

# The ExoMol database: molecular line lists for exoplanet and other hot atmospheres

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## Abstract

The ExoMol database ([www.exomol.com](http://www.exomol.com)) provides extensive line lists of molecular transitions which are valid over extended temperatures ranges. The status of the current release of the database is reviewed and a new data structure is specified. This structure augments the provision of energy levels (and hence transition frequencies) and Einstein  $A$  coefficients with other key properties, including lifetimes of individual states, temperature-dependent cooling functions, Landé  $g$ -factors, partition functions, cross sections,  $k$ -coefficients and transition dipoles with phase relations. Particular attention is paid to the treatment of pressure broadening parameters. The new data structure includes a definition file which provides the necessary information for utilities accessing ExoMol through its application programming interface (API). Prospects for the inclusion of new species into the database are discussed.

*Keywords:* infrared, visible, Einstein  $A$  coefficients, transition frequencies, partition functions, cooling functions, lifetimes, cross sections,  $k$  coefficients, Landé  $g$ -factors

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## 1. Introduction

Hot molecules exist in many environments in space including cool stars [1], failed stars generally known as brown dwarfs [1, 2] and exoplanets [3]. The atmospheric properties of these objects are known to be strongly influenced by the spectra of the molecules they contain. On Earth, spectra of hot molecules are observed in flames [4, 5], discharge plasmas [6, 7], explosions [8] and in the hot gases emitted, for example, from smoke stacks [9]. In addition, high-lying states can be important in non local thermodynamic equilibrium (LTE) environments both in space, for example, emissions observed from comets [10, 11], and on Earth. The spectra of key atmospheric molecules at room temperature have been the subject of systematically maintained databases such as HITRAN [12–14] and GEISA [15, 16]. As will be amply illustrated below, the spectra of hot molecules contain many, many more transitions and so far attempts to compile systematic databases have been limited. Databases of hot molecular spectra do exist for other specialized applications, such as the EM2C database for combustion applications [17] or one due to Parigger *et al* for studies of laser-induced plasmas [18].

Planets and cool stars share some common fundamental characteristics: they are faint, their radiation peaks in the infrared and their atmosphere is dominated by strong molecular absorbers. Modelling planetary and stellar atmospheres is therefore difficult as their spectra are extremely rich in structure with hundreds of thousands to many billions of spectral lines which may be broadened by high-pressure and temperature effects.

The ExoMol project [19] aims to provide the molecular line lists that astronomers need in order to understand the physics and chemistry of astronomical bodies cool enough to form molecules in their atmospheres. In particular for extrasolar planets, brown dwarfs and cool stars [2, 3, 20], as well as circumstellar

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structures such as planetary envelopes and molospheres [21]. In practice these data are also useful for a wide range of other scientific disciplines; examples include studies of the Earth’s atmosphere [22, 23], hypersonic non-equilibrium [24], analysis of laboratory spectra [25–27], measurements of hot reactive gases [28] and the proposed remote analysis of molecular composition using laser ablation [29].

The original ExoMol data structure was very focused in its goal, with the scope of the data limited to generating lists of transitions [30]. However, it has become obvious that the potential applications of the ExoMol spectroscopic data are much more diverse. For example, the ExoMol data can be used to compute partition functions [31], cross sections [32], lifetimes [33], Landé  $g$ -factors [34] and other properties. Our aim is to systematically provide this additional data to maximise its usefulness. To do this requires significant extension of the ExoMol data structure, which is the major purpose of this paper. At the same time this implementation should facilitate the adoption of an application programming interface (API) between the database and programs using it. Similar enhancements are actively being pursued by other related databases such as HITRAN [35–37]. A major new feature is the inclusion, albeit at a fairly crude level, of pressure-broadening parameters. These have been shown to be important for models of exoplanets [38, 39] and are known to be vital for many other applications. Section 2 summarizes current data coverage of the ExoMol database. Section 3 reviews the individual line lists included in the database with particular emphasis on any changes made to them since they were originally published. Section 4 describes the expansion of data that is now provided by the ExoMol database. Section 5 gives a formal description of the new data structures implemented to cover this provision and to provide the API functionality. Section 6 briefly describes the utility programs available as part of the ExoMol database and Section 7 the ExoMol website. Finally we discuss future prospects for the database.

## 2. Database coverage

The ExoMol project aims at complete coverage of the spectroscopic properties of molecules which are deemed to be important in hot astrophysical environments. Coverage concerns (a) the molecular species considered, including isotopologues; (b) the frequency range considered and (c) the upper temperature range for which the data is reasonably complete. Both the required temperature and frequency range completeness are to some extent a judgement on what is required for astronomical and other studies. For example, a molecule like nitric acid ( $\text{HNO}_3$ ) will dissociate at a relatively low temperature so coverage above about 500 K is unlikely to be important. On the other hand, several diatomics such SiO and CO are known to feature in stellar spectra, and coverage to temperatures over 6000 K is necessary.

The general ExoMol approach is molecule-by-molecule. That is, a comprehensive line list is created for a particular molecule which is then made available in the database. For more challenging larger systems such as  $\text{NH}_3$ ,  $\text{PH}_3$  and  $\text{HNO}_3$ , it has been our policy to produce initial, room-temperature line lists, see [40–42] respectively. This allows us to improve the model, validate the available experimental data and, in some cases, make spectral assignments. In all instances, the subsequent hot line lists [43–45] are both more complete and more accurate and should therefore be used even for studies at low temperatures.

Thus far ExoMol has not considered ultraviolet (UV) absorption; however, there is increasing interest in the consequences of UV radiation on exoplanets [46–48] so this may need to be reviewed in future. As discussed in the comments below, molecules are still being added to the database so that coverage by species is steadily increasing.

Tables 1 and 2 summarize the molecules for which the ExoMol database currently provides data. The division is between species that have explicitly

been studied as part of the ExoMol project (Table 1) and those for which the data has been taken from other studies (Table 2). The ExoMol project has a specific methodology based on the use of spectroscopically-determined potential energy surfaces and *ab initio* dipole surfaces, which are combined with explicit variational treatments of the nuclear motion problem. For open shell systems these treatments involve explicit inclusion of spin-orbit and related curve-coupling effects; the project has developed a nuclear motion program, DUO, for treating coupled diatomic curves [49]. For closed shell molecules a range of codes [50–55] are used. These are all essentially based on finding near-exact solutions of the ro-vibrational Schrödinger equation for a given potential energy surface but the level of approximation increases with the size of the molecule considered.

Data from other sources arise from a variety of methodologies which range from completely *ab initio*, appropriate for systems with very few electrons, to largely empirical. Table 2 gives a pointer to the method used in each case; for full details the reader should consult the cited reference.

The molecules listed in Tables 1 and 2 do not at present provide a comprehensive set of species. A number of key species are available from other sources. HITEMP [87] in principle provides a source of data for CO, OH, NO and CO<sub>2</sub>. In practice there are new hot line lists for CO [88], OH [89], and CO<sub>2</sub> [90, 91] which are more recent than the data given in the current release of HITEMP. High quality line lists for ozone are also available from elsewhere [92]. Line lists for some other missing species may also be found in the UGAMOP (<https://www.physast.uga.edu/ugamop/>) and Kurucz [93] databases. The coverage of these databases for a (possible) exoplanet characterization mission has recently been reviewed by two of us [94].

Finally we note that while the majority of exoplanet spectroscopy is performed at rather low resolution, very precise spectroscopic data can be required [95–97], which is proving an issue for particular key species [98]. Therefore, ExoMol aims

Molecule	$N_{\text{iso}}$	$T_{\text{max}}$	$N_{\text{elec}}$	$N_{\text{lines}}^a$	DSName	Reference
BeH	1	2000	1	16 400	Yadin	Yadin <i>et al</i> [56]
MgH	3	2000	1	10 354	Yadin	Yadin <i>et al</i> [56]
CaH	1	2000	1	15 278	Yadin	Yadin <i>et al</i> [56]
SiO	5	9000	1	254 675	EJBT	Barton <i>et al</i> [57]
HCN/HNC	2 <sup>a</sup>	4000	1	399 000 000	Harris	Barber <i>et al</i> [58]
CH <sub>4</sub>	1	1500	1	9 819 605 160	YT10to10	Yurchenko & Tennyson [59]
NaCl	2	3000	1	702 271	Barton	Barton <i>et al</i> [60]
KCl	4	3000	1	1 326 765	Barton	Barton <i>et al</i> [60]
PN	2	5000	1	142 512	YYLT	Yorke <i>et al</i> [61]
PH <sub>3</sub>	1	1500	1	16 803 703 395	SAITY	Sousa-Silva <i>et al</i> [62]
H <sub>2</sub> CO	1	1500	1	10 000 000 000	AYTY	Al-Refaie <i>et al</i> [62]
AlO	4	8000	3	4 945 580	ATP	Patrascu <i>et al</i> [63]
NaH	2	7000	2	79 898	Rivlin	Rivlin <i>et al</i> [64]
HNO <sub>3</sub>	1	500	1	6 722 136 109	AlJS	Pavlyuchko <i>et al</i> [45]
CS	8	3000	1	548 312	JnK	Paulose <i>et al</i> [65]
CaO	1	5000	5	21 279 299	VBATHY	Yurchenko <i>et al</i> [66]
SO <sub>2</sub>	1	2000	1	1 300 000 000	ExoAmes	Underwood <i>et al</i> [67]

$N_{\text{iso}}$  Number of isotopologues considered;

$T_{\text{max}}$  Maximum temperature for which the line list is complete;

$N_{\text{elec}}$  Number of electronic states considered;

$N_{\text{lines}}$  Number of lines: value is for the main isotope.

<sup>a</sup> A line list for H<sup>13</sup>CN/HN<sup>13</sup>C due to Harris *et al* [68] is also available.

Table 1: Datasets created by the ExoMol project and included in the ExoMol database.

to be both complete, and as accurate as possible.

### 3. Individual line lists

Below we consider some of the line lists presented in the ExoMol database and listed in Tables 1 and 2. We restrict our discussion to issues not covered in the original publication.

One general issue is that Medvedev and co-workers [88, 99] identified a numerical problem with the intensities of high overtone transitions computed with the standard compilation of the diatomic vibration-rotation program LEVEL [50]. Our line lists computed with LEVEL have been adjusted to remove transitions

which appeared to have been affected by this issue; such cases are noted below. The transitions removed are all very weak and it is anticipated that these changes will have very little effect on practical applications. We checked similar calculations performed with our in-house rovibronic program DUO [49] and found similar behaviour: when the value of the transition dipole moment becomes comparable with the double precision error ( $\approx 10^{-16}$  a.u.), the corresponding electric dipole intensities essentially represent numerical noise and have to be removed.

### 3.1. MgH

The ExoMol MgH line list only considers transitions within the X  $^2\Sigma^+$  ground electronic state [56]. A line list containing A  $^2\Pi - X ^2\Sigma^+$  and B'  $^2\Sigma^+ - X ^2\Sigma^+$  transitions has been given by GharibNezhad *et al* [74]. Both line lists are included in the ExoMol database.

Furthermore Szidarovszky and Császár [100] showed that, due to the relatively low dissociation energy of MgH, consideration of quasi-bound states in the partition sum significantly alters the partition function at higher temperatures. However, for reasons of self-consistency we recommend using the partition function given by Yadin *et al* [56].

### 3.2. CaH

The ExoMol CaH line list only considers transitions within the X  $^2\Sigma^+$  ground electronic state [56]. A line list containing E  $^2\Pi - X ^2\Sigma^+$  transitions has been given by Li *et al* [75]. Both line lists are included in the ExoMol database.

### 3.3. SiO

The original SiO ExoMol line lists of Barton *et al* [57] computed using LEVEL have been truncated by removing transitions with  $\Delta v \geq 6$ .

