Broadening of Lyα by neutral helium in DBA white dwarfs

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ABSTRACT

Traces of photospheric hydrogen are detected in at least half of all white dwarfs with helium-dominated atmospheres through the presence of Hα in high-quality optical spectroscopy. Previous studies have noted significant discrepancies between the hydrogen abundances derived from Hα and Lyα for a number of stars where ultraviolet spectroscopy is also available. We demonstrate that this discrepancy is caused by inadequate treatment of the broadening of Lyα by neutral helium. When fitting Hubble Space Telescope far-ultraviolet spectroscopy of 17 DB white dwarfs using our new line profile calculations, we find good agreement between log(N_H/N_H) measured from Lyα and Hα. Larger values of log(N_H/N_H) based on Lyα are still found for three stars, which are among the most distant in our sample, and we show that a small amount of interstellar absorption from neutral hydrogen can account for this discrepancy.

Key words: line: profiles – stars: atmospheres – white dwarfs – ultraviolet: stars.

1 INTRODUCTION

White dwarfs are the progeny of stars with initial masses \( \lesssim 8–10 \, M_\odot \) (e.g. Garcia-Berro, Ritossa & Iben 1997; Smartt et al. 2009; Doherty et al. 2015). The high surface gravities of these stars result in the chemical stratification of their thin non-degenerate envelopes (Schatzman 1948). Most white dwarfs have a sufficiently thick layer of hydrogen to result in optical and ultraviolet spectra dominated by Balmer and Lyman lines, respectively, and are spectroscopically classified as DA stars (Sion et al. 2003). However, a small fraction of white dwarfs have helium-dominated atmospheres, probably a result of a late thermal pulse removing the residual hydrogen (Iben et al. 1983; Althaus et al. 2015). For effective temperatures \( T_{\text{eff}} \gtrsim 10,000 \, \text{K} \), the spectra of these stars contain absorption lines of neutral helium (spectral type DB), at lower temperatures their spectra are featureless continua (DC white dwarfs). The relative fraction of white dwarfs with helium-dominated atmospheres is a function of effective temperature, and hence of cooling age. The changes in this fraction at the hotter end of the cooling sequence are generally understood by mixing of thin hydrogen layers in an increasingly deep convection zone. Variations of the ratio of helium versus hydrogen-dominated atmospheres near the cool end of the cooling sequence are less well understood (Bergeron, Ruiz & Leggett 1997; Chen & Hansen 2012; Giammichele, Bergeron & Dufour 2012).

Liebert et al. (1979) showed that the dichotomy between helium versus hydrogen-dominated atmospheres is not clean, identifying WD 1425+540 (G200-39) as a relatively cool (\( \sim 15,000 \, \text{K} \)) white dwarf exhibiting Balmer and helium absorption lines of roughly equal strength with an estimated \( \log(N_\text{H}/N_{\text{He}}) \sim -3.6 \). Shortly later, Liebert et al. (1984) discovered a hotter (\( \sim 30,000 \, \text{K} \)) example of a white dwarf with both Balmer and helium lines, showing that trace hydrogen in helium-atmosphere white dwarfs can be present across a wide range of effective temperatures. While these stars are classified as DAB or DBA white dwarfs depending on the relative strength of the Balmer and helium lines, it is important to note that the spectra of some helium atmosphere white dwarfs, in particular cooler ones, are dominated by Balmer lines, e.g. GD 362 (Ginninas, Dufour & Bergeron 2004; Kawka & Vennes 2005), GD 16 (Koester et al. 2005), and GD 17 (Gentile Fusillo et al. 2017). At higher temperatures, trace helium in hydrogen-dominated atmospheres becomes detectable via the presence of He H lines (Manseau, Bergeron & Green 2016). Whereas a small number of these mixed atmosphere stars turned out to be unresolved DA plus DB binaries (e.g. WD 1115+166, Bergeron & Liebert 2002), surveys with better instrumentation now demonstrate that a large fraction of helium atmosphere white dwarfs contain varying amounts of trace hydrogen (Voss et al. 2007; Bergeron et al. 2011; Koester & Kepler 2015; Rolland, Bergeron & Fontaine 2018).

