AVERAGES OF EIGENFUNCTIONS OVER HYPERSURFACES

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ABSTRACT. Let (M, g) be a compact, smooth, Riemannian manifold and $\{\phi_h\}$ an L^2-normalized sequence of Laplace eigenfunctions with defect measure μ. Let H be a smooth hypersurface with unit exterior normal ν. Our main result says that when μ is *not* concentrated conormally to H, the eigenfunction restrictions to H satisfy

$$\int_H \phi_h d\sigma_H = o(1) \qquad \text{and} \qquad \int_H h D_\nu \phi_h d\sigma_H = o(1),$$

$h \to 0^+$.

1. INTRODUCTION

On a compact Riemannian manifold (M, g), with no boundary, consider a sequence of Laplace eigenfunctions $\{\phi_h\}$,

$$-h^2 \Delta_g \phi_h = \phi_h,$$

normalized so that $\|\phi_h\|_{L^2(M)} = 1$. The goal of this article is to study the average oscillatory behavior of ϕ_h when restricted to a hypersurface $H \subset M$. Namely, the goal is to find a condition on the pair $(\{\phi_h\}, H)$ so that

$$\int_H \phi_h \, d\sigma_H = o(1), \tag{1}$$

as $h \to 0^+$, where σ_H denotes the hypersurface measure on H induced by the Riemannian structure.

It is important to point out that one cannot always expect to observe this oscillatory decay. For instance, on the round sphere, zonal harmonics of even degree integrate to a constant along the equator. Also, for any closed geodesic inside the square flat torus. there is a sequence of eigenfunctions that integrate to a non-zero constant.

Integrals of the form (1) have been studied for quite some time, going back to the work of Good [Goo83] and Hejhal [Hej82] that treated the case where H is a periodic geodesic inside a compact hyperbolic manifold. These authors proved that in such a case, $\int_H \phi_h \, d\sigma_H = O(1)$ as $h \to 0^+$. Zelditch [Zel92] generalized this to the case where H is any hypersurface inside a compact manifold, showing that for any hypersurface H,

$$\int_H \phi_h d\sigma_H = O(1). \tag{2}$$

In addition, it follows from [Zel92] that for a density one subsequence of eigenvalues $\{h_j\}_j$, one has $\lim_{j\to\infty} \int_H \phi_{h_j} \, d\sigma_H = 0$. Moreover, one can actually get an explicit polynomial bound of the form $O(h^{1/2-0})$ for the rate of decay of expectations for the density-one subsequence (see [JZ16]). However, the latter estimate is not satisfied for

all eigenfunctions and it is not clear which sequence of eigenfunctions must be removed for the estimate to hold. There are several articles that address this issue by restricting to special cases of Riemannian surfaces (M, g) and special curves $H \subset M$. Working on surfaces of strictly negative curvature, and choosing H to be a geodesic, Chen-Sogge [CS15] proved $\int_H \phi_h \, d\sigma_H = o(1)$. Subsequently, Sogge-Xi-Zhang [SXZ16] obtained a $O((\log h)^{-1/2})$ bound on the rate of decay under a relaxed curvature condition. Recently, working on surfaces of non-positive curvature Wyman [Wym17] obtained (1) when assuming curvature conditions on H. Finally, we remark that on average, one expects $\int_H \phi_h \, d\sigma_H \asymp h^{\frac{1}{2}}$ (see [Esw16]).

In this article we focus on establishing (1) given explicit conditions on the sequence of eigenfunctions $\{\phi_h\}$. We do not impose any geometric conditions on (M, g), nor do we assume it is a surface. Furthermore, we do not restrict our attention to geodesic curves and allow H to be any hypersurface in M. Instead, we prove that (1) holds provided that the sequence $\{\phi_h\}$ does not asymptotically concentrate in the conormal direction N^*H to H. One example where this holds is the case quantum ergodic sequences of eigenfunctions and any hypersurface H.

1.1. **Statements of the results.** Let $H \subset M$ be a closed smooth hypersurface, and write $S_H^*M \subset S^*M$ for the space of unit covectors with foot-points in H, and S^*H for the set of unit covectors tangent to H. We fix $t_0 > 0$ small enough and define a measure μ_H on $S_H^*M \subset S^*M$ by

$$\mu_H(A) := \frac{1}{2t_0} \mu\Big(\bigcup_{|s| \leq t_0} G^s(A) \Big), \tag{3}$$

where $G^t : S^*M \to S^*M$ denotes the geodesic flow. Remark 3 shows that if $A \subset S_H^*M$ is so that $\overline{A} \subset S_H^*M \backslash S^*H$, then $\mu_H(A)$ is independent of the choice of t_0 and it is natural to replace fixed t_0 with $\lim_{t_0 \to 0}$.

Definition 1. We say that μ is *conormally diffuse with respect to H* if

$$\mu_H(N^*H) = 0.$$

If $U \subset H$ is open, we say that μ is *conormally diffuse with respect to H over U* if

$$\mu_H(N^*H \cap S_U^*M) = 0.$$

As an example, this condition is satisfied when $\{\phi_h\}$ is a quantum ergodic (QE) sequence and $\mu = \mu_L$, the Liouville measure on S^*M. Note that the QE condition is much stronger than the assumption in Definition 1. In Section 5 we give examples of hypersurfaces and sequences of eigenfunctions for which the defect measure is conormally diffuse but is not absolutely continuous with respect to the Liouville measure. Our main result is the following.

Theorem 1. *Let $H \subset M$ be a closed hypersurface. Let $\{\phi_h\}$ be a sequence of eigenfunctions associated to a defect measure μ that is conormally diffuse with respect to H. Then,*

$$\int_H \phi_h \, d\sigma_H = o(1),$$

and

$$\int_H h\partial_\nu\phi_h \, d\sigma_H = o(1),$$

as $h \to 0^+$.

Remark 1. The proof of Theorem 1 actually shows that $\int_H \phi_h\chi d\sigma_H = o(1)$ for any $\chi \in C^\infty(H)$. We note also that the methods of this paper give another independent proof of (2).

As we have already pointed out, the Liouville measure $\mu = \mu_L$ is conormally diffuse. Consequently, the following result is a corollary of Theorem 1:

Theorem 2. *Le $H \subset M$ be a closed hypersurface and $\{\phi_h\}$ be any QE sequence sequence of eigenfunctions. Then,*

$$\int_H \phi_h \, d\sigma_H = o(1) \qquad and \qquad \int_H h\partial_\nu\phi_h \, d\sigma_H = o(1).$$

By Lindenstrauss' celebrated result [Lin06], Hecke eigenfunctions on compact, arithmetic hyperbolic surfaces are all QE (ie. they are quantum uniquely ergodic (QUE)). Together with Theorem 2 this yields

Theorem 3. *Let $(H/\Gamma, g)$ be a compact, arithmetic surface and $H \subset M$ be a closed, C^∞ curve. Then, for all Hecke eigenfunctions $\{\phi_h\}$,*

$$\int_H \phi_h \, d\sigma_H = o(1) \qquad and \qquad \int_H h\partial_\nu\phi_h \, d\sigma_H = o(1).$$

One can localize the results in Theorems 1-3. In the following, we write $d\sigma_H$ for the measure on H induced by the Riemannian structure.

Theorem 4. *Let (M, g) be a smooth, closed Riemannian manifold and $H \subset M$ be a closed hypersurface with $A \subset H$ a subset with piecewise C^∞ boundary and suppose $U \subset H$ is open with $\overline{A} \subset U$. Let $\{\phi_h\}$ be a sequence of eigenfunctions with defect measure μ conormally diffuse with respect to H over U. Then,*

$$\int_A \phi_h \, d\sigma_H = o(1),$$

and

$$\int_A h\partial_\nu\phi_h \, d\sigma_H = o(1),$$

as $h \to 0^+$.