### 3.4. HCN/HNC

The combined HCN/HNC ExoMol line list of Barber *et al* [58] used a calculated energy separation (isomerization energy) of  $5705\text{ cm}^{-1}$  between the ground states of HCN and HNC. Recent work by Nguyen *et al* [101] suggests that this value is too high. We have done some additional *ab initio* calculation using MOLPRO [102] at multi-reference configuration interaction (MRCI) level of theory. The complete basis set extrapolation (CBS) value CBS[56]z was found to equal  $5398.62\text{ cm}^{-1}$ . Subtracting from this value the zero point energy difference between HCN and HNC wells of  $88\text{ cm}^{-1}$ , gives the value for the energy separation of  $5310\text{ cm}^{-1}$ , closer to the value of  $5236 \pm 50\text{ cm}^{-1}$  recommended by Nguyen *et al* [101]. This value has been adopted by revising the states file and calculated partition function.

### 3.5. CH<sub>4</sub>

Although the YT10to10 line list [59] represents a major step forward in the modelling of hot methane spectra [103], it does not represent a complete solution to the problem. Rey *et al* [104] have also produced a line list, using similar procedures to those adopted by ExoMol, valid to higher temperatures (2000 K) but covering a more limited spectral range. The temperature coverage can be checked using available, high-temperature partition functions [105, 106].

Yurchenko *et al* [107] have extended the YT10to10 line list to higher temperatures (2000 K) and improved the predicted frequencies by replacing computed energy levels with empirical ones provided by Boudon [108]. The resulting line list has 35 billion transitions and Yurchenko *et al* propose that most of the lines in the line list can be represented by background cross sections. These temperature-dependent, but pressure-independent cross sections, generated using the vast majority of the lines, can then be used to supplement a reduced line list of 203 million lines selected to be the strongest in each spectral region. If



this methodology proves successful, it will be adopted by the ExoMol database for other very extensive line lists in future releases. At present these data are not included in the database.

### 3.6. *NaCl*

The original NaCl ExoMol line lists of Barton *et al* [60] computed using LEVEL have been truncated by removing transitions with  $\Delta v \geq 8$ .

### 3.7. *KCl*

The original KCl ExoMol line lists of Barton *et al* [60] computed using LEVEL have been truncated by removing transitions with  $\Delta v \geq 12$ .

### 3.8. *PN*

The original PN ExoMol line lists of Yorke *et al* [61] computed using LEVEL have been truncated by removing transitions with  $\Delta v \geq 6$ .

### 3.9. *AlO*

A number of ExoMol users pointed out that the vibrational labels used in the original AlO line list were not the best ones. Our analysis of AlO state lifetimes [109] reinforced this impression. The `.states` file has therefore been updated with revised vibrational state labels. We note that these labels are only approximate quantum numbers and this change should not affect any results obtained with the line list.

### 3.10. $H_3^+$

The  $H_3^+$  line list of Neal *et al* [69] (NMT) is the oldest in the ExoMol database. While it has continued to demonstrate its accuracy and predictive power [110, 111], perhaps surprisingly so, there are other issues with it. The line list was constructed using Jacobi coordinates which do not allow for a full treatment of

the symmetry. This led to the use of approximate nuclear spin statistical weights and, more problematically, to the removal of very weak transitions. Most of these transitions are symmetry-forbidden and should have zero dipole intensity. However “forbidden” rotational transitions [112, 113] have proved important for both astrophysical [114, 115] and laboratory [116, 117] studies. Methods are available which allow for the proper treatment of symmetry [118, 119]. A new line list [120] is nearly complete which extends the range of the NMT line list. The upper energy limit is increased to 25 000  $\text{cm}^{-1}$  and the highest rotational quantum number state considered in the calculations is  $J = 40$ . This line list will remove the symmetry problem as well as improve both the accuracy and coverage for  $\text{H}_3^+$ , thus making it useful for both higher and lower temperatures.

### 3.11. $\text{H}_2\text{O}$

The BT2 line list [70] has been outstandingly successful: it was used for the original detection of water in an exoplanet atmosphere [121], forms the basis of the well-used BT-Settl brown-dwarf model [122] and provides the hot water line list for HITEMP [87]. However, although it is more complete, for many levels it is less accurate than the Ames line list of Partridge and Schwenke [123]. Since the construction of BT2 there has been significant progress in deriving experimental energy levels for  $\text{H}_2^{16}\text{O}$  [124, 125], improvements in the *ab initio* dipole moment surface [126] which have been used to construct high accuracy room temperature line lists [127] and further improvements in representing the potential energy surface of the molecule [128]. A new  $\text{H}_2^{16}\text{O}$  line list, named POKAZATEL, [129] which builds on these advances will be released soon. The completeness of this line list is illustrated by the following figures: The highest rotational quantum number used in the calculations of the wavefunctions was  $J = 72$  — close to the highest  $J$  for which the bound states still do exist. The highest upper energy levels limit is equal to 40 000  $\text{cm}^{-1}$ , which is also close to the dissociation limit of water

[130].

While the BT2 and Ames line lists are seriously incomplete for temperatures above 3000 K, POKAZATEL considers every ro-vibrational transition in the molecule and therefore is appropriate for studies at higher temperatures.

Line lists for hot  $\text{H}_2^{17}\text{O}$  and  $\text{H}_2^{18}\text{O}$  will also be released very shortly [131]. Again these will take advantage of the available set of experimental energy levels [132, 133] obtained using the MARVEL procedure [134, 135]. These line lists will be as complete as BT2 and consider levels with  $J \leq 50$  and energies up to 30 000  $\text{cm}^{-1}$ .

### 3.12. $\text{NH}_3$

Ammonia is a ten electron system like water, so in terms of electronic calculations one might expect similar improvements here too.  $^{14}\text{NH}_3$  has also been the subject of a systematic study of experimental energy levels [136] and improvements in its *ab initio* treatment [137, 138]. New experimental data characterizing higher-lying energy levels is also available [26, 139]. Again the Ames group have performed high accuracy studies on this system [137, 140, 141], although there is no corresponding line list. Work is therefore in progress to compute a new ammonia line list which will replace the existing BYTe line list [43]; in particular this work the option of using the improved nuclear motion capability of the program TROVE working in curvilinear coordinates [54]. Recent, excellent *ab initio* calculations [138] which reproduce very highly excited energy levels of ammonia up to 18 000  $\text{cm}^{-1}$ , will be used as the starting point for a semi-empirical fit of the PES. This should give a line list which is both more accurate and more extensive than BYTe.

### 3.13. $\text{CrH}$

Chromium hydride is an important astronomical molecule for a variety of reasons, including classification of brown dwarfs [142, 143] and measurements of

magnetic fields [144]. The ExoMol database currently provides a substantially empirical line list due to Burrows *et al* [81]; we are in the process of computing a more complete CrH line list covering 8 states and the 4 main isotopologues [145].

### 3.14. *TiH*

The current titanium hydride line list due to Burrows *et al* [83] is in the process of being updated with a more extensive ExoMol one.

### 3.15. *Other species*

Besides the species listed above, line lists are in an advanced stage of construction for VO [146], H<sub>2</sub>O<sub>2</sub> [147] (see Al-Refaie *et al* [148] for a preliminary, room temperature version), SO<sub>3</sub> [149] (see also Underwood *et al* [150]), C<sub>2</sub>H<sub>4</sub>, CH<sub>3</sub>Cl (see also Owens *et al* [151]), C<sub>2</sub>, SiH, NS, NO, NaO, AlH and SH [152], MnH, and PO, PS and PH [153].

Finally we note that work is in progress looking at TiO [154]. TiO is a major absorber in cool stars [155, 156]. There are TiO line lists available from Schwenke [157] and Plez [158]. However it would appear that these line lists are incomplete; in particular there are missing bands that have been observed in the laboratory [159]. Furthermore, recent exoplanet studies suggest that the line frequencies are too inaccurate for high resolution studies [98].

## 4. Data Provided

The original aim of the ExoMol project was the provision of extensive line lists of energy levels, and hence transition frequencies, and Einstein *A* coefficients from which transition intensities and related information can be computed. The original ExoMol data structure [30] provided precisely for these data in a concise manner. The demands on and utility of the data provided by the project has led us to expand the scope of the database. In this section we outline the new data

provision and in the following sections we formally define the new data structures employed for this purpose.

#### *4.1. Extended states file*

There are other properties of a molecular state that can be computed by ExoMol which we now propose to store. The first of these is the radiative lifetime of each individual state, which can be computed in a straightforward fashion from the Einstein  $A$  coefficients [33]. Such lifetimes have been shown to be important in both laboratory [116, 117] and astronomical [114, 115] environments. Indeed the radiative lifetime of a state is a key component in determining the critical density of species in that state and hence whether it exists in local thermodynamic equilibrium in a given environment. Our recent study of lifetimes based on use of ExoMol data suggested that methane should show particularly interesting lifetime effects [33]. These computed lifetimes can be compared with experimental measurements, where available, providing a check on our calculated dipole moments. This is particularly useful for rovibronic spectra, where measurements of absolute transition intensities are unusual.

The second state-dependent property are Landé  $g$ -factors. These provide the behaviour of the states in the presence of a weak magnetic field as given by the Zeeman effect. Molecular spectra can provide important information on magnetic fields in stars, brown dwarfs and exoplanets [144, 160]. For open shell diatomics, Landé  $g$ -factors are given by a fairly straightforward formula in terms of standard Hund's case (a) quantum numbers [161]. The vibronic wavefunctions computed by DUO [49] contain the necessary information to compute these  $g$ -factors and this has currently recently been implemented within DUO [34]. The magnetic fields in some astronomical environments are strong enough that the Paschen-Back effect becomes important [144, 162]. In this case the shifts in the energy levels do not have a simple dependence on the magnetic field [163] and they will

have to be explicitly computed on a case-by-case basis. Landé  $g$ -factors are only given for systems whose spectroscopic model includes open shell electronic states (ie ones with unpaired electrons) since these are the systems whose states which show significant splitting in a magnetic field.

Besides (approximate) quantum numbers, which may appear in several forms, it is also often desirable to store more than one estimate of the energy level of the state. A general methodology, known as MARVEL, is available for obtaining self-consistent sets of empirical energy levels from networks of transitions [134, 135], which means that sets of empirical energy levels are available for key molecules [124, 132, 133, 164–166]. In this case it is preferable to replace computed energy levels with empirical ones, as has already been done for the HCN/HNC [58] and the CS [65] line lists. It is also recommended to keep the original computed levels as well as the uncertainties which are generally available for the empirical levels. The inclusion of approximate quantum numbers and extra energy levels included has been decided on a molecule-by-molecule basis; information on this is contained in the isotopologue definition file (see Section 5.3).

Molecule	$N_{\text{iso}}$	$T_{\text{max}}$	$N_{\text{elec}}$	$N_{\text{lines}}$	DSName	Reference	Methodology
$\text{H}_3^+$	2 <sup>a</sup>	4000	1	3 070 571	NMT	Neal <i>et al</i> [69]	ExoMol
$\text{H}_2\text{O}$	2 <sup>b</sup>	3000	1	505 806 202	BT2	Barber <i>et al</i> [70]	ExoMol
$\text{NH}_3$	2 <sup>c</sup>	1500	1	1 138 323 351	BYTe	Yurchenko <i>et al</i> [43]	ExoMol
$\text{HeH}^+$	4	10000	1	1 431	Engel	Engel <i>et al</i> [71]	Ab initio
$\text{HD}^+$	1	12000	1	10 119	CLT	Coppola <i>et al</i> [72]	Ab initio
$\text{LiH}$	1	12000	1	18 982	CLT	Coppola <i>et al</i> [72]	Ab initio
$\text{LiH}^+$	1	12000	1	332	CLT	Coppola <i>et al</i> [72]	Ab initio
$\text{ScH}$	1	5000	6	1 152 827	LYT	Lodi <i>et al</i> [73]	Ab initio
$\text{MgH}$	1		3	30 896	13GhShBe	GharibNezhad <i>et al</i> [74]	Empirical
$\text{CaH}$	1		2	6000	11LiHaRa	Li <i>et al</i> [75]	Empirical
$\text{NH}$	1		1	10 414	14BrRaWe	Brooke <i>et al</i> [76]	Empirical
$\text{CH}$	2		4	54 086	14MaPIVa	Masseron <i>et al</i> [77]	Empirical
$\text{CN}$	1		1	195 120	14BrRaWe	Brooke <i>et al</i> [78]	Empirical
$\text{CP}$	1		1	28 735	14RaBrWe	Ram <i>et al</i> [79]	Empirical
$\text{HCl}$	1		1	2588	11LiGoBe	Ram <i>et al</i> [80]	Empirical
$\text{CrH}$	1		2	13 824	02BuRaBe	Burrows <i>et al</i> [81]	Empirical
$\text{FeH}$	1		2	93 040	10WEReSe	Wende <i>et al</i> [82]	Empirical
$\text{TiH}$	1		3	181 080	05BuDuBa	Burrows <i>et al</i> [83]	Empirical

$N_{\text{iso}}$  Number of isotopologues considered;

$T_{\text{max}}$  Maximum temperature for which the line list is complete;

$N_{\text{elec}}$  Number of electronic states considered;

$N_{\text{lines}}$  Number of lines: value is for the main isotope.

