Thermal history modelling: HeFTy vs. QTQt

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A R T I C L E   I N F O

Article history:
Received 18 March 2014
Accepted 29 September 2014
Available online 7 October 2014

Keywords:
Thermochronology
Modelling
Statistics
Software
Fission tracks
(U-Th)/He

A B S T R A C T

HeFTy is a popular thermal history modelling program which is named after a brand of trash bags as a reminder of the 'garbage in, garbage out' principle. QTQt is an alternative program whose name refers to its ability to extract visually appealing ('cute') time–temperature paths from complex thermochronological datasets. This paper compares and contrasts the two programs and aims to explain the algorithmic underpinnings of these 'black boxes' with some simple examples. Both codes consist of 'forward' and 'inverse' modelling functionalities. The 'forward model' allows the user to predict the expected data distribution for any given thermal history. The 'inverse model' finds the thermal history that best matches some input data. HeFTy and QTQt are based on the same physical principles and their forward modelling functionalities are therefore nearly identical. In contrast, their inverse modelling algorithms are fundamentally different, with important consequences. HeFTy uses a 'Frequentist' approach, in which formalised statistical hypothesis tests assess the goodness-of-fit between the input data and the thermal model predictions. QTQt uses a Bayesian 'Markov Chain Monte Carlo' (MCMC) algorithm, in which a random walk through model space results in an assemblage of 'most likely' thermal histories. In principle, the main advantage of the Frequentist approach is that it contains a built-in quality control mechanism which detects bad data ('garbage') and protects the novice user against applying inappropriate models. In practice, however, this quality-control mechanism does not work for small or imprecise datasets due to an undesirable sensitivity of the Frequentist algorithm to sample size, which causes HeFTy to 'break' when datasets are sufficiently large or precise. QTQt does not suffer from this problem, as its performance improves with increasing sample size in the form of tighter credibility intervals. However, the robustness of the MCMC approach also carries a risk, as QTQt will accept physically impossible datasets and come up with 'best fitting' thermal histories for them. This can be dangerous in the hands of novice users. In conclusion, the name 'HeFTy' would have been more appropriate for QTQt, and vice versa.

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1. Introduction

Thermal history modelling is an integral part of dozens of tectonic studies published each year [e.g., Tian et al. (2014); Karlstrom et al. (2014); Cochrane et al. (2014)]. Over the years, a number of increasingly sophisticated software packages have been developed to extract time–temperature paths from fission track, U–Th–He, ⁴He/³He and vitrinite reflectance data [e.g., Corrigan (1991); Gallagher (1995); Willett (1997); Ketcham et al. (2000)]. The current ‘market leaders’ in inverse modelling are HeFTy (Ketcham, 2005) and QTQt (Gallagher, 2012). Like most well written software, HeFTy and QTQt hide all their implementation details behind a user friendly graphical interface. This paper has two goals. First, it provides a ‘glimpse under the bonnet’ of these two ‘black boxes’ and second, it presents an objective and independent comparison of both programs. We show that the differences between HeFTy and QTQt are significant and explain why it is important for the user to be aware of them. To make the text accessible to a wide readership, the main body of this paper uses little or no algebra (further theoretical background is deferred to the appendices). Instead, we illustrate the strengths and weaknesses of both programs by example. The physical models which geologists use to describe the diffusion of helium or the annealing of fission tracks are but approximations of reality. To simulate this fact in our linear regression example, we will now try to fit a linear model to a weakly non-linear dataset generated using Eq. (1) with a = 5, b = 2, c = 0 and σ = 1 (Fig. 2i). It is easy to fit a straight line model through these data and determine parameters A and B of Eq. (2) analytically by ordinary least squares regression. However, for the sake of illustrating the algorithms used by HeFTy and QTQt, it is useful to do the same exercise by numerical modelling. In the following, the words ‘HeFTy’ and ‘QTQt’ will be placed in inverted commas when reference is made to the underlying methods, rather than the actual computer programs by Ketcham (2005) and Gallagher (2012).

‘HeFTy’ explores the ‘model space’ by generating a large number (N) of independent random intercepts and slopes (Aᵢ, Bᵢ for j = 1 → N), drawn from a joint uniform distribution (Fig. 2ii). Each of these pairs corresponds to a straight line model, resulting in a set of residuals (yᵢ – Aᵢ – Bᵢxᵢ) which can be combined into a least-squares goodness-of-fit statistic:

\[ \chi^2 = \sum_{i=1}^{n} \frac{(y_i - A - B x_i)^2}{\sigma^2} \]  

(3)

Low and high \( \chi^2 \) values correspond to good and bad data fits, respectively. Under the ‘Frequentist’ paradigm of statistics (see Appendix A), \( \chi^2 \) can be used to formally test the hypothesis (‘H₀’) that the data were drawn from a straight line model with a = Aᵢ, b = Bᵢ and c = 0. Under this hypothesis, \( \chi^2 \) is predicted to follow a Chi-square distribution with n – 2 degrees of freedom:\n
\[ P(x, y|Aᵢ, Bᵢ) = P(\chi^2|\text{H₀}) \sim \chi^2_{n-2} \]  

(4)

where \( P(X|Y) \) stands for the ‘probability of X given Y’. The ‘likelihood function’ \( P(x, y|Aᵢ, Bᵢ) \) allows us to test how ‘likely’ the data are under the proposed model. The probability of observing a value at least as extreme as \( \chi^2 \) under the proposed (Chi-square) distribution is called the ‘p-value’. HeFTy uses cutoff-values of 0.05 and 0.5 to indicate ‘acceptable’ and ‘good’ model fits. Out of N = 1000 models tested in Fig. 2ii–ii, 50 fall in the first, and 180 in the second category.