Remark 2. We note that as a corollary of Theorem 4, the results in Theorems 2 and 3 for QE eigenfunctions extend to all smooth curve segments A.

2. Decomposition of defect measures

2.1. **Invariant Measures near transverse submanifolds.** Let N be a smooth manifold, $\mathcal{V}$ be a vector field on N and write $\varphi_t^{\mathcal{V}} : N \to N$ for the flow map generated by $\mathcal{V}$ at time t. Let $\Sigma \subset N$ be a smooth manifold transverse to $\mathcal{V}$. Then for $\epsilon > 0$ small enough, the map $\iota : (-2\epsilon, 2\epsilon) \times \Sigma \to N$

$$\iota(t, q) = \varphi_t^{\mathcal{V}}(q)$$

is a diffeomorphism onto its image and we may use $(-2\epsilon, 2\epsilon) \times \Sigma$ as coordinates on N near Σ.

Lemma 5. *Suppose that μ is a finite Borel measure on N and that $\mathcal{V}\mu = 0$ i.e. $(\varphi_t^{\mathcal{V}})_*\mu = \mu$. Then, for a Borel set $A \subset [-\varepsilon, \varepsilon] \times \Sigma$,*

$$\iota^*\mu(A) = dt d\mu_\Sigma(A)$$

where $d\mu_\Sigma$ is a finite Borel measure on Σ.

Proof. As above, we choose coordinates (t, q) so that $\iota^*\mathcal{V} = \partial_t$. Then, for all $F \in C_c^\infty(-2\epsilon, 2\epsilon) \times \Sigma$,

$$\int \partial_t F d\mu = 0.$$

Now, fix $\chi \in C_c^\infty((-2\epsilon, 2\varepsilon))$ with with $\int \chi dt = 1$. Let $f \in C_c^\infty((-2\epsilon, 2\varepsilon) \times \Sigma)$ and define

$$\bar{f}(q) := \int f(t, q) dt.$$

Then $f(t, q) - \chi(t)\bar{f}(q) = \partial_t F$ with

$$F(t, q) := \int_{-\infty}^t f(s, q) - \chi(s)\bar{f}(q) ds \in C_c^\infty((-2\varepsilon, 2\varepsilon) \times \Sigma).$$

Therefore, for all $f \in C_c^\infty((-2\varepsilon, 2\varepsilon) \times N)$ and $\chi \in C_c^\infty((-2\varepsilon, 2\varepsilon))$ with $\int \chi dt = 1$,

$$\int f(t, q) d\mu(t, q) = \int \chi(t)\bar{f}(q) d\mu(t, q) = \iiint f(s, q) ds \chi(t) d\mu(t, q).$$

Now, let $B \subset \Sigma$ be Borel and $I \subset (-2\varepsilon, \varepsilon)$ Borel and $f_n(t, q) \uparrow 1_I(t)1_B(q)$. Then by the dominated convergence theorem,

$$\mu(I \times B) = \iint |I|1_B(q)\chi(t) d\mu(t, q).$$

Next, let $\chi_n \uparrow (2\varepsilon)^{-1}1_{[-\varepsilon, \varepsilon]}$ with $\int \chi_n \equiv 1$. Then we obtain

$$\mu(I \times B) = \frac{|I|}{2\varepsilon}\mu([-\varepsilon, \varepsilon] \times B).$$

So, letting $\mu_\Sigma(B) := (2\varepsilon)^{-1}\mu([-\varepsilon, \varepsilon] \times B)$, we have that for rectangles $I \times B$, $\mu(I \times B) = dt d\mu_\Sigma(I \times B)$. But then, since these sets generate the Borel sigma algebra, the proof of the lemma is complete.

$\square$

2.2. **Fermi coordinates.** Throughout the remainder of the article we will work in the case that $H \subset M$ is a smooth, orientable, separating hypersurface. That is, $M \setminus H$ has two connected components. We then recover Theorem 1 for general H after proving Theorem 4 for such hypersurfaces. We then divide a given hypersurface into finitely many (possibly overlapping) subsets of separating orientable hypersurfaces and apply Theorem 4 to each. Let $H \subset M$ be a closed smooth hypersurface and let U_H be a Fermi collar neighborhood of H. In Fermi coordinates

$$U_H = \{(x', x_n) : x' \in H \text{ and } x_n \in (-c, c)\}$$

for some $c > 0$, and $H = \{(x', 0) : x' \in H\}$. Since H is a closed, separating hypersurface, it divides M into two connected components Ω_H and $M \backslash \Omega_H$. In the Fermi coordinates system, the point (x', x_n) is identified with the point $\exp_{x'}(x_n \nu_n) \in U_H$ where ν_n is the unit normal vector to Ω_H with base point at $x' \in H$.
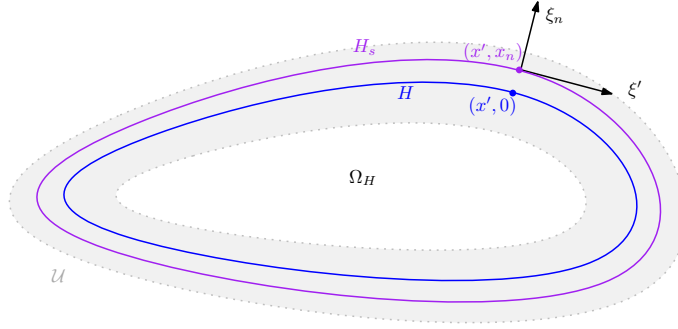
The Fermi coordinates on U_H induce coordinates (x', x_n, ξ', ξ_n) on $S^*_{\mathcal{U}} M = \{(x, \xi) \in S^* M : x \in \mathcal{U}\}$ with $(\xi', \xi_n) \in S^*_{(x', x_n)} M$. In these coordinates, ξ' is cotangent to H while ξ_n is conormal to H.

Note that in the Fermi coordinate system we have

$$|(\xi', \xi_n)|^2_{g(x', x_n)} = \xi_n^2 + R(x', x_n, \xi'), \tag{4}$$

where R satisfies that $R(x', 0, \xi') = |\xi'|^2_{g_H(x')}$ for all $(x', \xi') \in T^* H$ and g_H is the Riemannian metric induced on H by g.

2.3. **Transversals for defect measures.** We now apply Lemma 5 to the special case of defect measures, using the fact that they are invariant under the geodesic flow. In what follows we write $|\xi'|_{x'} := |\xi'|_{g_H(x')}$, where g_H is the Riemannian metric on H induced by g. Let

$$\mathcal{G}_H(\delta) := \{(x, \xi) \in S^*_H M : |\xi'|^2_{x'} \geq 1 - \delta^2\},$$

and define the set of non-glancing directions

$$\Sigma_\delta := S^*_H M \setminus \mathcal{G}_H(\delta).$$

Lemma 6. *Suppose μ is a defect measure associated to a sequence of Laplace eigenfunctions. Then, for all $\delta > 0$ there exists $\varepsilon > 0$ small enough so that*

$$\iota^* \mu = dt \, d\mu_{\Sigma_\delta} \qquad on \ \ (-\varepsilon, \varepsilon) \times \Sigma_\delta$$

where

$$\iota : (-\varepsilon, \varepsilon) \times \Sigma_\delta \to \bigcup_{|s| < \varepsilon} G^s(\Sigma_\delta), \qquad \iota(t, q) = G^t(q),$$

is a diffeomorphism and $d\mu_{\Sigma_\delta}$ is a finite Borel measure on Σ_δ.