<sup>a</sup> There is a  $\text{H}_2\text{D}^+$  line list available from Sochi and Tennyson [84].

<sup>b</sup> The VTT line list for  $\text{HDO}$  due to Voronin *et al* [85] is also available.

<sup>c</sup> There is a room temperature  $^{15}\text{NH}_3$  line list due to Yurchenko [86].

Table 2: Datasets not created as part of the ExoMol project but included in the ExoMol database.

#### 4.2. Cross sections, $k$ -coefficients and pressure effects

Many of the line lists stored in the database are huge. This makes finding other forms of storing the data desirable. ExoMol data has been used to generate temperature-dependent cross sections as a function of wavenumber [32]. Such cross sections now form part of the database which will also be expanded to contain pressure-dependent cross sections and  $k$ -coefficient tables generated from them.

Cross sections depend on pressure, as well as temperature, due to collisional broadening by the species in the atmosphere. The provision of temperature-, pressure- and broadening species-dependent cross sections would be both computationally very demanding and again result in very large datasets. Instead we have chosen to primarily provide parameters which characterise pressure profiles for key species plus software allowing cross sections to be generated for selected temperature-pressure parameters on the user's own computer. There is a widespread recognition that for detailed studies, such as those performed in the Earth's atmosphere at high resolution, Voigt profiles only provide an approximate solution to the problem [167]. However Voigt profiles are in widespread use and are easily computed [168]; they are therefore used to represent pressure broadening effects in ExoMol.

The new release of ExoMol includes, where possible, pressure broadening parameters. So far, these have not been considered for vibronic spectra. HITRAN provides a source of air-broadening and self-broadening parameters [169] so in the absence of other sources these are taken from HITRAN 2012 [170]. In gas giants  $\text{H}_2$  and He are the major broadeners. There is no systematic source of parameters for these broadeners so we have to treat each system on a case-by-case basis. While some work has been performed for broadening of hot water by  $\text{H}_2$  and He [171, 172], more data would clearly be helpful in this area. Hedges and Madhusudhan [39] recently put this in context by quantifying the effects of var-



ious factors involved in modelling line broadening in exoplanetary atmospheres, including completeness of broadening parameters, on molecular absorption cross sections.

ExoMol will provide Lorentzian half-widths and their temperature-exponents for key molecule-broadener systems. The availability of these parameters varies with species and broadener. In any case there are not individual values for every molecular line and often parameters have only been measured or calculated for a small fraction of transitions. Therefore, rather than pre-determine values for every molecular line and increase the size of the line lists, which are very large for some molecules, a separate pressure broadening parameters file is provided. This file contains three types of parameters: experimental, theoretical and semi-empirical. Experimental, theoretical and semi-empirical parameters from the literature are presented with their respective full or partial quantum number assignments. Additional semi-empirical parameters are determined by compiling all experimental and theoretical parameters as a function of  $J''$ , the total rotational quantum number of the lower level of the transition, and computing an average value for each  $J''$ . To avoid introducing additional error, no extrapolation is attempted beyond the  $J''_{\max}$  for which data are available, the parameters are simply assumed to be constant from this point. This is a very basic model similar to that used by other line-width studies [173–175]. Constructing the pressure broadening parameters file in this way ensures that parameters are provided for every spectral line. Table 3 lists the main sources used to provide pressure broadening parameters. Clearly further and improved parameters would be welcome.

Table 3: Sources of pressure broadening parameters for key molecule-broadener species currently being considered by ExoMol. Entries marked \* were extracted from HITRAN [170, 176].

Molecule	Broadener	Reference	Methodology
H <sub>2</sub> O	H <sub>2</sub>	Lavrentieva <i>et al</i> [177]	Semi-empirical
		Lavrentieva <i>et al</i> [178]	Semi-empirical
		Steyert <i>et al</i> [179]	Experimental
		Brown & Plymate [180]	Experimental
		Brown <i>et al</i> [181]	Experimental
		Gamache <i>et al</i> [182]	Theoretical
		Dick <i>et al</i> [183]	Experimental
		Faure <i>et al</i> [171]	Theoretical
		Langlois <i>et al</i> [184]	Experimental
		Dutta <i>et al</i> [185]	Experimental
		Golubiatnikov [186]	Semi-empirical
		Zeninari <i>et al</i> [187]	Semi-empirical
		Drouin and Wiesenfeld [188]	Theoretical
H <sub>2</sub> O	He	Lavrentieva <i>et al</i> [177]	Semi-empirical
		Lavrentieva <i>et al</i> [178]	Semi-empirical
		Steyert <i>et al</i> [179]	Experimental
		Gamache <i>et al</i> [182]	Theoretical
		Dick <i>et al</i> [183]	Experimental
		Lazarev <i>et al</i> [189]	Experimental
		Dutta <i>et al</i> [185]	Experimental
		Golubiatnikov [186]	Semi-empirical
Zeninari <i>et al</i> [187]	Semi-empirical		
H <sub>2</sub> O	Air	Mérienne <i>et al</i> [190]*	Experimental
		Gamache & Hartmann [191]*	Semi-empirical
		Gasster <i>et al</i> [192]*	Experimental
		Payne <i>et al</i> [193]*	Experimental
		Gamache [194]*	Semi-empirical
H <sub>2</sub> O	H <sub>2</sub> O	Gamache & Laraia [195]	Semi-empirical
		Mérienne <i>et al</i> [190]*	Experimental
		Gamache & Hartmann [191]*	Semi-empirical
		Markov [196]*	Experimental

*Continued on next page*

Table 3 – *Continued from previous page*

Molecule	Broadener	Reference	Methodology
CH <sub>4</sub>	H <sub>2</sub>	Golubiatnikov <i>et al</i> [197]*	Experimental
		Cazzoli <i>et al</i> [198]*	Semi-empirical
		Pine [199]	Experimental
		Margolis [200]	Experimental
		Fox <i>et al</i> [201]	Experimental
CH <sub>4</sub>	He	Strong <i>et al</i> [202]	Experimental
		Pine [199]	Experimental
		Varanasi & Chudamani[203]	Experimental
		Gabard <i>et al</i> [204]	Experimental
CH <sub>4</sub>	Air	Grigoriev <i>et al</i> [205]	Semi-empirical
		Fox <i>et al</i> [201]	Experimental
		Predoi-Cross <i>et al</i> [206]*	Experimental
		Smith <i>et al</i> [207]*	Experimental
		Brown <i>et al</i> [208]*	Semi-empirical
CH <sub>4</sub>	CH <sub>4</sub>	Anthony <i>et al</i> [209]	Theoretical
		Predoi-Cross <i>et al</i> [210]*	Experimental
		Smith <i>et al</i> [211]*	Experimental
NH <sub>3</sub>	H <sub>2</sub>	Brown <i>et al</i> [208]*	Semi-empirical
		Pine <i>et al</i> [212]*	Experimental
		Hadded <i>et al</i> [213]*	Experimental
NH <sub>3</sub>	He	Sharp & Burrows [214]*	Semi-empirical
		Pine <i>et al</i> [212]*	Experimental
		Hadded <i>et al</i> [213]*	Experimental
NH <sub>3</sub>	Air	Sharp & Burrows [214]*	Semi-empirical
NH <sub>3</sub>	NH <sub>3</sub>	Brown & Peterson [215]*	Experimental
PH <sub>3</sub>	H <sub>2</sub>	Brown & Peterson [215]*	Experimental
		Bouanich <i>et al</i> [216]	Semi-empirical
		Levy <i>et al</i> [217]	Experimental
		Sergent-Rozey <i>et al</i> [218]	Experimental
		Salem <i>et al</i> [219]	Experimental
		Pickett <i>et al</i> [220]	Experimental
PH <sub>3</sub>	He	Levy <i>et al</i> [221]	Experimental
		Salem <i>et al</i> [222]	Experimental
		Levy <i>et al</i> [217]	Experimental
		Sergent-Rozey <i>et al</i> [218]	Experimental
		Pickett <i>et al</i> [220]	Experimental

*Continued on next page*

Table 3 – *Continued from previous page*

Molecule	Broadener	Reference	Methodology
		Levy <i>et al</i> [221]	Experimental
PH <sub>3</sub>	Air	Butler <i>et al</i> [223]*	Experimental
		Kleiner <i>et al</i> [224]*	Semi-empirical
PH <sub>3</sub>	PH <sub>3</sub>	Butler <i>et al</i> [223]*	Experimental
		Kleiner <i>et al</i> [224]*	Semi-empirical
H <sub>2</sub> CO	H <sub>2</sub>	Nerf [225]	Experimental
H <sub>2</sub> CO	He	Nerf [225]	Experimental
H <sub>2</sub> CO	N <sub>2</sub>	Jacquemart <i>et al</i> [226]	Semi-empirical
H <sub>2</sub> CO	H <sub>2</sub> O	Jacquemart <i>et al</i> [226]*	Semi-empirical
HCN	Air	Yang <i>et al</i> [227]*	Experimental
		Devi <i>et al</i> [228]*	Semi-empirical
		Risland <i>et al</i> [229]*	Semi-empirical
HCN	HCN	Devi <i>et al</i> [228]*	Semi-empirical
		Devi <i>et al</i> [230]*	Semi-empirical
CS	Air	Blanquet <i>et al</i> [231]*	Semi-empirical
CS	CS	Misago <i>et al</i> [232]*	Semi-empirical

Another compact form of data input to exoplanet modelling codes involves the use of the  $k$ -coefficient approximation [233–235], for example, by the NEMESIS code [236] and cross sections, as used by Tau-REx [62, 237, 238]. Temperature- and pressure-dependent  $k$ -coefficients are also been provided for certain key species.

#### 4.3. Partition and cooling functions

In practice the ExoMol project always required partition functions, the study of which has often been performed independently [31, 239–241], and on occasion adopted from other studies [105]. Temperature-dependent partition functions are now formally included as part of the data structure.

Of course in various environments such as the early Universe and regions of star or planetary formation, the transformation of energy into radiation through molecular emissions provide an important source of cooling. Temperature-dependent cooling functions are important [72, 242, 243] and can also be computed from ExoMol line lists [33]. These are also now included in the data structure.

#### 4.4. *Dipoles*

While Einstein  $A$  coefficients are sufficient for the vast majority of radiative transport applications, the construction of these loses the phase information contained in the individual, complex-valued transition dipole moments. Incorporating the phase information in the ExoMol database will make it useful for theoretical modelling of the effects of an electric field on molecular structure. The areas of application includes cooling and trapping of molecular beams (see Ref. [244] and references therein), manipulating the long-range molecular interactions and collisional dynamics [245–247], molecular orientation and alignment [244, 248, 249], as well as rotational spectroscopy [250–253]. As shown in Section 5.7, the absolute values of transition dipole moments can be computed from the Einstein  $A$  coefficients and it is sufficient to complement this with only the sign of transition dipole moment for each molecular line.

#### 4.5. *Spectroscopic Models*

In addition to the spectroscopic data provided by the ExoMol database, each ExoMol line list has been constructed from a detailed spectroscopic model of the given molecule. In general this model includes a (spectroscopically-determined) potential energy surface and *ab initio* dipole moment surfaces, and input to the appropriate nuclear motion program. For systems with multiple electronic states, the spectroscopic model also includes coupling between electronic states, e.g. spin-orbit, electronic angular momentum etc.