QTQt also explores the ‘model space’ by random sampling, but it goes about this in a very different way than HeFTy. Instead of ‘carpet bombing’ the parameter space with uniformly distributed independent values, QTQt performs a random walk of serially dependent random values. Starting from a random guess anywhere in the parameter space, this ‘Markov Chain’ of random models systematically samples the model space so that models with high \( P(X, y|Aᵢ, Bᵢ) \) are more likely to be accepted than those with low values. Thus, QTQt bases the decision whether or not to accept or reject the jth model not on the absolute value of \( P(X, y|Aᵢ, Bᵢ) \), but on the ratio of \( P(X, y|Aᵢ, Bᵢ)/P(X, y|Aᵢ-1, Bᵢ-1) \). See Appendix B for further details about Markov Chain Monte Carlo (MCMC) modelling. The important thing to note at this point is that in well behaved systems like our linear dataset, QTQt’s MCMC approach yields identical results to HeFTy’s Frequentist algorithm (Fig. 2ii/iv).

2. Part I: linear regression

Before venturing into the complex multivariate world of thermochronology, we will first discuss the issues of inverse modelling in the simpler context of linear regression. The bivariate data in this problem \( \{xᵢ, yᵢ\} \) where \( xᵢ = \{x₁, ..., xᵢ, ..., xₙ\} \) and \( yᵢ = \{y₁, ..., yᵢ, ..., yₙ\} \) will be generated using a polynomial function of the form:

\[ yᵢ = a + b xᵢ + c xᵢ^2 + \epsilonᵢ \]  

(1)

where a, b and c are constants and \( \epsilonᵢ \) are the ‘residuals’, which are drawn at random from a Normal distribution with zero mean and standard deviation \( \sigma \). We will try to fit these data using a two-parameter linear model:

\[ y = A + B x \]  

(2)

On an abstract level, HeFTy and QTQt are two-way maps between the ‘data space’ \( \{xᵢ, yᵢ\} \) and ‘model space’ \( \{A, B\} \). Both programs comprise a ‘forward model’, which predicts the expected data distribution for any given set or parameter values, and an ‘inverse model’, which achieves the opposite end (Fig. 1). Both HeFTy and QTQt use a probabilistic approach to finding the set of models \( \{A, B\} \) that best fit the data \( \{xᵢ, yᵢ\} \), but they do so in very different ways, as discussed next.

2.1. Linear regression of linear data

For the first case study, consider a synthetic dataset of \( n = 10 \) data points drawn from Eq. (1) with \( a = 5, b = 2, c = 0 \) and \( \sigma = 1 \) (Fig. 2i). It is easy to fit a straight line model through these data and determine parameters A and B of Eq. (2) analytically by ordinary least squares regression. However, for the sake of illustrating the algorithms used by HeFTy and QTQt, it is useful to do the same exercise by numerical modelling. In the following, the words ‘HeFTy’ and ‘QTQt’ will be placed in inverted commas when reference is made to the underlying methods, rather than the actual computer programs by Ketcham (2005) and Gallagher (2012).

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2.2. Linear regression of weakly non-linear data

The physical models which geologists use to describe the diffusion of helium or the annealing of fission tracks are but approximations of reality. To simulate this fact in our linear regression example, we will now try to fit a linear model to a weakly non-linear dataset generated using Eq. (1) with \( a = 5, b = 2, c = 0.02 \) and \( \sigma = 1 \). First, we consider a small sample of \( n = 10 \) samples from this model (Fig. 2iii). The quadratic term (i.e., c) is so small that the naked eye cannot spot the non-linearity of these data, and neither can ‘HeFTy’. Using the same number of \( N = 1000 \) random guesses as before, ‘HeFTy’ finds 41 acceptable and

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1 The number of degrees of freedom is given by the number of measurements minus the number of fitted parameters, i.e., in this case the slope and intercept.
186 good fits using the $\chi^2$-test (Fig. 3ii). In other words, with a sample size of $n = 10$, the non-linearity of the input data is ‘statistically insignificant’ relative to the data scatter $\sigma$.

The situation is very different when we increase the sample size to $n=100$ (Fig. 4i–ii). In this case, ‘HeFTy’ fails to find even a single linear model yielding a $p$-value greater than 0.05. The reason for this is that the ‘power’ of statistical tests such as Chi-square increases with sample size (see Appendix C for further details). Even the smallest deviation from linearity becomes ‘statistically significant’ if a sufficiently large dataset is available. This is important for thermochronology, as will be illustrated in Section 3.1. Similarly, the statistical significance also increases with analytical precision. Reducing $\sigma$ from 1 to 0.2 has the same effect as increasing the sample size, as ‘HeFTy’ again fails to find any ‘good’ solutions (Fig. 5i–ii). ‘QTQt’, on the other hand, handles the large (Fig. 4iii–iv) and precise (Fig. 5iii–iv) datasets much better. In fact, increasing the quantity (sample size) or quality (precision) of data space model space

forward modelling
inverse modelling

$P(x,y | A,B)$ [Frequentist]
$P(A,B | x,y)$ [Bayesian]

Fig. 1. HeFTy and QTQt are ‘two-way maps’ between the ‘data space’ on the left and the ‘model space’ on the right. Inverse modelling is a two-step process. It involves (a) generating some random models from which synthetic data can be predicted and (b) comparing these ‘forward models’ with the actual measurements. HeFTy and QTQt fundamentally differ in both steps (Section 2.1).

