown using
 381 angle-dependent polarized photoluminescent emission spectroscopy that certain iridium complexes
 382 can be aligned in films.^{40,41} They suggested that the inherent asymmetry at the surface of a growing
 383 film promotes dopant alignment in these otherwise amorphous films. Alongside other investigations
 384 to investigate and control phosphor alignment, external quantum efficiencies in excess of 30% have
 385 been reported with projected efficiencies of up to 60% believed possible.^{39,42} Despite decades of
 386 research, new strategies for achieving efficient blue emission from Iridium complexes are still being
 387 sought with very impressive results being demonstrated recently. Isomerisation of a UV emitting N-
 388 heterocyclic carbene based Ir(III) complex from *fac* to *mer* subtly red-shifted its emission such that it
 389 displayed efficient deep blue phosphorescence with ~80% PLQY and EQE ~ 15%.⁴³ There have also
 390 been considerable progress in the development non-iridium based phosphorescent emitters. Recently,
 391 a new class of cyclometalated Au(III) complexes with tuneable emission spanning sky-blue to red were
 392 reported. These materials exhibited high photoluminescence quantum yields (PLQY) of up to 80% in
 393 solid-state thin films and impressive performance in OLEDs with EQEs up to 22% reported with
 394 excellent stability.³⁷ An extremely active area of research is the development of near-IR OLED emitters
 395 for security and communication applications. A new class of Pt(II) complexes that emit light at 740 nm
 396 with a PLQY of 81% were recently reported. In conjugation with a highly preferred horizontal dipole
 397 orientation OLED devices exhibited an external quantum efficiency of 24%.⁴⁴ The origin of the near-
 398 IR emission was attributed to the stacking of Pt(II) complexes in the solid state.

399 Almost certainly the biggest development in the field of OLEDs over the last 10 years has been the
400 demonstration of TADF using purely organic dyes. This work was pioneered by the group of Adachi
401 with the seminal publication demonstrating a series of purely organic materials with high PLQY
402 tuneable emission and OLED EQEs approaching 20%.⁴⁵ Within this work was detailed the simple design
403 concept that by spatially separating a molecule's HOMO and LUMO through a twisted donor-acceptor
404 structure, the energy gap between the first excited singlet and triplet could be substantially reduced.
405 This allows triplet states to be converted into emissive singlets through reverse intersystem crossing
406 (rISC) resulting in the phenomenon of thermally activated delayed fluorescence. By tuning the number,
407 twist and electron density of donors to acceptors, extremely impressive PLQY and OLED EQEs have
408 been obtained spanning blue to near-IR.^{46,47,48,49,50} Despite these extremely impressive results, the
409 exact mechanism by which rISC and TADF occur remains the subject of debate. It was rapidly
410 established that in order to design efficient TADF emitters it is vital to consider both locally excited
411 (LE) and charge-transfer (CT) states.⁵¹ However, the process by which the spin-flip between singlet and
412 triplet states is not clear. It has been suggested that 'hidden' n-pi mixed states may play an important
413 role in the process or that second order vibronic coupling between the various excited states is
414 required.⁵²⁻⁵⁴ Very recently, the important role of intermediate electronic states that mediate the spin-
415 flip process through molecular vibrations has been demonstrated using a new model which also
416 complements the notion of non-adiabatic coupling of triplet states promoting ISC.⁵⁵ Therefore, the
417 synthetic strategy of merely minimising the S_1 - T_1 energy gap requires refinement. One of the issues
418 with conventional TADF is the requirement of highly twisted D-A structures, which results in broad
419 spectral emission with low oscillator strength. This can potentially be overcome using a
420 "hyperfluorescence" strategy whereby a fluorescent guest is introduced into a TADF host albeit at the
421 expense of device simplicity.⁵⁶ It has also been suggested that using a 'multi-resonant' strategy, which
422 spatially separates HOMO and LUMO within a single planar molecule, narrow ΔE_{ST} gap materials with
423 high oscillator strength and narrow emission can be obtained.^{57,58} Indeed, extremely impressive
424 results using this strategy has recently been reported with a blue (469 nm) organoboron TADF emitter
425 based OLED with an EQE of 34%.⁵⁹

426 TADF has also been observed in organometallic systems, with Cu(I)-Carbene complexes demonstrating
427 outstanding optoelectronics properties in particular. Photoluminescence quantum yields of 100%
428 have been reported and OLED device efficiencies of >25% EQE.^{60,61} Emission in these systems appears
429 to stem from inter-ligand charge transfer states and importantly it appears that (as has also been
430 suggested in organic TADF⁶²) that there is considerable spin-mixing occurring, further indicating that
431 the "spin-pure" picture of TADF emitters may be overly simplistic. However, similarly to
432 phosphorescent OLED emitters, achieving highly stable blue OLEDs with good colour purity have
433 proved extremely challenging.