These spectroscopic models are a valuable product of the ExoMol procedure and there are several reasons for preserving them. Firstly, the spectroscopic model is ultimately the source of error in the final line list and is thus important when investigating discrepancies between experiment or astronomical observations and theoretical predictions. Secondly, they can be used as a starting point for future models or line lists. For example, new experimental data can be incorporated

through refinement of potential energy curves; improved *ab initio* methodologies can be incorporated by improvements to the dipole moment surface, or further electronic states can be added to the model. Finally, they can be useful for other applications besides calculation of the original line list. Current work within the ExoMol group is investigating the sensitivity of spectral lines to a possible variation in the proton-to-electron mass ratio [254, 255].

An important part of each spectroscopic model is the input file to the appropriate nuclear motion program (e.g. LEVEL [50], DUO [49], DVR3D [51], TROVE [53]). For the diatomic codes LEVEL and DUO, the program input generally includes full details of the appropriate (potential, dipole and coupling) curves. For the polyatomic nuclear motion programs DVR3D and TROVE, the PES and DMS are provided as standalone functions, usually written in Fortran, with the program inputs as separate files. These spectroscopic models are not part of the formal ExoMol data structure described in the next section but can be accessed from the webpage of the appropriate isotopologue in files labelled `.model`.

## 5. Data Structures

The data structure outlined below represents a significant extension of the original ExoMol data structure described by Tennyson *et al* [30]. However the core structure of the two files, the states and transitions files, used in the original specification remains unchanged meaning that utilities designed to work with the previous structure will still work.

A summary of the contents of the database is stored in a master file, `exo-mol.all`. This file points towards the files which store the actual data. Table 4 gives an overview of these files, which are available for each isotopologue under the updated ExoMol structure. Each of these files is specified in turn below.

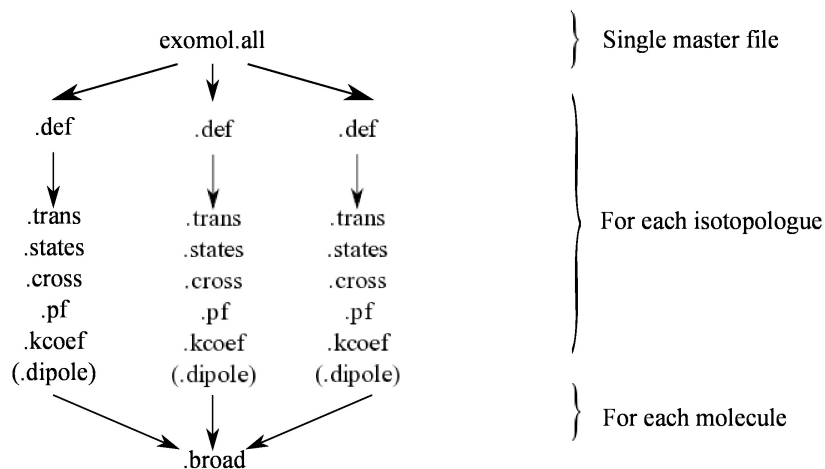


Figure 1: Structure of the ExoMol data files. There is a single master file describing the database,  $N_{\text{mol}}$  molecules and  $N_{\text{iso}}$  files per molecule.

Figure 1 gives a schematic representation of the file structure of the exomol data.

### 5.1. File naming convention

The naming convention for the files in the ExoMol database has the following specification. All file names, with the exception of the Pressure Broadening file, have a common part built using the template `Iso-slug__LineList__`. The “Iso-slug” is a machine-readable, plain ASCII-text “XML-safe” unique identifier for the molecular isotopologue, for example, `14N-1H3` is the iso-slug for  $^{14}\text{NH}_3$ . Ions are denoted by an appended `'_mZ'` or `'_pZ'` for charges  $-Z$  and  $+Z$  respectively (with the value of  $Z$  omitted if it is unity); so `1H3_p` is the iso-slug for  $^1\text{H}_3^+$ . “Linelist” is the line list publication name, see the column “DSName” in Tables 1 and 2; for example, the published names for the  $^{14}\text{NH}_3$  and  $\text{H}_3^+$  line lists are BYTe and NMT respectively. For external datasets that have been cast in ExoMol format, this name is an 8 character year-authors string which follows the convention of the IUPAC Task Group on water spectroscopy [132]. The extensive bibliography files, in BibTeX format, provided on the ExoMol website for each molecule in the database, uses the same naming convention.

File extension	$N_{\text{files}}$	File DSname	Contents
.all	1	Master	Single file defining contents of the ExoMol database..
.def	$N_{\text{tot}}$	Definition	Defines contents of other files for each isotopologue.
.states	$N_{\text{tot}}$	States	Energy levels, quantum numbers, lifetimes, (Landé $g$ -factors).
.trans	$N_{\text{tot}}$ <sup>a</sup>	Transitions	Einstein $A$ coefficients, (wavenumber).
.broad	$N_{\text{mol}}$	Broadening	Parameters for pressure-dependent line profiles.
.cross	$N_{\text{mol}}$ <sup>b</sup>	cross sections	Temperature or temperature and pressure-dependent cross sections.
.kcoef	$N_{\text{mol}}$ <sup>c</sup>	$k$ -coefficients	Temperature and pressure-dependent $k$ -coefficients.
.pf	$N_{\text{tot}}$	Partition function	Temperature-dependent partition function, (cooling function).
.dipoles	$N_{\text{tot}}$	Dipoles	Transition dipoles including phases.
.overview	$N_{\text{mol}}$	Overview	Overview of datasets available.
.readme	$N_{\text{iso}}$	Readme	Specifies data formats.
.model	$N_{\text{iso}}$	Model	Model specification.

$N_{\text{files}}$  total number of possible files.

$N_{\text{mol}}$  Number of molecules in the database

$N_{\text{tot}}$  is the sum of  $N_{\text{iso}}$  for the  $N_{\text{mol}}$  molecules in the database

$N_{\text{iso}}$  Number of isotopologues considered for the given molecule

<sup>a</sup> There are  $N_{\text{tot}}$  sets of .trans files but for molecules with large numbers of transitions the .trans files are subdivided into wavenumber regions.

<sup>b</sup> There are  $N_{\text{cross}}$  sets of .cross files for isotopologue.

<sup>c</sup> There are  $N_{\text{kcoef}}$  sets of .kcoef files for each isotopologue.

Table 4: Specification of the ExoMol file types. (Contents in brackets are optional.)

This common part of the file name can be followed by additional identifiers specific to different file types. For example, the transition file .trans for polyatomic molecules is usually split into smaller wavenumber ranges in order to simplify the manipulation of data. In this case the `Iso-slug__LineList__` is followed by `NUMIN-NUMAX`, which specifies the wavenumber range in  $\text{cm}^{-1}$ . The absorption cross section .cross file names also contain a wavenumber range, representing the full coverage of the line list, followed by the temperature in Kelvin, pressure in bar and the binning interval, or resolution (in  $\text{cm}^{-1}$ ), for which the cross sections have been computed. The file names have the form `NUMIN-NUMAX__Temperature [K] __Pressure [bar] __resolution` after the common part.

In case of the Pressure Broadening file .broad, the file name does not contain the string `LineList__` as the data in the broadening files is not line list specific and has been compiled from a number of sources. Here the string `Iso-slug__` is simply followed by `Broadener` identifying the broadener.



For example, the current BT2 line list package consists of the definition file

1H2-160\_\_BT2.def,

the States file

1H2-160\_\_BT2.states,

sixty Transitions files:

1H2-160\_\_BT2\_\_00000-00500.trans,

1H2-160\_\_BT2\_\_00500-01000.trans,

...

1H2-160\_\_BT2\_\_29500-30000.trans,

Partition function file 1H2-160\_\_BT2.pf,

Pressure Broadening files:

1H2-160\_\_H2.broad,

1H2-160\_\_He.broad,

1H2-160\_\_air.broad,

1H2-160\_\_self.broad

and eighteen cross section files:

1H2-160\_\_BT2\_\_00000-30000\_\_296K\_\_0bar\_\_0.01.cross,

1H2-160\_\_BT2\_\_00000-30000\_\_400K\_\_0bar\_\_0.01.cross,

...

1H2-160\_\_BT2\_\_00000-30000\_\_900K\_\_0bar\_\_0.01.cross,

1H2-160\_\_BT2\_\_00000-30000\_\_1000K\_\_0bar\_\_0.01.cross,

1H2-160\_\_BT2\_\_00000-30000\_\_1200K\_\_0bar\_\_0.01.cross,

...

1H2-160\_\_BT2\_\_00000-30000\_\_3000K\_\_0bar\_\_0.01.cross,

Here 1H2-160 is the iso-slug for H<sub>2</sub><sup>16</sup>O, BT2 is the name of the water line list [70], 00000-00500 is the wavenumber range 0-500 cm<sup>-1</sup>, H<sub>2</sub>, He, air and self are the broadeners, 296 K is the temperature, 0 bar is the pressure and 0.01 is the

Field	Fortran Format	C Format	Description
<b>Header Information</b>			
ID	A13	%13s	Always the ASCII string "EXOMOL.master"
V <sub>all</sub>	I8	%8d	Version number, format YYYYMMDD, recording last update of the whole database
N <sub>mol</sub>	I4	%4d	Number of molecules in the database
Followed by $N_{\text{mol}}$ entries of the form			
<b>Molecule Information</b>			
N <sub>name</sub>	I3	%3d	Number of molecule names listed.
Followed by $N_{\text{name}}$ entries of the form			
MolName	A27	%27s	(Common) name of the molecule
MolFormula	A27	%27s	Molecule chemical formula
N <sub>iso</sub>	I4	%4d	Number of isotopologues considered
Followed by $N_{\text{iso}}$ entries of the form			
MolKey	A27	%27s	Inchi key of the isotopologue
IsoFormula	A27	%27s	Isotopologue chemical formula
Iso-slug	A160	%160s	Isotopologue slug identifier, see text for details
DSName	A10	%10s	Isotopologue dataset name
V	I8	%8d	Version number with format YYYYMMDD
Metadata	Free format		

Table 5: Format of the ExoMol master file, `exomol.all`; each entry starts on a new line.

binning interval for the cross sections in wavenumbers.

For most cases a single `.broad` file is provided for a given molecule and designated for the parent (most abundant) isotopologue. Should broadening parameters be required for other isotopologues, this file should be employed. In a few cases where isotopic substitution lowers the symmetry, e.g. HDO as compared to H<sub>2</sub>O, then extra isotopologue-specific broadening files may be provided, although thus far there are no such files.

### 5.2. The master file

The file `exomol.all` is designed to be machine searchable making it easy to check the current contents of the database and when any of it was last updated. The file is structured as a list of molecules with the isotopologues associated with that molecule. The actual ExoMol data files are stored by isotopologue. Table 5 specifies the format of the file while Table 6 gives a portion of the current contents for water with two isotopologues,  $\text{H}_2^{16}\text{O}$  and  $\text{HD}^{16}\text{O}$ . The final entry gives web-searchable meta-data.

---

<code>EXOMOL.master</code>	<code># ID</code>
<code>20160315</code>	<code># Version number with format YYYYMMDD</code>
<code>30</code>	<code># Number of molecules in the database</code>
<code>1</code>	<code># Number of molecule names listed</code>
<code>water</code>	<code># Name of the molecule</code>
<code>H2O</code>	<code># Molecule chemical formula</code>
<code>2</code>	<code># Number of isotopologues considered</code>
<code>XLYOFNOQVPJJNP-DYCDLGHISA-N</code>	<code># Inchi key of molecule</code>
<code>1H-2H-16O</code>	<code># Iso-slug</code>
<code>(1H)(2H)(16O)</code>	<code># IsoFormula</code>
<code>VTT</code>	<code># Isotopologue dataset name</code>
<code>20160314</code>	<code># Version number with format YYYYMMDD</code>
<code>XLYOFNOQVPJJNP-UHFFFAOYSA-N</code>	<code># Inchi key of molecule</code>
<code>1H2-16O</code>	<code># Iso-slug</code>
<code>(1H)2(16O)</code>	<code># IsoFormula</code>
<code>BT2</code>	<code># Isotopologue dataset name</code>
<code>20160220</code>	<code># Version number with format YYYYMMDD</code>

---

Table 6: Extract from the `exomol.all` master file showing the three line header and the portion for the water molecule.