The origin of this trace hydrogen has been explored extensively (Beauchamp et al. 1996; Voss et al. 2007; Bergeron et al. 2011; Koester & Kepler 2015; Rolland et al. 2018), however, the two favoured explanations, a residual thin hydrogen layer or accretion from the interstellar medium face problems explaining the apparent change in \( \log(N_\text{H}/N_{\text{He}}) \) as a function of cooling age. Episodic accretion of water-bearing planetesimals has recently been proposed as an alternative source of hydrogen (Farihi et al. 2010; Farihi, Gänsicke & Koester 2013; Raddi et al. 2015; Gentile Fusillo et al. 2017), and
it is remarkable to note that Xu et al. (2017) showed that the prototypical DBA WD 1425+540 is accreting volatile-rich planetary material, with a composition similar to solar system comets.

In the vast majority of DBA and DAB stars, the hydrogen content has been derived from the analysis of optical spectroscopy, in particular the strength of the H α line (Voss et al. 2007; Bergeron et al. 2011; Koester & Kepler 2015; Rolland et al. 2018). Ultraviolet spectroscopy is available for a number of DB stars, allowing an independent assessment of log (N(H)/N(He)) from the Ly α absorption line, and Jura & Xu (2012) noticed significant discrepancies between the values of log (N(H)/N(He)) derived from H α and Ly α. The most extreme published example of differing log (N(H)/N(He)) derived from optical and ultraviolet spectroscopy is WD 1425+540 (Genest-Beaulieu & Bergeron 2017). However, so far no systematic comparison between the values derived from H α and Ly α has yet been attempted.

Here we present a homogeneous analysis of Hubble Space Telescope (HST) ultraviolet spectroscopy of 17 DB white dwarfs, making use of new calculations of the Ly α profile broadened by neutral helium.

### 2 OBSERVATIONS

The far-ultraviolet spectroscopy was obtained with the Cosmic Origin Spectrograph (COS) as part of three separate HST programs (#12474, #13453, and #14597) between 2011 October and 2017 October (Table 1). All observations were carried out using the G130M grating with a central wavelength of 1291 Å, except for WD 0125–236, which was observed with a central wavelength of 1327 Å. The 1291 Å observations covered the wavelength range 1130–1435 Å, with a gap at 1278–1288 Å due to the space between the two detector segments. For the 1327 Å observations, the wavelengths covered were 1170–1470 Å, with a gap 1318–1328 Å. We used all four FP-POS positions to mitigate the fixed pattern noise that is affecting the COS far-ultraviolet detector for the observations obtained in programs #13453 and #14597. The observations obtained in program #12474 were part of a snapshot survey of white dwarfs, and the short exposure times allowed only two FP-POS positions to be used (Gänsicke et al. 2012). To mitigate airglow emission, we replaced the region around O1 1302.04 Å with the spectrum extracted from the COS data obtained on the night side of the Earth.

The COS spectra of all 17 stars contain a broad Ly α line, indicating the presence of trace amounts of photospheric hydrogen. The strength of the Ly α line depends on the abundance of hydrogen, log (N(H)/N(He)), as well as on the effective temperature T eff, with the line weakening with increasing T eff for constant values of log (N(H)/N(He)). In the case of moderately weak Ly α absorption, the continuum can be smoothly interpolated across the Ly α line (WD 1107+265, top panel in Fig. 2). With increasing strength of the Ly α profile, the blue wing becomes increasingly suppressed, and a broad dip centred near 1150 Å becomes noticeable, with the most extreme example being WD 1425+540 (bottom panel in Fig. 2). With the exception of WD 0840+262, WD 1557+192, and WD 1940+374, the DBs in this sample also display absorption from photospheric metals, which will be discussed in a separate paper (Farhi et al., in preparation).

In addition to the HST observations, we obtained optical intermediate resolution spectroscopy of WD 0110–565 and WD 1349–230 using X-Shooter on the Very Large Telescope, which was used to determine the effective temperatures, surface gravities, as well as log (N(H)/N(He)) from the H α line.

### 3 SPECTROSCOPIC ANALYSIS

We modelled the HST/COS observations using theoretical spectra calculated with the physics and algorithms that are described in Koester (2010), and which were recently used for the analysis of DB white dwarfs in Koester & Kepler (2015). As the focus of this study is a comparison of the hydrogen abundances measured from H α and Ly α, we decided to adopt the atmospheric parameters from Rolland et al. (2018) who presented a homogeneous study of a large sample of DB white dwarfs. Two exceptions to this were WD 0110–565 and WD 1349–230, where T eff, log g, and log (N(H)/N(He)) based on H α were determined from our X-Shooter spectroscopy. We compared model spectra computed for these parameters (Table 2) with the PanSTARRS photometry available for all stars except WD 0110–565, and found good agreement when allowing for small amounts of reddening, E(B–V) < 0.02.