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Fig. 2. (i) — White circles show 10 data points drawn from a linear model (black line) with Normal residuals ($\sigma = 1$). Red and green lines show the linear trends that best fit the data according to the Chi-square test; (ii) — the ‘Frequentist’ Monte Carlo algorithm (‘HeFTy’) makes 1000 independent random guesses for the intercept (A) and slope (B) drawn from a joint uniform distribution. A Chi-square goodness-of-fit test is done for each of these guesses. $p$-Values $<0.05$ and $>0.5$ are marked as green (‘acceptable’) and red (‘good’), respectively. (iii) — White circles and black line are the same as in (i). The colour of the pixels (ranging from blue to red) is proportional to the number of ‘acceptable’ linear fits passing through them using the ‘Bayesian’ algorithm (‘QTQt’). (iv) — ‘QTQt’ makes a random walk (‘Markov Chain’) through parameter space, sampling the ‘posterior distribution’ and yielding an ‘assemblage’ of best fitting slopes and intercepts (black dots and lines). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)
the data only has beneficial effects as it tightens the solution space (Fig. 4iv vs. Fig. 5iv).

2.3. Linear regression of strongly non-linear data

For the third and final case study of our linear regression exercise, consider a pathological dataset produced by setting $a = 26, b = -10, c = 1$ and $\sigma = 1$. The resulting data points fall on a parabolic line, which is far removed from the 2-parameter linear model of Eq. (2) (Fig. 6). Needless to say, the HeFTy algorithm does not find any ‘acceptable’ models. Nevertheless, the QTQt-like MCMC algorithm has no trouble fitting a straight line through these data. Although the resulting likelihoods are orders of magnitude below those of Fig. 2, their actual values are not used to assess the goodness-of-fit, because the algorithm only evaluates the relative ratios of the likelihood for adjacent models in the Markov Chain (Appendix B). It is up to the subjective judgement of the user to decide whether to accept or reject the proposed inverse models. This is very easy to do in the simple regression example of this section, but may be significantly more complicated for high-dimensional problems such as the thermal history modelling discussed in the next section. In conclusion, the simple linear regression toy example has taught us that (a) the ability of a Frequentist algorithm such as HeFTy to find a suitable inverse model critically depends on the quality and quantity of the input data; while (b) the opposite is true for a Bayesian algorithm like QTQt, which always finds a suite of suitable models, regardless of how large or bad a dataset is fed into it. The next section of this paper will show that the same principles apply to thermochronology in exactly the same way.

3. Part II: thermal history modelling

The previous section revealed significant differences between ‘HeFTy-like’ and ‘QTQt-like’ inverse modelling approaches to a simple two-dimensional problem of linear regression. Both algorithms were shown to yield identical results in the presence of small and well-behaved datasets. However, their response differed in response to large or poorly behaved datasets. We will now show that exactly the same phenomenon manifests itself in the multi-dimensional context of thermal history modelling. Firstly, we will use a geologically straightforward thermochronological dataset to ‘break’ HeFTy (Section 3.1). Then, we will use a physically impossible dataset to demonstrate that it is impossible to break QTQt even when we want to (Section 3.2).

3.1. Large datasets ‘break’ HeFTy

We will investigate HeFTy with a large but otherwise unremarkable sample and using generic software settings like those used by the majority of published HeFTy applications. The sample (‘KL29’) was collected from a Mesozoic granite located in the central Tibetan Plateau (33.87N, 95.33E). It is characterised by a 102 ± 7 Ma apatite fission track (AFT) age and a mean (unprojected) track length of ~12.1 μm, which was calculated from a dataset of 821 horizontally confined fission tracks. It is the large size of our dataset that allows us to push HeFTy to its limits. In addition to the AFT data, we also measured five apatite U-Th–He (AHe) ages, ranging from 47 to 66 Ma. AFT and AHe analyses were done at the University of Melbourne and University College London using procedures outlined by Tian et al. (2014) and Carter et al. (2014) respectively. For the thermal history modelling, we used the multi-kinetic annealing model of Ketcham et al. (2007), employing Dpar as a kinetic parameter. Helium diffusion in apatite was modelled with the Radiation Damage Accumulation and Annealing Model (RDAAM) of Flowers et al. (2009). Goodness-of-fit requirements for ‘good’ and ‘acceptable’ thermal paths were defined as 0.5 and 0.05 (see Section 2.1) and the present-day mean surface temperature was set to 15 ± 15 °C. To speed up the inverse modelling, it was necessary to specify a number of ‘bounding boxes’ in time–temperature (t–T) space. The first of these t–T constraints was set at 140 °C/200 Ma–20 °C/180 Ma, i.e. slightly before the oldest AFT age. Five more equally broad boxes were used to guide the thermal history modelling (Fig. 7). The issue of ‘bounding boxes’ will be discussed in more detail in Section 5.