### 5.3. The definition file

A new addition to the ExoMol format is the inclusion of the definition file with extension `.def`. The definition file gives information on what ExoMol provides for a particular isotopologue and describes how the data can be used. The definition file serves multiple purposes:

- **Standardized ExoMol file usage:** The ExoMol format [30] `.states` file contains 4 standard fields: The ID of the state, the energy, the total degeneracy and the  $J$  quantum number. However additional fields such as the symmetry and vibrational quantum numbers are not standardized and differ between molecules. For usage that requires these additional quantum numbers, a subroutine or function must be created for each molecule to be read. The definition file therefore provides a means to overcome this by providing a single subroutine with the information to read any state file provided by the ExoMol project. This allows codes to easily integrate the ExoMol format only once and provide support for all molecules in the ExoMol project without additional changes.
- **Improved database function:** The definition file also provides detailed information on all fields in the ExoMol format. This means that database functionality such as sorting, filtering, splitting, selecting etc on specific or multiple fields can be easily performed on single or multiple molecules.
- **Facilitate updates:** Finally, updates of particular molecules can be easily propagated to any user of the ExoMol format. Each definition file is assigned a version number; this allows the user to check whether there are any updates to a particular molecule, with the possibility of an automatic download of any modified files and immediate usage without change to codes.

The format is defined in Table 7. The data format can be subdivided into different sections (highlighted in bold) describing a different aspect of the molecule in the ExoMol format. Each section may contain multiple subsections that define further aspects of that section. The file format therefore takes a tree structure. An example definition file for the CS line list [65] is given in Table. 8 and for the BT2 water line list [70] in Table. 9.

### 5.3.1. Header

The first section describes metadata related to internal molecular data.  $V$  is the version number of the definition file, it utilises a numerical dating format  $YYYYMMDD$  in order to allow for easy chronological ordering of different versions by integer comparison. The `MolKey` is the 27 character InchiKey [256] used to identify the molecule and  $N_{\text{atom}}$  describes the number of atoms in the molecule. This allows code to modify its behaviour based on whether it is dealing with diatomic molecules  $N_{\text{atom}} = 2$  or polyatomics  $N_{\text{atom}} > 2$ .

Table 7: Definition file format; each entry starts on a new line.

Field	Fortran Format	C Format	Description
<b>Header Information</b>			
ID	A10	%10s	Always the ASCII string "EXOMOL.def"
IsoFormula	A27	%27s	Isotopologue chemical formula
Iso-slug	A160	%160s	Isotopologue identifier, see text for details
DSName	A10	%10s	Isotopologue dataset name
$V$	I8	%8d	Version number with format YYYYMMDD
MolKey	A27	%27s	Standard inchi key of the molecule
$N_{\text{atom}}$	I4	%4d	Number of atoms
Atom definition (The following 2 lines occur $N_{\text{atom}}$ times)			
$I_{\text{atom}}$	I3	%3d	Isotope number
Atom	A3	%3s	Element symbol
<b>Isotopologue Information</b>			
$m_{\text{Da}}$ $m_{\text{kg}}$	F12.6,1X,ES14.8	%12.6f %14.8e	Isotopologue mass in Da and kg
$I_{\text{sym}}$	A6	%6s	Molecular symmetry Group (if $N_{\text{atom}} = 2$ then C or D)
$N_{\text{irrep}}$	I4	%4d	Number of irreducible representations
Symmetry definitions (The following 3 lines occur $N_{\text{irrep}}$ times)			
$I_{\text{irrep}}$	I3	%3d	Irreducible representation ID
Symmetry	A6	%6s	Irreducible representation label
$g_{\text{ns}}$	I3	%3d	Nuclear spin degeneracy
<b>ExoMol Information</b>			
$T_{\text{max}}$	F8.2	%8.2f	Maximum temperature of the line list
$N_{\text{broad}}$	I3	%3d	No. of pressure broadeners available
$D_{\text{avail}}$	I2	%2d	Dipole availability (1=Yes 0=No)

Continued on next page  
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Table 7 – Continued from previous page

Field	Fortran Format	C Format	Description
$N_{\text{cross}}$	I3	%3d	No. of cross section files available
$N_{\text{kcoef}}$	I3	%3d	No. of $k$ -coefficient files available
<b>State file information</b>			
$\text{Life}_{\text{avail}}$	I2	%2d	Flag denoting lifetime availability (1=Yes 0=No)
$g_{\text{avail}}$	I2	%2d	Flag denoting Landé $g$ -factor availability (1=Yes 0=No)
$N_{\text{states}}$	I10	%10d	No. of states in <code>.states</code> file
$N_{\text{cases}}$	I3	%3d	No. of quanta cases
Quanta case definition (The following line occurs $N_{\text{cases}}$ times)			
Case Label	A8	%8s	Label of the quanta case
$N_{\text{quanta}}$	I3	%3d	No. of quanta defined
Quanta definition (The following 3 lines occur $N_{\text{quanta}}$ times)			
Quanta Label	A3	%3s	
$F_{\text{Fortran}}$ $F_{\text{C}}$	A8,1x,A8	%8s %8s	Fortran and C format for quanta
Description	A40	%40s	Short description of quanta
<b>Trans file definition</b>			
$N_{\text{trans}}$	I15	%15d	Total number of transitions
$N_{\text{files}}$	I4	%4d	Number of transition files
$\tilde{\nu}_{\text{max}}$	F8.2	%8.2f	Maximum wavenumber (in $\text{cm}^{-1}$ )
<b>Partition function information</b>			
$T_{\text{max}}^Q$	F8.2	%8.2f	Maximum temperature of partition functions
$T_{\text{Step}}$	F5.2	%5.2f	Temperature step size
$C_{\text{avail}}$	I2	%2d	Cooling function availability (1=Yes 0=No)
<b>Pressure Broadening Information</b> (default values)			
$\gamma_D$	F6.4	%6.4f	Default value of Lorentzian half-width for all lines (in $\text{cm}^{-1}/\text{bar}$ )
$n_D$	F5.3	%5.3f	Default value of the temperature exponent for all lines
<b>Pressure Broadening Information</b> ( $N_{\text{broad}} > 0$ , the following 6 lines occur $N_{\text{broad}}$ times)			
Broadener Label	A8	%8s	Label for particular broadener
Broadener File Name	A20	%20s	File name of particular broadener
$J_{\text{max}}$	I7/F7.1	%7d/%7.1f	Maximum $J''$ for which pressure broadened parameters are provided
$\gamma_L$	F6.4	%6.4f	Value of Lorentzian half-width for lines with $J'' > J_{\text{max}}$ (in $\text{cm}^{-1}/\text{bar}$ )
$n_L$	F5.3	%5.3f	Value of the temperature exponent for lines with $J'' > J_{\text{max}}$
$N_Q$	I4	%4d	Number of defined quantum number sets
Quantum Sets (The following 3 lines occur $N_Q$ times)			
Set Code	A2	%2s	A code that defines this set of quantum numbers
$N_{\text{lines}}$	I6	%6d	No. of lines in the broad that contain this code
$N_{\text{Quanta}}$	I4	%4d	No. of quantum numbers defined

Continued on next page

Table 7 – Continued from previous page

Field	Fortran Format	C Format	Description
Quantum Numbers (The following line occurs $N_{\text{Quanta}}$ times)			
Quantum Label	A4	%4s	

### 5.3.2. Atom Definition

This section describes each of the atoms in the molecule. The  $N_{\text{atom}}$  field describes how many of them are defined. An integer mass number and a string of the molecule allows a definite description of the isotopologue.

### 5.3.3. Isotopologue Information

This section gives mass and symmetry information about the molecules.  $m$  is the molecular mass given in both Da and kg; the mass is used for modeling Doppler broadening. Next is the symmetry information of the molecule. The molecular symmetry group [257] and number of irreducible representations (irreps),  $N_{\text{irrep}}$ , are described here, after which  $N_{\text{irrep}}$  lines defining each irrep in that group are given. The  $I_{\text{sym}}$  field gives the ID of the symmetry used in the `.states` file. However, for linear molecules including diatomics the symmetry group label will only take values of  $C$  or  $D$  corresponding to  $C_{\infty v}$  or  $D_{\infty h}$  respectively and  $N_{\text{sym}}$  is set to 0.  $g_{\text{ns}}$  corresponding to nuclear spin degeneracy will be provided for  $C$ . For  $D$ ,  $g_{\text{even}}$  and  $g_{\text{odd}}$  are provided; these correspond to the nuclear-spin degeneracy factor needed for when the rovibronic wavefunction (i.e. the wavefunction without nuclear spin) is respectively even or odd with respect to interchange of identical atoms. Otherwise we follow symmetry labels and conventions as set out by Bunker and Jensen [257]. In this case  $g_{\text{ns}}$  is provided for each irreducible representation.

### 5.3.4. ExoMol Information

This section relates to the availability of certain aspects of the line list. First and foremost it describes the maximum temperature for which the line list is

applicable ( $T_{\max}$ ). Secondly, for molecules with pressure broadening parameters, it describes the number of broadeners ( $N_{\text{broad}}$ ) available; this is used later on to determine what additional structures should be parsed in the file. Finally, there is a boolean field which describes whether we provide dipoles for the particular molecule.

### 5.3.5. State file information

This section is the most important as it describes how the `.states` file format is structured. The first two fields describe whether lifetime information and Landé  $g$ -factors are available. These will be placed after the initial ‘standard’ fields in the states file. Then there is the  $N_{\text{states}}$  field which gives the number of states in the file. Finally the number of quanta cases,  $N_{\text{cases}}$ . A quanta case describes what the following quantum numbers represent. For example, commonly for polyatomics the ExoMol `.states` file provides two representations of quantum numbers; the standard normal mode quantum numbers and the local mode ‘TROVE’ quantum numbers. For this example  $N_{\text{cases}} = 2$ ; see Down *et al* [14] for instance. Each case is given a label and must describe the number of quanta defined  $N_{\text{quanta}}$ , after which  $N_{\text{quanta}}$  lines will follow, each describing a particular quantum number. Quantum numbers are described by a string label comprising a string format identifier and a short, human-readable description. The string format identifier utilises the FORTRAN string formatting to allow for the construction of formats more easily for FORTRAN which has stricter string handling than C/C++. An overview of the new `.states` format is given in Section 5.4.

### 5.3.6. Transition and partition information

This section simply describes the availability of `.trans` and `.pf` files. In particular, for the partition function file (`.pf`) it describes the maximum temperature provided ( $T_{\max}^Q$ ), the step size of the temperature ( $T_{\text{Step}}$ ) and whether a cooling function is also present.



### 5.3.7. Pressure broadening information

Finally, the pressure broadening information is provided here. If  $N_{\text{broad}} = 0$  this simply contains default values of Lorentzian half-width  $\gamma_L$  and temperature exponent  $n_L$ . If  $N_{\text{broad}} > 0$  this section has the following structure: the broadener is defined first, immediately followed by the associated file name,  $J_{\text{max}}$ , and the asymptotic values of the Lorentzian half-width  $\gamma_L$  and temperature exponent  $n_L$  for the  $a0$  quantum number set (see Section 5.6).  $N_Q$  quantum number sets are explicitly defined by a label (e.g.  $a0$ ) that corresponds to codes in the broadener file, the number of lines in the broadener file with that code and the number of quantum numbers in that set. After this each additional quantum number is also described; the total angular momentum quantum number of the lower state  $J''$  is compulsory. The quantum numbers relate to the quanta defined previously in the states section followed by either a prime or double prime. Further details relating to pressure broadening are given in Section 5.6.

The final field of the definition file is used for keywords to facilitate the search of the database by web search engines.