Using the values of log (N(H)/N(He)) listed in Table 2, we computed synthetic Ly α profiles using a broadening theory for Ly α which considers Stark broadening by electrons, H+, and He+ ions (Tremblay & Bergeron 2009; Tremblay, private communication), convolved with van der Waals broadening by neutral helium and hydrogen atoms. These predicted Ly α lines are significantly narrower than the observed ones for all objects (Fig. 3 and A1–A3, right-hand panels) except in WD 0840+262, WD 1940+374, and WD 2144–079. An attempt to match the width of the observed Ly α profiles requires hydrogen abundances up to a factor 10 larger than those derived from H α (Fig. 3 and A1–A3, middle panels), and still fails to reproduce the shape of the blue wing of Ly α. This problem was already noted in WD 1425+540 by Xu et al. (2017) and studied in detail by Genest-Beaulieu & Bergeron (2017), who suggested that an inhomogeneous distribution of hydrogen in the atmosphere, with log (N(H)/N(He)) increasing in the outer layers, might explain the strong observed Ly α profile. Here we argue that the most likely origin of this discrepancy is that an inadequate broadening theory of Ly α by neutral helium is used in the current model atmosphere code.

We had already developed a unified profile calculation (Koester & Wolff 2000), but that study was concerned with the analysis of cool (T eff < 9000 K) helium-rich white dwarfs with practically no flux below 1500 Å. We therefore used a simplified theory for only one temperature and one perturber density and scaled this profile.
Christova (2009), who predicted a small satellite feature in the blue wing of Ly\(\alpha\) near 1150 Å, coinciding with the depression detected near 1150 Å, coinciding with the depression detected

Improved line profile calculations were performed by Allard & Christova (2009), who predicted a small satellite feature in the blue wing of Ly\(\alpha\) near 1150 Å, coinciding with the depression detected near 1150 Å, coinciding with the depression detected

Table 2. Atmospheric parameters of the 17 DB white dwarfs. Effective temperatures, surface gravities, and hydrogen abundances measured from H\(\alpha\) are taken from Rolland et al. (2018; R18) or derived from our X-Shooter spectra (XS). Hydrogen abundances were measured from the Ly\(\alpha\) absorption line in the COS spectra using Ly\(\alpha\) line profiles that include Stark broadening by electrons, H\(^+\), and He\(^+\) and van der Waals broadening by neutral helium and hydrogen (old) and using our improved unified broadening calculations (new). Stars with photospheric metals are flagged as \(z\) and interstellar absorption is likely to contribute to the observed Ly\(\alpha\) profile of WD 1349–230, WD 1557+192, and WD 2354+159, flagged as \(i\).