Proof. In what follows we use Lemma 5 with $N = S^*M$, $\mathcal{V} = H_p$ the Hamiltonian flow for $p = |\xi|_g$, and $\varphi_t^{\mathcal{V}} = G^t$ the geodesic flow. Note that since μ is a defect measure for a sequence of Laplace eigenfunctions, it is invariant under the geodesic flow G^t. Then, for $q \in \Sigma_\delta$,

$$|H_p x_n(q)| > c\,\delta > 0$$

and hence Σ_δ is transverse to G^t. Therefore, there exists $\varepsilon > 0$ so that $\iota : (-2\varepsilon, 2\varepsilon) \times \Sigma_\delta \to S^*M$, with $\iota(t, q) = G^t(q)$, is a coordinate map. $\square$

Remark 3. For each $A \subset S_H^*M$ with $\overline{A} \subset S_H^*M \setminus S^*H$, there exists $\delta_0 > 0$ so that

$$d\mu_{\Sigma_\delta}(A) = \lim_{t \to 0} \frac{1}{2t} \mu\Big(\bigcup_{|s| \leq t} G^s(A) \Big)$$

for all $0 < \delta \leq \delta_0$. Indeed, since $\overline{A}$ is compact, there exists $\delta_0 = \delta_0(A) > 0$ so that $\overline{A} \subset \Sigma_{\delta_0}$. Then, by Lemma 6, there exists $\varepsilon = \varepsilon(A) > 0$ so that if $|t| \leq \varepsilon$, then

$$\mu\Big(\bigcup_{|s| \leq t} G^s(A) \Big) = 2t\, d\mu_{\Sigma_\delta}(A).$$

In particular, we conclude that the quotient $\frac{1}{2t}\mu\Big(\bigcup_{|s| \leq t} G^s(A) \Big)$ is independent of t as long as $|t| \leq \varepsilon$.

We also need the following description of μ.

Lemma 7. *Suppose μ is a defect measure associated to a sequence of Laplace eigenfunctions, and let $\delta > 0$. Then, in the notation of Lemma 6, there exists $\varepsilon_0 > 0$ small enough so that*

$$\mu = |\xi_n|^{-1} d\mu_{\Sigma_\delta}(x', \xi', \xi_n) dx_n,$$

for $(x', x_n, \xi', \xi_n) \in \iota((-\varepsilon_0, \varepsilon_0) \times \Sigma_\delta)$.

Remark 4. Notice that $|H_p x_n| > \gamma$ on $\Sigma_\gamma = S_H^*M \setminus \mathcal{G}(\gamma)$. Therefore, there exists $c_0, c_1 > 0$ so that

$$\{(x', x_n, \xi', \xi_n) : \ |x_n| \leq c_0\gamma, \ |\xi'|_{x'}^2 \leq 1 - c_0^{-1}\gamma^2\} \subset \bigcup_{|t| \leq c_1\gamma} G^t(\Sigma_\gamma).$$

Proof. By Lemma 6,

$$\iota^*\mu = d\mu_{\Sigma_\delta}(x', \xi', \xi_n) dt \quad \text{on } (-\varepsilon, \varepsilon) \times \Sigma_\delta.$$

Then, for $q \in \Sigma_\delta$

$$|\partial_t x_n(\iota(0, q))| = |H_p x_n(\iota(0, q))| = \frac{|\xi_n(\iota(0, q))|}{|\xi|_g} > \delta$$

and hence for $\varepsilon_0 > 0$ small enough and $q \in \Sigma_\delta$, $t \in (-\varepsilon_0, \varepsilon_0)$,

$$|\partial_t x_n(\iota(t,q))| = |H_p x_n(\iota(t,q))| = \frac{|\xi_n(\iota(t,q))|}{|\xi|_g} > \frac{\delta}{2}.$$

Therefore, $dt = f(x', x_n, \xi', \xi_n) dx_n$ where

$$f(x', x_n, \xi', \xi_n) = |H_p x_n(\iota^{-1}(x_n, (x', \xi', \xi_n)))|^{-1} = \frac{|\xi|_g}{|\xi_n|} = |\xi_n|^{-1}$$

where in the last equality, we use that $|\xi|_g = 1$. In particular,

$$\mu = |\xi_n|^{-1} d\mu_{\Sigma_\delta}(x', \xi', \xi_n) dx_n.$$

$\square$

Before proceeding to the proof of Theorem 1 we note that Lemma 7 implies that for all $\delta > 0$,

$$\mu(S_H^* M \setminus \mathcal{G}_H(\delta)) = 0. \tag{5}$$

Remark 5. Notice that the measure

$$|\xi_n|^{-1} d\mu_{\Sigma_\delta}(x', \xi', \xi_n) = \frac{1}{\sqrt{1 - |\xi'|_{x'}^2}} d\mu_{\Sigma_\delta}(x', \xi', \xi_n)$$

is hypersurface measure on $S_H^* M \setminus \mathcal{G}(\delta)$ induced by μ where we take ∂_{x_n} to be the normal vector field to $S_H^* M$. For example, if μ_L is Liouville measure, then, parametrizing $S_H^* M \setminus \mathcal{G}(\delta)$ by (x', ξ')

$$d(\mu_L)_{\Sigma_\delta} = c\mathbf{1}_{\{S_H^* M \setminus \mathcal{G}(\delta)\}}(x', \xi', \xi_n) dx' d\xi'$$

for some $c > 0$.

3. Proof of Theorem 1

Consider the cut-off function $\chi_\alpha \in C^\infty(\mathbb{R}, [0, 1])$ with

$$\chi_\alpha(t) = \begin{cases} 0 & |t| \geq \alpha \\ 1 & |t| \leq \frac{\alpha}{2}, \end{cases}$$

with $|\chi_\alpha'(t)| \leq 3/\alpha$ for all $t \in \mathbb{R}$.

For $\delta > 0$ consider the symbol

$$\beta_\delta(x', \xi') = \chi_\delta(|\xi'|_{x'}) \in S^0(T^* H) \tag{6}$$

where we continue to write $|\xi'|_{x'} := |\xi'|_{g_H(x')}$. We refer the reader to the Appendix where the semiclassical notation used in this section is introduced. The operator $Op_h(\beta_\delta) \in \Psi^0(H)$ microlocalizes near the conormal direction in $T^* H$ which is identified with $\xi' = 0$ via the orthogonal projection. The first step towards the proof of Theorem 1 is to reduce the problem to study averages over H of the functions ϕ_h and $h\partial_\nu \phi_h$ when microlocalized near the conormal direction.