In a first experiment, we modelled a small subset of our data comprising just the first 100 track length measurements. After one million iterations, HeFTy returned 39 ‘good’ and 1373 ‘acceptable’ thermal histories, featuring a poorly resolved phase prior to 120 Ma, followed by rapid cooling to ~60 °C, a protracted isothermal residence in the upper part of the AFT partial annealing zone from 120 to 40 Ma, and ending with a phase of more rapid cooling from 60 to 15 °C since 40 Ma. This is in every way an unremarkable thermal history, which correctly reproduces the negatively skewed (c-axis projected) track length distribution, and predicts AFT and AHe ages of 102 and 59 Ma,
respectively, well within the range of the input data (Fig. 7i). Next, we move on to a larger dataset, which was generated using the same AFT and AHe ages as before, but measuring an extra 269 confined fission tracks in the same slide as the previously measured 100 tracks. Despite the addition of so many extra measurements, the resulting length distribution looks very similar to the smaller dataset. Nevertheless, HeFTy struggles to find suitable thermal histories. In fact, the program fails to find a single ‘good’ t–T path even after a million iterations, and only comes up with a measly 109 ‘acceptable’ solutions. A closer look at the model predictions reveals that HeFTy does a decent job at modelling the track length distribution, but that this comes at the expense of the AFT and AHe age predictions, which are further removed from the measured values than in the small dataset (Fig. 7ii). In a final experiment, we prepared a second fission track slide for sample KL29, yielding a further 452 fission track length measurements. This brings the total tally of the length distribution to an unprecedented 821 measurements, allowing us to push HeFTy to its breaking point. After one million iterations, HeFTy does not manage to find even a single ‘acceptable’ t–T path (Fig. 7iii).

It is troubling that HeFTy performs worse for large datasets than it does for small ones. It seems unfair that the user should be penalised for the addition of extra data. The reasons for this behaviour will be discussed in Section 4. But first, we shall have a closer look at QTQt, which has no problems in finding a tight fit.

3.2. ‘Garbage in, garbage out’ with QTQt

Contrary to HeFTy, QTQt does not mind large datasets. In fact, its inverse modelling results improve with the addition of more data. This is because large datasets allow the ‘reversible jump MCMC’ algorithm (Appendix B) to add more anchor points to the candidate models, thereby improving the resolution of the t–T history. Thus, QTQt does not punish but reward the user for adding data. For the full 821-length dataset of KL29, this results in a thermal history similar to the HeFTy model of Fig. 7i. We therefore conclude that QTQt is much more robust than HeFTy in handling large and complex datasets. However, this greater robustness also carries a danger with it, as will be shown next. We now apply QTQt to a semi-synthetic dataset generated by arbitrarily changing the AHe age of sample KL29 from 55 ± 5 Ma to 102 ± 7 Ma, i.e. identical to its AFT age. As discussed in Section 3.1, the sample has a short (~12.1 μm) mean (unprojected) fission track length, indicating slow cooling through the AFT partial annealing zone. The identical AFT and AHe ages, however, imply infinitely rapid cooling. The combination of
Fig. 5. (i)–(ii) — As Figs. 3i–ii and 4i–ii but with higher precision data ($\sigma = 0.2$ instead of 1). Again, 'HeFTy' fails to find any 'acceptable' solutions. (iii)–(iv) — The same data analysed by 'QTQt', which works fine.

Fig. 6. (i) — White circles show 10 data points drawn from a strongly non-linear model (black line). (ii) — Although it clearly does not make any sense to fit a straight line through these data, 'QTQt' nevertheless manages to do exactly that. 'HeFTy' (not shown), of course, does not.
the AFT and AHe data is therefore physically impossible and, not surprisingly, HeFTy fails to find a single ‘acceptable’ fit even for a moderate sized dataset of 100 track lengths. QTQt, however, has no problem finding a ‘most likely’ solution (Fig. 8ii).

The resulting assemblage of models is characterised by a long period of isothermal holding at the base of the AFT partial annealing zone (~120 °C), followed by rapid cooling at 100 Ma, gentle heating to the AHe partial retention zone (~60 °C) until 20 Ma and rapid cooling to the surface thereafter (Fig. 8ii). This assemblage of thermal history models is largely unremarkable and does not, in itself, indicate any problems with the input data. These problems only become clear when we compare the measured with the modelled data. While the fit to the track length measurements is good, the AFT and AHe ages are off by 20%. It is then up to the user to decide whether or not this is ‘significant’ enough to reject the model results. This is not necessarily as straightforward as it may seem. For instance, the original QTQt paper by Gallagher (2012) presents a dataset in which the measured and modelled values for the kinetic parameter DPar differ by 25%. In this case, the author has made a subjective decision to attach less credibility to the DPar measurement. This may very well be justified, but nevertheless requires expert knowledge of thermochronology while remaining, once again, subjective. This subjectivity is the price of Bayesian MCMC modelling.