<table>
<thead>
<tr>
<th>WD</th>
<th>(T_{\text{eff}}) (K)</th>
<th>log (g) c.g.s.</th>
<th>H(\alpha)</th>
<th>log ((N_\text{H}/N_{\text{He}}))</th>
<th>Ly(\alpha) new</th>
<th>Ly(\alpha) old</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>0100–068(^{2})</td>
<td>19820 ± 531</td>
<td>8.06 ± 0.04</td>
<td>−5.14 ± 1.06</td>
<td>−4.80 ± 0.15</td>
<td>−4.80</td>
<td>R18</td>
<td></td>
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<tr>
<td>0110–565(^{3})</td>
<td>19124 ± 16</td>
<td>8.17 ± 0.01</td>
<td>−4.20 ± 0.20</td>
<td>−4.10 ± 0.10</td>
<td>−4.10</td>
<td>XS</td>
<td></td>
</tr>
<tr>
<td>0125–236(^{3})</td>
<td>16550 ± 436</td>
<td>8.24 ± 0.07</td>
<td>−5.21 ± 0.32</td>
<td>−5.00 ± 0.15</td>
<td>−4.80</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>0435+410(^{3})</td>
<td>16790 ± 408</td>
<td>8.18 ± 0.08</td>
<td>−4.21 ± 0.07</td>
<td>−4.10 ± 0.15</td>
<td>−4.00</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>0437+138(^{3})</td>
<td>15120 ± 361</td>
<td>8.25 ± 0.07</td>
<td>−4.68 ± 0.06</td>
<td>−4.60 ± 0.10</td>
<td>−4.10</td>
<td>R18</td>
<td></td>
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<tr>
<td>0840+262</td>
<td>17700 ± 863</td>
<td>8.28 ± 0.05</td>
<td>−4.18 ± 0.06</td>
<td>−4.15 ± 0.15</td>
<td>−3.95</td>
<td>R18</td>
<td></td>
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<tr>
<td>1107+265(^{3})</td>
<td>15130 ± 357</td>
<td>8.11 ± 0.06</td>
<td>−5.77 ± 0.46</td>
<td>−5.40 ± 0.15</td>
<td>−5.00</td>
<td>R18</td>
<td></td>
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<tr>
<td>1349–230(^{3})</td>
<td>17905 ± 124</td>
<td>8.05 ± 0.01</td>
<td>−4.90 ± 0.20</td>
<td>−3.90 ± 0.20</td>
<td>−3.90</td>
<td>XS</td>
<td></td>
</tr>
<tr>
<td>1352+004(^{3})</td>
<td>13980 ± 340</td>
<td>8.05 ± 0.09</td>
<td>−5.31 ± 0.17</td>
<td>−5.10 ± 0.15</td>
<td>−4.70</td>
<td>R18</td>
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<tr>
<td>1425+540(^{3})</td>
<td>14410 ± 341</td>
<td>7.89 ± 0.06</td>
<td>−4.26 ± 0.03</td>
<td>−4.00 ± 0.20</td>
<td>−3.70</td>
<td>R18</td>
<td></td>
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<tr>
<td>1557+192(^{2})</td>
<td>19510 ± 546</td>
<td>8.15 ± 0.05</td>
<td>−4.30 ± 0.26</td>
<td>−3.40 ± 0.30</td>
<td>−3.20</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>1644+198(^{3})</td>
<td>15210 ± 360</td>
<td>8.14 ± 0.06</td>
<td>−5.68 ± 0.39</td>
<td>−5.20 ± 0.10</td>
<td>−4.90</td>
<td>R18</td>
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<td>1822+410(^{3})</td>
<td>16230 ± 383</td>
<td>8.00 ± 0.06</td>
<td>−4.45 ± 0.06</td>
<td>−4.40 ± 0.20</td>
<td>−4.20</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>1940+374</td>
<td>16850 ± 406</td>
<td>8.07 ± 0.09</td>
<td>−5.97 ± 1.50</td>
<td>−5.60 ± 0.15</td>
<td>−5.10</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>2144+079(^{3})</td>
<td>16340 ± 408</td>
<td>8.18 ± 0.05</td>
<td>−6.22 ± 0.35</td>
<td>−6.35 ± 0.20</td>
<td>−6.20</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>2229+130(^{3})</td>
<td>14870 ± 352</td>
<td>8.15 ± 0.06</td>
<td>−4.91 ± 0.08</td>
<td>−4.70 ± 0.15</td>
<td>−4.30</td>
<td>R18</td>
<td></td>
</tr>
<tr>
<td>2354+159(^{3})</td>
<td>24830 ± 1670</td>
<td>8.15 ± 0.04</td>
<td>−4.59 ± 0.36</td>
<td>−3.60 ± 0.20</td>
<td>−3.50</td>
<td>R18</td>
<td></td>
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Figure 1. Ly\(\alpha\) line profiles computed for a range of effective temperatures and values of log \((N_\text{H}/N_{\text{He}})\), using van der Waals broadening by neutral hydrogen and helium (right, red lines) and using our improved unified line broadening theory (left, blue lines). All models were computed with log \(g\) = 8 fixed, and the normalised spectra were offset vertically by multiples of 0.6 units.

with the neutral helium density. This treatment is not appropriate to describe the whole Ly\(\alpha\) profile, including the blue wing visible in the spectra of the present sample of hotter DB white dwarfs. Improved line profile calculations were performed by Allard & Christova (2009), who predicted a small satellite feature in the blue wing of Ly\(\alpha\) near 1150 Å, coinciding with the depression detected