Lemma 8. For any $\delta > 0$ and $u \in L^2(H)$,

$$\int_H u \, d\sigma_H = \int_H Op_h(\beta_\delta) u \, d\sigma_H + O_\delta(h^\infty) \|u\|_{L^2(H)}.$$

Proof. We wish to show that

$$\langle (1 - Op_h(\beta_\delta)) u, 1 \rangle_{L^2(H)} = \langle u, (1 - Op_h(\beta_\delta))^* 1 \rangle_{L^2(H)} = O_\delta(h^\infty).$$

To prove this, we simply note that in local coordinates

$$(1 - Op_h(\beta_\delta))^* 1(x) = \frac{1}{(2\pi h)^{n-1}} \iint e^{\frac{i}{h} \langle x - x', \xi' \rangle} a_\delta(x, \xi'; h)(1 - \chi_{2\delta})(|\xi'|_x) \, d\xi' dx',$$

for some symbol $a_\delta \in S^0$. The phase function $\Phi(x', \xi'; x) = \langle x - x', \xi' \rangle$ has critical points in (x', ξ') given by

$$(x', \xi') = (x, 0).$$

By repeated integration by parts with respect to the operator

$$L := \frac{1}{|x - x'|^2 + |\xi'|^2} \left(\sum_{j=1}^n \xi'_j h D_{x'_j} + \sum_{j=1}^n (x'_j - x_j) h D_{\xi'_j} \right),$$

using that $L(e^{i\Phi/h}) = e^{i\Phi/h}$, one gets

$$(1 - Op_h(\beta_\delta))^* 1(x) = \frac{1}{(2\pi h)^{n-1}} \iint e^{i \frac{(x - x')\xi'}{h}} a_\delta(x, \xi')(1 - \chi_\delta)(|\xi'|_x) \chi_1(|x - x'|) \, d\xi' dx' + O_\delta(h^\infty)$$

$$= O_\delta(h^\infty),$$

uniformly in $x \in H$. The last line follows by repeated integrations by parts with respect to L using the fact that $(1 - \chi_\delta)^{(k)}(0) = 0$ for all $k \geq 0$.   □

3.1. **Proof of Theorem 1.** We wish to show that for any $\varepsilon > 0$ there exists $h_0(\varepsilon) > 0$ so that

$$\left| \int_H \phi_h \, d\sigma_H \right| \leq \varepsilon \qquad \text{and} \qquad \left| \int_H h \partial_\nu \phi_h d\sigma_H \right| \leq \varepsilon, \tag{7}$$

for all $h \leq h_0$.

In view of Lemma 8, we can microlocalize the problem to the conormal direction; that is, the claim in (7) follows provided we prove that given $\varepsilon > 0$ there exist $\delta(\varepsilon) > 0$ and $h_0(\varepsilon) > 0$ so that

$$\left| \int_H Op_h(\beta_\delta) \phi_h \, d\sigma_H \right| \leq \varepsilon \qquad \text{and} \qquad \left| \int_H Op_h(\beta_\delta) h \partial_\nu \phi_h d\sigma_H \right| \leq \varepsilon, \tag{8}$$

for all $h \leq h_0(\varepsilon)$.

To prove (8), by Cauchy-Schwarz, it clearly suffices to establish the stronger bounds

$$\| Op_h(\beta_\delta) \phi_h \|_{L^2(H)} \leq \varepsilon \qquad \text{and} \qquad \| Op_h(\beta_\delta) h \partial_\nu \phi_h \|_{L^2(H)} \leq \varepsilon, \tag{9}$$

for all $h \leq h_0(\varepsilon)$ and $\delta(\varepsilon) > 0$ sufficiently small.

From now on, we fix $\varepsilon > 0$. Using Green's formula [CTZ13], it is straightforward to check that for any operator $A : C^\infty(M) \to C^\infty(M)$ one has the Rellich Identity

$$\frac{i}{h} \int_{\Omega_H} [-h^2 \Delta_g, A] \phi_h \, \overline{\phi_h} \, dv_g = \int_H A \phi_h \, \overline{h D_\nu \phi_h} \, d\sigma_H + \int_H h D_\nu(A\phi_h) \, \overline{\phi_h} \, d\sigma_H, \tag{10}$$

where $D_\nu = \frac{1}{i} \partial_\nu$, with ν being the unit outward vector normal to Ω_H.

Let $\delta > 0$ and $\alpha > 0$ be two real valued parameters to be specified later and consider the operator

$$A_{\delta,\alpha}(h) := Op_h(\beta_\delta^2) \circ Op_h(\chi_\alpha(x_n)) \circ hD_\nu,$$

where β_δ is defined in (6).

Remark 6. We note that when we write $Op_h(\beta_\delta^2)$ above, we are actually considering the operator $Op_h(\beta_\delta^2) \otimes \mathrm{Id}_{x_n}$. That is, for $u \in C^\infty(M)$,

$$[Op_h(\beta_\delta^2)u](x', y_n) = [Op_h(\beta_\delta^2)u|_{x_n=y_n}](x').$$

The operator $A_{\delta,\alpha}(h)$ is the semiclassical normal derivative operator h-microlocalized to a neighbourhood of the conormal direction to H over the collar neighbourhood U_H.

We note that

$$\int_H A_{\delta,\alpha}(h)\phi_h\,\overline{hD_\nu\phi_h}\,d\sigma_H = \left\langle Op_h(\beta_\delta^2)hD_\nu\phi_h\,,\,hD_\nu\phi_h\right\rangle_{L^2(H)}, \qquad (11)$$

since $\chi_\alpha(x_n) = 1$ for $x_n \in [-\frac{\alpha}{2}, \frac{\alpha}{2}]$. Without loss of generality, we may assume that

$$\Omega_H \cap U_H = \{(x', x_n) : x' \in H \text{ and } x_n < 0\}.$$

With this choice, $D_\nu = D_{x_n}$. We next recall that

$$\gamma_H(h^2 D_\nu^2 \phi_h) = (I + h^2\Delta_{g_H})\gamma_H(\phi_h) + ha_1\gamma_H(\phi_h) + ha_2\gamma_H(hD_\nu\phi_h),$$

where $\gamma_H : M \to H$ is the restriction map to H, and $a_1, a_2 \in C^\infty(H)$. Since $\chi'_\alpha(0) = 0$ it follows from the restriction upper bounds $\|\phi_h\|_H = O(h^{-1/4})$ [BGT07, HT12, Tac10, Tat98] and $\|hD_\nu\phi_h\|_H = O(1)$ [CHT15, Tac14] that

$$\langle (hD_\nu)^2\phi_h, \phi_h\rangle_{L^2(H)} - \langle (1 + h^2\Delta_{g_H})\phi_h, \phi_h\rangle_{L^2(H)} = O_{L^2}(\sqrt{h}).$$

Consequently,

$$\begin{aligned}
\int_H hD_\nu(A_{\delta,\alpha}(h)\phi_h)\,\overline{\phi_h}\,d\sigma_H &= \left\langle hD_\nu Op_h(\beta_\delta^2)\chi_\alpha(x_n)hD_\nu\phi_h\,,\,\phi_h\right\rangle_{L^2(H)} \\
&= \left\langle Op_h(\beta_\delta^2)(hD_\nu)^2\phi_h\,,\,\phi_h\right\rangle_{L^2(H)} \\
&= \left\langle Op_h(\beta_\delta^2)(1 + h^2\Delta_{g_H})\phi_h\,,\,\phi_h\right\rangle_{L^2(H)} + O(h^{\frac{1}{2}}). \quad (12)
\end{aligned}$$

Substitution of (11) and (12) in (10) gives

$$\frac{i}{h}\int_{\Omega_H}[-h^2\Delta_g, A_{\delta,\alpha}(h)]\phi_h\,\overline{\phi_h}\,dv_g =$$

$$= \left\langle Op_h(\beta_\delta^2)hD_\nu\phi_h\,,\,hD_\nu\phi_h\right\rangle_H + \left\langle Op_h(\beta_\delta^2)(1 + h^2\Delta_{g_H})\phi_h\,,\,\phi_h\right\rangle_H + O(h^{\frac{1}{2}}). \quad (13)$$