4. Discussion

The behaviour shown by HeFTy and QTQt in a thermochronological context (Section 3) is identical to the toy example of linear regression...
is that both the Laslett et al. (1987) and Ketcham et al. (2007) models are incorrect, albeit to different degrees. As sample size increases, the problem is that Geology itself imposes unrealistic assumptions on our thermal modelling efforts. Our understanding of diffusion and annealing kinetics is based on short term experiments carried out in completely different environments than the geological processes which we aim to understand. For example, helium diffusion experiments are done under ultra-high vacuum at temperatures of hundreds of degrees over the duration of at most a few weeks. These are very different conditions than those found in the natural environment, where diffusion takes place under hydrostatic pressure at a few tens of degrees over millions of years (Villa, 2006). But even if we disregard this problem, and imagine a utopian scenario in which our annealing and diffusion models are an exact description of reality, the p-value conundrum would persist, because there are dozens of other experimental factors that can go wrong, resulting in dozens of reasons for K–S and χ² to reject the data. Examples are observer bias in fission track length measurements to our dataset, HeFTy fails to yield any ‘acceptable’ solutions.

(Section 2). HeFTy is ‘too picky’ when it comes to large datasets and QTQt is ‘not picky enough’ when it comes to bad datasets. These opposite types of behaviour are a direct consequence of the statistical underpinnings of the two programs. The sample size dependence of HeFTy is caused by the fact that it judges the merits of the trial models by means of formalised statistical hypothesis tests, notably the Kolmogorov–Smirnov (K–S) and χ²-tests. These tests are designed to make a black or white decision as to whether the hypothesis is right or wrong. However, as stated in Section 2.2, the physical models produced by Science (including Geology) are “but approximations of reality” and are therefore always “somewhat wrong”. This should be self-evident from a brief look at the Settings menu of HeFTy, which offers the user the choice between, for example, the kinetic annealing model of Laslett et al. (1987) or Ketcham et al. (2007). Surely it is logically impossible for both models to be correct. Yet for sufficiently small samples, HeFTy will find plenty of ‘good’ t–T paths in both cases. The truth of the matter is that both the Laslett et al. (1987) and Ketcham et al. (2007) models are incorrect, albeit to different degrees. As sample size increases, the ‘power’ of statistical tests such as K–S and χ² to detect the ‘wrongness’ of the annealing models increases as well (Appendix C). Thus, as we keep adding fission track length measurements to our dataset, HeFTy will find it more and more difficult to find ‘acceptable’ t–T paths. Suppose, for the sake of the argument, that the Laslett et al. (1987) annealing model is ‘more wrong’ than the Ketcham et al. (2007) model. This will manifest itself in the fact that beyond a critical sample size, HeFTy will fail to find even a single ‘acceptable’ model using the Laslett et al. (1987) model, while the Ketcham et al. (2007) model will still yield a small number of ‘non-disprovable’ t–T paths. However, if we further increase the sample size beyond this point, then even the Ketcham et al. (2007) model will eventually fail to yield any ‘acceptable’ solutions.

The problem is that Geology itself imposes unrealistic assumptions on our thermal modelling efforts. Our understanding of diffusion and annealing kinetics is based on short term experiments carried out in completely different environments than the geological processes which we aim to understand. For example, helium diffusion experiments are done under ultra-high vacuum at temperatures of hundreds of degrees over the duration of at most a few weeks. These are very different conditions than those found in the natural environment, where diffusion takes place under hydrostatic pressure at a few tens of degrees over millions of years (Villa, 2006). But even if we disregard this problem, and imagine a utopian scenario in which our annealing and diffusion models are an exact description of reality, the p-value conundrum would persist, because there are dozens of other experimental factors that can go wrong, resulting in dozens of reasons for K–S and χ² to reject the data. Examples are observer bias in fission track length measurements to our dataset, HeFTy fails to yield any ‘acceptable’ solutions.

One apparent solution to this problem is to adjust the p-value cutoffs for ‘good’ and ‘acceptable’ models from their default values of 0.5 and 0.05 to another value, in order to account for differences in sample size. Thus, large datasets would require lower p-values than small ones. The aim of such a procedure would be to objectively accept or reject models based on a sample-independent ‘effect size’ (see Appendix C). Although this sounds easy enough in theory, the implementation details...
are not straightforward. The problem is that HeFTy is very flexible in accepting many different types of data and it is unclear how these can be normalized in a common reference frame. For example, one dataset might include only AFT data, a second AFT and AHe data, while a third might throw some vitrinite reflectance data into the mix as well. Each of these different types of data is evaluated by a different statistical test, and it is unclear how to consistently account for sample size in this situation. On a related note, it is important to discuss the current way in which HeFTy combines the p-values for each of the previously mentioned hypothesis tests. Sample KL29 of Section 3, for example, yields three different p-values: one for the fission track lengths, one for the AFT ages and one for the AHe ages. HeFTy bases the decision whether to reject or accept a t–T path based on the lowest of these three values (Ketcham, 2005). This causes a second level of problems, as the chance of erroneously rejecting a correct null hypothesis (a so-called ‘Type-I error’) increases with the number of simultaneous hypothesis tests. In this case we recommend that the user adjusts the p-value cutoff by dividing it by the number of datasets (i.e., use a cutoff of 0.05/3 = 0.17 for ‘good’ and 0.05/3 = 0.017 for ‘acceptable’ models). This is called the ‘Bonferroni correction’ [e.g., p. 424 of Rice (1995)].