434

435 **Beyond the state of the art**

436 A typical chemical design approach for semiconductor optimisation involves systematic variations in
437 the aromatic conjugated units comprising the core or backbone, undertaken in order to isolate and
438 study the molecular design effect on both morphology and frontier molecular orbital energy levels. As
439 a justification of the design concept, results are often compared with previous publications. In the
440 transistor field, unreliable charge carrier mobility values, extracted from invalid model equations and
441 assumptions, has both inflated expectations and led to incorrect structure-property relationships
442 being proposed, with molecular design conclusions reached on tenuous evaluations. This has seriously
443 impeded progress in new materials for transistor applications. However, there are clearly some

444 obvious opportunities for further exploration. The presence of aliphatic chains in these materials is
445 primarily to impart solubility and secondarily control solid-state structure. However, both alkyl chain
446 and lamellar crystallisation can come at the *expense* of conjugated backbone planarity and crystallinity.
447 It is imperative that new semiconductor design is done in a holistic fashion whereby the self-assembly
448 of the system works with, rather than against optimum performance. It may be possible to utilise side
449 chains in creative new ways such as the suppression of various undesirable phonon modes which have
450 recently been shown to severely affect charge transport.⁶³ Alternatively, entirely removing the alky
451 chains should fundamentally improve charge transport. In-situ side chain detachment methods to
452 increase backbone densification, without compromising microstructural order is a promising approach
453 to performance improvement, yet currently most methods remove the inert chains at the expense of
454 order and percolation, thus reducing transport properties. Molecular order however has proven to be
455 key in optimising charge transport and finding ways to further reduce the disorder in conjugated
456 systems or alternatively design materials with a higher tolerance towards rotational fluctuations, will
457 be essential to improve the charge transport further. To date this approach has proven challenging as
458 it is difficult to predict solid state packing of organic semiconductors and the different vibrational
459 modes associated with a particular unit cell. With improving computational capabilities however
460 material property predictions are getting ever more sophisticated and accurate, providing important
461 insight into material packing motifs and phonon coupling modes. This opens up the opportunity to not
462 only calculate the many phonon coupling modes in a conjugated system, but most importantly to
463 predict and identify the modes most detrimental to charge carrier transport.

464 Besides the electronic charge transport, mixed conduction, i.e. the ability of organic conductors and
465 semiconductors to facilitate both electronic and ionic charge transport in the solid state, is achieving
466 a growing amount of attention due to potential use for biological interfacing.^{64,65} In this context, the
467 organic electrochemical transistor (OECT), a three-terminal device that can be gated through an
468 aqueous electrolyte, is emerging as a powerful tool for bioelectronic applications. In an OECT, the bulk
469 active layer contributes to the current modulation through doping or dedoping. Development of n-
470 and p-type semiconducting polymers for OECTs therefore necessitate new design rules for frontier
471 energy level optimisation. For stable operation in aqueous environments, as a first prerequisite, the
472 HOMO level of p-type materials must be electrochemically accessible before water electrolysis takes
473 place at 1.23 V (versus the normal hydrogen electrode (NHE)). High performance mixed conductors
474 have much shallower HOMO levels around -4.4 to -4.6 eV depending on the measurement
475 technique.^{66,67} Shallow HOMO levels in that range are obtained with glycolated polythiophenes similar
476 to the poly(3-alkoxythiophene) depicted in Figure 1 where the oxygen atoms mesomerically donate
477 electron density to the polymer backbone while the polar nature of the oligoether sidechain facilitates
478 ion transport. The development of n-type semiconductors is essential for sensing of biologically
479 relevant metabolites and cations, such as sodium, potassium and calcium. To be operationally stable
480 in water, the main challenge centres around the stability of the radical anions towards water and/or
481 oxygen.⁶⁸ Here, LUMO levels deeper than -4.0 eV are required.^{69,70} Currently the performance of
482 electron transport OECT semiconductors significantly lags behind their hole transporting counterparts.
483 Although this has been tentatively attributed to coulombic charge pinning from counterions, charge
484 trapping from water, and other phenomena, new molecular designs are required to overcome these
485 limitations.