Next, we observe that

$$\|Op_h(\beta_\delta)hD_\nu\phi_h\|_H^2 = \left\langle Op_h(\beta_\delta^2)hD_\nu\phi_h\,,\,hD_\nu\phi_h\right\rangle_H + O(h) \qquad (14)$$

since $\|hD_\nu\phi_h\|_H = O(1)$ [CHT15]. On the other hand, for $(x', \xi') \in \mathrm{supp}\,\beta_\delta$ we have $|\xi'|_x^2 \le \delta^2$ and so,

$$\beta_\delta^2 \cdot (1 - |\xi'|_x^2) - \beta_\delta^2 \cdot (1 - 2\delta^2) = \beta_\delta^2\left(2\delta^2 - |\xi'|_x^2\right) \ge \beta_\delta^2\delta^2 \ge 0.$$

Therefore, combining the sharp Garding inequality with the bound $\|\phi_h\|_H = O(h^{-1/4})$ gives

$$
\begin{aligned}
(1 - 2\delta^2)\|Op_h(\beta_\delta)\phi_h\|_H^2 &= \langle Op_h(\beta_\delta^2(1 - 2\delta^2))\phi_h, \phi_h\rangle_H + O(h^{\frac{1}{2}}) \\
&\leq \langle Op_h(\beta_\delta^2 \cdot (1 - |\xi'|_x^2))\phi_h, \phi_h\rangle_H + O(h^{\frac{1}{2}}) \\
&= \langle Op_h(\beta_\delta^2)(1 + h^2\Delta_{g_H})\phi_h, \phi_h\rangle_H + O(h^{\frac{1}{2}}).
\end{aligned}
\tag{15}
$$

Substitution of (14) and (15) into (13) gives

$$
\|Op_h(\beta_\delta)hD_\nu\phi_h\|_H^2 + (1 - 2\delta^2)\|Op_h(\beta_\delta)\phi_h\|_H^2 \leq \\
\frac{i}{h}\int_{\Omega_H}[-h^2\Delta_g, A_{\delta,\alpha}(h)]\phi_h\,\overline{\phi_h}\,dv_g + O(h^{\frac{1}{2}}). \tag{16}
$$

The claim in (9) follows at once from (16) provided we show that for any $\varepsilon > 0$ there exist $\delta, \alpha > 0$ and $h_0 > 0$ (all possibly depending on ε) such that

$$
\left|\left\langle \frac{i}{h}[-h^2\Delta_g, A_{\delta,\alpha}(h)]\phi_h\,,\,\phi_h\right\rangle_{L^2(\Omega_H)}\right| \leq \varepsilon^2 \qquad \forall h \leq h_0(\varepsilon) \tag{17}
$$

To prove (17) we note that

$$
\left\langle \frac{i}{h}[-h^2\Delta_g, A_{\delta,\alpha}(h)]\phi_h\,,\,\phi_h\right\rangle_{L^2(\Omega_H)} = \\
= \left\langle Op_h\Big(\big\{\sigma(-h^2\Delta_g), \sigma(A_{\delta,\alpha}(h))\big\}\Big)\phi_h\,,\,\phi_h\right\rangle_{L^2(\Omega_H)} + O(h), \tag{18}
$$

where $\sigma(A_{\delta,\alpha}(h))(x,\xi) = \beta_\delta^2(x',\xi')\,\chi_\alpha(x_n)\xi_n$, and according to (4), the Poisson bracket

$$
\big\{|(\xi',\xi_n)|_x^2\,,\,\sigma(A_{\delta,\alpha}(h))\big\} = 2\chi_\alpha'(x_n)\beta_\delta^2(x',\xi')\,\xi_n^2 + \chi_\alpha(x_n)q_\delta(x',x_n,\xi',\xi_n) \tag{19}
$$

where,

$$
q_\delta(x,\xi) := \xi_n\partial_{\xi'}R \cdot \partial_{x'}\beta_\delta^2 - \xi_n\partial_{x'}R \cdot \partial_{\xi'}\beta_\delta^2 - \partial_{x_n}R \cdot \beta_\delta^2.
$$

We now estimate each term in the RHS of (19) separately.

Lemma 9. *Let $\{\phi_h\}$ be an L^2-normalized eigenfunction sequence with defect measure μ. Then,*

(i) $|\langle Op_h(\chi_\alpha(x_n)q_\delta)\phi_h\,,\,\phi_h\rangle_{L^2(\Omega_H)}| \leq R_{\alpha,\delta} + o(1),$
where

$$
R_{\alpha,\delta} := \|q_\delta\|_{L^\infty} \cdot \mu(\{(x',x_n,\xi',\xi_n) \in S_{U_H}^*M : |x_n| \leq \alpha, |\xi'| < \delta\})^{\frac{1}{2}}.
$$

In addition,

(ii) $\langle Op_h(2\chi_\alpha'(x_n)\beta_\delta^2(x',\xi')\,\xi_n^2)\phi_h\,,\,\phi_h\rangle_{L^2(\Omega_H)} = \int_{S_{\Omega_H}^*M} 2\chi_\alpha'(x_n)\beta_\delta^2(x',\xi')\,\xi_n^2\,d\mu + o(1).$

In both (i) and (ii), $o(1)$ denotes a term that vanishes as $h \to 0^+$.

We postpone the proof of Lemma 9 until the end of this section. Assuming this result for the moment, we now conclude the proof of the theorem. From Lemma 9 and (19), it follows that

$$\left\langle \frac{i}{h}[-h^2\Delta_g, A_{\delta,\alpha}(h)]\phi_h \, , \, \phi_h \right\rangle_{L^2(\Omega_H)} = \int_{S^*_{\Omega_H} M} 2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi') \, \xi_n{}^2 \, d\mu + R_{\alpha,\delta} + o(1).$$
(20)

Since μ is a Radon measure, and hence monotone,

$$\lim_{\alpha\to 0} R_{\alpha,\delta} = \|q_\delta\|_{L^\infty} \cdot \mu(\{(x',0,\xi) \in S^*_H M; \ |\xi'| < \delta\})^{\frac{1}{2}}.$$
(21)

Thus, using Lemma 7 (or more precisely (5)) gives

$$\lim_{\alpha\to 0} R_{\alpha,\delta} = 0.$$
(22)

Moreover, since the LHS of (16) is independent of α, we are free to take the $\alpha \to 0$ limit of both sides. In view of (20) and (21), it follows that after taking $h \to 0^+$ and then $\alpha \to 0^+$,

$$\limsup_{h\to 0} \left(\|Op_h(\beta_\delta)hD_\nu\phi_h\|_H^2 + (1-2\delta^2)\|Op_h(\beta_\delta)\phi_h\|_H^2 \right) \leq$$

$$\leq \limsup_{\alpha\to 0}\limsup_{h\to 0^+} \frac{i}{h}\int_{\Omega_H} [-h^2\Delta_g, A_{\delta,\alpha}(h)]\phi_h \, \overline{\phi_h} \, dv_g$$

$$= \limsup_{\alpha\to 0^+}\int_{S^*_{\Omega_H} M} 2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi') \, \xi_n^2 \, d\mu.$$
(23)

The last line in (23) follows from (22).