---

EXOMOL.def	#	ID
(12C)(32S)	#	IsoFormula
12C-32S	#	Iso-slug
JnK	#	Isotopologue dataset name
20160217	#	Version number with format YYYYMMDD
DXHPZXWIPWDXHJ-UHFFFAOYSA-N	#	Inchi key of molecule
2	#	Number of atoms
12	#	Isotope number 1
C	#	Element symbol 1
32	#	Isotope number 2
S	#	Element symbol 2
43.972071 7.30173406e-26	#	Isotopologue mass (Da) and (kg)
C	#	Symmetry group
1	#	Number irreducible representations
1	#	Irreducible representation ID
0	#	Irreducible representation label

1	# Nuclear spin degeneracy
3000	# Maximum temperature of linelist
2	# No. of pressure broadeners available
0	# Dipole availability (1=yes, 0=no)
0	# No. of cross section files available
0	# No. of k-coefficient files available
0	# Lifetime availability (1=yes, 0=no)
0	# Lande g-factor availability (1=yes, 0=no)
11497	# No. of states in .states file
1	# No. of quanta cases
dcs	# Quantum case label
5	# No. of quanta defined
i	# Quantum label 1
I12 %12d	# Format quantum label 1
State ID	# Description quantum label 1
E	# Quantum label 2
F12.6 %12.6f	# Format quantum label 2
State energy in cm-1	# Description quantum label 2
g	# Quantum label 3
I6 %6d	# Format quantum label 3
State degeneracy	# Description quantum label 3
J	# Quantum label 4
I7 %7d	# Format quantum label 4
Total rotational angular momentum	# Description quantum label 4
v	# Quantum label 5
I4 %4d	# Format quantum label 5
State vibrational quantum number	# Description quantum label 5
199045	# Total number of transitions
1	# No. of transition files
10996.09	# Maximum wavenumber (in cm-1)
3000.0	# Maximum temperature of partition function
1.00	# Temperature step size
0	# Cooling function availability (1=yes, 0=no)
0.0700	# Default value of Lorentzian half-width for all lines (in cm-1/bar)
0.500	# Default value of temperature exponent for all lines
air	# Label for a particular broadener
12C-32S__air.broad	# Filename of particular broadener
46	# Maximum J for which pressure broadening parameters provided
0.0698	# Value of Lorentzian half-width for J" > Jmax

0.750	# Value of temperature exponent for lines with J" > Jmax
1	# Number of defined quantum number sets
a0	# A code that defines this set of quantum numbers
47	# No. of lines in the broad that contain this code
0	# No. of quantum numbers defined
self	# Label for a particular broadener
12C-32S__self.broad	# Filename of particular broadener
46	# Maximum J for which pressure broadening parameters provided
0.0620	# Value of Lorentzian half-width for J" > Jmax
0.500	# Value of temperature exponent for lines with J" > Jmax
1	# Number of defined quantum number sets
a0	# A code that defines this set of quantum numbers
47	# No. of lines in the broad that contain this code
0	# No. of quantum numbers defined

Table 8: File 12C-32S\_\_JnK.def: the definition file for  $^{12}\text{C}^{32}\text{S}$ .

EXOMOL.def	# ID
(1H)2(16O)	# IsoFormula
1H2-16O	# Iso-slug
BT2	# Isotopologue dataset name
20160220	# Version number with format YYYYMMDD
XLYOFNOQVPJJNP-UHFFFAOYSA-N	# Inchi key of molecule
3	# Number of atoms
1	# Isotope number 1
H	# Element symbol 1
16	# Isotope number 2
0	# Element symbol 2
18.010565 2.99072463e-26	# Isotopologue mass (Da) and (kg)
C2v	# Symmetry group
4	# Number of irreducible representations
1	# Irreducible representation ID
A1	# Irreducible representation label
1	# Nuclear spin degeneracy
2	# Irreducible representation ID
A2	# Irreducible representation label
1	# Nuclear spin degeneracy
3	# Irreducible representation ID
B1	# Irreducible representation label

3	# Nuclear spin degeneracy
4	# Irreducible representation ID
B2	# Irreducible representation label
3	# Nuclear spin degeneracy
3000.0	# Maximum temperature of linelist
0	# No. of pressure broadeners available
0	# Dipole availability (1=yes, 0=no)
0	# No. of cross section files available
0	# No. of k-coefficient files available
0	# Lifetime availability (1=yes, 0=no)
0	# Lande g-factor availability (1=yes, 0=no)
221097	# No. of states in .states file
1	# No. of quanta cases
nltcs	# Quantum case label
14	# No. of quanta defined
i	# Quantum label 1
I12 %12d	# Format quantum label 1
State ID	# Description quantum label 1
E	# Quantum label 2
F12.6 %12.6f	# Format quantum label 2
State energy in cm-1	# Description quantum label 2
g_tot	# Quantum label 3
I6 %6d	# Format quantum label 3
Total degeneracy of the level	# Description quantum label 3
J	# Quantum label 4
I7 %7d	# Format quantum label 4
Total rotational angular momentum	# Description quantum label 4
+/-	# Quantum label 5
A1 %1s	# Format quantum label 5
Total parity	# Description quantum label 5
Gamma	# Quantum label 6
I2 %2d	# Format quantum label 6
Symmetry block number (1)	# Description quantum label 6
N	# Quantum label 7
I10 %10d	# Format quantum label 7
Reference within the symmetry block	# Description quantum label 7
nucspin	# Quantum label 8
A1 %1s	# Format quantum label 8
Nuclear spin isomer label(2)	# Description quantum label 8

Gamma_rve	# Quantum label 9
A2 %2s	# Format quantum label 9
Rovibrational symmetry label	# Description quantum label 9
v1	# Quantum label 10
I2 %2d	# Format quantum label 10
v1 symmetric stretch quantum number(3	# Description quantum label 10
v2	# Quantum label 11
I2 %2d	# Format quantum label 11
v2 bend quantum number	# Description quantum label 11
v3	# Quantum label 12
I2 %2d	# Format quantum label 12
v2 asymmetric stretch quantum number	# Description quantum label 12
Ka	# Quantum label 13
I2 %2d	# Format quantum label 13
Ka rotational quantum number	# Description quantum label 13
Kc	# Quantum label 14
I2 %2d	# Format quantum label 14
Kc rotational quantum number	# Description quantum label 14
505806255	# Total number of transitions
16	# No. of transition files
29971.78	# Maximum wavenumber (in cm-1)
3000.0	# Maximum temperature of partition function
1.0	# Step size of temperature
1	# Cooling function availability (1=yes, 0=no)
0.0700	# Default value of Lorentzian half-width for all lines (in cm-1/ba
0.500	# Default value of temperature exponent for all lines

Table 9: File BT2.def: the definition file for H<sub>2</sub><sup>16</sup>O.

#### 5.4. States file

The basic ExoMol data structure is a comprehensive list of the molecular states involved. This file contains a numbered list of energy terms, in cm<sup>-1</sup> with the zero defined by the lowest level, immediately followed by the total statistical weight  $g_{\text{tot}} = (2J + 1) \times g_{\text{ns}}$  in the HITRAN convention [258] which retains the full nuclear-spin degeneracy of all atoms. These three columns are sufficient, in conjunction with the Transitions file, to model absorption and/or emission in-

tensities. The degeneracy is followed by the total angular momentum quantum number  $J$  (integer or half-integer) The next columns are used for the state lifetimes (s) and/or Landé  $g$ -factors with availability determined in the definition file. This is followed by good (rigorous) quantum number(s) such as the state symmetry, parity or a total symmetry (irreducible representation). Then follows any approximate quantum numbers that are specified. There is some flexibility over this portion of the file which is defined on a case-by-case basis, see Ref. [259] for example.

For polyatomic molecules the approximate quantum numbers include, for example, the vibrational normal mode  $v_i$ , local mode  $n_i$  quantum numbers, rotational quantum numbers  $K, K_a, K_c$ , polyad numbers, vibrational and rotational symmetries etc. For the diatomic molecules these are the spin components  $F_i, \Sigma$  or  $\Omega$ , electronic term designations, e.g  $^3\Pi_g$ , vibrational quantum numbers  $v$ , rotational angular momentum  $N$  (Hund's case (b)), projection of the electronic angular momentum,  $\Lambda$ , etc. The spin components quantum numbers can be integral, half-integral or even strings F1, F2, ... If the theoretical energy is replaced by an experimental counterpart, the former is usually kept as an extra field.

Regardless of molecule, the structure of any ExoMol state file can be determined from its complimentary definition file. The structure in order is:

1. Mandatory fields
2. Lifetime if  $\text{Life}_{\text{avail}} = 1$
3. Landé  $g$ -factor if  $g_{\text{avail}} = 1$
4. Quanta Cases in order of definition

The detailed specification of the mandatory section of the states file is given in Table 10. Using the state file information from Table 8, the mandatory definition in the FORTRAN string format can easily be programmatically generated:

```
"(I12,1x,F12.6,1x,I6,1x,I7,1x,ES12.4,1x,A1,1x,A1,1x,A9,1x,I2,1x,I2,1x,I2,1x,
```

Field	Fortran Format	C Format	Description
$i$	I12	%12d	State ID
$E$	F12.6	%12.6f	State energy in $\text{cm}^{-1}$
$g_{\text{tot}}$	I6	%6d	State degeneracy
$J$	I7/F7.1	%7d/%7.1f	$J$ -quantum number (integer/half-integer)
$(\tau)$	ES12.4	%12.4e	Lifetime in s (optional)
$(g)$	F10.6	%10.6f	Landé $g$ -factor (optional)

ID: state identifier: a non-negative integer index, starting at 1

$J$ : total angular momentum quantum, excluding nuclear spin

Fortran format,  $J$  integer: (I12,1x,F12.6,1x,I6,I7,1x,ES12.4,1x,F10.6)

or  $J$  half-integer: (I12,1x,F12.6,1x,I6,F7.1,1x,ES12.4,1x,F10.6)

Table 10: Specification of the mandatory part of the states file and extra data options.

I2)"

This format corresponds to the sample state file given in Table 11.

Note that for lifetimes it is necessary to deal with states with infinite radiative lifetimes (e.g. the ground state) and for which lifetimes are not available: for highly lying states not all downward transitions are considered so it is not possible to compute the lifetime. Here, and elsewhere in the database, infinity is specified by the string 'INF' and unknown numbers by 'NaN'.

### 5.5. Transitions file

The transitions file has a simple structure, see Table 12. Two pointers,  $i$  and  $f$  point to rows in the `.states` file to identify the upper and lower states involved plus the full information characterising these states. The  $A$  gives the Einstein  $A$  coefficient for this transition. This file can be very large but, optionally, there is a fourth column  $\tilde{\nu}$  which gives the transition wavenumber. Otherwise this must be computed from the states file as the difference of the upper and lower energy levels. A sample transitions file is given in Table 13.

Many of the molecules in the database are characterised by a very large number of transitions. To make the use of `.trans` tractable these files are often split

$i$	$\tilde{E}$	$g$	$J$	$\tau$	$g$	$+/-$	$e/f$	State	$v$	$\Lambda$	$\Sigma$	$\Omega$
4051	8242.235601	15	7	6.509E-01	0.106733	+	f	a3Pi	0	1	1	2
4052	8294.404892	15	7	2.106E-03	0.018230	+	f	a3Pi	0	1	0	1
4053	8369.059680	15	7	4.594E-01	0.000045	+	f	a3Pi	0	1	-1	0
4054	8627.261523	15	7	6.479E-05	0.017893	+	f	Ap1Pi	0	1	0	1
4055	8781.692232	15	7	1.881E-01	0.106728	+	f	a3Pi	1	1	1	2
4056	8833.629438	15	7	1.466E-03	0.018239	+	f	a3Pi	1	1	0	1
4057	8907.535275	15	7	1.595E-01	0.000048	+	f	a3Pi	1	1	-1	0
4058	9167.417952	15	7	4.406E-05	0.017996	+	f	Ap1Pi	1	1	0	1
4059	9314.686162	15	7	1.019E-01	0.106678	+	f	a3Pi	2	1	1	2
4060	9365.515048	15	7	1.139E-03	0.018390	+	f	a3Pi	2	1	0	1
4061	9434.128287	15	7	7.821E-02	0.000113	+	f	a3Pi	2	1	-1	0
4062	9533.168336	15	7	2.160E-03	0.029567	+	f	b3Sigma+	0	0	1	1
4063	9545.896549	15	7	1.114E-02	0.005414	+	f	b3Sigma+	0	0	0	0
4064	9709.239641	15	7	3.342E-05	0.018256	+	f	Ap1Pi	2	1	0	1
4065	9841.636252	15	7	8.905E-02	0.106743	+	f	a3Pi	3	1	1	2
4066	9893.716709	15	7	1.039E-03	0.018284	+	f	a3Pi	3	1	0	1
4067	9966.036778	15	7	6.823E-02	0.000061	+	f	a3Pi	3	1	-1	0
4068	10103.239512	15	7	1.106E-03	0.026224	+	f	b3Sigma+	1	0	1	1
4069	10115.858820	15	7	2.887E-03	0.008900	+	f	b3Sigma+	1	0	0	0
4070	10240.534926	15	7	2.729E-05	0.018317	+	f	Ap1Pi	3	1	0	1

$i$ : State counting number.