In summary, the very idea to use statistical hypothesis tests to evaluate the model space is problematic. Unfortunately, we cannot use p-values to make a reliable decision to find out whether a model is ‘good’ or ‘acceptable’, independent of sample size. QTQt avoids this problem by ranking the models from ‘bad’ to ‘worse’, and then selecting the ‘most likely’ ones according to the posterior probability (Appendix A). Because the MCMC algorithm employed by QTQt only determines the posterior probability up to a multiplicative constant, it does not care ‘how bad’ the fit to the data is. The advantage of this approach is that it always produces approximately the same number of solutions, regardless of sample size. The disadvantage is that the ability to automatically detect and reject faulty datasets is lost. This may not be a problem, one might think, if sufficient care is taken to ensure that the analytical data are sound and correct. However, that does not exclude the possibility that there are flaws in the forward modelling routines. For example, recall the two fission track annealing models previously mentioned in Section 4. Although the Ketcham et al. (2007) model may be a better representation of reality than the Laslett et al. (1987) model and, therefore, yield more ‘good’ fits in HeFTy, the difference would be invisible to QTQt users. The program will always yield an assemblage of t–T models, regardless of the annealing model used. As a second example, consider the poor age reproducibility that characterizes many U–Th–He datasets and which has long puzzled geochronologists (Fitzgerald et al., 2006). A number of explanations have been proposed to explain this dispersion over the years, ranging from invisible and insoluble actinide-rich mineral inclusions (Vermeesch et al., 2007), α-implantation by ‘bad neighbours’ (Spiegel et al., 2009), fragmentation during mineral separation (Brown et al., 2013) and radiation damage due to α- recoil (Flowers et al., 2009). The latter two hypotheses are linked to precise forward models which can easily be incorporated into inverse modelling software such as HeFTy and QTQt. Some have argued that dispersed data are to be preferred over non-dispersed measurements because they offer more leverage for t–T modelling (Beucher et al., 2013). However, all this assumes that the physical models are correct, which, given the fact that there are so many competing ‘schools of thought’, is unlikely to be true in all situations. Nevertheless, QTQt will take whatever assumption specified by the user and run with it. It is important to note that HeFTy is not immune to these problems either. Because sophisticated physical models such as RDAAM comprise many additional parameters and, hence, ‘degrees of freedom’, the statistical tests used by HeFTy are easily underpowered (Appendix C), yielding many ‘good’ solutions and producing a false sense of confidence in the inverse modelling results.

In conclusion, the evaluation of whether an inverse model is physically sound is more subjective in QTQt than it is in HeFTy. There is no easy way to detect analytical errors or invalid model assumptions other than by subjectively comparing the predicted data with the input measurements. Note that it is possible to ‘fix’ this limitation of QTQt by explicitly evaluating the multiplicative constant given by the denominator in Bayes’ Theorem (Appendix A). We could then set a cutoff value for the posterior probability to define ‘good’ and ‘acceptable’ models, just like in HeFTy. However, this would cause exactly the same problems of sample size dependency as we saw earlier. Conversely, HeFTy could be modified in the spirit of QTQt, by using the p-values to rank models from ‘bad’ to ‘worse’, and then simply plotting the ‘most likely’ ones. The problem with this approach is the sensitivity of HeFTy to the dimensionality of the model space. In order to be able to objectively compare two samples using the proposed ranking algorithm, the parameter space should be devoid of ‘bounding boxes’, and be fixed to a constant search range in time and temperature. This would make HeFTy unreasonably slow, for reasons explained in Section 5.

5. On the selection of time–temperature constraints

As we saw in Section 3.1, HeFTy allows, and generally even requires, the user to constrain the search space by means of ‘bounding boxes’. Often these boxes are chosen to correspond to geological constraints, such as known phases of surface exposure inferred from independently dated unconformities. But even when no formal geological constraints are available, the program often still requires bounding boxes to speed up the modelling. This is a manifestation of the so-called ‘curse of dimensionality’, which is a problem caused by the exponential increase in ‘volume’ associated with adding extra dimensions to a mathematical space. Consider, for example, a unit interval. The average nearest neighbour distance between 10 random samples from this interval will be 0.1. To achieve the same sampling density for a unit square requires not 10 but 100 samples, and for a unit cube 1000 samples. The parameter space explored by HeFTy comprises not two or three but commonly dozens of parameters (i.e., anchor points in time–temperature space), requiring tens of thousands of uniformly distributed random sets to be explored in order to find the tiny subset of statistically plausible models. Furthermore, the ‘sampling density’ of HeFTy’s randomly selected t–T paths also depends on the allowed range of time and temperature. For example, keeping the temperature range equal, it takes twice as long to sample a t–T space spanning 200 Myr than one spanning 100 Myr. Thus, old samples tend to take much longer to model than young ones. The only way for HeFTy to get around this problem is by shrinking the search space. One way to do this is to only permit monotonically rising t–T paths. Another is to use ‘bounding boxes’, like in Section 3.1 and Fig. 7. It is important not to make these boxes too small, especially when they are derived from geological constraints. Otherwise the set of ‘acceptable’ inverse models may simply connect one box to the next, mimicking the geological constraints without adding any new geological insight.