To analyze the RHS of (23), fix $\gamma > 0$ small. By Lemma 7 there exists $\varepsilon_\gamma > 0$ and a measure μ_{Σ_γ} on $\Sigma_\gamma = \{(x,\xi) \in S^*_H M : \ |\xi'|^2_{x'} \leq 1 - \gamma^2\}$ so that

$$\mu(x,\xi) = f(x',x_n,\xi',\xi_n)d\mu_{\Sigma_\gamma}(x',\xi',\xi_n)dx_n, \qquad (x,\xi) \in \bigcup_{|t|\leq\varepsilon_\gamma} G^t(\Sigma_\gamma).$$

By Remark 4 we may assume that we work with α,δ small enough so that

$$\operatorname{supp}(\chi'_\alpha \cdot \beta_\delta^2) \subset \bigcup_{|t|\leq\varepsilon_\gamma} G^t(\Sigma_\gamma).$$
(24)

Since $\operatorname{supp}(\chi'_\alpha) \subset (-\alpha, 0)$, by the Fubini theorem we have

$$\int_{S^*_{\Omega_H} M} 2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2 \, d\mu =$$

$$= \int_{-c}^0 2\chi'_\alpha(x_n) \left(\int_{S^*_H M} \beta_\delta^2(x',\xi')\xi_n^2|\xi_n|^{-1}d\mu_{\Sigma_\gamma}(x',\xi',\xi_n) \right) dx_n$$
(25)

$$= \int_{S^*_H M}\int_{-c}^0 2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')|\xi_n|dx_n d\mu_{\Sigma_\gamma}(x',\xi',\xi_n).$$

Sending $\alpha \to 0$ gives

$$\lim_{\alpha \to 0^+} \int_{S^*_{\Omega_H} M} 2\chi'_\alpha(x_n)\beta^2_\delta(x',\xi')\xi^2_n \, d\mu = \int_{S^*_H M} 2\beta^2_\delta(x',\xi')|\xi_n|d\mu_{\Sigma_\gamma}(x',\xi',\xi_n).$$

Sending $\delta \to 0$ and using that $\beta_\delta \equiv 1$ on N^*H, $|\beta_\delta| \leq C$ we obtain

$$\lim_{\delta \to 0} \lim_{\alpha \to 0^+} \int_{S^*_{\Omega_H} M} 2\chi'_\alpha(x_n)\beta^2_\delta(x',\xi')\xi^2_n \, d\mu = \int_{N^*H} 2 \, d\mu_{\Sigma_\gamma}(x',\xi',\xi_n) = 2\,\mu_{\Sigma_\gamma}(N^*H). \tag{26}$$

Since μ is conormally diffuse, we have by Remark 3 that $\mu_{\Sigma_\gamma}(N^*H) = 0$ and so (9) follows from (26) and (23). $\square$

3.2. **Proof of Lemma 9.**

Proof. First, we use the standard fact that $\{\phi_h\}$ are microsupported on S^*M [CHT15] to h-microlocally cut them off near S^*M. More precisely, for $r > 0$ small, consider the annular shell

$$A(r) := \{(x,\xi) \in T^*M : 1 - r < |\xi|_{g(x)} < 1 + r\}.$$

Let $\tilde\chi \in C^\infty_c(T^*M)$ be a cutoff function equal to 1 on $A(r)$ and zero on $T^*M \setminus A(2r)$. Then, [CHT15]

$$\|\phi_h - Op_h(\tilde\chi)\phi_h\|_{L^2(M)} = O(h^\infty). \tag{27}$$

Proof of (i): Since $\|\phi_h\|_{L^2(M)} = 1$, by Cauchy-Schwarz,

$$\begin{aligned}
\big|\langle Op_h(\chi_\alpha(x_n)q_\delta)\phi_h\,,\,\phi_h\rangle_{L^2(\Omega_H)}\big|^2 &\leq \|Op_h(\chi_\alpha(x_n)q_\delta)\phi_h\|^2_{L^2(M)} \\
&= \big\langle [Op_h(\chi_\alpha(x_n)q_\delta)]^*[Op_h(\chi_\alpha(x_n)q_\delta)]\phi_h, \phi_h\big\rangle_{L^2(M)} \\
&= \big\langle [Op_h(\chi_\alpha(x_n)q_\delta)]^*[Op_h(\chi_\alpha(x_n)q_\delta)]\phi_h, Op_h(\tilde\chi)\phi_h\big\rangle_{L^2(M)} + O(h^\infty) \\
&= \big\langle Op_h(\tilde\chi \cdot \chi^2_\alpha(x_n) \cdot |q_\delta|^2)\phi_h, \phi_h\big\rangle_{L^2(M)} + O(h) \\
&= \int_{S^*M} \tilde\chi \cdot \chi^2_\alpha(x_n) \cdot |q_\delta|^2 \, d\mu + o(1) \\
&\leq \|q_\delta\|^2_{L^\infty} \cdot \mu\big(\{(x',x_n,\xi',\xi_n) \in S^*_{U_H} M : |x_n| \leq \alpha, |\xi'| < \delta\}\big) + o(1),
\end{aligned}$$

where the penultimate identity follows from the fact that μ is the defect measure associated to $\{\phi_h\}$ and the symbol $\tilde\chi \cdot \chi^2_\alpha(x_n) \cdot |q_\delta|^2 \in C^\infty_c(T^*U_H)$.

Proof of (ii): Let $\rho \in C^\infty_c(\mathbb{R})$ be a smooth cut-off function with $\rho(x_n) = 0$ for $x_n \geq 0$ and $\rho(x_n) = 1$ for $x_n \leq -\alpha/2$. Then, since $\Omega_H \cap U_H$ is identified with the set of points on which $x_n < 0$, and $\text{supp}(\chi'_\alpha) \subset (-\infty, -\alpha/2] \cup [\alpha/2, +\infty)$, we have

$$\rho(x_n)\chi'_\alpha(x_n) = \begin{cases} 0 & \text{on } \Omega^c_H, \\ \chi'_\alpha(x_n) & \text{on } \Omega_H. \end{cases}$$

Note that since $\chi'_\alpha(x_n) = 0$ for $x_n \in [-\alpha/2, \alpha/2]$, we may regard $\rho\chi'_\alpha$ as a smooth function defined on all of M. We then have that

$$\langle Op_h(2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2)\phi_h\,,\,\phi_h\rangle_{L^2(\Omega_H)} =$$
$$= \langle Op_h(2\rho(x_n)\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2)\phi_h\,,\,\phi_h\rangle_{L^2(M)}.$$

Microlocalizing the eigenfunctions near S^*M by using the cut-off $\tilde{\chi}$ we obtain

$$\langle Op_h(2\,\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2)\phi_h\,,\,\phi_h\rangle_{L^2(\Omega_H)} =$$
$$= \langle Op_h(\tilde{\chi}\rho(x_n)2\,\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2)\phi_h\,,\,\phi_h\rangle_{L^2(M)} + O(h).$$

Using that μ is the defect measure associated to $\{\phi_h\}$, and that the symbol $\tilde{\chi}\beta_\delta^2\,\xi_n^2 \in C_c^\infty(T^*M)$, we obtain

$$\langle Op_h(\tilde{\chi}\rho(x_n)2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2)\phi_h\,,\,\phi_h\rangle_{L^2(M)} =$$
$$= \int_{S^*M} 2\rho(x_n)\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2\,d\mu + o(1)$$
$$= \int_{S^*_{\Omega_H}M} 2\chi'_\alpha(x_n)\beta_\delta^2(x',\xi')\xi_n^2\,d\mu + o(1),$$

as claimed. $\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\square$

Remark 7. By replacing the test operator $A_{\delta,\alpha}(h)$ with

$$\tilde{A}_{\delta,\alpha}(h) := Op_h(\beta_\delta^2(x',\xi')) \circ f(x') \circ Op_h(\chi_\alpha(x_n)) \circ hD_\nu,$$

where $f \in C^\infty(H)$ and carrying out the same argument as in the proof of Theorem 1, it is easy to see that under the assumption $\mu_H(\pi^{-1}(\operatorname{supp} f) \cap N^*H) = 0$,

$$\int_H f\,\phi_h\,d\sigma_H = o(1) \quad \text{and} \quad \int_H f\,hD_\nu\phi_h\,d\sigma_H = o(1).$$

4. PROOF OF THEOREM 4

To prove Theorem 4 we need the following result.