Figure 2. Ly\(\alpha\) absorption line profiles based on three different atmosphere structures, using the new broadening theory with (blue) and without (green) metals in the calculation of the equation of state, and using the old line broadening including metals in the equation of state (red). Synthetic spectra were computed from these three atmosphere structures without metals and using the new line broadening to facilitate the comparison of the resulting Ly\(\alpha\) profiles. Whereas the photospheric metals do not affect the structure of the atmosphere, the stronger Ly\(\alpha\) absorption resulting from the improved line broadening results in a somewhat stronger line blanketing, which will slightly affect the derived effective temperature, if included in a detailed fit. For the purpose of this paper, we adopt \(T_{\text{eff}}\) and log \(g\) from Table 2.
Figure 3. HST/COS spectra of three DB white dwarfs (grey) with the strength of Ly$\alpha$ increasing from top to bottom. The region around O I 1302.04 Å is noisier as we only used the night-side data in this wavelength range. Right-hand panels: model spectra (red) computed with log ($N_{\text{H}}/N_{\text{He}}$) based on a fit to the H$\alpha$ line, and using Ly$\alpha$ line profiles that include Stark broadening by electrons, H$^+$, and He$^+$ and van der Waals broadening by neutral helium and hydrogen. Middle panels: using the same broadening theory, we varied log ($N_{\text{H}}/N_{\text{He}}$) to reproduce the red wing of Ly$\alpha$. The models (red) reproduce well the strength of the observed Ly$\alpha$ line for WD 1107+265, but clearly fail to fit the observations of WD 0437+138 and WD 1425+540. Left-hand panels: fits to the COS spectra using models (blue) including our updated treatment of Ly$\alpha$ broadening by neutral helium result in much better agreement with the observations, and reproduce the depression near 1150 Å, a satellite feature of Ly$\alpha$. The values of log ($N_{\text{H}}/N_{\text{He}}$) used for the model calculations are given in each panel.

in the COS spectra of the white dwarfs with the strongest Ly$\alpha$ absorption. Since these authors did not provide numerical tables, we have repeated the profile calculation based on the unified theory by Allard and co-workers (Allard & Kielkopf 1982; Allard & Koester 1992; Allard et al. 1999), but using our own complete new implementation described in Hollands et al. (2017). Adiabatic potential curves and dipole moments were obtained from Theodorakopoulos et al. (1984, 1987) and Belyaev (2015). We calculated tables covering the whole temperature, density, and wavelength ranges where this mechanism is important. Fig. 1 compares these new line profiles to our old treatment of Ly$\alpha$ broadening described above for a range of effective temperatures and hydrogen abundances, and illustrates the increasing importance of broadening by neutral helium both for decreasing $T_{\text{eff}}$ and increasing log ($N_{\text{H}}/N_{\text{He}}$). Given the strength of this additional broadening of the blue wing of Ly$\alpha$, and the fact that most of the DBs also have photospheric metals, we evaluated their influence on the atmospheric structure. Using WD 1425+540 as test case, we computed the structure of the atmosphere, adopting log ($N_{\text{H}}/N_{\text{He}}$) = −4.0, (1) using the new Ly$\alpha$ line profile and including the best-fitting metal abundances in the equation of state, (2) using the old Ly$\alpha$ line profile and including the best-fitting metal abundances in the equation of state, and (3) using the new Ly$\alpha$ line profile but no metals. We then computed synthetic spectra without metals using the new line broadening profiles, and compared the resulting Ly$\alpha$ absorption lines (Fig. 2) and illustrate that the inclusion of metals has no noticeable effect on the atmospheric structure, but that the stronger Ly$\alpha$ line resulting from the new broadening theory causes a small amount of additional line blanketing. We conclude that the atmospheric parameters of the white dwarfs in our sample can potentially be improved making use of the improved Ly$\alpha$ profiles, though including the Gaia Data Release 2 parallaxes (Gaia Collaboration 2018) in the fit may have a stronger effect. However,
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4 CONCLUSIONS

We have developed new line profile calculations for Ly $\alpha$ that account for broadening by neutral helium. When fitting HST/COS spectroscopy of 17 DB white dwarfs using these new line profiles, we find good agreement between log ($N_{H}/N_{H_{\alpha}}$) measured from Ly $\alpha$ and H $\alpha$, resolving discrepancies that were present in analyses that did not properly account for the additional line broadening. Three stars still show a larger value of log ($N_{H}/N_{H_{\alpha}}$) measured from Ly $\alpha$ compared to that derived from H $\alpha$, and we argue that a small contribution of interstellar absorption from neutral hydrogen is a plausible cause for this discrepancy.

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Figure A1. Same as Fig. 3 for the full DB sample analysed in this paper.
Figure A2. Same as Fig. 3 for the full DB sample analysed in this paper.
Figure A3. Same as Fig. 3 for the full DB sample analysed in this paper.

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