$\tilde{E}$ : State energy in  $\text{cm}^{-1}$ .

$g$ : State degeneracy.

$J$ : Total angular momentum.

$\tau$ : Lifetime in s.

$g$ : Landé  $g$ -factor

$+/-$ : Total parity.

$e/f$ : rotationless-parity.

$v$ : State vibrational quantum number.

$\Lambda$ : Projection of the electronic angular momentum.

$\Sigma$ : Projection of the electronic spin.

$\Omega$ :  $\Omega = \Lambda + \Sigma$ , projection of the total angular momentum.

Table 11: Extract from the state file for  $^{40}\text{Ca}^{16}\text{O}$ , 40Ca-160\_VBATHY.state.



Field	Fortran Format	C Format	Description
$i$	I12	%12d	Upper state ID
$f$	I12	%12d	Lower state ID
$A$	ES10.4	%10.4e	Einstein $A$ coefficient in $\text{s}^{-1}$
$\tilde{\nu}_{fi}$	E15.6	%15.6e	Transition wavenumber in $\text{cm}^{-1}$ (optional).

Fortran format: (I12,1x,I12,1x,ES10.4,1x,ES15.6)

Table 12: Specification of the transitions file.

into wavenumber regions, as discussed above. The `.trans` files are also compressed using `.bz2` format; a utility is provided for reading this format without needing to uncompress the file.

### 5.6. Pressure broadening file

Like the `.states` file, the first four fields of the `.broad` file are mandatory for all records, see Table 14. This includes one quantum number,  $J''$ , which is always known and hence guarantees that at least a semi-empirical Lorentzian half-width ( $\gamma_{\text{ref}}$ ) and temperature dependence, represented by exponent  $n$  (see Eq. 1), can be generated for every molecular line. Additional upper and lower state labels, which give molecule and quantum number dependent behavior, follow the compulsory fields. These are given in the definitions file along with the maximum  $J''$  for which approximate parameters have been generated,  $J''_{\text{max}}$ , and the respective values of Lorentzian half-width and temperature exponent,  $\gamma_{\text{max}}$  and  $n_{\text{max}}$ . When  $J'' > J''_{\text{max}}$ ,  $\gamma_{\text{max}}$  and  $n_{\text{max}}$  should be used for the Lorentzian half-width and temperature exponent respectively.

The `.broad` file has a hierarchical structure; values of  $\gamma_{\text{ref}}$  and  $n$  with full quantum assignments are presented first, followed by values with partial quantum assignments, then finally values with  $J''$  dependence only. This represents the preferential order in which the values of  $\gamma_{\text{ref}}$  and  $n$  should be used. This hierarchical structure is demonstrated in Table 16 where a portion of the  $\text{H}_2\text{O}-\text{H}_2$  `.broad` file is presented.

$i$	$f$	$A_{if}$	$\tilde{\nu}_{if}$
10571	10884	9.5518E-06	120.241863
21053	21375	1.9515E-05	120.242886
8726	9672	1.8658E-04	120.243522
11655	11950	5.0065E-06	120.243733
93209	93967	5.7055E-03	120.244192
2228	3175	7.3226E-07	120.244564
46727	46432	1.0599E-04	120.244658
44436	44774	1.4626E-04	120.245583
29037	28723	1.8052E-04	120.245669
4458	4805	1.0431E-08	120.246396
69313	68434	5.0531E-06	120.248178
22640	22985	1.1281E-07	120.248891
57027	56721	7.1064E-06	120.250180

$i$ : Upper state counting number;

$f$ : Lower state counting number;  $A_{if}$ : Einstein  $A$  coefficient in  $s^{-1}$ ;  $\tilde{\nu}_{if}$ : transition wavenumber in  $cm^{-1}$ .

Table 13: Extract from the transitions file for  $^{40}Ca^{16}O$ , 40Ca-160\_VBATHY.trans.

Field	Fortran Format	C format	Description
code	A2	%2s	Code identifying quantum number set following $J''$
$\gamma_{ref}$	F6.4	%6.4f	Lorentzian half-width at reference temperature and pressure in $cm^{-1}/atm$
$n$	F6.3	%6.3f	Temperature exponent
$J''$	I7/F7.1	%7d/%7.1f	Lower $J$ -quantum number integer/half-integer

Fortran format,  $J$  integer: (A2,1x,F6.4,1x,F6.3,1x,I7)

or  $J$  half-integer: (A2,1x,F6.4,1x,F6.3,1x,F7.1)

Table 14: Specification of the mandatory part of the pressure broadening parameters file.

a0	0.0860	0.096	0
a0	0.0850	0.093	1
a0	0.0840	0.091	2
a0	0.0840	0.089	3
a0	0.0830	0.087	4
...			
a0	0.0720	0.067	35
a0	0.0720	0.066	36
...			
Field	Fortran Format	C format	Description
code	A2	%2s	Code identifying quantum number set following $J''^*$
$\gamma_{\text{ref}}$	F6.4	%6.4f	Lorentzian half-width at reference temperature and pressure in $\text{cm}^{-1}/\text{bar}$
$n$	F5.3	%5.3f	Temperature exponent
$J''$	I7/F7.1	%7d	Lower $J$ -quantum number

\*Code definition: a0 = none

Table 15: File `12C-32S__air.broad`: Air `.broad` file for  $^{12}\text{C}^{32}\text{S}$ : portion of the file (upper part); field specification (lower part).

Once a suitable  $\gamma_{\text{ref}}$  and  $n$  have been identified, the Lorentzian half-width of a spectral line at temperature  $T$  and pressure  $P$  can be calculated as:

$$\gamma(T) = \gamma_{\text{ref}} \times \left(\frac{T_{\text{ref}}}{T}\right)^n \times \left(\frac{P}{P_{\text{ref}}}\right). \quad (1)$$

Here,  $T_{\text{ref}} = 296$  K and  $P_{\text{ref}} = 1$  bar.

A separate `.broad` file is supplied for each molecule-broadener system and the broadener is specified in the file name. Example `.broad` files for CS-Air and  $\text{H}_2\text{O}-\text{H}_2$  are given in Table 15 and Table 16. For particular molecule-broadener systems where no pressure broadening information is available, the default values of  $\gamma_{\text{ref}}$  and  $n$  given in the isotopologue's definition file may be used.

...
b2 0.0356 0.300 9 8 1 9 0 8 0 0 0 0 1 0
b2 0.0522 0.300 9 8 1 9 0 8 0 0 0 1 0 0
b2 0.0521 0.300 7 8 1 6 1 7 0 0 0 0 0 1
...
a5 0.0600 0.546 3 4 0 3 1 4
a5 0.0618 0.551 3 4 2 1 1 4
a5 0.0569 0.525 2 3 2 1 3 0
...
a1 0.0301 0.268 14 15
a1 0.0291 0.230 15 16
a1 0.0282 0.218 16 17
...
a0 0.0242 0.165 24
a0 0.0239 0.160 25
a0 0.0236 0.150 26
...

Field	Fortran Format	C format	Description
code	A2	%2s	Code identifying quantum number set following $J''^*$
$\gamma_{\text{ref}}$	F2.4	%2.4f	Lorentzian half-width at reference temperature and pressure in $\text{cm}^{-1}/\text{bar}$
$n$	F2.3	%2.3f	Temperature exponent
$J''$	I7	%7d	Lower $J$ -quantum number
$J'$	I7	%7d	Upper $J$ -quantum number
$K''_a$	I2	%2d	Lower rotational quantum number
$K''_c$	I2	%2d	Lower rotational quantum number
$K'_a$	I2	%2d	Upper rotational quantum number
$K'_c$	I2	%2d	Upper rotational quantum number
$v''_1$	I2	%2d	Lower vibrational quantum number
$v''_2$	I2	%2d	Lower vibrational quantum number
$v''_3$	I2	%2d	Lower vibrational quantum number
$v'_1$	I2	%2d	Upper vibrational quantum number
$v'_2$	I2	%2d	Upper vibrational quantum number
$v'_3$	I2	%2d	Upper vibrational quantum number

\*Code definitions: b2 =  $J', K''_a, K''_c, K'_a, K'_c, v''_1, v''_2, v''_3, v'_1, v'_2, v'_3$ ; a5 =  $J', K''_a, K''_c, K'_a, K'_c$ ;  
a1 =  $J'$ ; a0 = no further quantum numbers.

Table 16: 1H2-160\_H2.broad: H<sub>2</sub>O - H<sub>2</sub> broad file: portion of the file (upper part); field specification (lower part).

Field	Fortran Format	C Format	Description
<i>i</i>	I12	%12d	Upper state ID
<i>f</i>	I12	%12d	Lower state ID
<i>D</i>	F12.8	%12.8d	Dipole moment in Debye

Fortran format: (I12,1x,I12,1x,F12.8)

Table 17: Specification of the dipole file.

### 5.7. *Dipoles file*

The dipole file mirrors the simple structure of the transitions file, see Table 17. Two pointers, *i* and *f* point to rows in the `.states` file to identify the upper and lower states involved, as specified by the state ID variable, plus the full information characterising these states. *D* gives the signed transition dipole (in Debye). In this form the combination of the `.states` and `.dipole` files can be used to simulate the effect of weak linear polarized electric fields (Stark) on the molecule, both static and time-dependent, where the polarization of the field is along *Z*. The  $|D|$  value is related to the Einstein *A* coefficient from the transitions file via a simple transformation:

$$A_{if} = \frac{64 \times 10^{-36} \pi^4}{3h} |D|^2 (2J'' + 1) \tilde{\nu}_{if}^3 \quad (2)$$

where *h* is Planck's constant, and *J''* is the total angular momentum of the lower state. However the phase of the (potentially complex) *D*, which is important for calculating the Stark effect, is lost in this transformation. Therefore, in the dipole file we provide the actual value of *D*.

The ExoMol format of the `.states/.dipole` files combination can then also be used for modelling higher-order field effects, including optical activity, polarization phenomena, nonlinear optical properties, and effects of strong fields including the effects of strong magnetic fields [260]. In this case the matrix elements of the corresponding properties will be stored as additional columns after the *D* column. Table 18 gives an example of a dipole file for CaO.

$i$	$f$	$\mu_{if}$
33	1	0.52209235E+01
34	1	0.14271796E+00
35	1	0.50123653E-01
36	1	0.16344445E-01
37	1	-0.48627678E-03
38	1	-0.86965402E-03
39	1	0.31660861E-03
40	1	0.21892847E-03
41	1	-0.15794191E-03

$i$ : Upper state counting number;  $f$ : Lower state counting number;  $\mu_{if}$ : transition dipole in Debye.

Table 18: Extract from the dipole file for  $^{40}\text{Ca}^{16}\text{O}$ , `40Ca-16O_VBATHY.dipole`.

### 5.8. Cross section file

Absorption cross sections are provided by ExoMol in files with the extension `.cross`. The two columns of this file are described in Table 19 and are wavenumber ( $\tilde{\nu}$ , in  $\text{cm}^{-1}$ ) and cross section value ( $\sigma$ , in  $\text{cm}^2\text{molec}^{-1}$ ). The two fields are separated by a single space. The stored cross sections may be temperature or temperature and pressure dependent as denoted by the header which includes the temperature in K and the pressure in bar.