Lemma 10. *Suppose $A \subset H$ has piecewise smooth boundary. Then for all $\epsilon > 0$*

$$\|(1 - Op_h(\beta_\delta))^* 1_A\|_{L^2(H)} = O_\epsilon(h^{\frac{1}{2}-\epsilon}).$$

Proof. To prove this result we first introduce a cut-off function χ_h so that $(1 - \chi_h)1_A$ is smooth and close to 1_A. Let $\chi_h \in C_c^\infty(H)$ satisfy

    i) $\chi_h \equiv 1$ on $\{x \in H : d(x, \partial A) \le h^{1-\epsilon}\}$.

    ii) $\operatorname{supp}\chi \subset \{x \in H : d(x, \partial A) \le 2h^{1-\varepsilon}\}$.

    iii) $|\partial_x^\alpha \chi| \le C_\alpha h^{|\alpha|(1-\varepsilon)}$.

Then, $(1 - \chi_h)1_A$ satisfies the same bound as in (iii), and hence integrating by parts as in Lemma 8 i.e. with $L := \frac{1}{|x-x'|^2+|\xi'|^2}\left(\sum_{j=1}^{n}\xi'_j h D_{x'_j} + \sum_{j=1}^{n}(x'_j - x_j)h D_{\xi'_j}\right)$, gives

$$
\begin{aligned}
&[(1 - Op_h(\beta_\delta))^*(1-\chi_h)1_A](x) \\
&= \frac{1}{(2\pi h)^{n-1}} \iint e^{\frac{i}{h}\langle x-x',\xi'\rangle}(1 - \beta_\delta(x',\xi'))(1-\chi_h(x'))1_A(x')dx'd\xi' \\
&= \frac{1}{(2\pi h)^{n-1}} \iint e^{\frac{i}{h}\langle x-x',\xi'\rangle}(L^*)^N\big[(1 - \beta_\delta(x',\xi'))(1-\chi_h(x'))1_A(x')\big]dx'd\xi' \\
&= O_N(h^{1-n+N(1-\varepsilon)}).
\end{aligned}
$$

In particular,

$$\|(1 - Op_h(\beta_\delta))^*(1-\chi_h)1_A\|_{L^\infty} = O_\epsilon(h^\infty). \tag{28}$$

On the other hand

$$\|\chi_h 1_A\|_{L^2(H)} = O(h^{\frac{1-\epsilon}{2}}). \tag{29}$$

Combining (28) and (29) together with L^2 boundedness of $Op_h(\beta_\delta)$ proves the lemma.
$\square$

4.1. **Proof of Theorem 4.** Let $A \subset H$ be an open subset with piecewise C^∞ boundary and indicator function χ_A. Suppose that $U \subset H$ is open with $\overline{A} \subset U$. Then since $C^\infty(H)$ is dense in $L^2(H)$, for any $\varepsilon > 0$, we can find $f \in C^\infty(H)$

$$\|f - 1_A\|_{L^2(H)} \le \varepsilon, \qquad \operatorname{supp} f \subset U.$$

Now,

$$
\left|\int_H 1_A \phi_h d\sigma_H\right| \le
$$
$$
\le \left|\int_H 1_A Op_h(\beta_\delta)\phi_h d\sigma_H\right| + \left|\langle(1 - Op_h(\beta_\delta))\phi_h, 1_A\rangle_H\right|
$$
$$
\le \left|\int_H (1_A - f)Op_h(\beta_\delta)\phi_h d\sigma_H\right| + \left|\int_H fOp_h(\beta_\delta)\phi_h d\sigma_H\right| + \left|\langle\phi_h, (1 - Op_h(\beta_\delta))^*1_A\rangle_H\right|
$$
$$
\le \left|\int_H (1_A - f)Op_h(\beta_\delta)\phi_h d\sigma_H\right| + o(1). \tag{30}
$$

The last line follows by applying Lemma 10, the universal upper bound $\|\phi_h\|_{L^2(H)} \le Ch^{-\frac{1}{4}}$ [BGT07] and Cauchy-Schwarz to the third term, and by applying Remark 7 to the second term.

Now, since β_δ is supported away from

$$S^*H := \{(x',\xi') \in T^*H : |\xi'|_{x'} = 1\},$$

we have that $\|Op_h(\beta_\delta)\phi_h\|_{L^2(H)} \le C$ [BGT07, Tac10] and hence applying Cauchy–Schwarz to (30)

$$\left|\int_H 1_A \phi_h d\sigma_H\right| \le C\varepsilon + o(1).$$

Since $\varepsilon > 0$ was arbitrary, the theorem follows. $\square$

Remark 8. It is clear from the proof of Theorem 4 that one can decrease the regularity assumption on ∂A and only assume that ∂A has Minkowski box dimension $< n - \frac{3}{2}$ where $n = \dim M$. However, we do not pursue this here.

5. Examples

5.1. Non vanishing averages on the torus.
Let $\mathbb{T}^2$ be the 2-dimensional square flat torus. We identify $\mathbb{T}^2$ with $\{(x_1, x_2) : (x_1, x_2) \in [0,1) \times [0,1)\}$. Consider the sequence of normalized eigenfunctions

$$\phi_h(x_1, x_2) = e^{\frac{i}{h} x_1}.$$

Consider the curve $H \subset \mathbb{T}^2$ defined as $H = \{(x_1, x_2) : x_1 = 0\}$. Then, since $\phi_h|_H \equiv 1$, we have

$$\int_H \phi_h d\sigma_H = 1, \quad h^{-1} \in 2\pi\mathbb{Z}^+.$$

We claim that in this case the measure μ associated to $\{\phi_h\}$ is *not* conormally diffuse with respect to H. Actually, we next prove that

$$\mu(x_1, x_2, \xi_1, \xi_2) = \delta_{(1,0)}(\xi_1, \xi_2) \cdot dx_1\, dx_2, \qquad (x, \xi) \in S^*\mathbb{T}^2. \tag{31}$$

Given (31), it follows that

$$\mu_H = \delta_{(1,0)}(\xi_1, \xi_2), \qquad (x, \xi) \in S^*\mathbb{T}^2.$$

In particular,

$$\mu_H(N^*H) = 1,$$

so the measure μ is *not* conormally diffuse with respect to H.

To see that (31) holds, fix any $a \in C_c^\infty(T^*\mathbb{T}^2)$. Then,

$$\langle Op_h(a)\phi_h, \phi_h \rangle = \frac{1}{(2\pi h)^n} \int_{\mathbb{T}^2} \int_{\mathbb{T}^2} \int_{\mathbb{R}^2} a(x, \xi) e^{\frac{i}{h} \psi(x,y,\xi)} d\xi dy dx$$

for the phase function

$$\psi(x, y, \xi) := \langle x - y, \xi \rangle + y_1 - x_1.$$

We next do Stationary Phase in (y, ξ). The critical points for the phase are $(y, \xi) = (x, (1, 0))$. Also,

$$\mathrm{Hess}_{(y,\xi)}\psi = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}.$$

It follows that

$$\langle Op_h(a)\phi_h, \phi_h \rangle = \int_{\mathbb{T}^2} a(x, (1,0)) dx = \int_{S^*\mathbb{T}^2} a(x, \xi)\, \delta_{(1,0)}(\xi) dx,$$

as claimed.