486 For photovoltaic applications on the other hand, the focus should not be on optimising the electronic
487 and solid state properties of a single material, but rather on understanding the design criteria to
488 control the morphology of the active layer, perhaps incorporating multiple materials, chosen to create
489 an energy cascade with large charge transfer (CT) state energies. While the frontier energy levels and
490 the isolated molecular packing of both donor and acceptor is generally well controlled, morphological

491 control through molecular design is still lacking, both on microscopic and macroscopic length scales.
492 Although morphological stability of the BHJ blend has been improved significantly by the introduction
493 of low diffusive NFAs, it would be highly advantageous if the tertiary structure could be controlled to
494 a greater extent for instance by choice of heteroatoms, functional groups and sidechains. Ideally, such
495 approaches should also provide a means to firstly generate much thicker photoactive layers without
496 unwanted vertical phase separation and excessive charge recombination and secondly scale up the
497 cell's active area more effortlessly. At much shorter length scales, being able to control the donor-
498 acceptor interfaces in the blend, for instance through molecular shape and dipole moments, would
499 allow for greater understanding and control of the energy and charge transfer processes. Another
500 important aspect that require further attention is the role of the CT state, how the CT state can be
501 manipulated through chemical design and how this can be utilised to suppress recombination
502 processes. Recent work indicates that electrostatic potential mapping for instance could be a useful
503 tool to help understand how molecular structure and intermolecular interactions influence the charge
504 separation processes.⁷¹ The recent reports of Y6, a curved NFA structure markedly different from the
505 previous best NFA performers in terms of molecular design, has to a large extent taken the OPV
506 community by surprise. While it promisingly shows that there is still ample room for device
507 performance improvements, it also underlines that a unified picture on molecular design criteria has
508 not yet been developed. The understanding of NFAs, their detailed structure-property relations and
509 the dominating reasons for the excellent performance continue to be explored. Factors such as off-
510 axis dipole, high electron mobility, facile energy transfer processes and the potential for good Forster
511 energy overlap with the donor could potentially play significant roles. As the community now embarks
512 on in-depth fundamental studies of Y6-based devices as well as the synthesis of Y6 analogues, further
513 pieces of this puzzle will emerge, which in turn will push the device performance closer towards the
514 20% efficiency mark and hopefully add to the understanding of molecular and frontier energy level
515 design criteria.

516 The commercialisation of OLED technology has already been a monumental success. However, a
517 detailed understanding of the exact photophysical processes occurring during charge recombination
518 and emission is required alongside a clear relationship between chemical structure and photophysical
519 properties for continuing progress to be achieved. It is clearly necessary to further understand the
520 origin of stability/instability of excited states in particular in the presence of charges. Identification of
521 degradation pathways may allow suitable chemical modification to prevent the degradation of these
522 materials. Despite this, current generations of OLED materials are operating close to their theoretical
523 maxima and for a step-change in performance, a new generation of materials must be developed.
524 Some encouraging signs of emerging 4th generation, which warrant exploration, are as follows. There
525 have been recent reports of extremely efficient OLED devices using radical based fluorophores.^{72,73}
526 These compounds, based on the well-known luminescent tris(2,4,6-trichlorophenyl)methyl (TTM)
527 radical emit from a doublet excited state entirely circumventing spin dependent recombination.
528 However, thus far these materials have only demonstrated electroluminescence in the red/near-IR
529 and it remains to be seen whether it is possible to tune them to cover the entire visible spectrum.
530 Another interesting new approach by Adachi *et al.* involves the use of host materials capable of
531 undergoing singlet fission which, in principle, could allow for IQEs of 200%.⁷⁴ Very low EQEs for this
532 device were observed due to the low PLQY of the emitting guest erbium (III) complex. In order for this
533 strategy to succeed it will be necessary to develop wide band-gap singlet fission hosts and/or highly
534 emissive near-IR emitters capable of harvesting triplet excitons. There have been several recent
535 reports of utilising triplet fusion for OLEDs whereby two triplets annihilate and produce an emissive
536 singlet.⁷⁵⁻⁷⁷ Although this limits IQE to a maximum of 50% (if only triplets are injected) it could
537 potentially allow for low driving voltages and may offer stability advantages for blue emitters.

538 There is also a need to improve the underpinning chemistry used to synthesise organic
539 semiconductors. The presence of (often hard to detect) residual impurities can result in batch-to-batch
540 variations or unintentional doping. The availability and compatibility of chemical building blocks
541 dictates the choice of chemistry employed in synthesis. Polymerisations often relies on palladium
542 catalysed cross-coupling methodology (i.e. Yamamoto, Suzuki, and particularly Stille) and while these
543 coupling reactions are extremely efficient, their atom economy and use of chlorinated solvents does
544 not comply with sustainability requirements, nor are they particularly suitable for large scale
545 production. C-H activation and condensation chemistry circumvents some of these problems, by
546 shortening the synthetic pathway and making the extensive functionalisation of the monomers
547 obsolete.⁷⁸ The narrow scope and imposed restrictions on employable building blocks however leave
548 ample room for improvement. Developing new chemical polymerisation protocols that allow for
549 control over molecular weight and polydispersity whilst reducing chemical defects (i.e. homocoupling,
550 partial oxidation) will be an integral part in advancing the quality of the synthesised materials.

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552

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