The absorption cross section is given as a sequence of average values,  $\sigma_i$  within wavenumber bins,  $\tilde{\nu}_i$  of specified width,  $\Delta\tilde{\nu}$ . That is, in general, cross section values are given by

$$\sigma_i = \sum_j \frac{S_j}{\Delta\tilde{\nu}} \int_{\tilde{\nu}_i - \Delta\tilde{\nu}/2}^{\tilde{\nu}_i + \Delta\tilde{\nu}/2} f(\tilde{\nu}; \tilde{\nu}_{0,j} \dots) d\tilde{\nu}, \quad (3)$$

where the sum is taken over individual absorption lines of strength  $S_j$  and central wavenumber  $\tilde{\nu}_{0,j}$ , and  $f(\tilde{\nu}; \tilde{\nu}_{0,j} \dots)$  is a normalized lineshape function which is calculated using transition-dependent parameters,  $(\tilde{\nu}_{0,j} \dots)$ . In the case that the bin width is much smaller than the line-width,  $\Delta\tilde{\nu} \ll \text{HWHM}$ , this formula

Field	Fortran Format	C Format	Description
$\tilde{\nu}_i$	F12.6	%12.6f	Central bin wavenumber, $\text{cm}^{-1}$
$\sigma_i$	ES14.8	%14.8e	Absorption cross section, $\text{cm}^2\text{molec}^{-1}$

Fortran format: (F12.6,1x,ES14.8)

Table 19: Specification of the `.cross` cross section file format

$\tilde{\nu}/\text{cm}^{-1}$	$\sigma/\text{cm}^2\text{molec}^{-1}$
1525.530000	8.58106599E-36
1525.540000	4.67629889E-37
1525.550000	4.53342436E-36
1525.560000	6.55402974E-34
1525.570000	2.08219623E-31
1525.580000	1.90580250E-33
1525.590000	4.18283497E-32
1525.600000	1.30145956e-26
1525.610000	5.31710192E-27
1525.620000	7.96854826E-28
1525.630000	1.04806968E-18
1525.640000	7.80260131E-18
1525.650000	3.18257478E-22
1525.660000	1.27578728E-34
1525.670000	6.09403363E-30
1525.680000	1.50281502E-29
1525.690000	2.79935737E-29

Table 20: Extract from `1H2-160_BT2_0-30000_296K_0bar_0.01.cross`: Cross section function file for  $\text{H}_2^{16}\text{O}$ .

reduces to

$$\sigma_i = \sum_j S_j f(\tilde{\nu}; \tilde{\nu}_{0,j} \dots). \quad (4)$$

See Hill *et al* [32] for further details. Table 20 gives an example of a cross section file for water.

### 5.9. *k*-coefficients file

*k*-coefficients are provided by ExoMol in files with the extension `.kcoef`. The format of this file is described in Table 21. The *k*-coefficients are provided in

Field	Fortran Format	C Format	Description
$P$	F12.6	%12.6f	Pressure in bar
$T$	F8.2	%8.2f	Temperature in K
$\lambda_{\min}$	F12.6	%12.6f	Minimum central wavelength of bins, in $\mu\text{m}$ ,
$\lambda_{\max}$	F12.6	%12.6f	Maximum central wavelength of bins, in $\mu\text{m}$ ,
$\Delta\lambda$	F12.6	%12.6f	Spacing of wavelength bin centres, in $\mu\text{m}$ ,
$\delta\lambda$	F12.6	%12.6f	Bin widths, in $\mu\text{m}$ ,
$N_{\text{bins}}$	I6	%6d	No. of wavelength bins provided
$N_g$	I6	%6d	No. of $g$ -ordinate values at which $k$ is provided within each bin
$g_{\text{ord}}$	$N_g*(F10.8, 1X)$	$N_g* \%10.8f$	$N_g$ $g$ -ordinate values, $0 \leq g \leq 1$ to be read on a single line.

Table 21: Specification of the `.kcoef` cross section file format; the  $k$ -coefficients are in  $\text{cm}^2$  molecule $^{-1}$ .

$N_{\text{bins}}$  wavelength bins spanning a specified wavelength range  $\lambda_{\min}$  to  $\lambda_{\max}$ . The spacing between the bin centers is  $\Delta\lambda$  while each bin is  $\delta\lambda$  wide. Clearly, for the case that the bins do not overlap  $\Delta\lambda = \delta\lambda$ . Each bin contains  $N_g$  values of  $k$  corresponding to the  $g$ -ordinate values (which will be an increasing sequence of  $k$ -coefficients).

Figure 22 gives an example of a  $k$ -coefficient file for water.

---

```

EXOMOL.def begin{figure}
  0.005000
  240.00
  0.300000
  30.000000
  0.025000
  0.025000
  1189
  20
  0.00343570 0.01801404 0.04388279 ... 0.98198600 0.99656430
      0.3 6.41554300e-23 6.42944800e-23 6.60117000e-23 ... 6.32264600e-21
      0.3025 6.67216472e-23 6.68662592e-23 6.86521680e-23 ... 7.09400881e-21
  ...
  30.000000 1.47557489e-24 1.47877304e-24 1.51826910e-24 ... 1.45420858e-22

```

---

Table 22: Extract from `1H2-16O_BT2_0-30000_296K_0.01.kcoef`:  $k$ -coefficient file for  $\text{H}_2^{16}\text{O}$ .

### 5.10. Partition function file

Temperature-dependent properties such as the partition function and the cooling function are provided by the partition function file with extension `.pf`. This file gives partition function,  $Q$ , and, if available ( $C_{\text{avail}} = 1$ ), cooling function,  $W$ , values as a function of temperature,  $T$ . Usually  $T$  is given in steps of  $T_{\text{Step}}$  (usually 1 K) up to  $T_{\text{max}}^Q$ . Table 23 gives the specification for the `.pf` file



Field	Fortran Format	C Format	Description
$T$	F8.1	%8.1d	Temperature in K
$Q(T)$	F15.4	%15.4d	Partition function (dimensionless).
$W(T)$	ES12.4	%12.4e	Cooling function in ergs s <sup>-1</sup> molecule <sup>-1</sup> (if available).

Fortran format: (F8.1,1x,F15.4,1x,ES12.4) or (F8.1,1x,F15.4)

Table 23: Specification of the .pf partition function file.

$T/K$	$Q(T)$	$W(T)$
1.0	4.0001	1.5662E-24
2.0	4.0274	6.9307E-22
3.0	4.2080	5.1097E-21
4.0	4.5752	1.3731E-20
5.0	5.0664	2.6214E-20
6.0	5.6252	4.3743E-20
7.0	6.2198	6.7576E-20
8.0	6.8342	9.8328E-20
9.0	7.4602	1.3618E-19
10.0	8.0937	1.8114E-19
11.0	8.7323	1.0279E-18
12.0	9.3745	2.6263E-18
13.0	10.0193	5.0095E-18
14.0	10.6663	8.1986E-18
15.0	11.3148	1.2208E-17
16.0	11.9646	1.7050E-17

Table 24: Extract from the partition function file for <sup>40</sup>Ca<sup>1</sup>H, 40Ca-1H\_\_Yadin.pf, including cooling function,  $W(T)$ .

example, for H<sub>2</sub><sup>17</sup>O, H has spin  $\frac{1}{2}$  so has degeneracy  $(2I_H + 1) = 2$  and <sup>17</sup>O has spin  $\frac{5}{2}$  so  $(2I_O + 1) = 6$ . This means that para water has degeneracy 1 and ortho water has degeneracy 18. This is in contrast to standard compilations of astronomical partition functions such as those of Irwin [261], Sauval and Tatum [262], and Barklem and Collet [263] who normalise to unit atomic nuclear spin degeneracy factors. Final results are independent of which convention is chosen provided they are used self-consistently.

## 6. Utility Programs

As part of the database, ExoMol provides a number of sample utilities for manipulating data presented in the ExoMol format. This covers the most common applications of the line lists, such as computing absorption/emission cross sections [32] or coefficients (‘stick’ intensities), partition functions, cooling functions and lifetimes. These utilities are available as Fortran 95 and Python programs combined into an EXOCROSS package. ExoCross allows users to utilise different line profiles to generate cross-sections, including the Voigt profile, using the line parameters compiled in the ExoMol database (file `.broad`) when available. Other standard line profiles include the temperature-dependent Doppler profile, rectangular-shape profiles (i.e. averaged cross-sections over the wavenumber bin), as well as general Gaussian and Lorentzian profiles with line-widths as free input parameters. The utilities can be obtained directly from the ExoMol website together with examples of input files.

As mentioned above a Python utility, `extract_trans.py` is provided which reads the `.trans` in `.bz2` format without requiring it to be uncompressed.

The molecular data produced by ExoMol is increasingly being used in a wide variety of applications. In many of these applications, pre-existing software packages utilise different data formats for their input files. The detailed specification of our format above should be sufficient to allow conversion to any required format using a suitable parser program. We have written a number of utility programs (in Python) that convert the ExoMol format to other common formats. These can be used to produce the full line list in different formats such as HITRAN, for which the Python utility `exomol2hitran.py` is available. Oscillator strengths can be generated by the utility `exomol2gf.py`. However, we should stress that the ExoMol format is much more compact than most other alternatives (notably HITRAN format), primarily due to the separation of information into the states

and trans file. The size of the ExoMol line lists means that it can be prohibitive to store the full data in another format; for example, the methane 10to10 line list would require 14 TB to be stored in HITRAN format. It is therefore recommended that the utility programs provided be integrated within the larger program, such that the new format data is produced for a small number of lines at a time which are immediately utilised by the program, thus avoiding the storage of data.

## 7. Website

The ExoMol website ([www.exomol.com](http://www.exomol.com)) is the main source of the ExoMol data. Data can be accessed two ways: (a) by first selecting a molecule and then a specific isotopologue-dataset name combination or (b) by selecting a particular data type, such as partition function. We note that cross sections are also being made available via the Virtual Atomic and Molecular Data Centre (VAMDC) [264, 265]. This API takes the form of a “Table Access Protocol” service at [www.exomol.com/tap/sync](http://www.exomol.com/tap/sync).

The `exomol.all` file described in Section 5.2 is available at [www.exomol.com/exomol.all](http://www.exomol.com/exomol.all) and the data files comprising each dataset can be accessed at a URL of the form [www.exomol.com/db/<MolFormula>/<IsoSlug>/<DSName>/<Filename>](http://www.exomol.com/db/<MolFormula>/<IsoSlug>/<DSName>/<Filename>), for example [www.exomol.com/db/Si0/28Si-160/EBJT/28Si-160\\_\\_EBJT.pf](http://www.exomol.com/db/Si0/28Si-160/EBJT/28Si-160__EBJT.pf).

A separate API for machine-access of the ExoMol data in the format described in this paper will be available at [www.exomol.com/api](http://www.exomol.com/api) and documented on the website. This will allow the automatic retrieval of data files by the HTTP GET method.

## 8. Conclusions

The ExoMol database presented here is a molecule-by-molecule set of comprehensive line lists for modelling spectra and other properties of hot gases. The choice of molecules is dictated by the need to model the atmospheres of exoplanets and other hot astronomical objects, but the spectroscopic data have much wider applications than this. We are still in the process of adding molecules to the database and are receptive to suggestions of other key species to include. In particular the discovery [266, 267] and observation of spectra [268] of very hot, rocky planets will undoubtedly lead to the consideration of novel species [269, 270].

In this paper we have significantly expanded the flexibility and scope of the data model used by ExoMol. The original ExoMol format of a `.states` and `.trans` file has been already adopted as a possible input format by some programs, such as PGOPHER [271]. The extended format proposed here is both backward-compatible and designed to allow direct extraction of data from the database using an application programming interface (API). We believe that this will be helpful to scientists and relieve them of the need to manually deal with datasets containing many billions of spectral lines. Our new data structure allows the inclusion of important new molecular parameters including those to determine pressure broadening, state-resolved radiative lifetimes and Landé  $g$ -factors.

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