5.2. Defect measures that are not Liouville.
As we already pointed out in the Introduction, the assumptions on μ for being conormally diffuse are much weaker than asking μ to be absolutely continuous with respect to the Liouville measure on S^*M. In these examples we build a defect measure μ that is not absolutely continuous with respect to the Liouville measure but still satisfies the hypothesis of Theorem 1 for a suitable choice of curve H.

5.2.1. *Toral Eigenfunctions.* Let $\mathbb{T}^2$ be the 2-dimensional square flat torus. We identify $\mathbb{T}^2$ with $\{(x_1, x_2) : (x_1, x_2) \in [0, 1) \times [0, 1)\}$. Consider the sequence of eigenfunctions

$$\phi_h(x_1, x_2) = e^{\frac{i}{h}x_1}, \quad h^{-1} \in 2\pi\mathbb{Z}.$$

As shown in Section 5.1, the associated defect measure is

$$\mu(x_1, x_2, \xi_1, \xi_2) = \delta_{(1,0)}(\xi_1, \xi_2) dx_1 \, dx_2.$$

Next, consider the curve $H \subset \mathbb{T}^2$ defined as $H = \{(x_1, x_2) : x_2 = 0\}$. Since $N^*H = \{(x_1, x_2, \xi_1, \xi_2) \in S^*\mathbb{T}^2 : \xi_1 = 0\}$, we have for $\delta > 0$ sufficiently small,

$$\mu_H(N^*H) = 0.$$

Theorem 1 therefore implies that

$$\lim_{h \to 0^+} \int_H \varphi_h d\sigma = \lim_{h \to 0^+} \int_0^1 e^{\frac{ix_1}{h}} dx_1 = 0.$$

Of course, in this case the much stronger result $\int_0^1 e^{\frac{ix_1}{h}} dx_1 = 0$ holds for all $h^{-1} \in 2\pi\mathbb{Z}$.

5.2.2. *Gaussian Beams.* Consider the two dimensional sphere S^2 equipped with the round metric, and use coordinates

$$(\theta, \omega) \mapsto (\cos\theta \cos\omega, \sin\theta \cos\omega, \sin\omega) \in S^2,$$

with $[0, 2\pi) \times [-\pi/2, \pi/2]$. For each of the frequencies $h^{-1} = \sqrt{l(l+1)}$ with $\ell \in \mathbb{N}$ we associate the Gaussian beam

$$\phi_h(\theta, \omega) = \frac{1}{2^l l!} \Big(\frac{2l+1}{4\pi(2l)!}\Big)^{\frac{1}{2}} e^{-il\theta}(\cos\omega)^l.$$

It is normalized so that

$$\|\phi_h\|_{L^2(S^2)} = 1, \qquad (-h^2\Delta_{S^2} - 1)\phi_h = 0.$$

Then, let $\chi \in C_c^\infty(-1, 1)$ with $\chi \equiv 1$ on $[-\frac{1}{2}, \frac{1}{2}]$ and define

$$u_h(\theta, \omega) = \frac{1}{2^l l!} \Big(\frac{2l+1}{4\pi(2l)!}\Big)^{\frac{1}{2}} e^{-il\theta}\chi(\omega)e^{-l\omega^2/2}.$$

Observe that

$$u_h - \phi_h = o_{L^2}(1),$$

so for the purposes of computing the defect measure, we may compute with u_h. Using this, by an elementary stationary phase argument, (see e.g. [Zwo12, Section 5.1]) the defect measure associated to ϕ_h is

$$\mu = \frac{1}{2\pi}\delta_{\{\omega=0,\xi=-1,\zeta=0\}}d\theta$$

where ξ is dual to θ and ζ is dual to ω. Let $H = \{(\theta, \omega) : \omega = 0\}$ be the equator. In particular, $N^*H = \{(\theta, \omega, \xi, \zeta) \in S^*S^2 : \omega = 0, \xi = 0, \zeta = \pm 1\}$. Then,

$$\mu_H(N^*H) = \mu(\{\omega \in (-t_0, t_0), \xi = 0, \zeta = \pm 1\}) = 0$$

and Theorem 1 implies

$$\int_H \phi_h(\theta, 0) d\theta = o(1).$$

6. Appendix on Semiclassical notation

We next review the notation used for semiclassical operators and symbols and some of the basic properties. First, recall that for a compact manifold M of dimension n, we write

$$S^m(T^*M) := \{a(\cdot; h) \in C^\infty(T^*M) : \ |\partial_x^\alpha \partial_\xi^\beta a(x, \xi; h)| \le C_{\alpha\beta}(1 + |\xi|)^{m-|\beta|}\}.$$

We write $\Psi^m(M)$ for the semiclassical pseudodifferential operators of order m on M and

$$Op_h : S^m(T^*M) \to \Psi^m(M)$$

for a quantization procedure with $Op_h(1) = \mathrm{Id} + O_{\mathcal{D}' \to C^\infty}(h^\infty)$ and for u supported in a coordinate patch, $\varphi \in C_c^\infty(M)$ with $\varphi \equiv 1$ on $\mathrm{supp}\, u$ we have

$$Op_h(a)u(x) = \frac{1}{(2\pi h)^n} \iint e^{\frac{i}{h}\langle x-y, \xi\rangle} \varphi(x) a(x, \xi) u(y) d\xi dy + O_{\mathcal{D}' \to C^\infty}(h^\infty)u.$$

Then there exists a principal symbol map

$$\sigma : \Psi^m(M) \to S^m(T^*M)/hS^{m-1}(T^*M)$$

so that

$$Op_h \circ \sigma(A) = A + O_{\Psi^{m-1}}(h), \quad A \in \Psi^m, \qquad \sigma \circ Op_h = \pi : S^m \to S^m/hS^{m-1},$$

where π is the natural projection map. Moreover, for $A \in \Psi^{m_1}$, $B \in \Psi^{m_2}$,

- $\sigma(AB) = \sigma(A)\sigma(B) \ \in S^{m_1+m_2}/hS^{m_1+m_2-1}$,

- $\sigma([A, B]) = \frac{h}{i}\{\sigma(A), \sigma(B)\} \ \in hS^{m_1+m_2-1}/h^2 S^{m_1+m_2-2}$,

where $\{\cdot, \cdot\}$ denotes the poisson bracket. For more details on the semiclassical calculus see e.g. [Zwo12, Chapters 4,14] [DZ16, Appendix E].

Finally, we recall the for any $\{u(h)\}_{0<h<h_0} \subset L^2(M)$ a bounded family of functions, we may extract a subsequence $h_k \to 0$ so that for $a \in C_c^\infty(T^*M)$,

$$\langle Op_h(a)u_{h_k}, u_{h_k}\rangle_{L^2(M)} \underset{h_k \to 0}{\to} \int a(x, \xi) d\mu$$

for a positive Radon measure μ. We call μ a *defect measure for u_{h_k}*. For $p \in S^m(T^*M)$ real valued, if $u(h)$ solves

$$Op_h(p)u = o(h), \qquad \|u(h)\|_{L^2} = 1,$$

then for any defect measure μ associated to $u(h)$,

$$\mathrm{supp}\, \mu \subset \{p(x, \xi) = 0\}, \qquad \exp(tH_p)_* \mu = \mu$$

where H_p denotes the Hamiltonian vector field associated to p. See e.g. [Zwo12, Chapter 5] for more details.

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