The curse of dimensionality affects QTQt in a different way than HeFTy. As explained in Section 2, QTQt does not explore the multi-dimensional parameter space by means of independent random uniform guesses, but by performing a random walk which explores just a small subset of that space. Thus, an increase in dimensionality does not significantly slow down QTQt. However, this does not mean that QTQt is insensitive to the dimensionality of the search space. The ‘reversible jump MCMC’ algorithm allows the number of parameters to vary from one trial model to the next (Appendix B). To prevent spurious overfitting of the data, this number of parameters is usually quite low. Whereas HeFTy commonly uses ten or more anchor points (i.e., >20 parameters) to define a t–T path, QTQt uses far fewer than that. For example, the maximum likelihood models in Fig. 8 use just three and six t–T anchor points for the datasets comprising 100 and 821 track lengths, respectively. The crudeness of these models is masked by averaging, either through the graphical trick of colour-coding the number of intersecting t–T paths, or by integrating the model assemblages into ‘maximum mode’ and ‘expected’ models (Sambridge et al., 2013).
HeFTy is named after a well known brand of waste disposal bags, as a welcome reminder of the ‘garbage in, garbage out’ principle. QTQt, on the other hand derives its name from the ability of thermal history modelling software to extract colourful and easily interpretable time-temperature histories from complex analytical datasets. In light of the observations made in this paper, it appears that the two programs have been ‘exchanged at birth’, and that their names should have been swapped. First, HeFTy is an arguably easier to use and visually more appealing (‘cute’) piece of software than QTQt. Second, and more importantly, QTQt is more prone to the ‘garbage in, garbage out’ problem than HeFTy. By using p-values, HeFTy contains a built-in quality control mechanism which can protect the user from the worst kinds of ‘garbage’ data. For example, the physically impossible dataset of Section 3.2 was ‘blocked’ by this safety mechanism and yielded no ‘acceptable’ thermal history models in HeFTy. However, in normal to small datasets, the statistical tests used by HeFTy are often underpowered and the ‘garbage in, garbage out’ principle remains a serious concern. Nevertheless, HeFTy is less susceptible to overinterpretation than QTQt, which lacks an ‘objective’ quality control mechanism. It is up to the expertise of the analyst to make a subjective comparison between the input data and the model predictions made by QTQt.

Unfortunately, and this is perhaps the most important conclusion of our paper, HeFTy’s efforts in dealing with the ‘garbage’ data come at a high cost. In its attempt to make an ‘objective’ evaluation of candidate models, HeFTy acquires an undesirable sensitivity to sample size. HeFTy’s power to resolve even the tiniest violations of the model assumptions increases with the amount and the precision of the input data. Thus, as was shown in a regression context (Section 2.2 and Fig. 3) as well as thermochronology (Section 3 and Fig. 7), HeFTy will fail to come up with even a single ‘acceptable’ model if the analytical precision is very high or the sample size is very large. Put in another way, the ability of HeFTy to extract thermal histories from AFT and AHe (Tian et al., 2014), apatite U–Pb (Cochrane et al., 2014) or 4He/3He data (Karlstrom et al., 2014) only exists by virtue of the relative sparsity and low analytical precision of the input data. It is counter-intuitive and unfair that the user should be penalised for acquiring large and precise datasets. In this respect, the MCMC approach taken by QTQt is more sensible, as it does not punish but reward large and precise datasets, in the form of more detailed and tightly constrained thermal histories. Although the inherent subjectivity of QTQt’s approach may be perceived as a negative feature, it merely reflects the fact that thermal history models should always be interpreted in a wider geological context. What is ‘significant’ in one geological setting may not necessarily be so in another, and no computer algorithm can reliably make that call on behalf of the geologist. As George Box famously said, “all models are wrong, but some are useful”.

Acknowledgements

The authors would like to thank James Schwanenthal and Martin Rittner (UCL) for assistance with the U–Th–He measurements, Guangwei Li (Melbourne) for measuring some of sample KL29’s fission track lengths, and Ed Sobel and an anonymous reviewer for feedback on the submitted manuscript. This research was funded by ERC grant #259505 and NERC grant #NE/K003232/1.

Appendix A. Frequentist vs. Bayesian inference

HeFTy uses a ‘Frequentist’ approach to statistics, which means that all inferences about the unknown model [A,B] are based on the known data [x,y] via the likelihood function P(x,y|A,B). In contrast, QTQt follows the ‘Bayesian’ paradigm, in which inferences are based on the so-called ‘posterior probability’ P(A,B|x,y). The two quantities
are related through Bayes’ Rule:

\[ P(A, B | x, y) \propto P(x, y | A, B) \cdot P(A, B) \]  

(5)

where \( P(A, B) \) is the ‘prior probability’ of the model \((A,B)\). If the latter follows a uniform distribution (i.e., \( P(A,B) = \) constant for all \((A,B)\)), then \( P(A, B | x, y) \propto P(x, y | A, B) \) and the posterior is proportional to the likelihood (as in Section 2). Note that the constant of proportionality is not specified, reflecting the fact that the absolute values of the posterior probability are not evaluated. Bayesian credible intervals comprise those models yielding the (typically 95%) highest posterior probabilities, without specifying exactly how high these should be. How this is done in practice is discussed in Appendix B.

Appendix B. A few words about MCMC modelling

Appendix A explained that HeFTy evaluates the likelihood \( P(x,y|A,B) \) whereas QTQt evaluates the posterior \( P(A,B|x,y) \). A more important difference is how this evaluation is done. As explained in Section 2.1, HeFTy considers a large number of independent random models and judges whether or not the data could have been derived from these based on the actual value of \( P(x,y|A,B) \). QTQt, on the other hand, generates a ‘Markov chain’ of serially dependent models in which the jth candidate model is generated by randomly modifying the \((j-1)\)th model, and is accepted or rejected at random with probability \( \alpha \):

\[ \alpha = \min \left( \frac{P(A_j, B_j | x, y) \cdot P(A_j | A_{j-1}, B_{j-1})}{P(A_{j-1}, B_{j-1} | x, y) \cdot P(A_{j-1} | A_{j-2}, B_{j-2})} \right) \]  

(6)

where \( P(A_j, B_j | A_{j-1}, B_{j-1}) \) and \( P(A_{j-1}, B_{j-1} | A_{j-2}, B_{j-2}) \) are the ‘proposal probabilities’ expressing the likelihood of the transition from model state \((j-1)\) to model state \(j\) and vice versa. It can be shown that, after a sufficiently large number of iterations, this routine assembles a representative collection of models from the posterior distribution so that those areas of the parameter space for which \( P(A,B|x,y) \) is high are more densely sampled than those areas where \( P(A,B|x,y) \) is low. The collection of models covering the 95% highest posterior probabilities comprises a 95% ‘credible interval’. For the thermochronological applications of Section 3, QTQt uses a generalised version of Eq. (6) which allows a variable number of model parameters. This is called ‘reversible jump MCMC’ (Green, 1995). For the linear regression problem of Section 2, the proposal probabilities are symmetric so that \( P(A_j, B_j | A_{j-1}, B_{j-1}) = P(A_{j-1}, B_{j-1} | A_j, B_j) \) and the prior probabilities are constant (see Appendix A) so that Eq. (6) reduces to a ratio of likelihoods. The crucial point to note here is that the MCMC algorithm does not use the actual value of the posterior, only relative differences. This is the main reason behind the different behaviours of HeFTy and QTQt exhibited in Sections 2 and 3.

Appendix C. A primer for thermochronologists

Sections 2.2 and 3.1 showed how HeFTy inevitably ‘breaks’ when it is fed with too much data. This is because (a) no physical model of Nature is ever 100% accurate and (b) the power of statistical tests such as Chi-square to resolve even the tiniest violation of the model assumption monotonically increases with sample size. To illustrate the latter point in more detail, consider the linear regression exercise of Section 2.2, which tested a second order polynomial dataset against a linear null hypothesis. Under this null hypothesis, the Chi-square statistic (Eq. (3)) was predicted to follow a Chi-square distribution with \( n-2 \) degrees of freedom. Under this ‘null distribution’, \( \chi^2_{\text{null}} \) is 95% likely to take on a value of \( \sim 15.5 \) for \( n = 10 \) and of \( \sim 122 \) for \( n = 100 \). If the null hypothesis was correct, and we were to accidently observe a value greater than these, then this would have amounted to a so-called ‘Type I’ error. In reality, however, we know that the null hypothesis is false due to the fact that \( c = 0.02 \) or \( 0 \) in Eq. (1). It turns out that in the simple case of linear regression, we can actually predict the expected distribution of \( \chi^2_{\text{null}} \) under this ‘alternative hypothesis’. It can be shown that in this case, the statistic does not follow an ordinary (‘central’) Chi-square distribution, but a ‘non-central’ Chi-square distribution (Cohen, 1977) with \( n-2 \) degrees of freedom and a ‘noncentrality parameter’ \( \lambda \) given by:

\[ \lambda = \sum_{i=1}^{n} \left( a_i + bx_i + c \right)^2 / \sigma^2 \]  

(7)

Using this ‘alternative distribution’, it is easy to show that \( \chi^2_{\text{null}} \) is 50.1% likely to fall below the cutoff value of 15.5 for \( n = 10 \), thus failing to reject the wrong null hypothesis and thereby committing a ‘Type II’ error. By increasing the sample size to \( n = 100 \), the probability (\( \beta \)) of committing a Type II error decreases to a mere 0.32% (Table 1). The ‘power’ of a statistical test is defined as \( 1 - \beta \). It is a universal property of statistical tests that this number increases with sample size. As a second example, the case of the t-test is discussed in Appendix B of Vermeesch (2013). Because Frequentist algorithms such as HeFTy are intimately linked to statistical tests, their power to resolve even the tiniest deviation from linearity, the slightest inaccuracy in our annealing models, or any bias in the \( \alpha \)-ejection correction will eventually result in a failure to find any ‘acceptable’ solution.

Appendix D. Supplementary data

Supplementary data to this article can be found online at http://dx.doi.org/10.1016/j.ejarscirev.2014.09.010.

References


http://dx.doi.org/10.1016/j.gca.2005.01.024.

Table 1

<table>
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<tr>
<th>( n = 10 )</th>
<th>( n = 100 )</th>
<th>( n = 1000 )</th>
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<tbody>
<tr>
<td>95</td>
<td>61.2</td>
<td>0.72</td>
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\( \beta = 0.05 \) for \( n = 10 \) and \( 0.001 \) for \( n = 100 \) and \( 0.02 \) for \( n = 1000 \).


